



Neutrino factory optimization for non-standard interactions

Toshihiko Ota

J. Kopp, TO, and W. Winter
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B. Gavela, D. Hernandez, TO, and W. Winter
Phys. Rev. **D79** (2009) 013007

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Preface

Within the current precision — Leading Order (LO)

Oscillation probabilities for $\nu_\mu \rightarrow \nu_\alpha$ (@atmospheric region $\Delta m_{31}^2 L/E \sim 1$)

$$\underbrace{P_{\nu_\mu \rightarrow \nu_e}}_0 + P_{\nu_\mu \rightarrow \nu_\mu} + \underbrace{P_{\nu_\mu \rightarrow \nu_\tau}}_{1 - P_{\nu_\mu \rightarrow \nu_\mu}} = 1 \quad (\text{unitarity})$$



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Future experiments are sensitive to the Next LO

$$P_{\nu_\mu \rightarrow \nu_e} = 0 \quad \boxed{\text{Leading Order}}$$

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$$P_{\nu_\mu \rightarrow \nu_e} = 0$$

Leading Order

$$+ \mathcal{O}(s_{13}^2)$$

Mass-Texture, LFV Prediction...

$$+ \mathcal{O}(s_{13} \Delta m_{21}^2 / \Delta m_{31}^2)$$

CP violation (Leptogenesis)...



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- | | |
|---|--|
| $P_{\nu_\mu \rightarrow \nu_e} = 0$ | Leading Order |
| $+ \mathcal{O}(s_{13}^2)$ | Mass-Texture, LFV Prediction... |
| $+ \mathcal{O}(s_{13} \Delta m_{21}^2 / \Delta m_{31}^2)$ | CP violation (Leptogenesis)... |
| $+$ | Direct evidence of New Physics |



Outline

- 1 Introduction
- 2 NSI search in a neutrino factory
 - Neutrino factory for standard oscillation parameters
 - Neutrino factory for NSI
 - Correlations between NSIs
 - Silver detector for NSI
 - Optimization of Golden detector baselines
- 3 Model building for large NSI [short comment]
- 4 Summary



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Introduction — NSI in oscillation

- NSI — Exotic interactions with neutrinos which are parametrized as four-Fermi interactions:

Standard oscillation

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \langle \nu_\beta | e^{-i H L} | \nu_\alpha \rangle \right|^2$$



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$$P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \langle \nu_\beta | e^{-iHL} | \nu_\alpha \rangle \right|^2$$

With NSI in source and detection

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \langle \nu_\beta^d | e^{-iHL} | \nu_\alpha^s \rangle \right|^2$$

- CC type NSI — flavour mixture states at source and detection
Grossman PLB359 (1995) 141.

$$|\nu_\alpha^s\rangle = |\nu_\alpha\rangle + \sum_{\gamma=e,\mu,\tau} \epsilon_{\alpha\gamma}^s |\nu_\gamma\rangle, \quad \text{e.g., } \pi^+ \xrightarrow{\epsilon_{\mu e}^s} \mu^+ \nu_e$$

$$\langle \nu_\alpha^d | = \langle \nu_\alpha | + \sum_{\gamma=e,\mu,\tau} \epsilon_{\gamma\alpha}^d \langle \nu_\gamma |, \quad \text{e.g., } \nu_\tau N \xrightarrow{\epsilon_{\tau e}^d} e^- X$$



Introduction — NSI in oscillation

- NSI — Exotic interactions with neutrinos which are parametrized as four-Fermi interactions:

Standard oscillation

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \langle \nu_\beta | e^{-iH L} | \nu_\alpha \rangle \right|^2$$

With NSI in propagation

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \langle \nu_\beta | e^{-i(H + V_{\text{NSI}})L} | \nu_\alpha \rangle \right|^2$$

- NC type NSI — extra matter effect in propagation

e.g., Wolfenstein PRD17 (1978) 2369. Valle PLB199 (1987) 432. Guzzo Masiero Petcov PLB260 (1991) 154.
Roulet PRD44 (1991) R935.

$$(V_{\text{NSI}})_{\beta\alpha} = \sqrt{2} G_F N_e \begin{pmatrix} \epsilon_{ee}^m & \epsilon_{e\mu}^m & \epsilon_{e\tau}^m \\ \epsilon_{e\mu}^{m*} & \epsilon_{\mu\mu}^m & \epsilon_{\mu\tau}^m \\ \epsilon_{e\tau}^{m*} & \epsilon_{\mu\tau}^{m*} & \epsilon_{\tau\tau}^m \end{pmatrix}, \quad \text{e.g., } \nu_e \xrightarrow{\epsilon_{e\tau}^m} \nu_\tau \text{ in propagation}$$

- We will focus on NSI in the propagation in this talk.

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2 NSI search in a neutrino factory

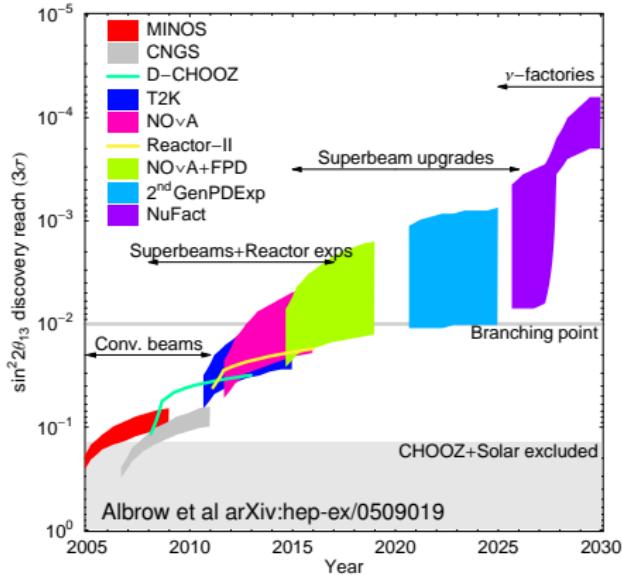
- Neutrino factory for standard oscillation parameters
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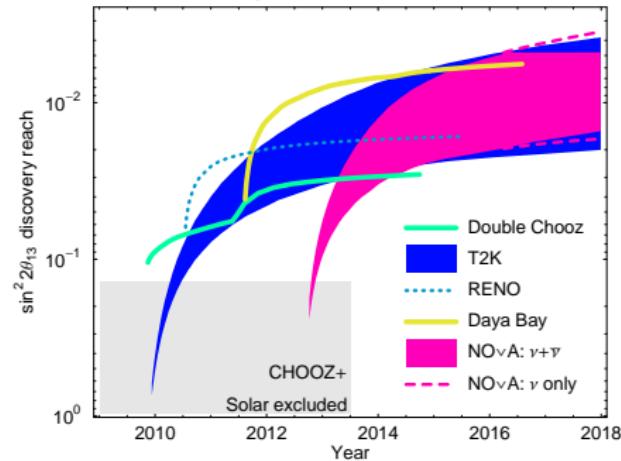


