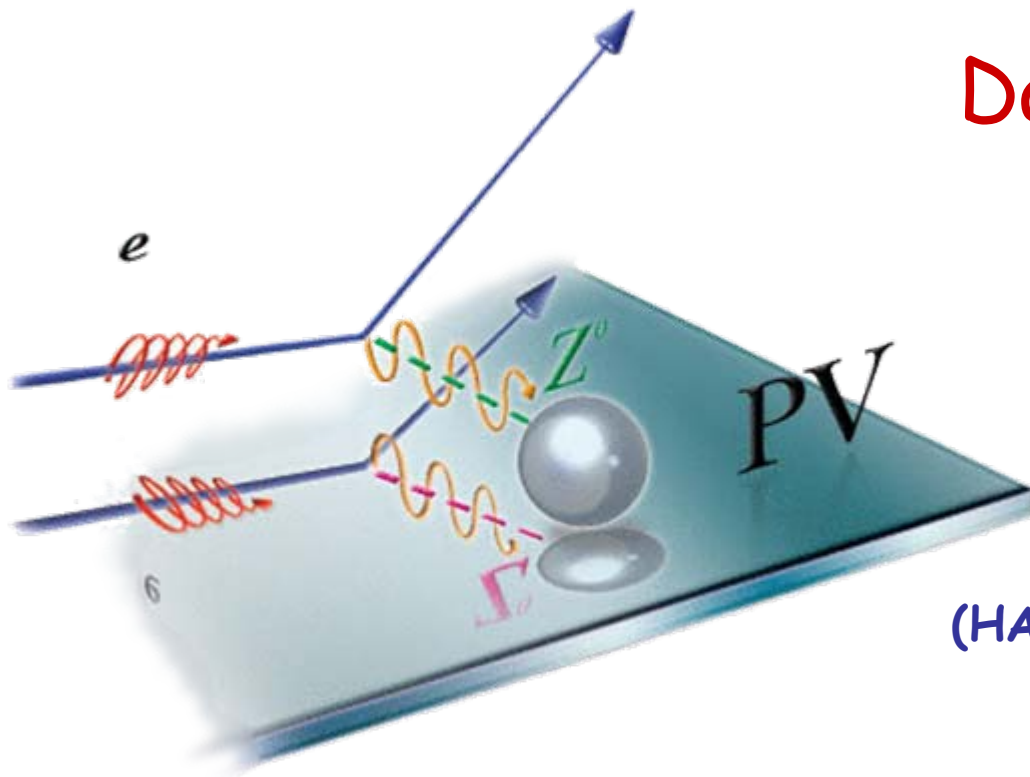


Parity-Violating Electron Scattering & Strangeness in the Nucleon

David S. Armstrong

College of William & Mary

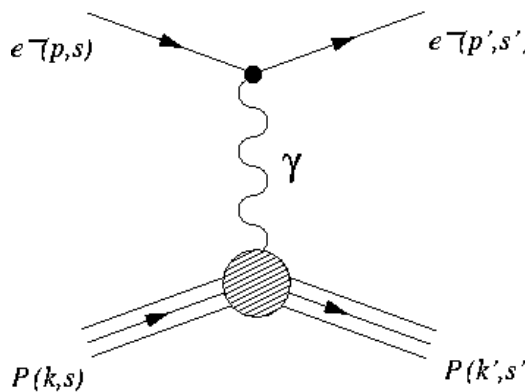


(HAPPEX and G0 Collaborations)

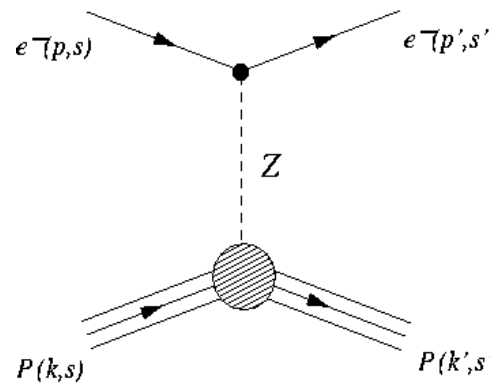
Outline

- Parity-violation in electron scattering
- Elastic Vector Strange Form Factors: G_E^s and G_M^s
- First generation results: HAPPEX-I, SAMPLE, PV-A4
- Latest results:
 - GO (forward-angle)
 - HAPPEX-II and HAPPEX-Helium
- The present situation at $Q^2 = 0.1 \text{ (GeV/c)}^2$
- The future...

Parity Violating Electron Scattering → Weak NC Amplitudes



$$M^{EM} = \frac{4\pi\alpha}{Q^2} Q_\ell \ell^\mu J_\mu^{EM}$$



$$M_{PV}^{NC} = \frac{G_F}{2\sqrt{2}} \left[g_A \ell^{\mu 5} J_\mu^{NC} + g_V \ell^\mu J_{\mu 5}^{NC} \right]$$

Interference: $\sigma \sim |M^{EM}|^2 + |M^{NC}|^2 + 2\text{Re}(M^{EM*})M^{NC}$

Interference with EM
amplitude makes Neutral
Current (NC) amplitude
accessible

$$\Rightarrow A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{|M_{PV}^{NC}|}{|M^{EM}|} \sim \frac{Q^2}{(M_Z)^2}$$

Tiny ($\sim 10^{-6}$) cross section asymmetry isolates weak interaction

Form Factors

$$J_{\mu}^{EM} = \sum_q Q_q \langle \bar{N} | \bar{u}_q \gamma_{\mu} u_q | N \rangle = \bar{N} \left[\gamma_{\mu} F_1^{\gamma} + \frac{i \sigma_{\mu\nu} q^{\nu}}{2M_N} F_2^{\gamma} \right] N$$

Adopt the Sachs FF: $G_E^{\gamma} = F_1^{\gamma} + \tau F_2^{\gamma}$ $G_M^{\gamma} = F_1^{\gamma} + F_2^{\gamma}$
(Roughly: Fourier transforms of charge and magnetization)

NC probes **same** hadronic flavor structure, with different couplings:

$$G_{E/M}^{\gamma} = \frac{2}{3} G_{E/M}^u - \frac{1}{3} G_{E/M}^d - \frac{1}{3} G_{E/M}^s$$

$$G_{E/M}^Z = \left(1 - \frac{8}{3} \sin^2 \theta_W \right) G_{E/M}^u - \left(1 - \frac{4}{3} \sin^2 \theta_W \right) G_{E/M}^d - \left(1 - \frac{4}{3} \sin^2 \theta_W \right) G_{E/M}^s$$

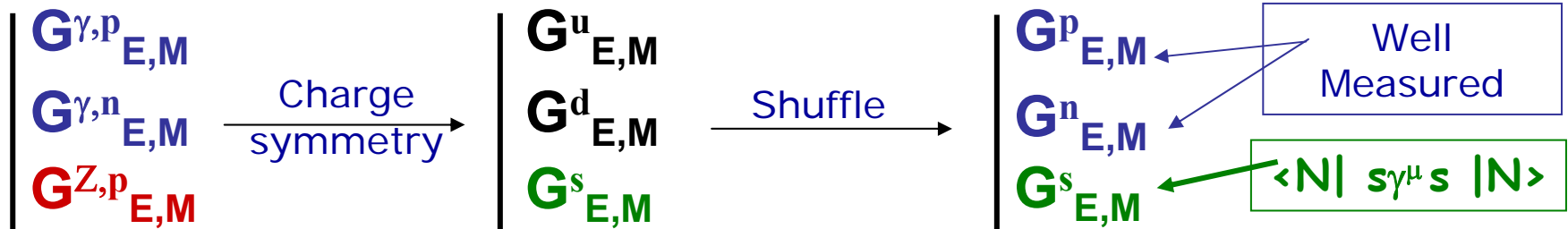
$G_{E/M}^Z$ provide an important new benchmark for testing non-perturbative QCD structure of the nucleon

Charge Symmetry

One expects the neutron to be an isospin rotation of the proton*:

$$G_{E/M}^{p,u} = G_{E/M}^{n,d}, \quad G_{E/M}^{p,d} = G_{E/M}^{n,u}, \quad G_{E/M}^{p,s} = G_{E/M}^{n,s}$$

$$G_{E/M}^{\gamma,p} = \frac{2}{3} G_{E/M}^u - \frac{1}{3} G_{E/M}^d - \frac{1}{3} G_{E/M}^s \longrightarrow G_{E/M}^{\gamma,n} = \frac{2}{3} G_{E/M}^d - \frac{1}{3} G_{E/M}^u - \frac{1}{3} G_{E/M}^s$$



$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \propto \frac{M_Z M_\gamma}{|M_\gamma|^2} = -\frac{G_F Q^2}{\sqrt{2} \pi \alpha} F(G_{E/M}^p, G_{E/M}^n, G_{E/M}^s, G_A)$$

*Neglecting trivial breaking due to Coulomb force

Isolating the form factors: vary the kinematics or target

For a proton:

$$A = \left[\frac{-G_F Q^2}{4\pi\alpha\sqrt{2}} \right] \frac{A_E + A_M + A_A}{\sigma_p} \quad \sim \text{few parts per million}$$

$$A_E = \varepsilon G_E^p G_E^Z, \quad A_M = \tau G_M^p G_M^Z, \quad A_A = -(1 - 4\sin^2 \theta_W) \varepsilon' G_M^p G_A^e$$

Forward angle

Backward angle

$$G_{E,M}^Z = (1 - 4\sin^2 \theta_W)(1 + R_V^p)G_{E,M}^p - (1 + R_V^n)G_{E,M}^n - G_{E,M}^s$$

$$G_A^e = -G_A + \Delta s + \eta F_A + R^e$$

For ${}^4\text{He}$: G_E^s alone (but
only available at low Q^2)

$$A_{PV} = \frac{G_F Q^2}{\pi\alpha\sqrt{2}} \left[\sin^2 \theta_W + \frac{G_E^s}{2(G_E^p + G_E^n)} \right]$$

For deuterium:

enhanced G_A^e sensitivity

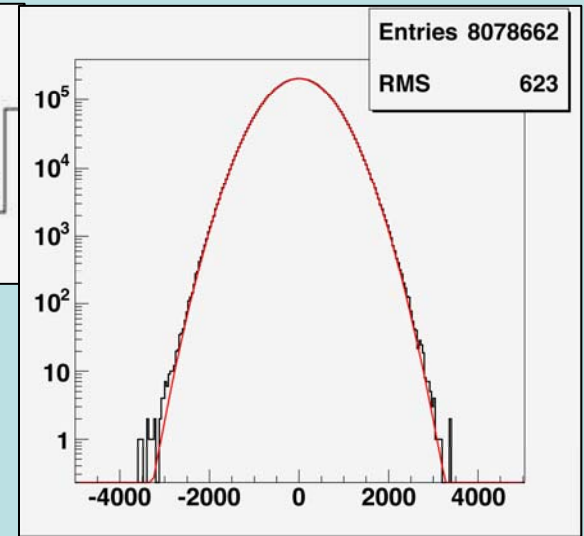
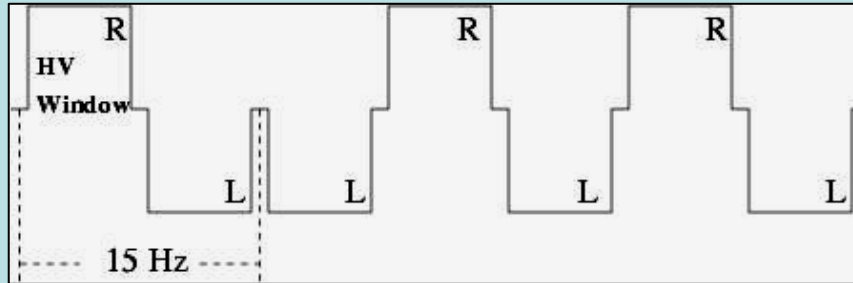
Measurement of P-V Asymmetries

$$A_{LR} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \approx 10^{-6}$$

5% Statistical Precision on 1 ppm
→ requires 4×10^{14} counts

Rapid Helicity Flip: Measure the asymmetry at 10^{-4} level, 10 million times

$$A_{LR} = \frac{N_R - N_L}{N_R + N_L}$$



- High luminosity: thick targets, high beam current
- Control noise (target, electronics)
- Polarized source uses optical pumping of strained photocathode: high polarization and rapid flip

Statistics: high rate, low noise

Systematics: beam asymmetries, backgrounds, Helicity correlated DAQ

Normalization: Polarization, Linearity, Background

Early History: Tests of Weinberg-Salam-Glashow

- C. Prescott, *et al.* SLAC E122 DIS on deuterium
Phys. Lett. 77B, 3 47 (1978), Phys. Lett. 84B, 524 (1979)
- W. Heil, *et al.* Mainz quasielastic ${}^9\text{Be}$
Nucl. Phys. B327, 1 (1989)
- P. Souder, *et al.* MIT/Bates ${}^{12}\text{C}$ elastic
PRL 65, 694 (1990)

HAPPEX (first generation)

Hydrogen Target: $E=3.3 \text{ GeV}$, $\theta=12.5^\circ$, $Q^2=0.48 \text{ (GeV/c)}^2$

$$A^{PV} = \left[\frac{-G_F M_p^2 \tau}{\pi \alpha \sqrt{2}} \right] \left\{ (1 - 4 \sin^2 \theta_W) - \frac{[\varepsilon G_E^{p\gamma} (G_E^{n\gamma} + G_E^s) + \tau G_M^{p\gamma} (G_M^{n\gamma} + G_M^s)]}{\varepsilon (G_E^{p\gamma})^2 + \tau (G_M^{p\gamma})^2} \right\} - A_A$$

$$A^{PV} = -14.92 \text{ ppm} \pm 0.98 \text{ (stat) ppm} \pm 0.56 \text{ (syst) ppm}$$

