

SPACE CHARGE NUMERICAL SIMULATION EXPERIMENTS

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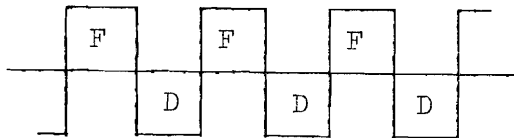
Comments are presented here on the work contained in a CERN report.^(a)

These comments are mainly oriented toward to comparison with the work of I Haber.^(b)

The CERN work was performed about 10 years ago, at the same time as both numerical computations and experimental measurements of emittance growth in a linac were being carried out, in the hope of contributing to the understanding of the phenomena or of being the start of some theory.

1. Focusing System Considered

FD Focusing:



is defined by the betatron phase shift per focusing period at zero beam intensity (represented by μ and expressed in radians in the report; $\mu(\text{radians}) = \sigma_0$ (degrees)).

The intensity of space charge is expressed by the dimensionless parameter:

$$\delta \cong \frac{q}{mv} \frac{\lambda_0 I}{2 \epsilon \beta \gamma}$$

where q is the charge, m is the mass, v is the velocity of the particle, λ_0 is the betatron wavelength and I is the current (in coulombs per unit time).

a) Etude numerique d'effets de charge d'espace en focalisation periodique, by P. M. Lapostolle, CERN report ISR/78-13.

b) Presented at HIF Workshop, Oct. 29-Nov. 9, 1979, Claremont Hotel, Oakland, CA (See present proceedings).

There is a direct connection between δ and σ/σ_0 :

$$\text{for } \delta = 1.5 \quad \sigma / \sigma_0 \approx 0.5,$$

$$\text{for } \delta = 2.0 \quad \sigma / \sigma_0 \approx 0.4$$

(The latter value is the threshold of the first Gluckstern mode for a round beam, K-V distribution.)

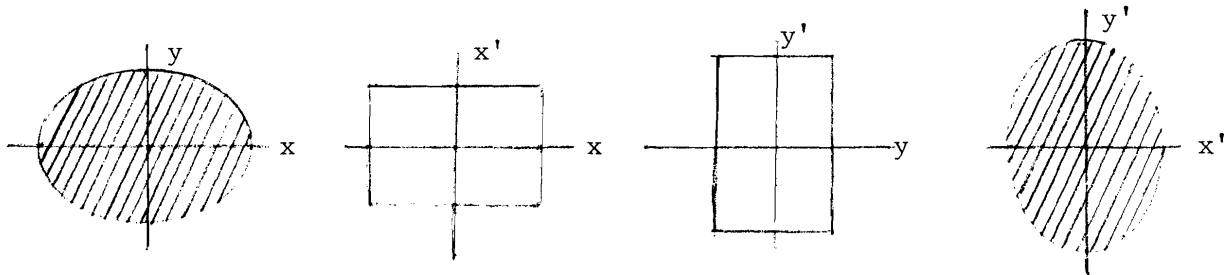
2. Particle Filling at the Input

In most of the runs, 2,500 macroparticles were used with a random uniform distribution inside a fixed boundary.

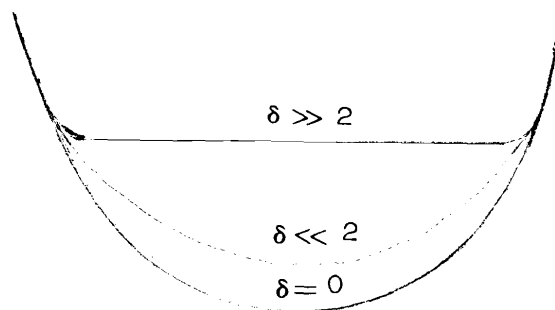
Such a fixed boundary only exists for continuous focusing and is then defined through a Bessel function. The boundary was approximated in actual cases by a properly sized (final adjustment empirical):

hyperellipsoid for $\delta < 2$,

uniform distribution in x, y and x', y' for $\delta > 2$.



In the case, $\delta > 2$, the potential well in which particles move is no longer paraboloidal, but tends to present a flat bottom where no restoring



force exists, except on the edges where the force extends over a thickness equal to the Debye length.