Oscillation experiments in next generation



Huber Lindner Schwetz Winter arXiv:0907.1896

$\sin^2 2\theta_{13}$ discovery (NH, 90% CL)

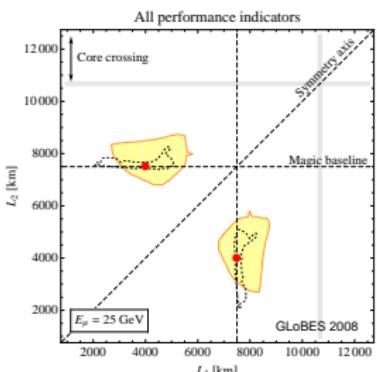
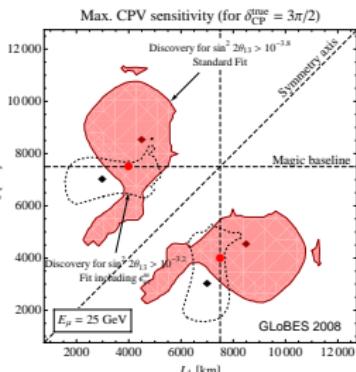
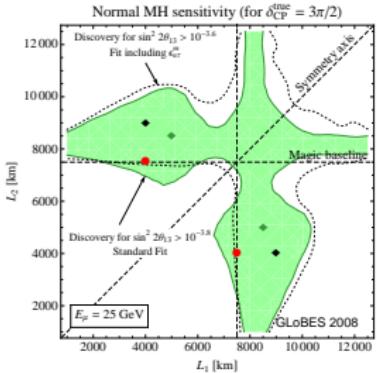
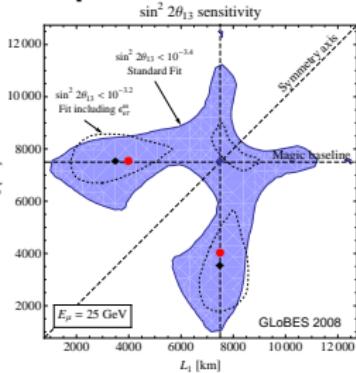


Central theme of this talk

What is the optimal setup (E_μ , L , and type of detectors) for measuring θ_{13} , δ_{CP} , $\text{sign}[\Delta m_{31}^2]$, and NSI parameters $\epsilon_{\alpha\beta}^m$?

Kopp O Winter

Optimization for standard oscillation parameters



For θ_{13} , δ_{CP} , and $\text{sign}[\Delta m_{31}^2]$

Optimum at 4000+7500 km
Optimization does not change in presence of NSI

- The qualitatively different observations (matter osc. max and magic baseline) help resolve the parameter degeneracies.
- If we have two Golden dets at 4000km+7500km, Silver det (included in IDS-NF) does not contribute improving sensitivities.



Relevant NSI in a neutrino factory experiment

Relevant NSI in each channel

e.g., Kikuchi Minakata Uchinami JHEP0903 (2009) 114, Kopp Lindner O Sato PRD77 (2008) 013007

- Appearance channel: $\epsilon_{e\mu}^m$ and $\epsilon_{e\tau}^m$

On bounds on ϵ : Biggio Blennow Fernandez-Martinez JHEP0903 (2009) 139 and arXiv:0907.0097



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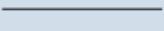
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- Silver channel: $\epsilon_{e\tau}^m$ and $\epsilon_{e\mu}^m$  ECC (Silver) det

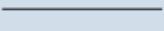
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Tau-associated NSI $\epsilon_{e\tau}^m$ $\epsilon_{\mu\tau}^m$, and $\epsilon_{\tau\tau}^m$

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Central questions for NSI search

- Correlation between NSIs and its resolution



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- Correlation between NSIs and its resolution
- Necessity of Silver channel



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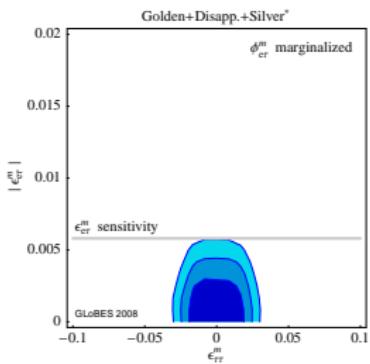
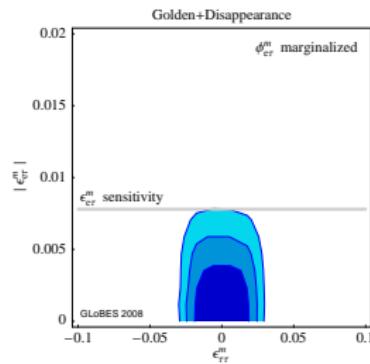
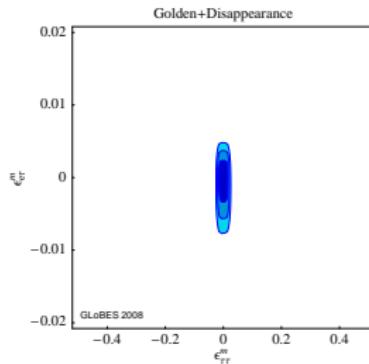
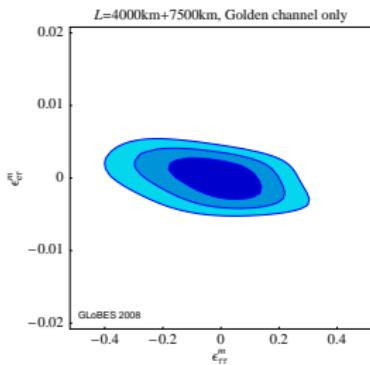
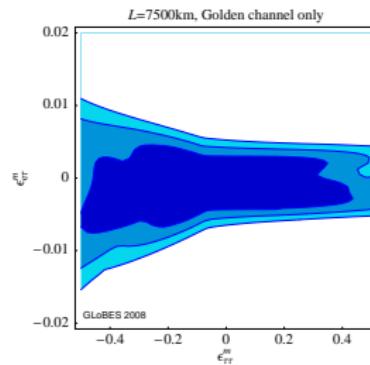
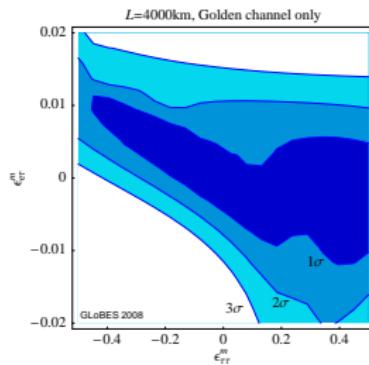
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We address them with full simulations powered by GLoBES