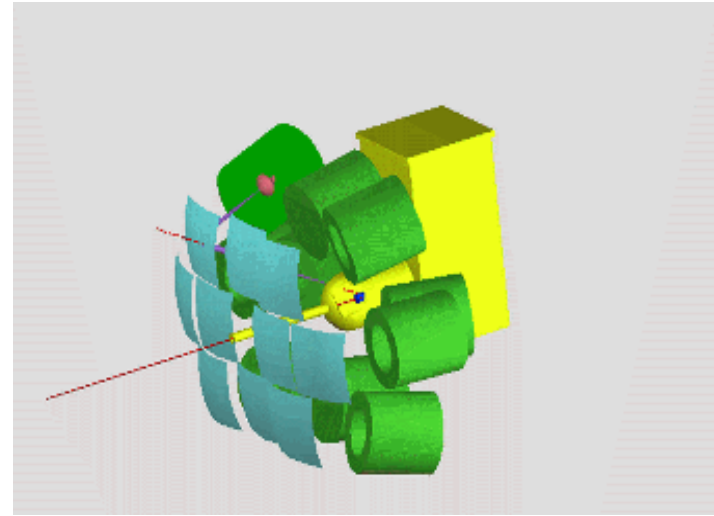
$$G_E^s + 0.39 G_M^s = 0.01 \pm 0.020 \text{ (exp)} \pm 0.010 \text{ (FF)}$$

"Parity Quality" Beam @ JLab

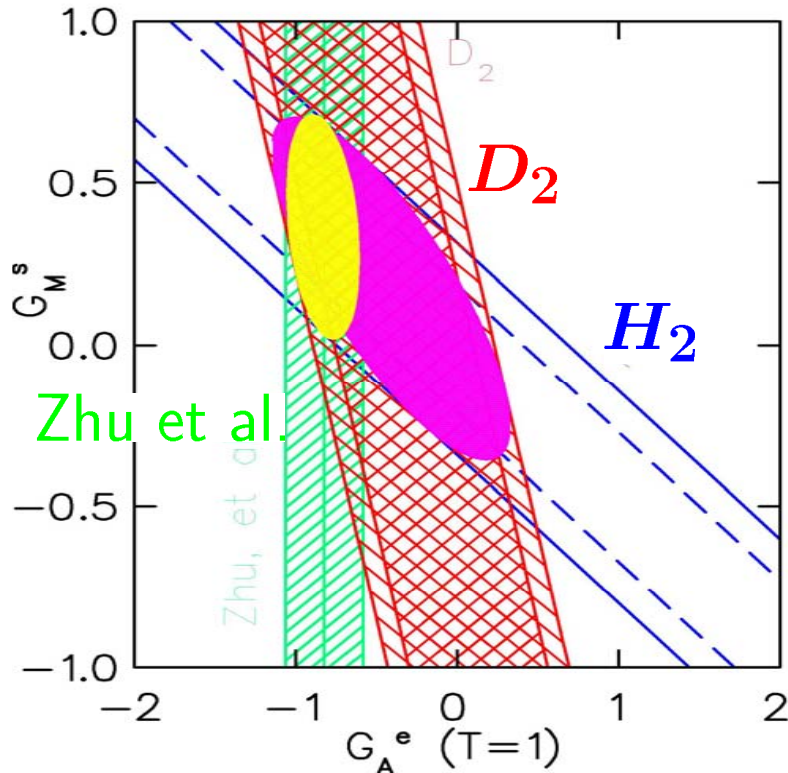
Phys. Rev. Lett. 82, 1096 (1999);
Phys. Lett. B509, 211 (2001);
Phys. Rev. C 69, 065501 (2004)

A_A suppressed by $\varepsilon' (1 - 4 \sin^2 \theta_W)$ where $\varepsilon' = [\tau(1 + \tau)(1 - \varepsilon^2)]^{\frac{1}{2}} \approx (0.08)(0.08)$ here.

SAMPLE (MIT/Bates)



$Q^2(\text{GeV}^2)$	A_{PV} (ppm)	$A_0 + \alpha G_M^s + \beta G_A^e(T=1)$
0.1, LH_2	$-5.61 \pm 0.67 \pm 0.88$	$-5.56 + 3.37 G_M^s + 1.54 G_A^e$
0.1, LD_2	$-7.06 \pm 0.73 \pm 0.72$	$-7.06 + 0.72 G_M^s + 1.66 G_A^e$
0.03, LD_2	$-3.51 \pm 0.57 \pm 0.58$	$-2.14 + 0.27 G_M^s + 0.76 G_A^e$



$$G_M^s = 0.23 \pm 0.36 \pm 0.40$$

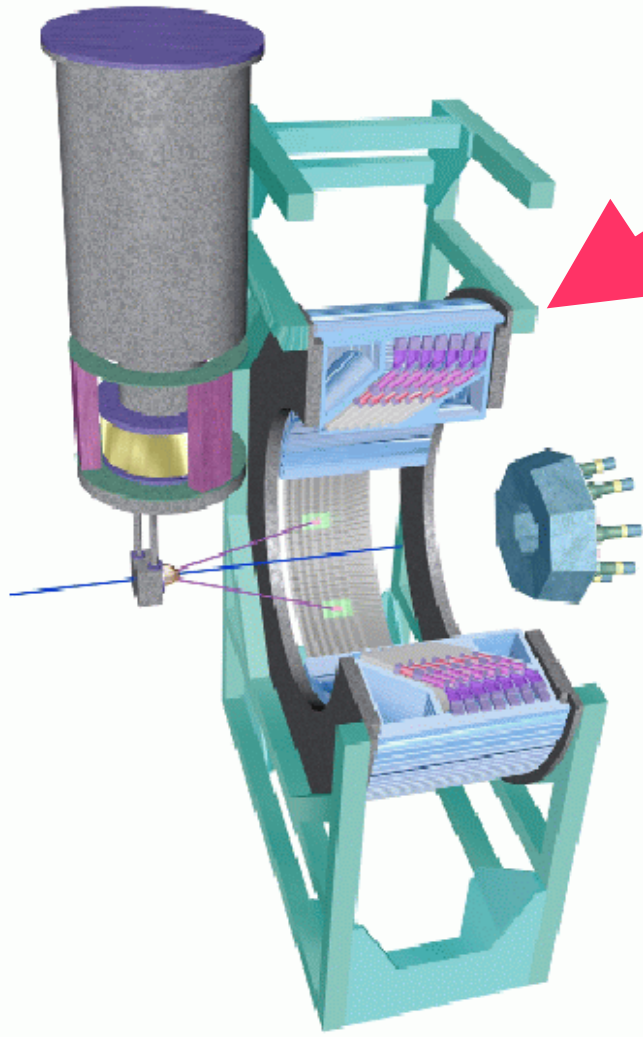
$$G_A^e(T=1) = -0.53 \pm 0.57 \pm 0.50$$

E.J. Beise *et al.*, Prog Nuc Part Phys 54 (2005)

Results of Zhu *et al* commonly used to constrain G_M^s result:

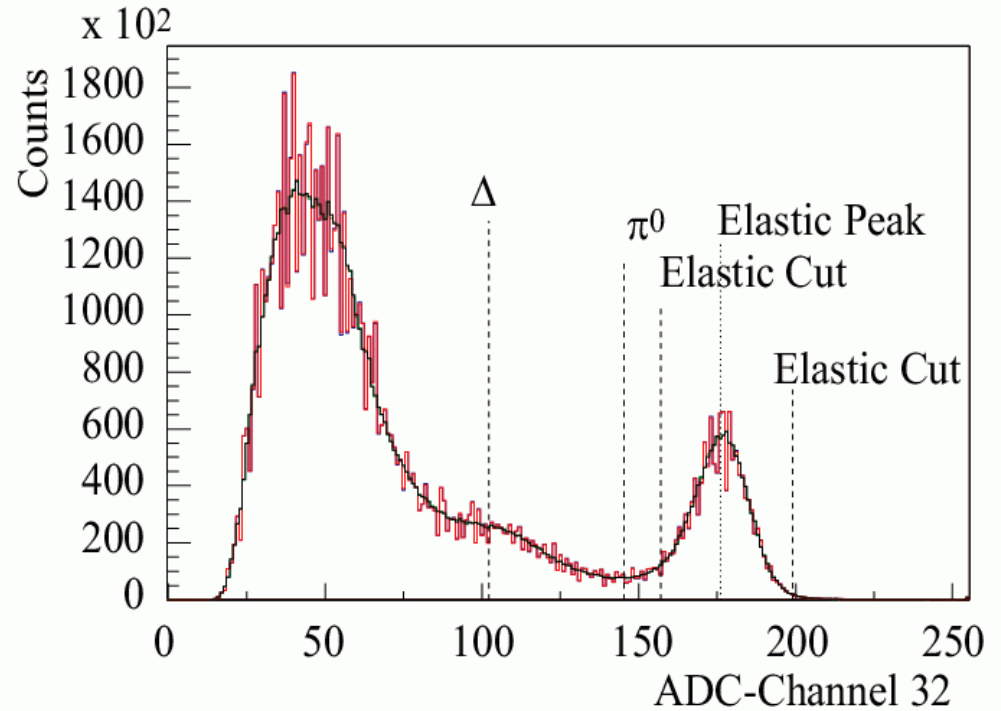
$$G_M^s = 0.37 \pm 0.20_{Stat} \pm 0.36_{Syst} \pm 0.07_{FF}$$

PV-A4 at Mainz



Fast PbF_2 calorimeter

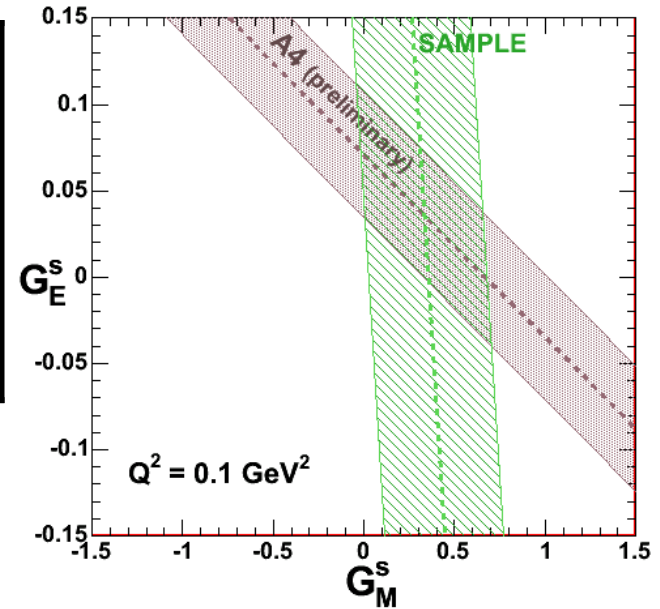
- 1022 crystals in 146 frames
- $\theta = 30\text{-}40^\circ$, $\varphi = 0\text{-}360^\circ$



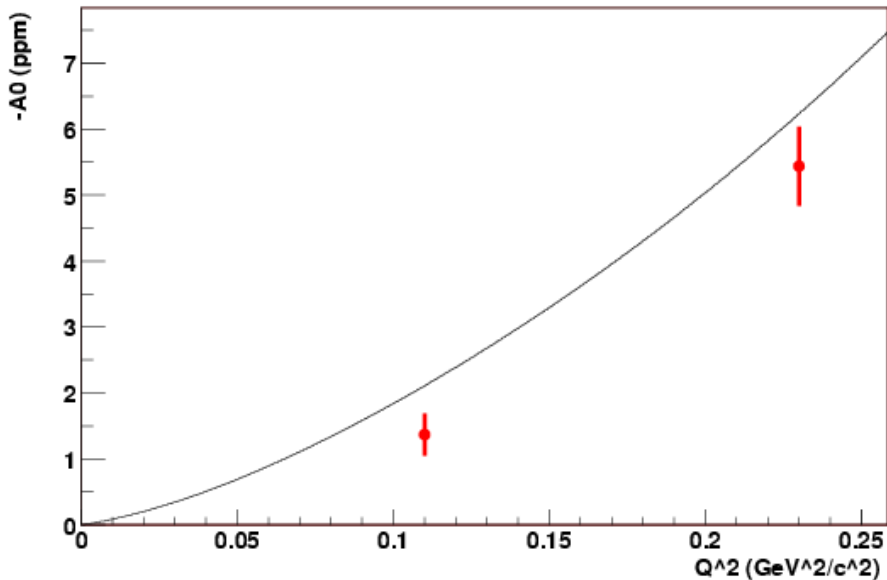
For $Q^2=0.108 \text{ (GeV/c)}^2$, 16×10^6 histograms $\rightarrow 10^{13}$ elastic scattering events!

PV-A4 (MAMI/Mainz)

Q^2 (GeV^2)	$A \pm \text{stat} \pm \text{syst}$ (ppm)	$G_E^s + \alpha G_M^s$
0.230	$-5.44 \pm 0.54 \pm 0.26$	$G_E^s + 0.225 G_M^s$ $= 0.039 \pm 0.034$
0.101	$-1.36 \pm 0.29 \pm 0.13$	$G_E^s + 0.106 G_M^s$ $= 0.071 \pm 0.036$



A0 at theta=35.0 deg

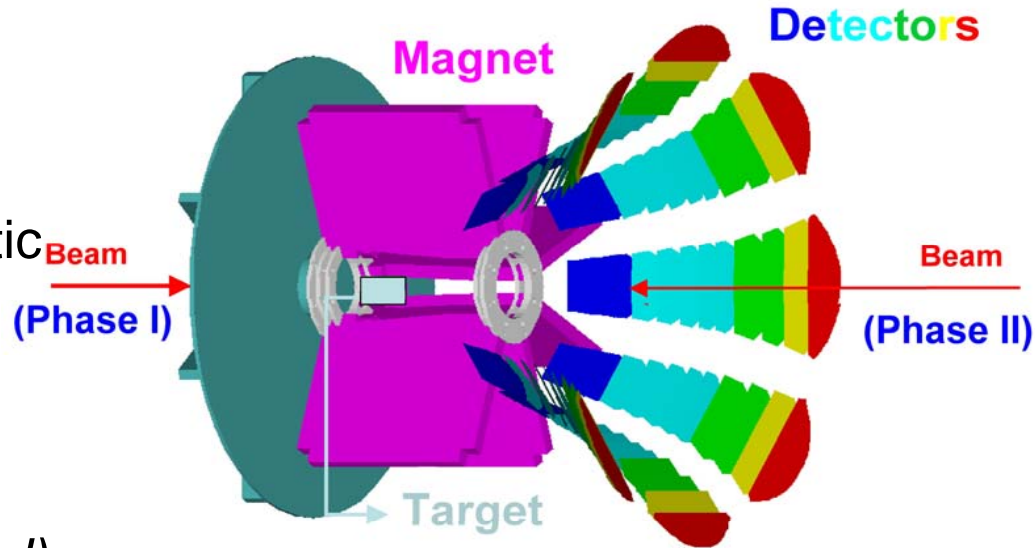


"Evidence for Strange Quark Contributions to the Nucleon's Form Factors at $Q^2 = 0.1 \text{ GeV}^2$ "
PRL 94, 152001 (2006)

Back Angle runs underway to separate G_M^s , G_A at additional points...