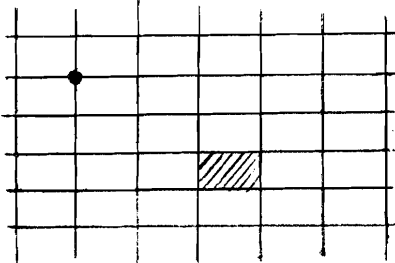
For $\delta = 2$, the Debye length is about equal to the beam radius.

For large δ , individual particles do not perform sinusoidal oscillations and their frequency spectrum may become wide. (This was one reason not to take, during this work, σ/σ_0 as a parameter but, rather, δ , along with the fact that it was intended to keep the possibility of considering unequal sizes in x and y).

Some other distributions have been tested: a) KV, to check the stability limit (when stable, the 4-D surface immediately become slightly wrinkled due to statistical fluctuations, but with a constant thickness). b) Gaussian and uniform-gaussian have also been tested to represent the output of a source.

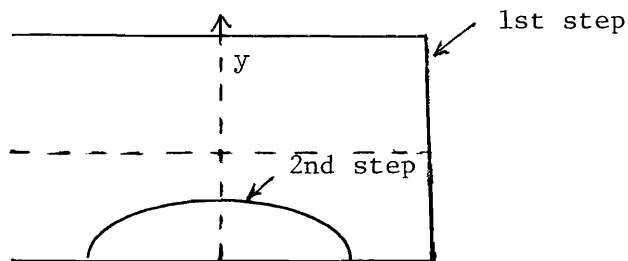
3. Methods of Computation

Space charge is computed in free space, but assuming a 4-fold symmetry (2-fold only in certain cases).



direct 10 x 10 box-to-nodes computation
or
FFT (16 x 16 or 32 x 11) } + interpolation

The rectangular mesh is variable (sized to 1.1 outer particles) in order to keep maximum accuracy. (The potential on the boundary is computed from a series expansion of the outer potential: with a rectangular mesh of aspect ratio larger than 1.1, a 2-step computation is performed, the first is on a square mesh to compute the potential on a close boundary which would enter the circle of convergence of the expansion. For example, consider the 2-fold symmetry case illustrated here:



4. Results

Detailed results appear in the report quoted (a). Only a summary and some remarks are given here:

- For $\sigma_0 > 90^\circ$, even for $0.4 < \sigma/\sigma_0 < 0.6$, large emittance increases appear (as for KV). (The emittance level reached indicates an overshoot effect in this case).

- For $\sigma_0 < 90^\circ$ and large enough tune depression ($\sigma/\sigma_0 < 0.4$) slower increases are seen: the smaller the tune depression, the slower the increase, the smaller the σ_0 , the slower the increase (at least in the cases computed and with 4-fold symmetry assumed in the computations; on the CERN linac, there was some experimental evidence of this effect).

- For continuous focusing, even with tune depressions as large as 95%, no deviation from the original distribution was observed (apart from the immediate statistical effect). However, no way was found to specify the absence of oscillations.

- It is mainly the outer part of the beam which is affected by higher order space charge effects; the central core almost keeps its 2-D phase space density. (The 4-D is obviously conserved.)

- It is not obvious how to determine from the simulation experiments whether or not a more stable distribution exists.

A few additional tests were performed with the simulation program:

1. Filamentation due to mismatch in a continuous focusing system, even

for cases with very large space charge,

$$x_{\min}^2 + x_{\max}^2 = ct$$

$$y_{\min}^2 + y_{\max}^2 = ct$$

The above leads, after damping has taken place, to an emittance increase. (Similar behavior was observed with a non-linear third-order field and no space charge.)

2. Some early simulation work with non-circular beams which was not followed up:

Starting with $\sigma_{ox} = \sigma_{oy}$ and unequal $\epsilon_{ox}, \epsilon_{oy}$, it was found that ϵ_x and ϵ_y become equal with their sum remaining constant (constant energy). Starting with $\sigma_{ox} \neq \sigma_{oy}$ (by a factor of ≈ 2) nothing was seen, neither with equal initial energies in x and y, nor with equal emittances. (These runs were made when the program was not optimized, interpolations were crude and the accuracy was such that small variations were considered irrelevant.)