GLoBES Website: <http://www.mpi-hd.mpg.de/lin/globes/>

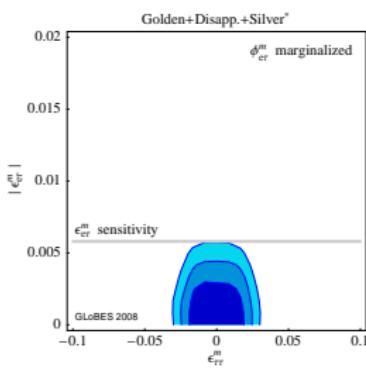
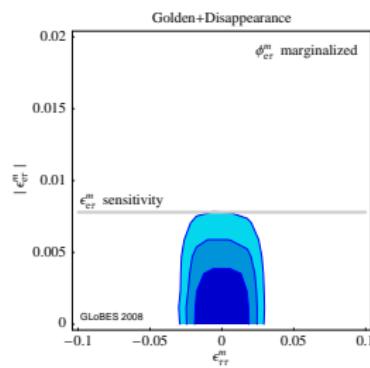
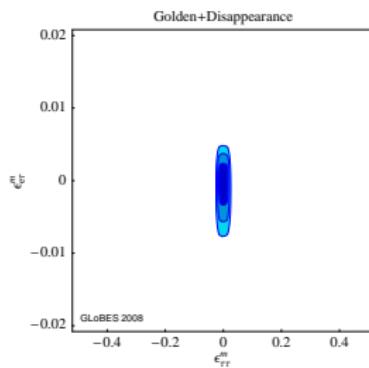
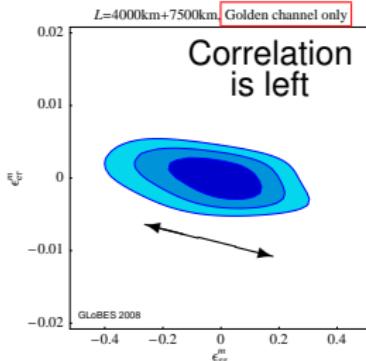
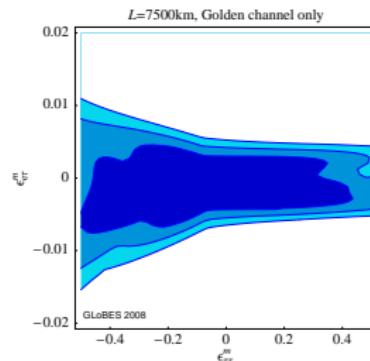
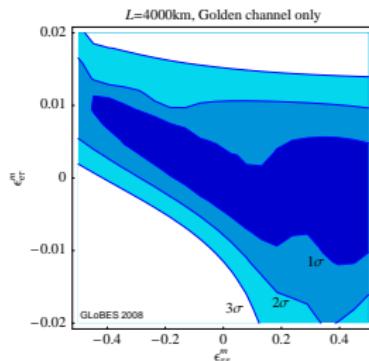
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Correlation between $\epsilon_{e\tau}^m$ and $\epsilon_{\tau\tau}^m$