G^0 (JLab - Hall C)

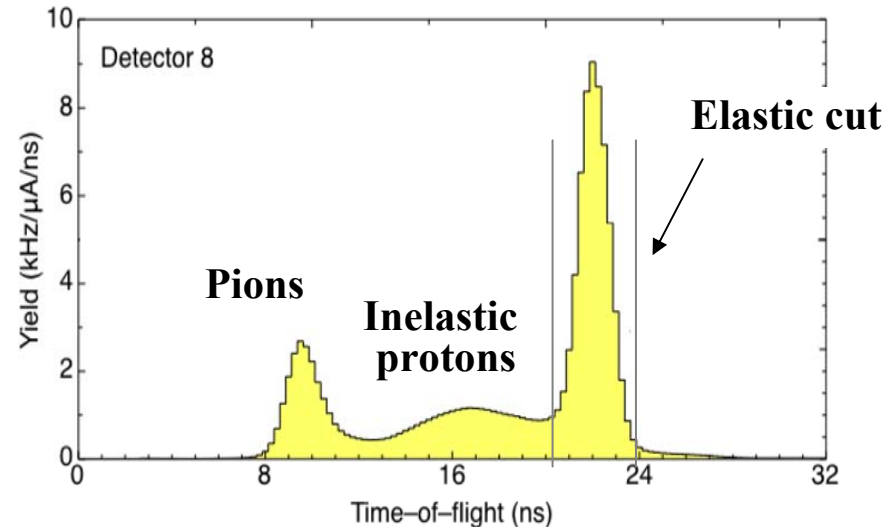
- LH₂/LD₂ target (20 cm)
 $L = 2 \cdot 10^{38} \text{ cm}^{-2} \text{ s}^{-1}$
- Superconducting toroidal magnetic spectrometer
- 16 “Rings” divided into 8 octants



Forward angle mode (completed):

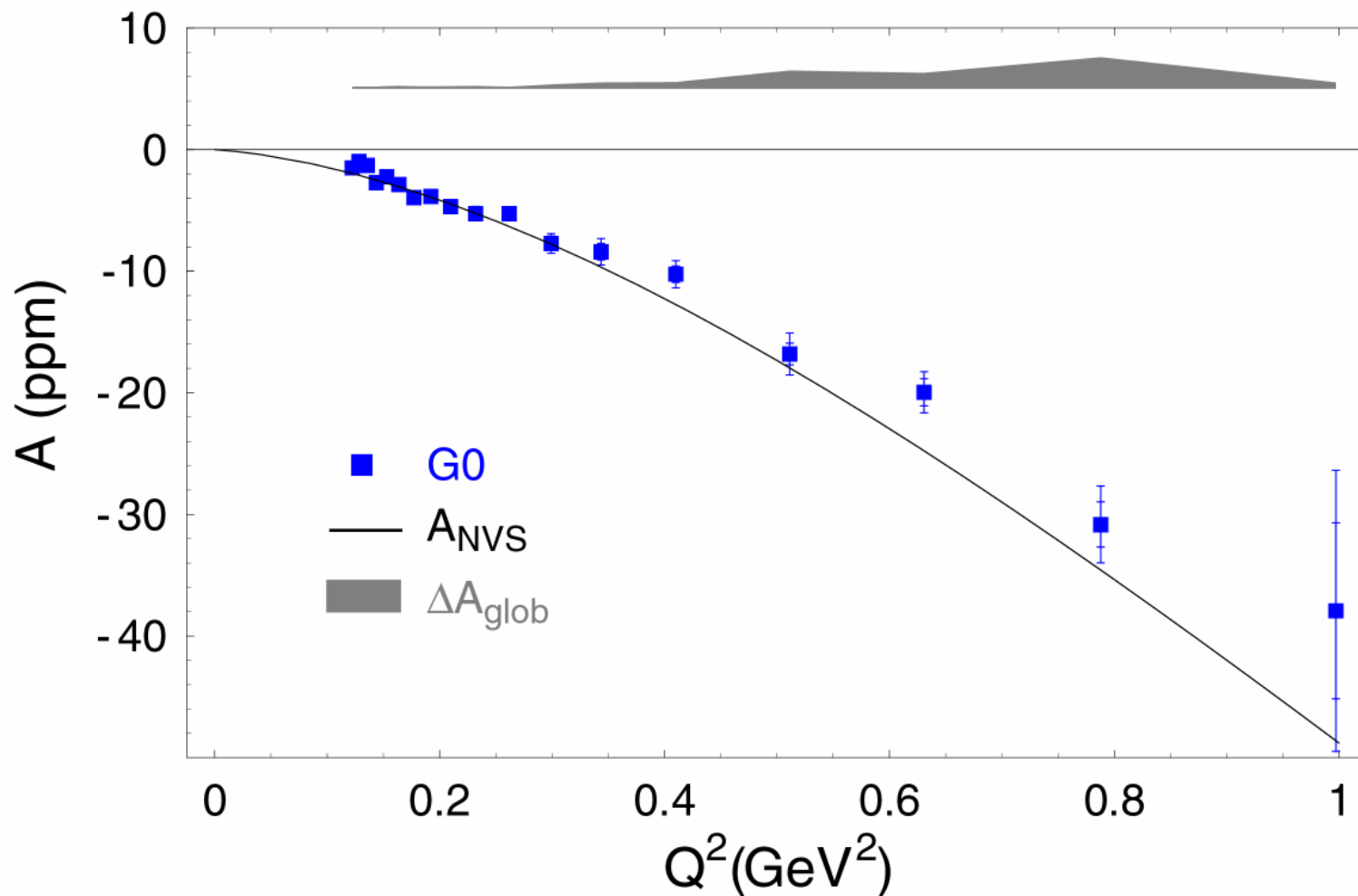
- LH₂: $E_e = 3.0 \text{ GeV}$
Recoil proton detection ($52^\circ < \theta_p < 76^\circ$)
↳ $0.12 \leq Q^2 \leq 1.0 \text{ (GeV/c)}^2$
- Counting experiment – separate backgrounds via time-of-flight
Histograms built each 33 ms

E99-016, E01-115 and E01-116

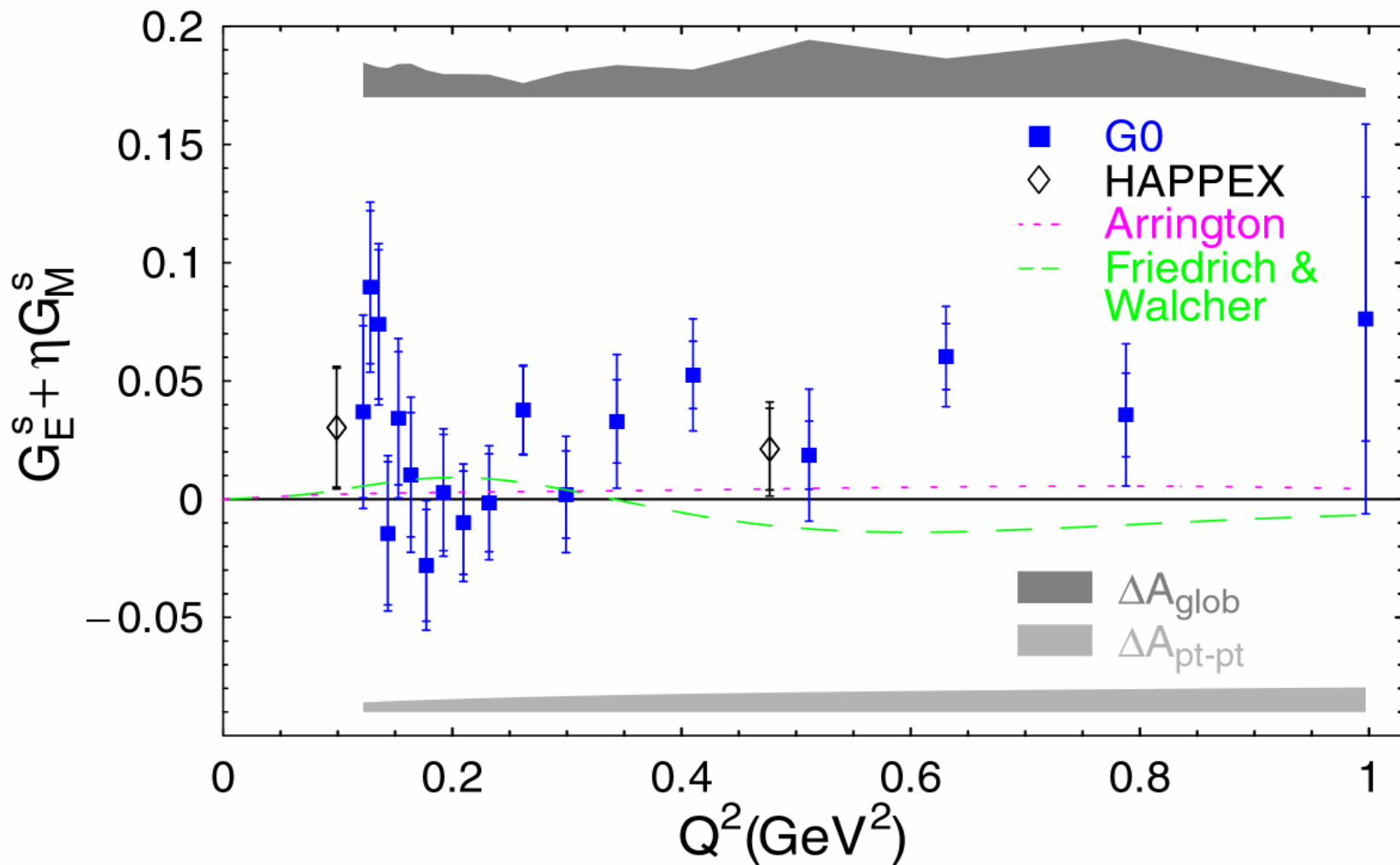


GO Asymmetries (Forward-Angle)

- EM form factors: Kelly PRC 70 (2004) 068202
- A_{NVS} = "no vector strange" asymmetry = $A(G_E^s, G_M^s = 0)$
- inside error bars: *stat*, outside: *stat & pt-pt*



G0: Forward-angle results



$G_E^s = G_M^s = 0$: Hypothesis excluded at 89% C.L.

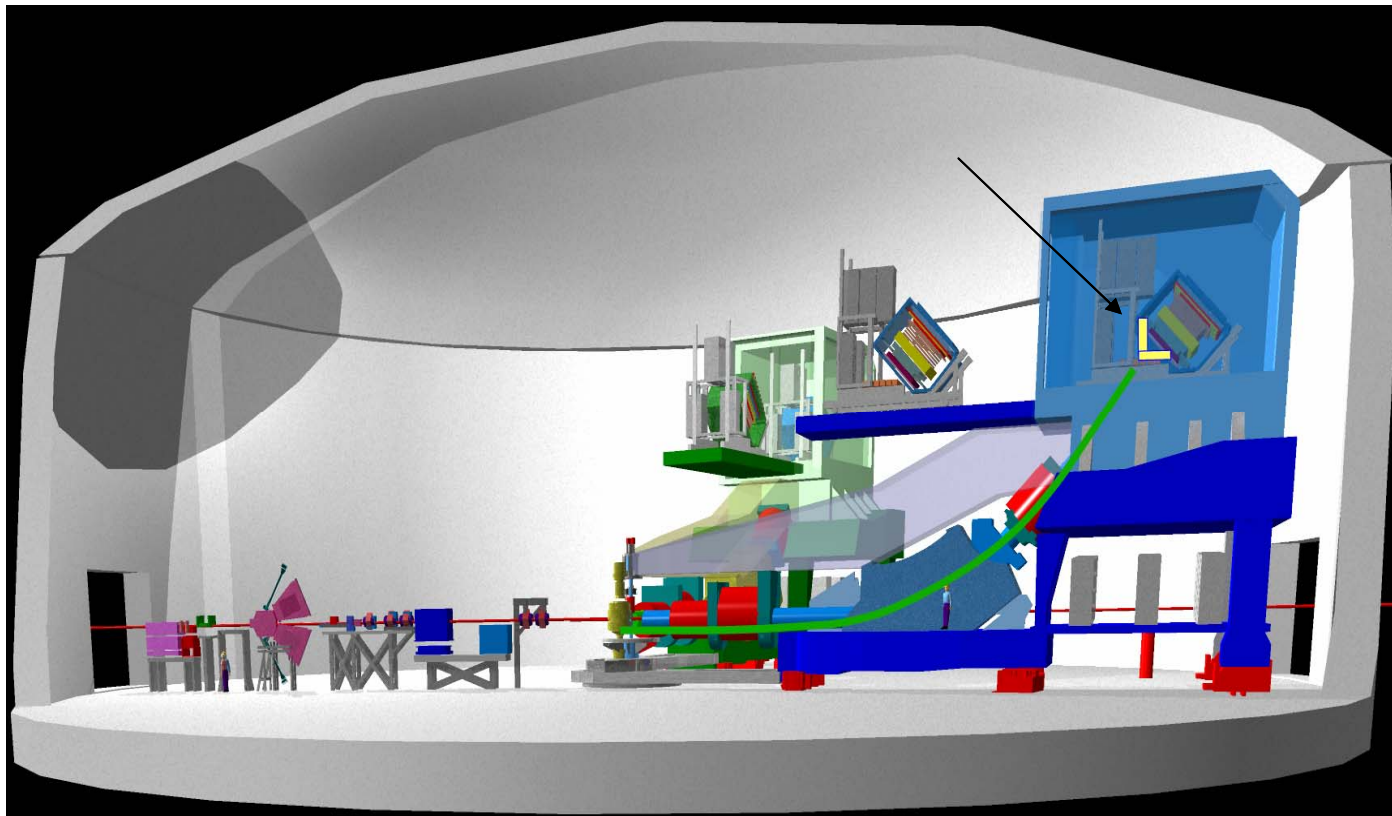
HAPPEX (second generation)

$$E=3 \text{ GeV} \quad \theta=6^\circ \quad Q^2=0.1 \text{ (GeV/c)}^2$$

•Hydrogen : $G_E^s + \alpha G_M^s$

• ^4He : Pure G_E^s : $A^{PV} = -\frac{A_0}{2} \left(2 \sin^2 \theta_W + \frac{G_E^s}{G_E^{p\gamma} + G_E^{n\gamma}} \right)$