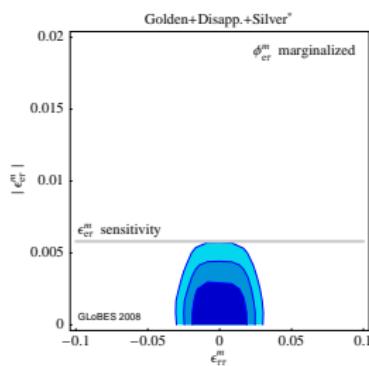
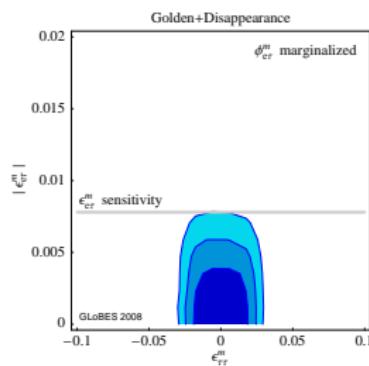
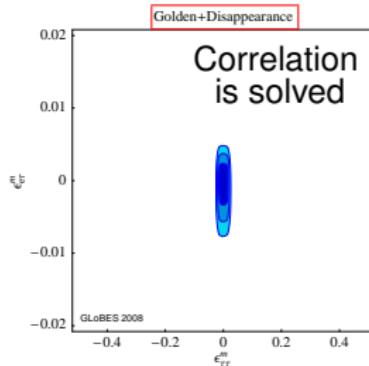
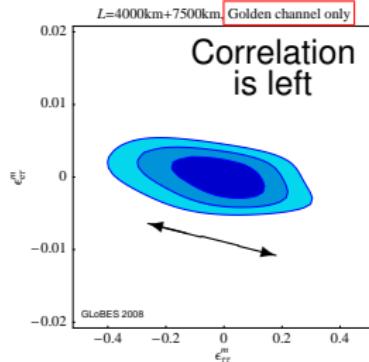
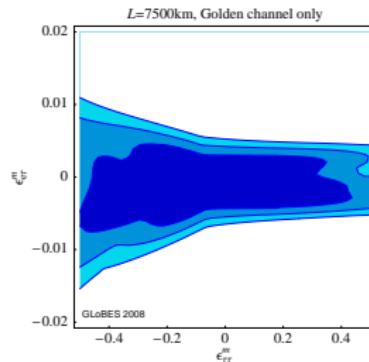
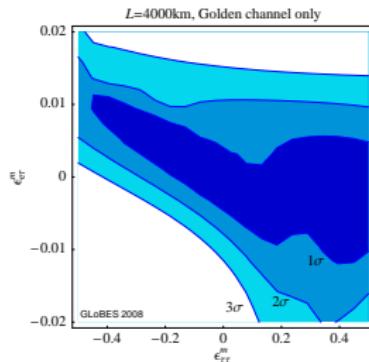
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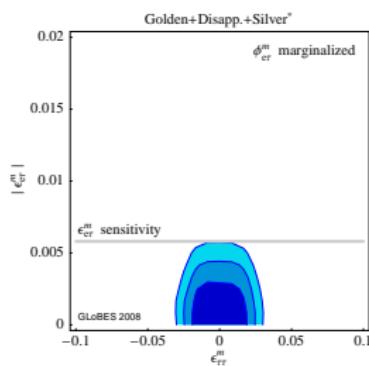
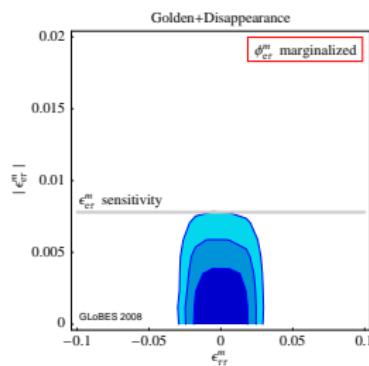
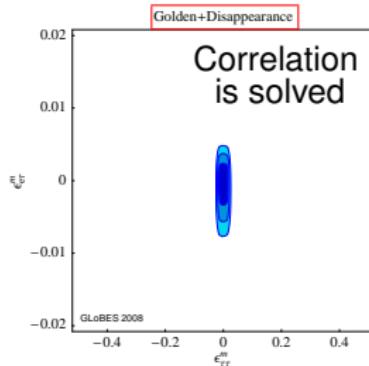
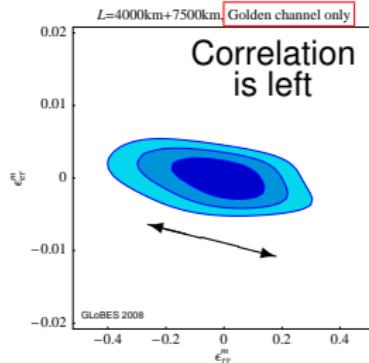
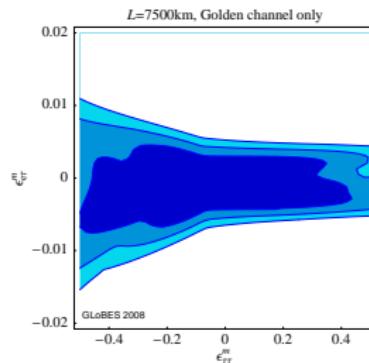
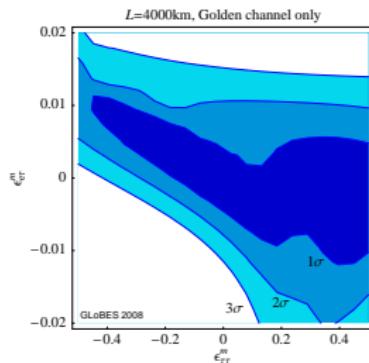
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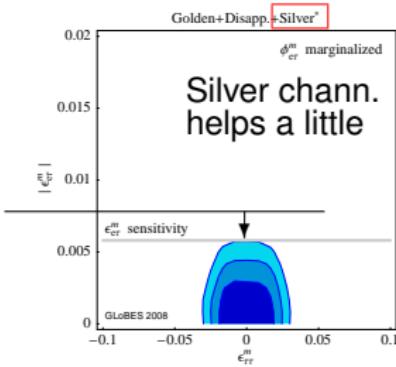
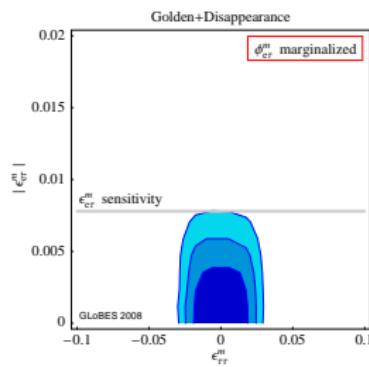
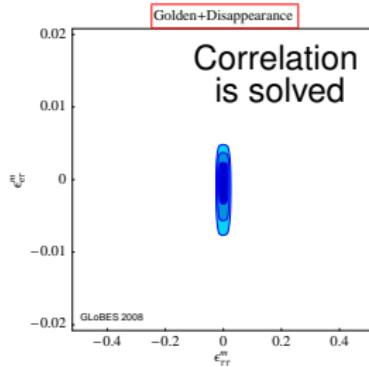
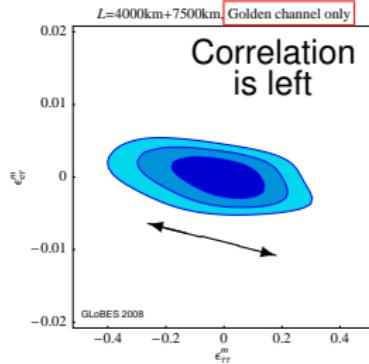
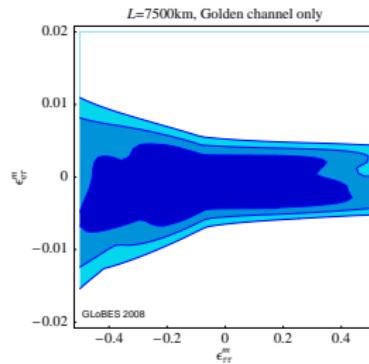
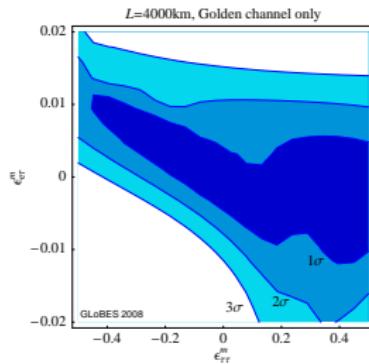
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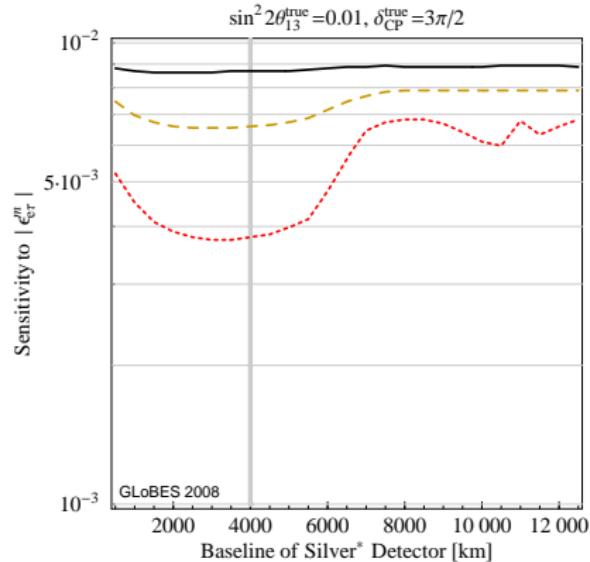
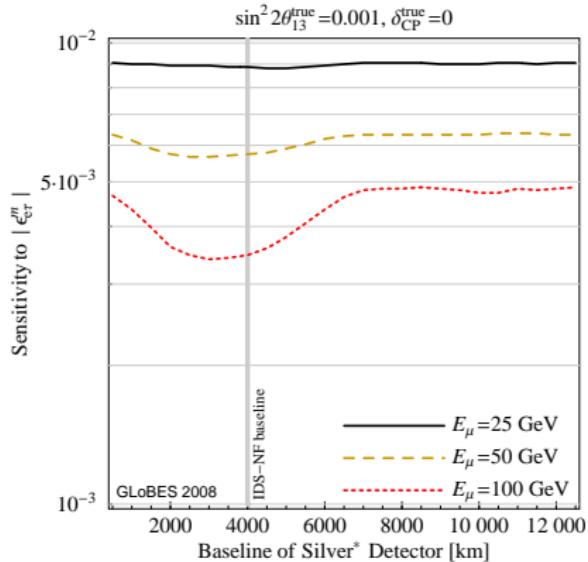
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Kopp O Winter