New results: just released (P. Souder at Dallas APS meeting)

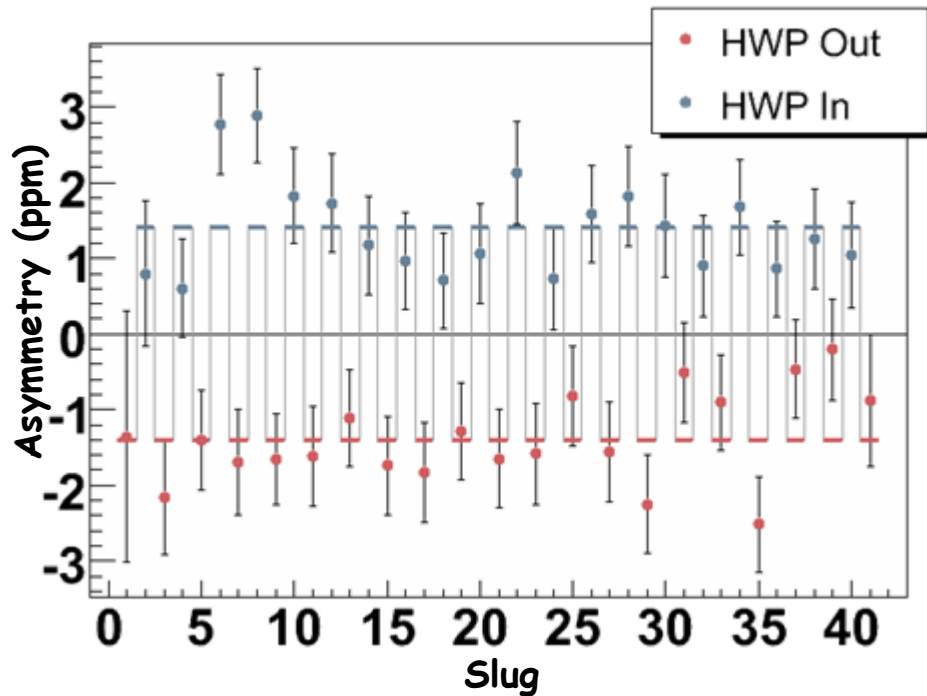
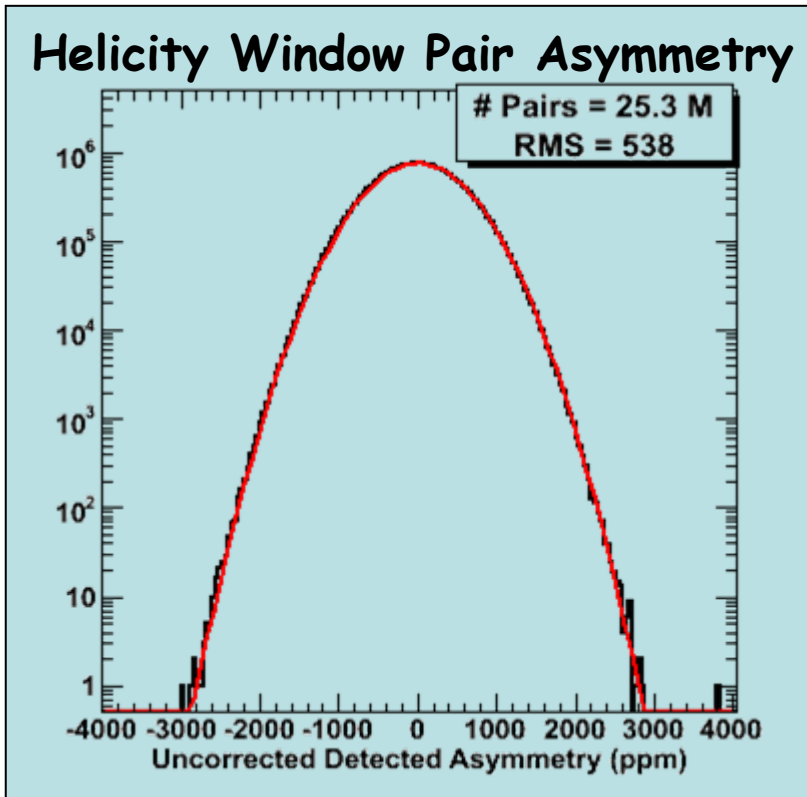


^1H Preliminary Results

Raw Parity Violating Asymmetry

A_{raw} correction ~ 11 ppb

Helicity Window Pair Asymmetry



$$Q^2 = 0.1089 \pm 0.0011 \text{ GeV}^2$$

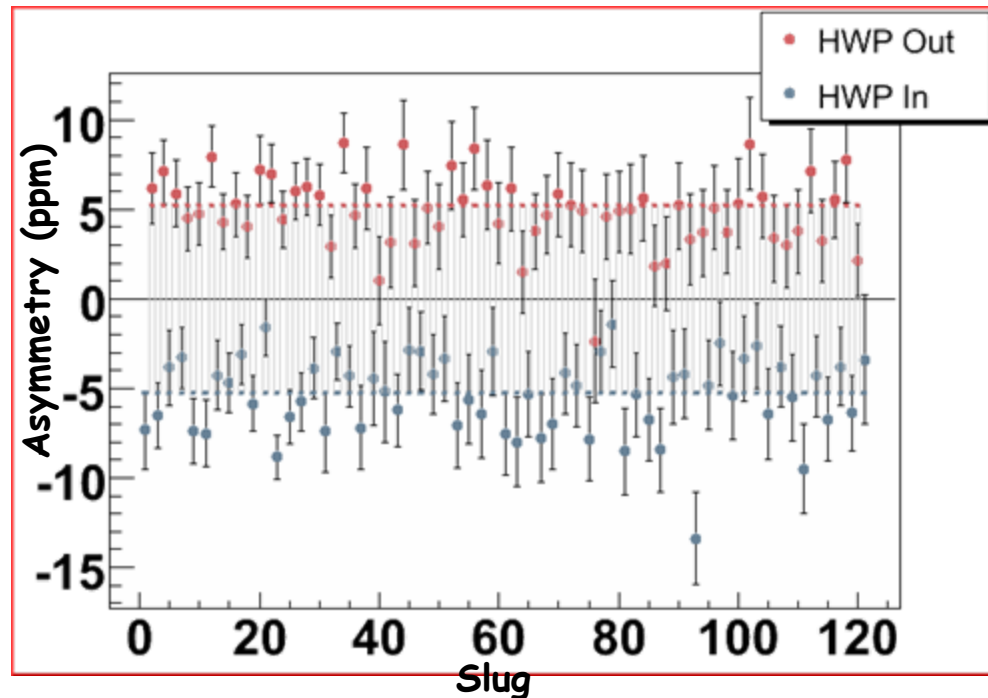
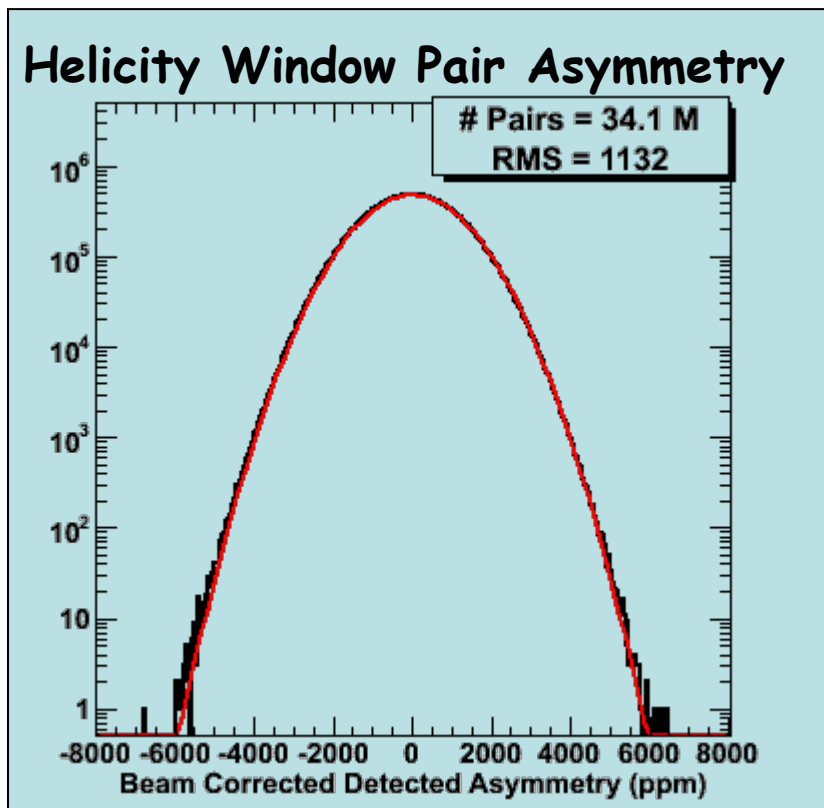
$$A_{\text{raw}} = -1.418 \text{ ppm} \pm 0.105 \text{ ppm (stat)}$$

^4He Preliminary Results

Raw Parity Violating Asymmetry

A_{raw} correction ~ 0.12 ppm

Helicity Window Pair Asymmetry



$$Q^2 = 0.07725 \pm 0.0007 \text{ GeV}^2$$

$$A_{\text{raw}} = 5.253 \text{ ppm} \pm 0.191 \text{ ppm (stat)}$$

HAPPEX-II 2005 Preliminary Results

HAPPEX-⁴He:

$$Q^2 = 0.0772 \pm 0.0007 \text{ (GeV/c)}^2$$
$$A_{PV} = +6.43 \pm 0.23 \text{ (stat)} \pm 0.22 \text{ (syst) ppm}$$

$$A(G^s=0) = +6.37 \text{ ppm}$$

$$G^s_E = 0.004 \pm 0.014_{\text{(stat)}} \pm 0.013_{\text{(syst)}}$$

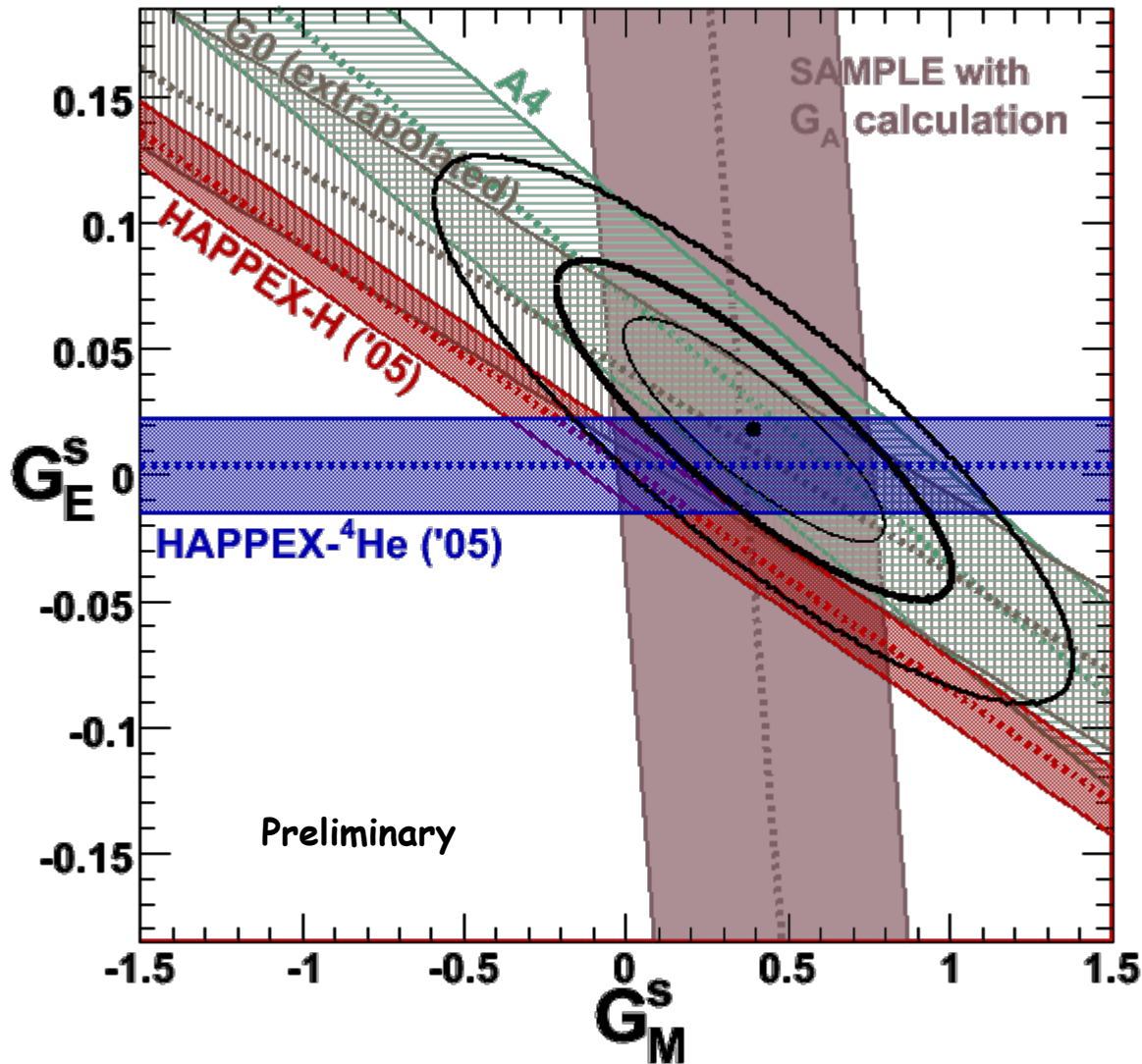
HAPPEX-H:

$$Q^2 = 0.1089 \pm 0.0011 \text{ (GeV/c)}^2$$
$$A_{PV} = -1.60 \pm 0.12 \text{ (stat)} \pm 0.05 \text{ (syst) ppm}$$

$$A(G^s=0) = -1.640 \text{ ppm} \pm 0.041 \text{ ppm}$$

$$G^s_E + 0.088 G^s_M = 0.004 \pm 0.011_{\text{(stat)}} \pm 0.005_{\text{(syst)}} \pm 0.004_{\text{(FF)}}$$