Optimization of silver detector baseline

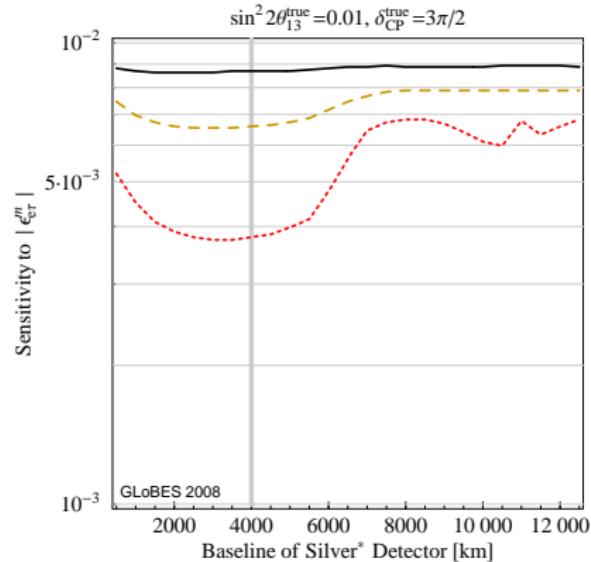
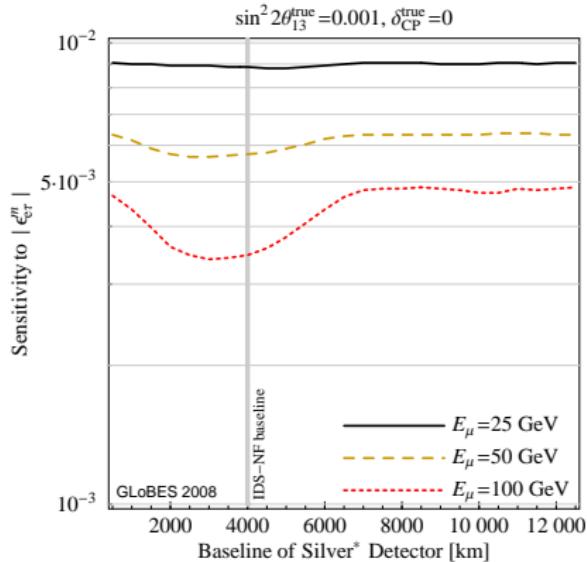


- Fix two Golden dets at $L = 4000 + 7500 \text{ km}$.



Kopp O Winter

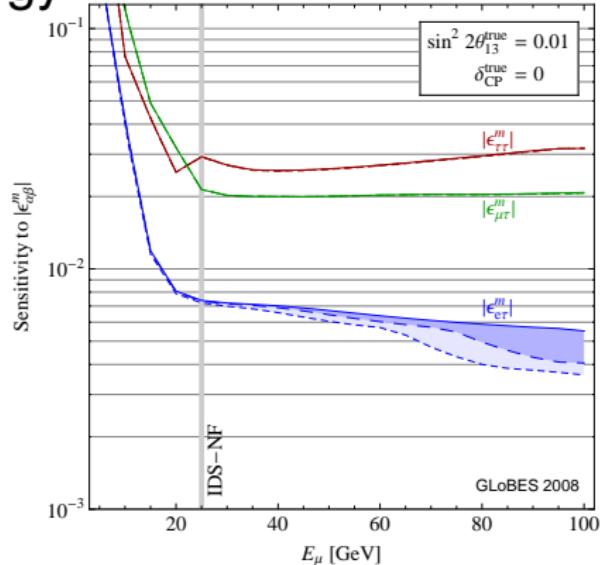
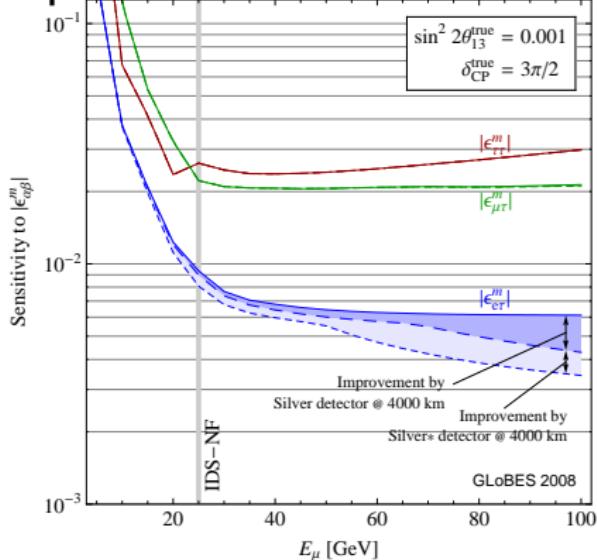
Optimization of silver detector baseline



- Fix two Golden dets at $L = 4000 + 7500 \text{ km}$.
- Silver detector only relevant $L \sim 4000 \text{ km}$ and $E_\mu \gg 25 \text{ GeV}$.
→ Fix Silver det at $L = 4000 \text{ km}$...

Kopp O Winter

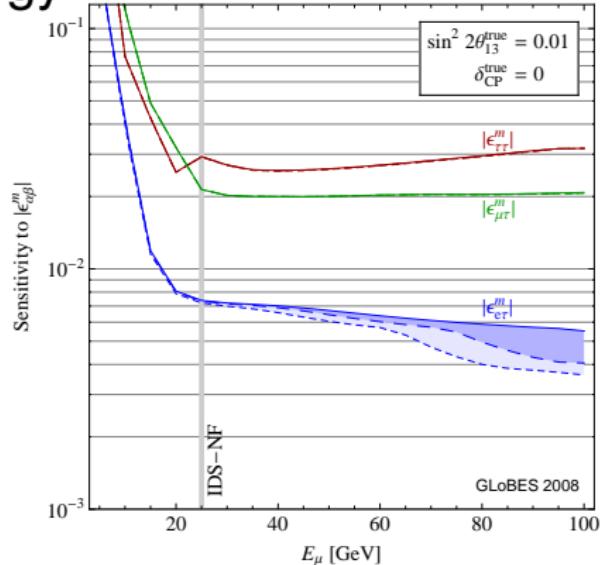
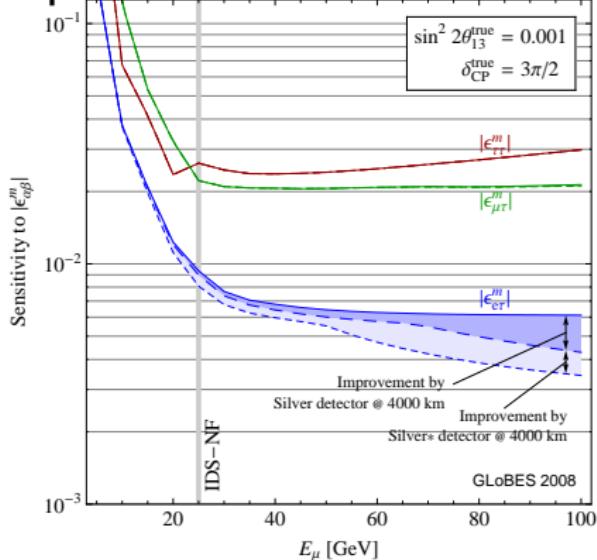
Optimization of muon energy



- Fix Goldens at $L = 4000 + 7500 \text{ km}$ and Silver at $L = 4000 \text{ km}$

Kopp O Winter

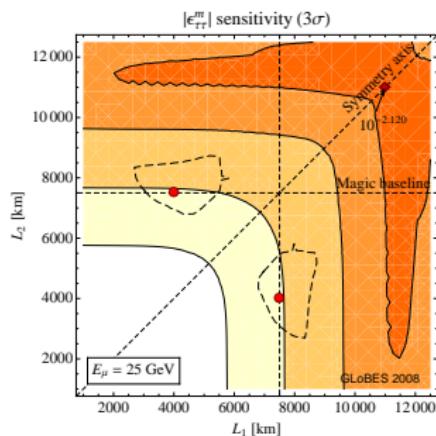
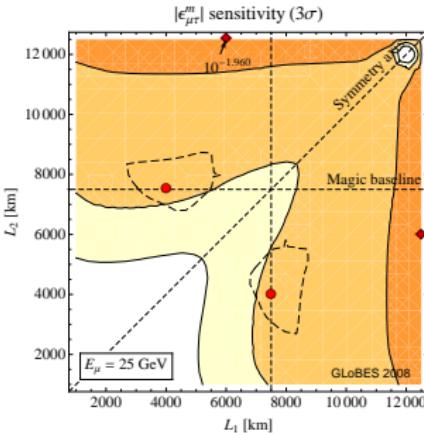
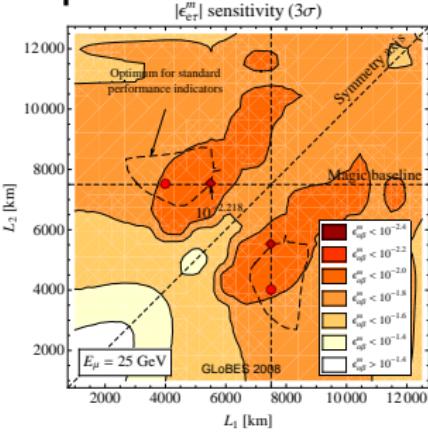
Optimization of muon energy



- Fix Goldens at $L = 4000 + 7500 \text{ km}$ and Silver at $L = 4000 \text{ km}$
- Silver detector only useful at $E_\mu \gg 50 \text{ GeV}$.
→ Fix $E_\mu = 25 \text{ GeV}$ (IDS-NF baseline), omit Silver detector...

Kopp O Winter

Optimization of Golden detector baselines



- For $\epsilon_{e\tau}^m$: $L = 4000\text{km} + 7500\text{km}$ is almost optimal.
- For $\epsilon_{\mu\tau}^m$ and $\epsilon_{\tau\tau}^m$: longer baseline is preferred.
Sensitivity is dominated by the longer baseline, and it is simply proportional to the baseline length.

Note For NSI search, higher E_μ (and longer L) is preferred.

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NSI from Dim.6 and Dim.8

Dim.6 op — 4-Fermi

$\nu_\alpha \xrightarrow{\epsilon_{\alpha\beta}^m} \nu_\beta$ is parametrized as 4-Fermi interaction

$$\frac{1}{\Lambda^2} [\bar{\nu}_\beta \gamma^\rho P_L \nu_\alpha] [\bar{f} \gamma_\rho P_L f]$$

NSI from Dim.6 and Dim.8

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Dim.6 op — 4-Fermi

$\nu_\alpha \xrightarrow{\epsilon_{\alpha\beta}^m} \nu_\beta$ is parametrized as 4-Fermi interaction

$$\begin{aligned} & \frac{1}{\Lambda^2} [\bar{L}_\beta \gamma^\rho L_\alpha] [\bar{f} \gamma_\rho P_L f] \\ &= \frac{1}{\Lambda^2} [[\bar{\nu}_\beta \gamma^\rho P_L \nu_\alpha] + [\bar{\ell}_\beta \gamma^\rho P_L \ell_\alpha]] [\bar{f} \gamma_\rho P_L f] \end{aligned}$$

includes also SU(2) counter process Exception $[\bar{L}^c L][\bar{L}^c L]$

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includes also SU(2) counter process Exception $[\bar{L}^c L][\bar{L}^c L]$

Dim.8 op — 4-Fermi and 2-Higgs

SM gauge invariant form

$$\frac{1}{\Lambda^4} [(\bar{L}_\beta H) \gamma^\rho (H^\dagger L_\alpha)] [\bar{f} \gamma_\rho P_L f]$$

NSI from Dim.6 and Dim.8

Dim.6 op — 4-Fermi

$\nu_\alpha \xrightarrow{\epsilon_{\alpha\beta}^m} \nu_\beta$ is parametrized as 4-Fermi interaction

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includes also SU(2) counter process Exception $[\bar{L}^c L][\bar{L}^c L]$

Dim.8 op — 4-Fermi and 2-Higgs

SM gauge invariant form and after EWSB

$$\begin{aligned} & \frac{1}{\Lambda^4} [(\bar{L}_\beta H) \gamma^\rho (H^\dagger L_\alpha)] [\bar{f} \gamma_\rho P_L f] \\ &= \frac{v^2}{2\Lambda^4} [\bar{\nu}_\beta \gamma^\rho \nu_\alpha] [\bar{f} \gamma_\rho P_L f] + (\text{Higgs ints.}) \end{aligned}$$

$SU(2)$ relation is broken with Higgs vev.