HAPPEX-II 2005 Preliminary Results

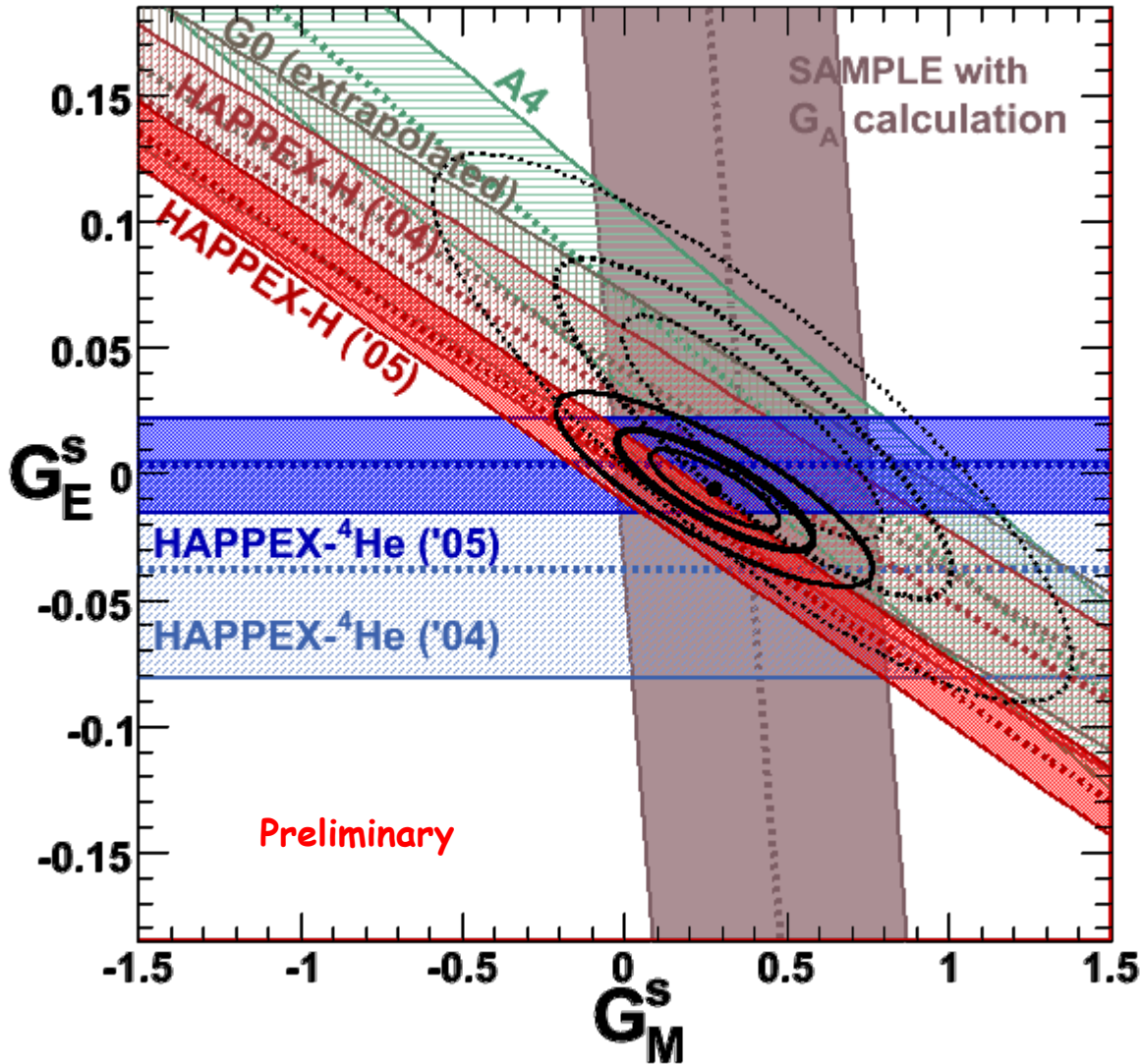


Three bands:

1. Inner: Project to axis for 1-D error bar
2. Middle: 68% probability contour
3. Outer: 95% probability contour

Caution: the combined fit is approximate. Correlated errors and assumptions not taken into account

World Data near $Q^2 \sim 0.1 \text{ GeV}^2$



$$G_M^s = 0.28 \pm 0.20$$

$$G_E^s = -0.006 \pm 0.016$$

$\sim 3\% \pm 2.3\%$ of proton magnetic moment

$\sim 0.2 \pm 0.5\%$ of electric distribution

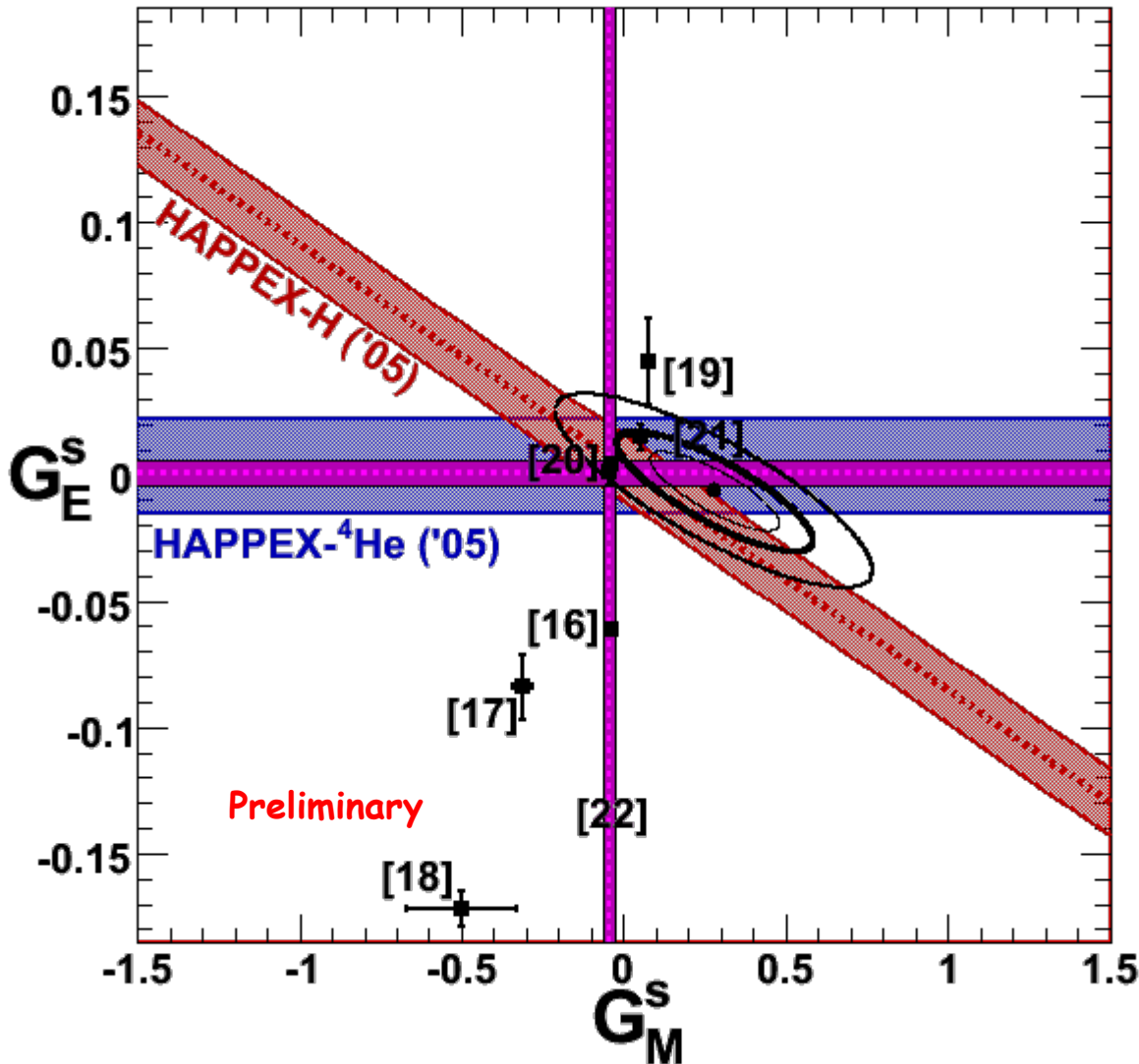
HAPPEX-only fit suggests something even smaller:

$$G_M^s = 0.12 \pm 0.24$$

$$G_E^s = -0.002 \pm 0.017$$

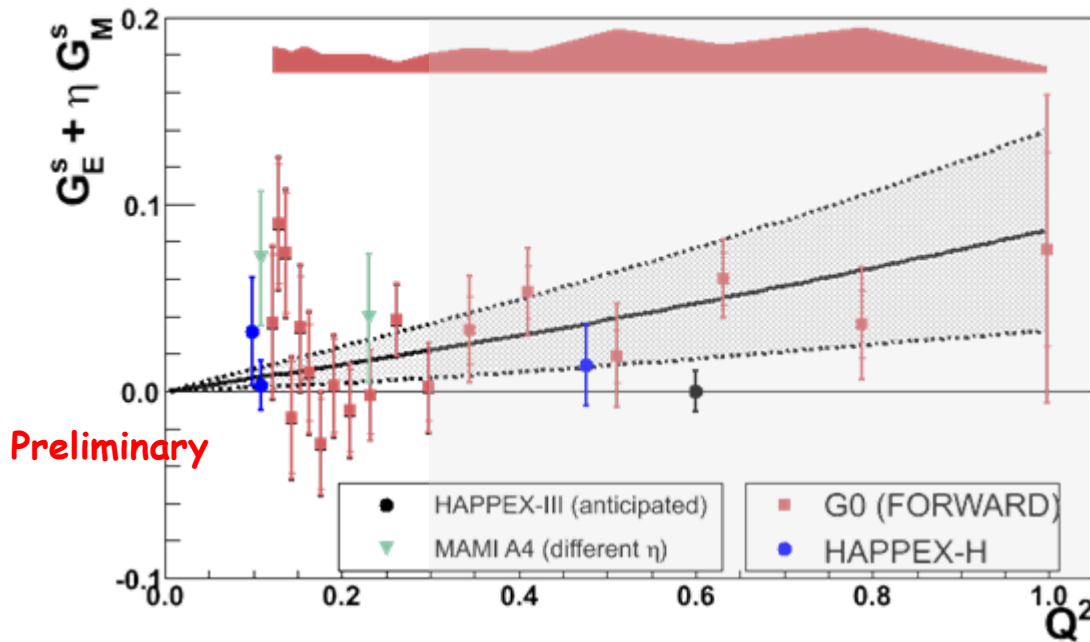
Caution: the combined fit is approximate. Correlated errors and assumptions not taken into account

World data consistent with state of the art theoretical predictions



16. **Skyrme Model** - N.W. Park and H. Weigel, Nucl. Phys. A **451**, 453 (1992).
17. **Dispersion Relation** - H.W. Hammer, U.G. Meissner, D. Drechsel, Phys. Lett. B **367**, 323 (1996).
18. **Dispersion Relation** - H.-W. Hammer and Ramsey-Musolf, Phys. Rev. C **60**, 045204 (1999).
19. **Chiral Quark Soliton Model** - A. Sliva *et al.*, Phys. Rev. D **65**, 014015 (2001).
20. **Perturbative Chiral Quark Model** - V. Lyubovitskij *et al.*, Phys. Rev. C **66**, 055204 (2002).
21. **Lattice** - R. Lewis *et al.*, Phys. Rev. D **67**, 013003 (2003).
22. **Lattice + charge symmetry** - Leinweber *et al.*, Phys. Rev. Lett. **94**, 212001 (2005) & hep-lat/0601025

A Simple Fit (for a simple point)



Simple fit:

$$G_E^s = r_s \tau$$

$$G_M^s = \mu_s$$

Includes only data $Q^2 < 0.3 \text{ GeV}^2$

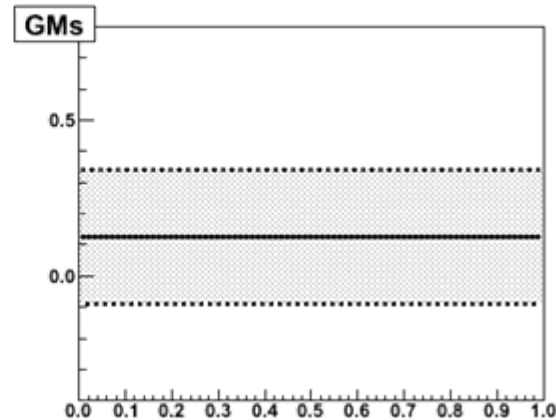
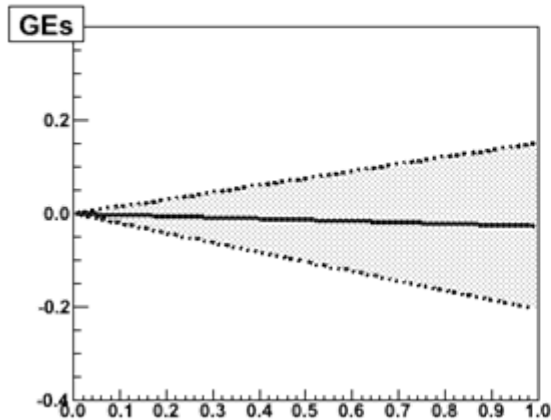
Includes SAMPLE constrained with G_A theory and HAPPEX-He 2004, 2005

G0 Global error allowed to float with unit constraint

Nothing intelligent done with form factors, correlated errors, etc.

Quantitative values should NOT be taken very seriously, but some clear, basic points:

- The world data are consistent.
- Rapid Q^2 dependence of strange form-factors is not required.
- Sizeable contributions at higher Q^2 are not definitively ruled out. (To be tested by HAPPEX-III, G0 and A4 backangle.)