Gavela Hernandez O Winter

NSI from Dim.8

- NSI from Dim.6 (4-Fermi) are strongly constrained

No $SU(2)$ violationBergmann Grossman Pierce PRD**61** (2000) 053005

- NSI from Dim.8 (4-Fermi+2Higgs) have some chance

 $SU(2)$ relation is broken with Higgs vev

Berezhiani Rossi PLB**535** (2002) 207, Davidson Pena-Garay Rius Santamaria JHEP**0303** (2003) 011, Biggio Blennow Fernandez-Martinez JHEP**0903** (2009) 139 and arXiv:0907.0097

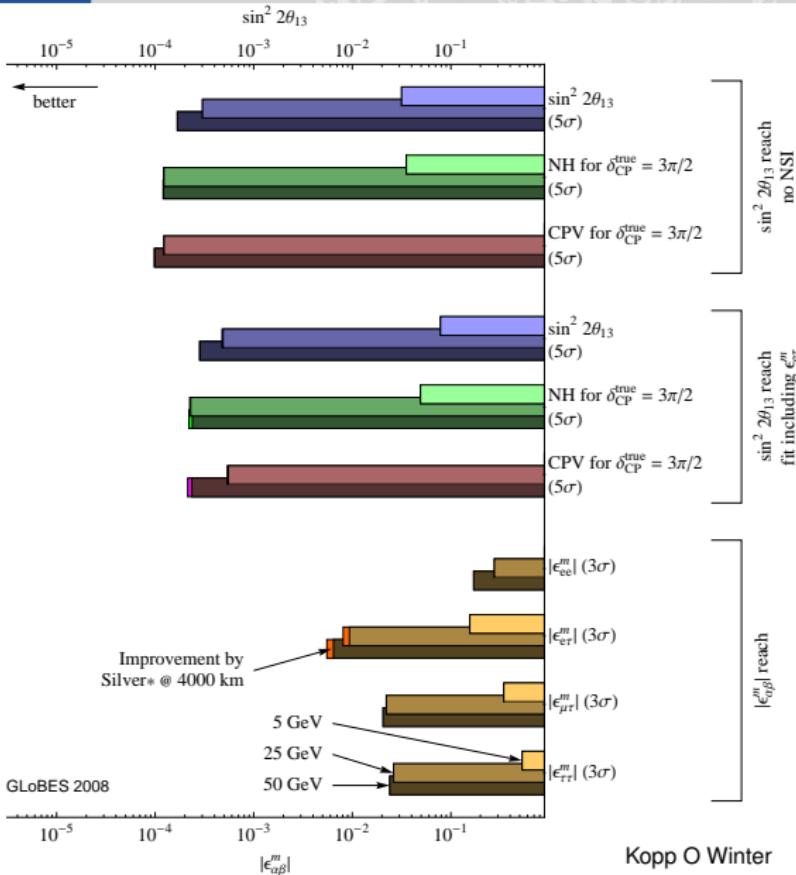
- NSI from a Dim.8 diagram is always constrained
- Antusch Baumann Fernandez-Martinez NPB**810** (2009) 369
- Combination of diagrams can help to obtain large NSI avoiding constraints.