A Global Fit: R.D. Young, et al. nucl-ex/0604010

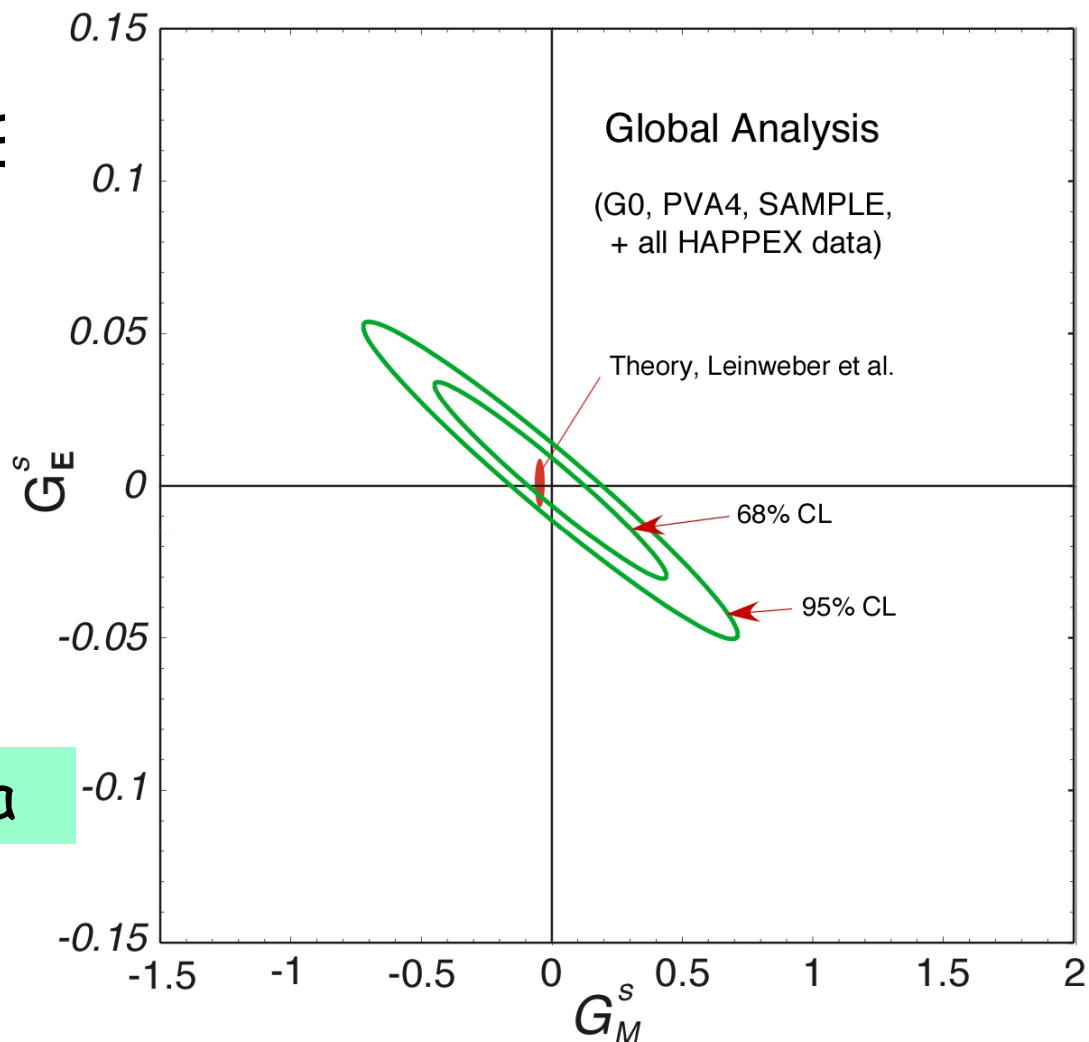
- all data $Q^2 < 0.3$, leading moments of G_E^s , G_M^s

- Reanalysis of SAMPLE nuclear corrections

- Kelly's EMF

- Float G_A^e separately for neutron and proton

With HAPPEX-2005 data



Figures: courtesy of R. Carlini

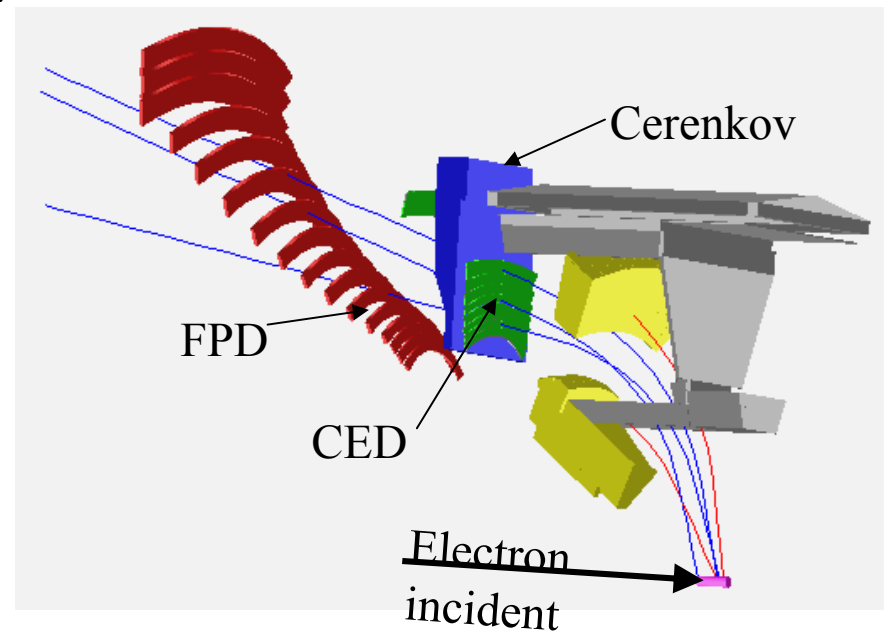
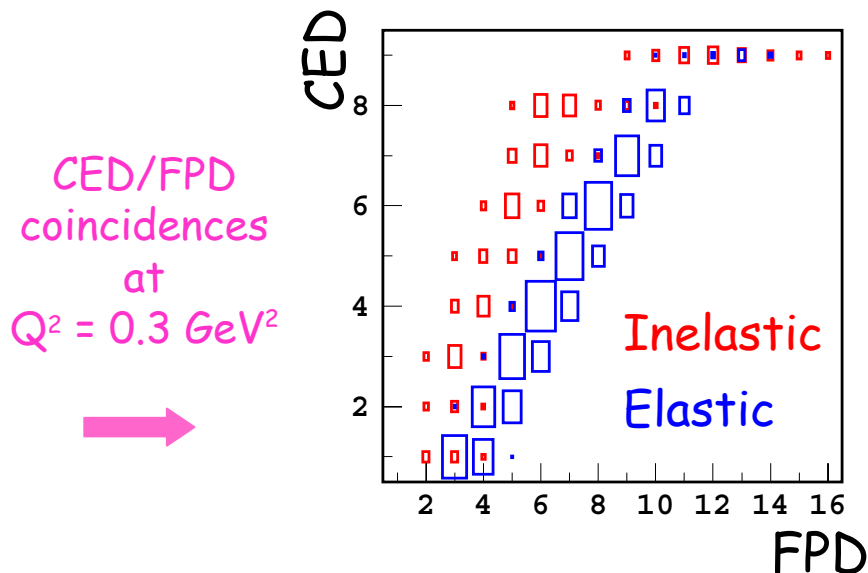
G^0 : Backward Angle

- Detect scattered electrons at $\theta_e \sim 110^\circ$
 - Need separate runs at $E = 362, 687$ MeV
for $Q^2 = 0.23, 0.63$ $(\text{GeV}/c)^2$
for both LH_2 and LD_2 targets
- Get G_M^S and G_A

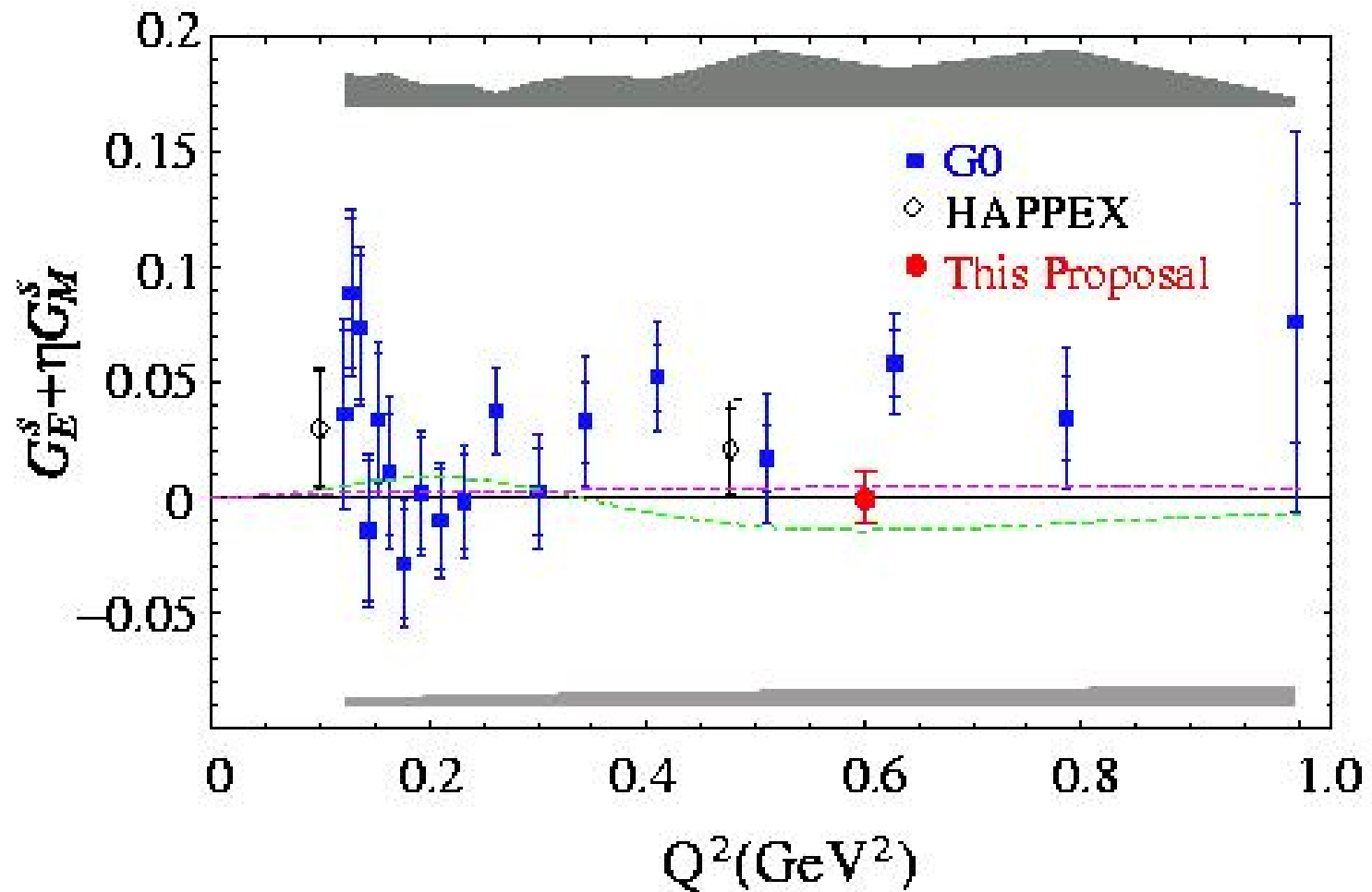
Additional detectors:

- Cryostat Exit Detectors (CED) to separate elastic/inelastic e^-
- Cerenkov detectors for π rejection

First Run - just completed (last weekend)



HAPPEX-III (2008)



Conclusions

- *Marvelous* consistency of data, *esp.* at $Q^2=0.1 \text{ GeV}^2$.
- $Q^2 = 0.1 \text{ GeV}^2$ data: G_M^s and G_E^s consistent with zero; constraining axial FF to Zhu *et al.* theory favors positive G_M^s
- Still room (& hints?) for non-zero values at higher Q^2

Future:

- GO Backward: will allow G_M^s and G_E^s separation at two Q^2
- Mainz: PV-A4 backward-angle program underway
- HAPPEX-III: high precision forward-angle @ $Q^2 = 0.6 \text{ GeV}^2$
- Qweak: Standard Model test at low Q^2 (2009)

Backup Slides

Two Photon Exchange

1. Beyond single boson exchange in electroweak interference:
 - $\gamma\gamma$ and γZ box and crossing diagrams.
 - effects appear small at large ε and small Q^2
 - not a concern at present experimental precision.
2. Electromagnetic Form Factors used to extract strange form factors:
 - which form factors to use?
3. Transverse Asymmetry/Beam normal asymmetry/Vector analyzing power:
 - ☹️ "background" to PV measurements, if electron beam not 100% longitudinal and detectors not perfectly symmetric.
 - 😊 interesting in its own right - imaginary parts of TPE.

Validity of charge symmetry assumption

$$u \leftrightarrow d$$

$$G_{E,M}^{u,p} = G_{E,M}^{d,n}$$

$$G_{E,M}^{d,p} = G_{E,M}^{u,n}$$

$$G_{E,M}^{s,p} = G_{E,M}^{s,n}$$

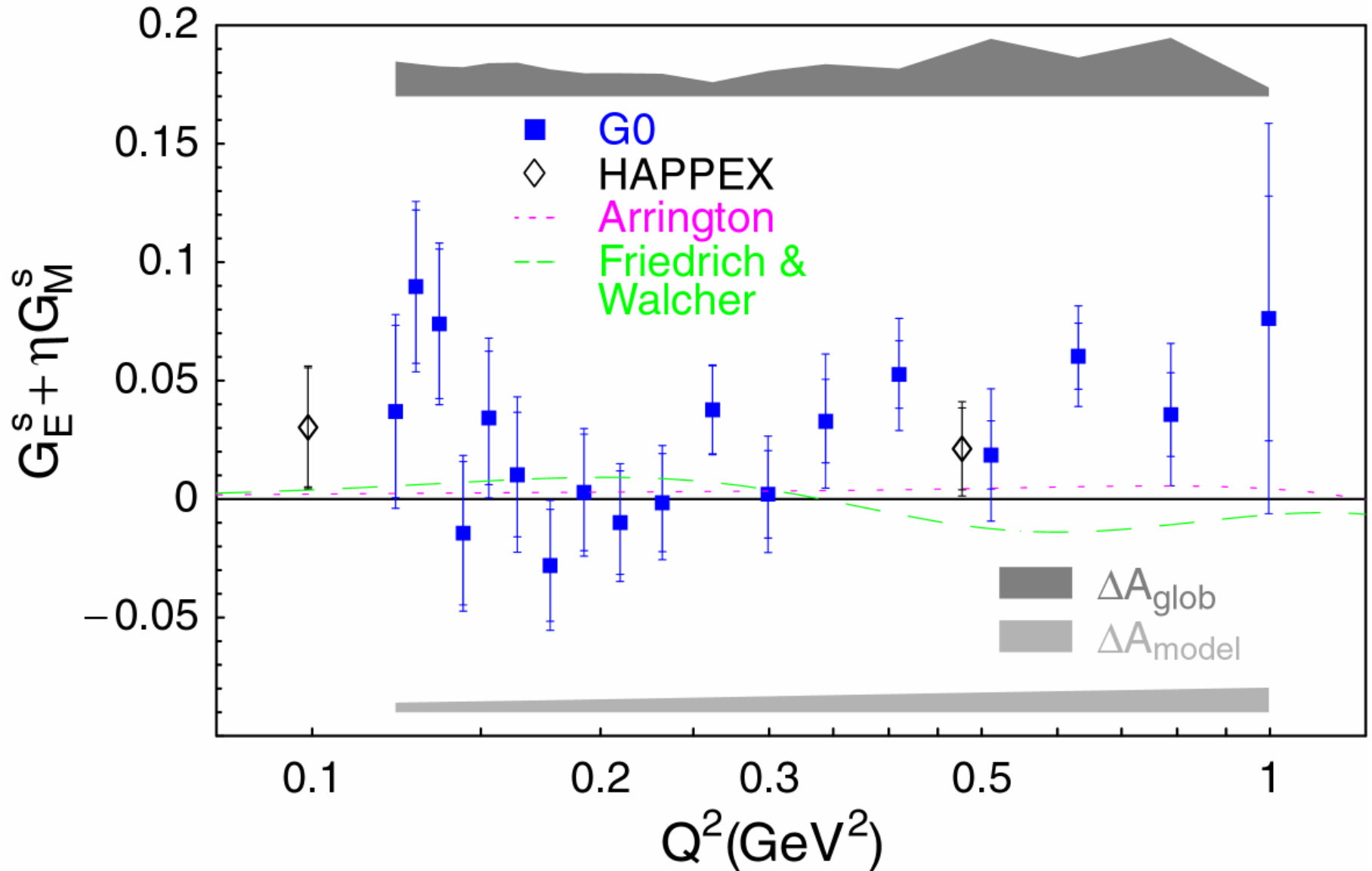
Size of charge symmetry breaking effects in some n,p observables:

- n - p mass difference $\rightarrow (m_n - m_p)/m_n \sim 0.14\%$
- polarized elastic scattering $\vec{n} + p, \vec{p} + n$ $\Delta A = A_n - A_p = (33 \pm 6) \times 10^{-4}$
Vigdor et al, PRC 46, 410 (1992)
- Forward backward asymmetry $n + p \rightarrow d + \pi^0$ $A_{fb} \sim (17 \pm 10) \times 10^{-4}$
Opper et al., nucl-ex 0306027 (2003)

\rightarrow For vector form factors theoretical CSB estimates indicate $< 1\%$ violations (unobservable with currently anticipated uncertainties)

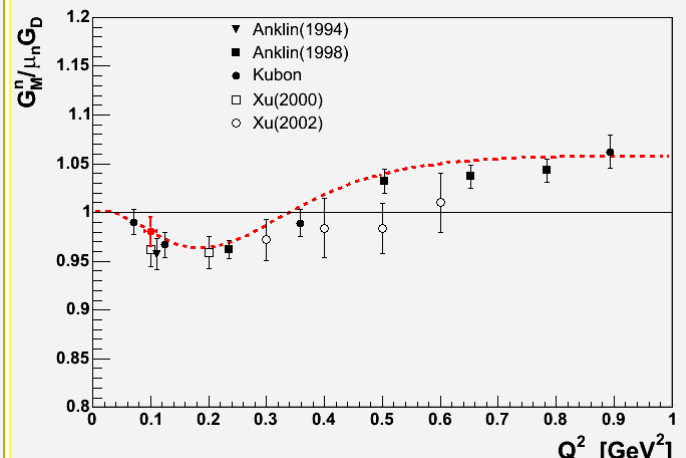
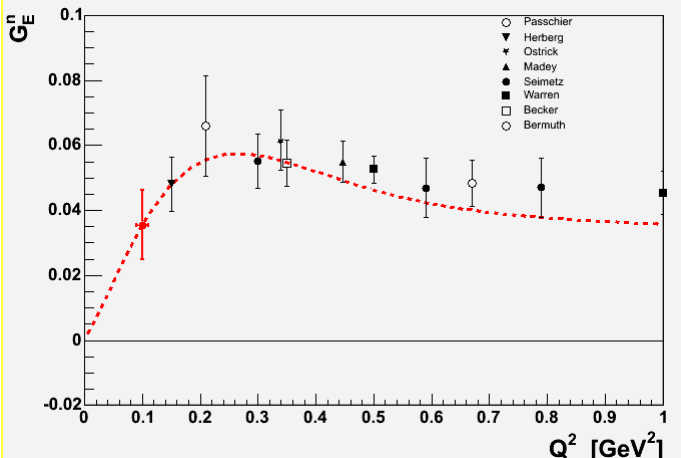
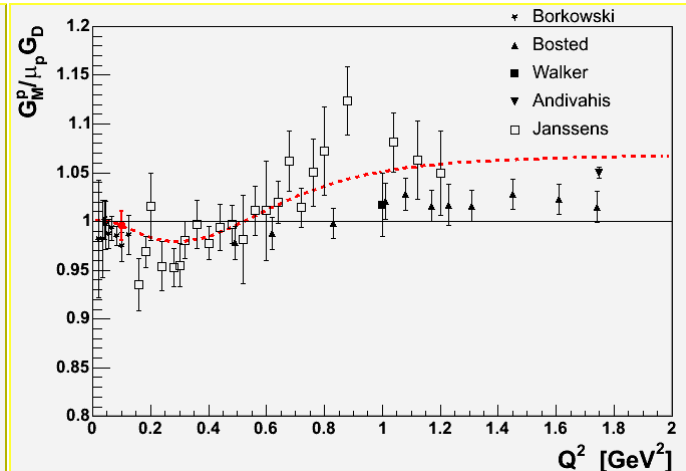
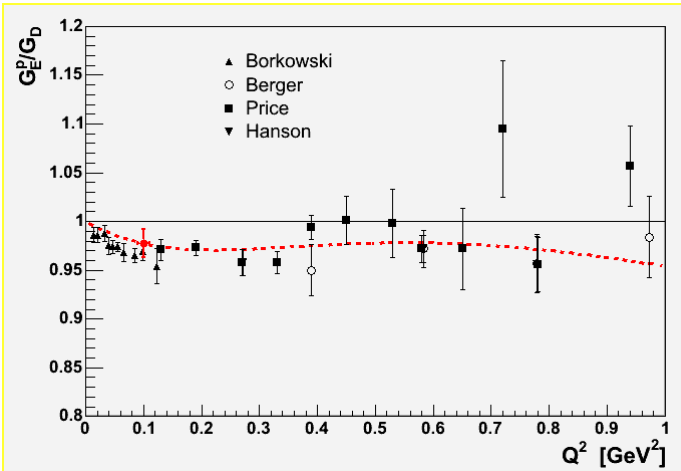
Miller PRC 57, 1492 (1998) Lewis & Moberg, PRD 59, 073002(1999)

Strange Quark Contribution to Proton



EM Form Factors

Electromagnetic form factors parameterized as by:
Friedrich and Walcher, Eur. Phys. J. A, 17, 607 (2003)



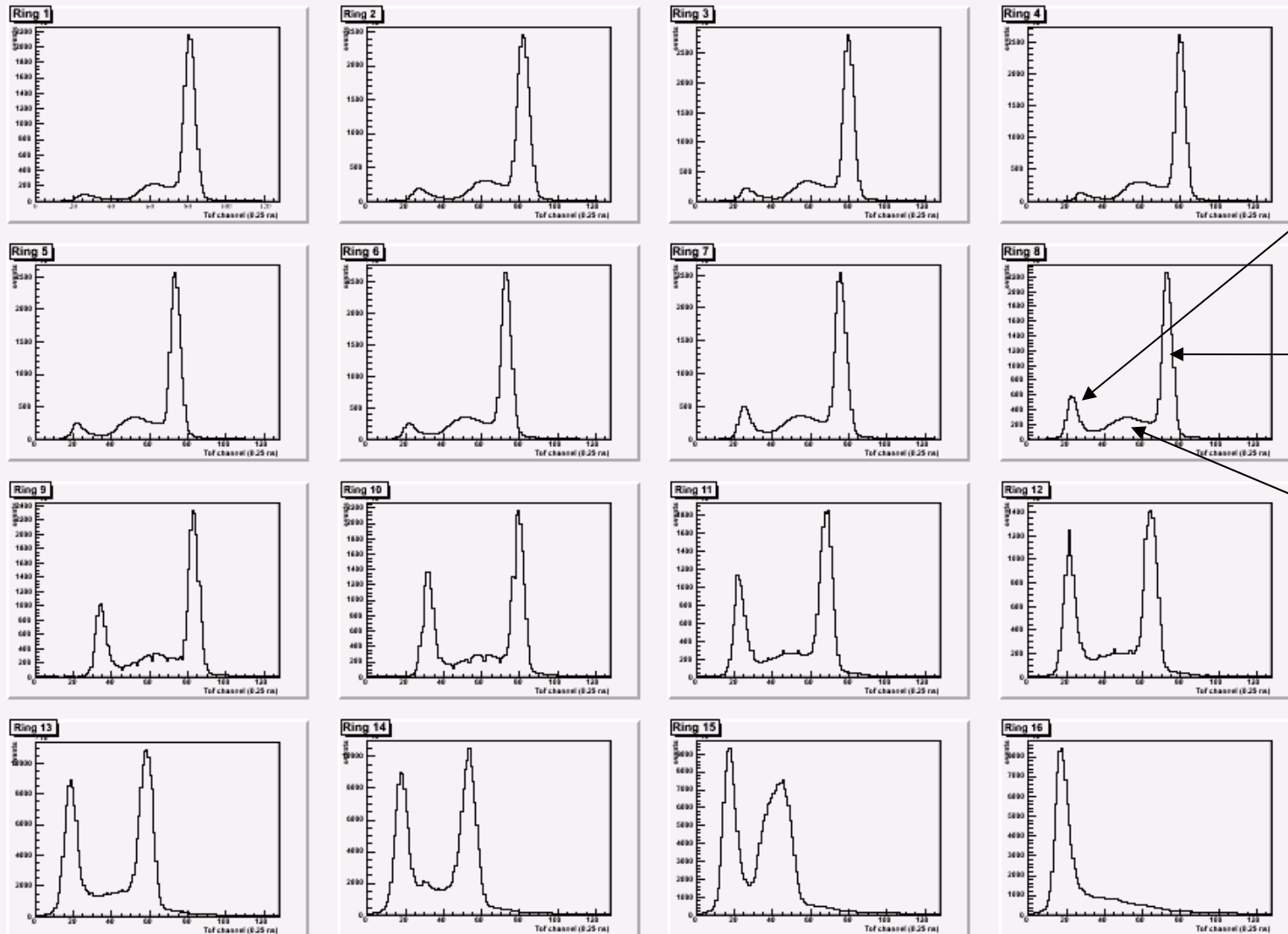
G_E^n from BLAST:
Claimed uncertainty
at 7-8%

FF	Error
G_E^p	2.5%
G_M^p	1.5%
G_E^n	10%
G_M^n	1.5%
$G_A^{(3)}$	-
$G_A^{(8)}$	-

Time of Flight Spectra

Time of flight spectra for all
16 detectors of a single octant
- recorded every 33 ms

Time of Flight Run 22140 Octant 2



pions

elastic
protons
(signal)

inelastic
protons

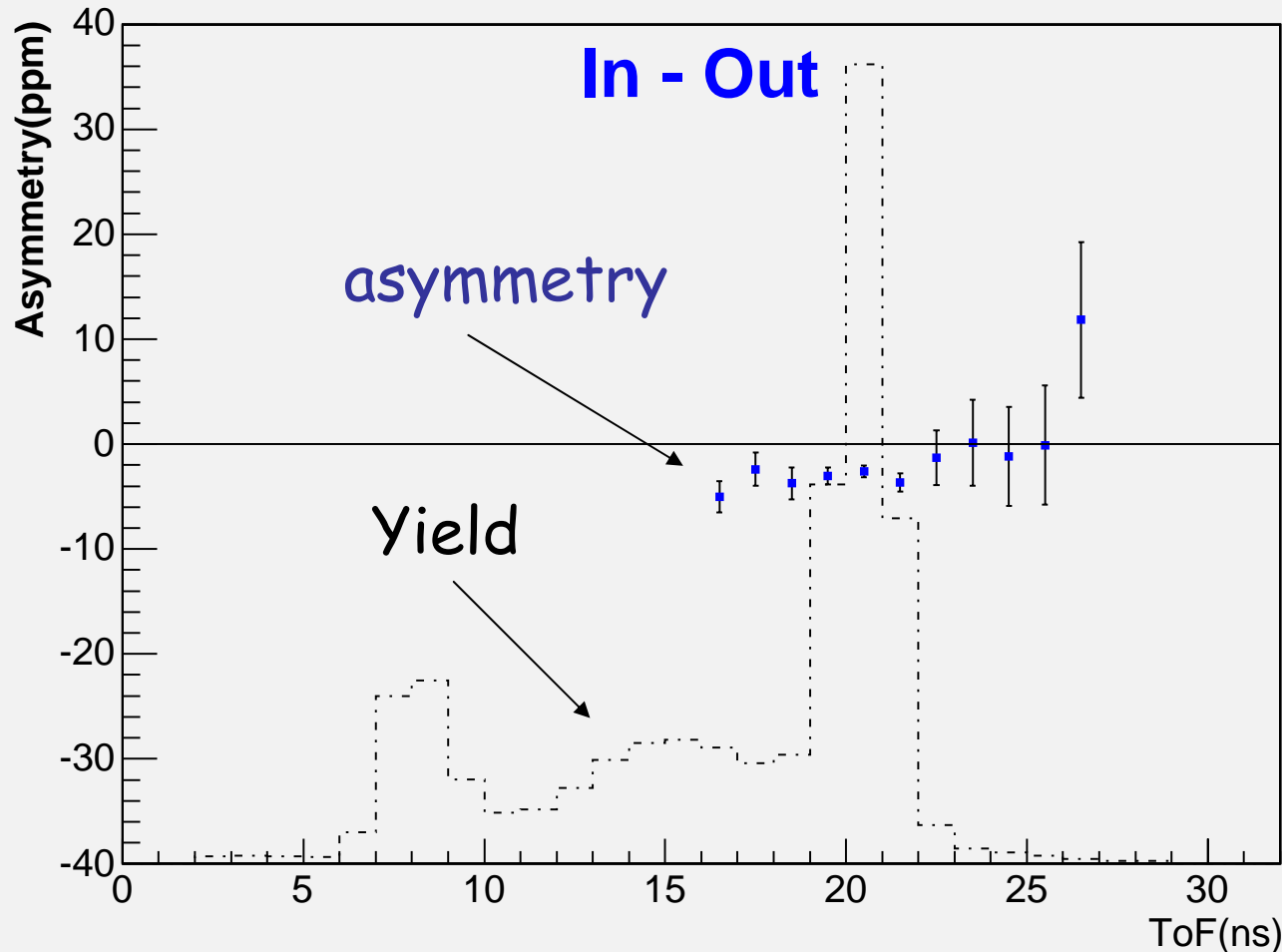
Positive Background Asymmetries

- Det. 12-16 see smoothly varying peak in background asymmetries
 - maximum magnitude $\sim +45$ ppm
- Source is protons from hyperon weak decay scattering inside spectrometer
 - GEANT simulation with generator for hyperon production based on CLAS data
 - simulate both Λ and $\Sigma^{+,0}$ decays
 - polarization transfer for Λ 100%
 - assume 70% for Σ^+
 - Σ^0 asymmetry scaled by further factor of $-1/3$ (CG coefficient)
 - simulation explains source; use measured data for actual analysis