#	Dim. eight operator	\mathcal{C}_{LEH}^1	\mathcal{C}_{LEH}^3	$\mathcal{O}_{\text{NSI?}}$	Mediators
Combination $\bar{L}L$					
1	$(\bar{L}\gamma^\rho L)(\bar{E}\gamma_\rho E)(H^\dagger H)$	1			$\mathbf{1}_0^v$
2	$(\bar{L}\gamma^\rho L)(\bar{E}H^\dagger)(\gamma_\rho)(HE)$	1			$\mathbf{1}_0^v + 2_{-3/2}^{L/R}$
3	$(\bar{L}\gamma^\rho L)(\bar{E}H^T)(\gamma_\rho)(H^*E)$	1			$\mathbf{1}_0^v + 2_{-1/2}^{L/R}$
4	$(\bar{L}\gamma^\rho \bar{\tau}L)(\bar{E}\gamma_\rho E)(H^\dagger \bar{\tau}H)$		1		$\mathbf{3}_0^v + \mathbf{1}_0^R$
5	$(\bar{L}\gamma^\rho \bar{\tau}L)(\bar{E}H^\dagger)(\gamma_\rho \bar{\tau})(HE)$		1		$\mathbf{3}_0^v + 2_{-3/2}^{L/R}$
6	$(\bar{L}\gamma^\rho \bar{\tau}L)(\bar{E}H^T)(\gamma_\rho \bar{\tau})(H^*E)$		1		$\mathbf{3}_0^v + 2_{-1/2}^{L/R}$
Combination $\bar{E}L$					
7	$(\bar{L}E)(\bar{E}L)(H^\dagger H)$	-1/2			$\mathbf{2}_{+1/2}^s$
8	$(\bar{L}E)(\bar{\tau})(\bar{E}L)(H^\dagger \bar{\tau}H)$	-1/2			$\mathbf{2}_{+1/2}^s$
9	$(\bar{L}H)(H^\dagger E)(\bar{E}L)$	-1/4	-1/4	✓	$\mathbf{2}_{+1/2}^s + \mathbf{1}_0^R + 2_{-1/2}^{L/R}$
10	$(\bar{L}FH)(H^\dagger E)(\bar{\tau})(\bar{E}L)$	-3/4	1/4		$\mathbf{2}_{+1/2}^s + \mathbf{3}_0^R + 2_{-1/2}^{L/R}$
11	$(\bar{L}i\tau^2 H^*)(H^T E)(i\tau^2)(\bar{E}L)$	1/4	-1/4		$\mathbf{2}_{+1/2}^s + \mathbf{1}_{-1}^R + 2_{-3/2}^{L/R}$
12	$(\bar{L}\bar{\tau}i\tau^2 H^*)(H^T E)(i\tau^2 \bar{\tau})(\bar{E}L)$	3/4	1/4		$\mathbf{2}_{+1/2}^s + \mathbf{3}_{-1}^{L/R} + 2_{-3/2}^{L/R}$
Combination $\bar{E}^c L$					
13	$(\bar{L}\gamma^\rho E^c)(\bar{E}^c \gamma_\rho L)(H^\dagger H)$	-1			$\mathbf{2}_{-3/2}^v$
14	$(\bar{L}\gamma^\rho E^c)(\bar{\tau})(\bar{E}^c \gamma_\rho L)(H^\dagger \bar{\tau}H)$	-1			$\mathbf{2}_{-3/2}^v$
15	$(\bar{L}H)(\gamma^\rho)(H^\dagger E^c)(\bar{E}^c \gamma_\rho L)$	-1/2	-1/2	✓	$\mathbf{2}_{-3/2}^v + \mathbf{1}_0^R + 2_{+3/2}^{L/R}$
16	$(\bar{L}FH)(\gamma^\rho)(H^\dagger E^c)(\bar{\tau})(\bar{E}^c \gamma_\rho L)$	-3/2	1/2		$\mathbf{2}_{-3/2}^v + \mathbf{3}_0^R + 2_{+3/2}^{L/R}$
17	$(\bar{L}i\tau^2 H^*)(\gamma^\rho)(H^T E^c)(i\tau^2)(\bar{E}^c \gamma_\rho L)$	-1/2	1/2		$\mathbf{2}_{-3/2}^v + \mathbf{1}_{-1}^R + 2_{+1/2}^{L/R}$
18	$(\bar{L}\bar{\tau}i\tau^2 H^*)(\gamma^\rho)(H^T E^c)(i\tau^2 \bar{\tau})(\bar{E}^c \gamma_\rho L)$	-3/2	-1/2		$\mathbf{2}_{-3/2}^v + \mathbf{3}_{-1}^{L/R} + 2_{+1/2}^{L/R}$
Combination $\bar{E}^T L$					
19	$(\bar{L}E)(\bar{E}H)(H^\dagger L)$	-1/4	-1/4	✓	$\mathbf{2}_{+1/2}^s + \mathbf{1}_0^R + 2_{-1/2}^{L/R}$
20	$(\bar{L}E)(\bar{\tau})(\bar{E}H)(H^\dagger \bar{\tau}L)$	-3/4	1/4		$\mathbf{2}_{+1/2}^s + \mathbf{3}_0^R + 2_{-1/2}^{L/R}$
21	$(\bar{L}H)(\gamma^\rho)(H^\dagger L)(\bar{E}\gamma_\rho E)$	1/2	1/2	✓	$\mathbf{1}_0^v + \mathbf{1}_0^R$
22	$(\bar{L}FH)(\gamma^\rho)(H^\dagger \bar{\tau}L)(\bar{E}\gamma_\rho E)$	3/2	-1/2		$\mathbf{1}_0^v + \mathbf{3}_0^R$
23	$(\bar{L}i\tau^2 E^c)(\bar{E}^c H)(\gamma^\rho)(H^T L)$	-1/2	-1/2	✓	$\mathbf{2}_{-3/2}^v + \mathbf{1}_0^R + 2_{+3/2}^{L/R}$
24	$(\bar{L}\gamma^\rho E^c)(\bar{E}^c H)(\gamma^\rho)(H^\dagger L)$	-3/2	1/2		$\mathbf{2}_{-3/2}^v + \mathbf{3}_{-1}^R + 2_{+3/2}^{L/R}$
Combination $\bar{H}L$					
25	$(\bar{L}E)(i\tau^2)(\bar{E}H^*)(H^T i\tau^2 L)$	1/4	-1/4		$\mathbf{2}_{+1/2}^s + \mathbf{1}_{-1}^R + 2_{+2/2}^{L/R}$
26	$(\bar{L}E)(\bar{\tau}i\tau^2)(\bar{E}H^*)(H^T i\tau^2 \bar{\tau}L)$	3/4	1/4		$\mathbf{2}_{+1/2}^s + \mathbf{3}_{-1}^R + 2_{-3/2}^{L/R}$
27	$(\bar{L}i\tau^2 H^*)(\gamma^\rho)(H^T i\tau^2 L)(\bar{E}\gamma_\rho E)$	-1/2	1/2		$\mathbf{1}_0^v + 2_{-1}^{L/R}$
28	$(\bar{L}\bar{\tau}i\tau^2 H^*)(\gamma^\rho)(H^T i\tau^2 \bar{\tau}L)(\bar{E}\gamma_\rho E)$	-3/2	-1/2		$\mathbf{1}_0^v + 3_{-1}^{L/R}$
29	$(\bar{L}\gamma^\rho E^c)(i\tau^2)(\bar{E}^c H^*)(\gamma_\rho)(H^T i\tau^2 L)$	1/2	-1/2		$\mathbf{2}_{-3/2}^v + \mathbf{1}_{-1}^R + 2_{+1/2}^{L/R}$
30	$(\bar{L}\gamma^\rho E^c)(\bar{\tau}i\tau^2)(\bar{E}^c H^*)(\gamma_\rho)(H^T i\tau^2 \bar{\tau}L)$	3/2	1/2		$\mathbf{2}_{-3/2}^v + \mathbf{3}_{-1}^{L/R} + 2_{+1/2}^{L/R}$



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Optimal Nufact setup for NSI

$E_\mu = 25 \text{ GeV OK}$
 $L = 4000\text{km} + 7500\text{km OK}$

Same for the standard oscillation parameters.

- Silver channel ($\nu_e \rightarrow \nu_\tau$) does not help much.
- Combination of Golden and Disapp. channels, and also combination of two baselines help to resolve parameter correlations.



Backup: Experimental setup

Our setup is based on IDS-NF baseline:

- 50 kton Magnetized Iron Detector (MIND) for Golden and Disapp. channels.
- 10 kton emulsion cloud chamber for Silver channel
- Silver* = signal $\times 5$ and background $\times 3$.
- $2.5 \cdot 10^{21}$ useful muon decays per baseline and polarity
- Charge ID in Golden and Silver channels,
but not in Disapp.

We can count out charge missID background.

Huber Lindner Rolinec Winter PRD74 (2006) 073003.

Kopp O Winter

Backup: higher muon energy $E_\mu = 50 \text{ GeV}$