"Side-band" background correction

- Asymmetry and yield measured on either side of elastic peak
- > smooth interpolation is simple

Detector 8 - typical for low to mid Q^2

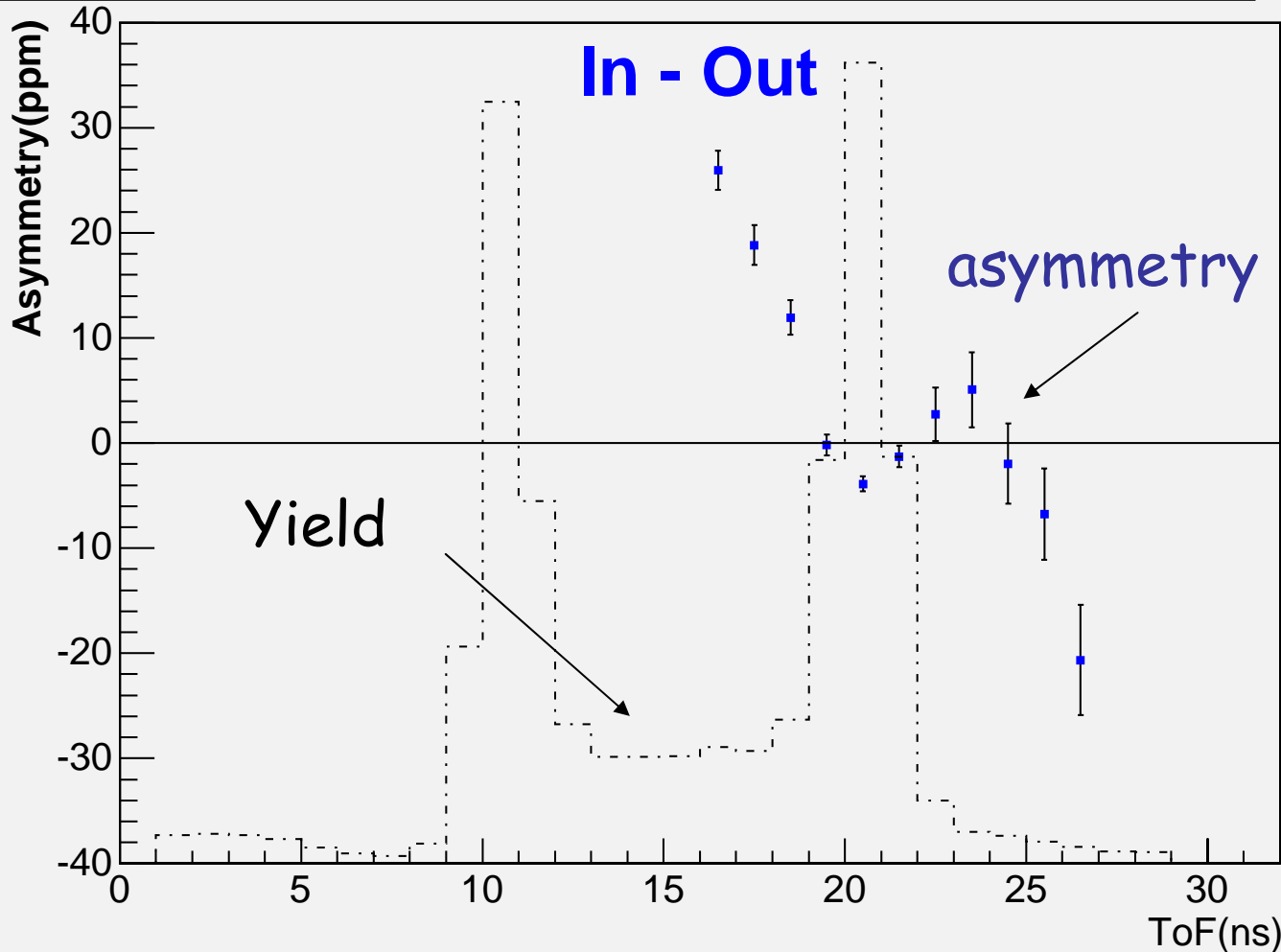


Error in elastic asymmetry due to background is
2% - 5%
for these detectors

"Side-band" background correction @ larger Q^2

- Background asymmetry 'large' & varying significantly under elastic peak

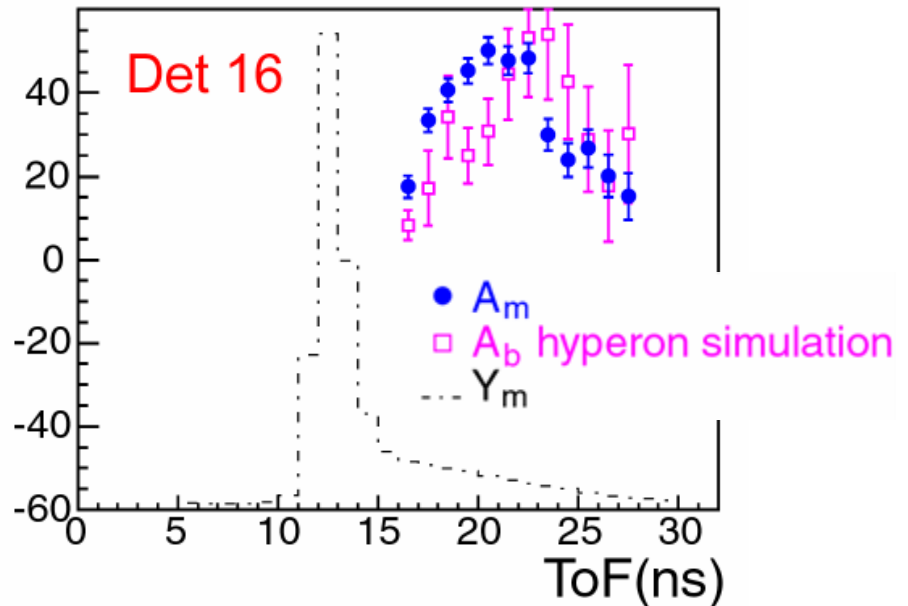
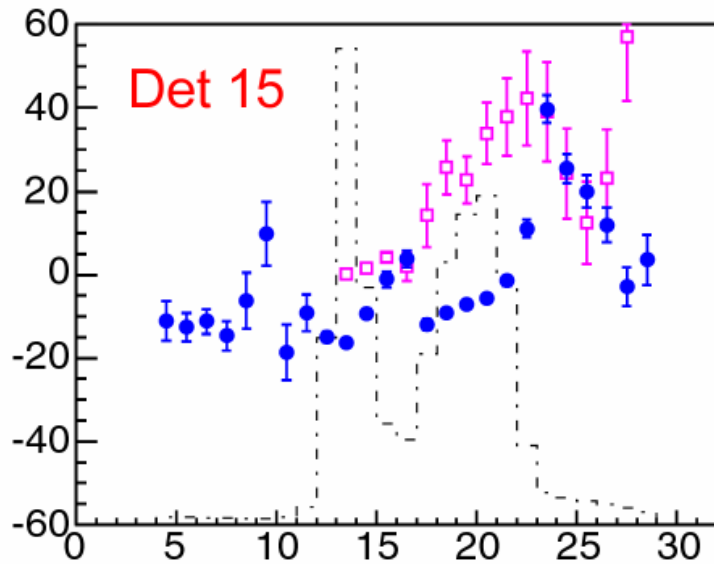
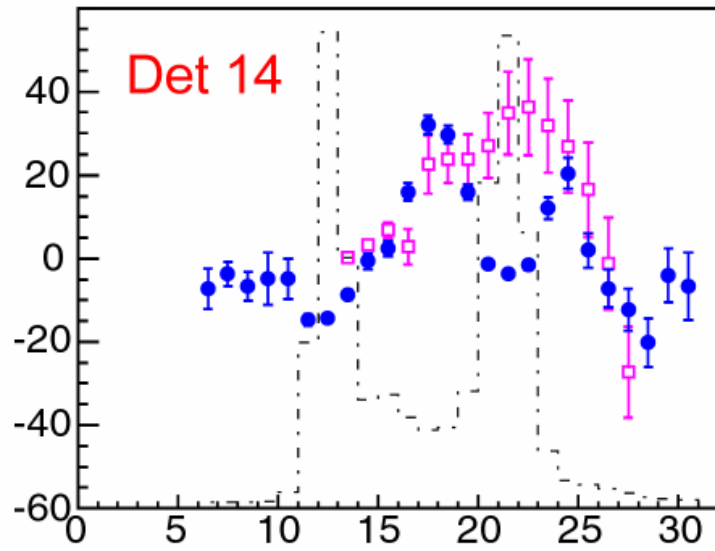
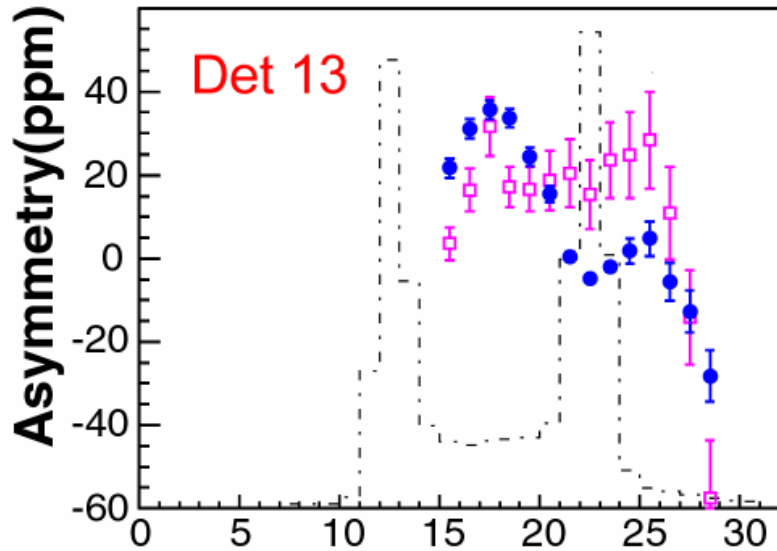
Detector 13 - typical for higher Q^2



Error in elastic asymmetry due to background dominates here

7% - 20% error
(depending on detector)

Positive Background Asymmetries: GEANT



G0: Asymmetry with EW Radiative Corrections

- Full form of asymmetry used to extract $G_E^s + \eta G_M^s$

$$A = -\frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \frac{1}{\varepsilon G_E^{p2} + \tau G_M^{p2}} \left\{ (1 - 4\sin^2 \theta_W) \left(\varepsilon G_E^{p2} + \tau G_M^{p2} \right) (1 + R_V^p) - \right. \\ \left. \left(\varepsilon G_E^p G_E^n + \tau G_M^p G_M^n \right) (1 + R_V^n) - \left(\varepsilon G_E^p G_E^s + \tau G_M^p G_M^s \right) (1 + R_V^{(0)}) - \right. \\ \left. \varepsilon' (1 - 4\sin^2 \theta_W) G_M^p G_A^e \right\}$$

where

$$G_A^e = -G_A^p (1 + R_A^{T=1}) + \left[\frac{1}{2} (3F - D) R_A^{T=0} + \Delta s (1 + R_A^{(0)}) \right] G_A^{dip}$$

and

$$G_A^p = g_A G_A^{dip} = (F + D) G_A^{dip} = \frac{g_A}{(1 + Q^2 / \Lambda_A^2)^2}$$

Simple Fits to World Hydrogen Data

- Fit $G_E^s(Q^2) + \eta(Q^2, E_i) G_M^s(Q^2) = \frac{4\pi\alpha\sqrt{2}}{G_F Q^2} \frac{\varepsilon G_E^{p^2} + \tau G_M^{p^2}}{\varepsilon G_E^p (1 + R_V^{(0)})} (A_{phys} - A_{NVS}(Q^2, E_i))$

with simple forms for G_M^s , G_E^s

$$G_E^s(Q^2) = \frac{c_2 Q^4}{1 + d_1 Q^2 + d_2 Q^4 + d_3 Q^6} \quad \text{à la Kelly}$$

$$G_M^s(Q^2) = \frac{G_M^s(Q^2 = 0)}{\left(1 + Q^2 / \Lambda_M^s\right)^2}$$

with

$$G_M^s(Q^2 = 0) = 0.81 \quad \text{from } Q^2 = 0.1 \text{ GeV}^2 \text{ plot, dipole ff}$$

Error Budget-Helium

2005

False Asymmetries	48 ppb
Polarization	192 ppb
Linearity	58 ppb
Radiative Corrections	6 ppb
Q ² Uncertainty	58 ppb
Al background	32 ppb
Helium quasi-elastic background	24 ppb
Total	216 ppb

2004

False Asymmetries	103 ppb
Polarization	115 ppb
Linearity	78 ppb
Radiative Corrections	7 ppb
Q ² Uncertainty	66 ppb
Al background	14 ppb
Helium quasi-elastic background	86 ppb
Total	205 ppb

Error Budget-Hydrogen

2005

False Asymmetries	17 ppb
Polarization	37 ppb
Linearity	15 ppb
Radiative Corrections	3 ppb
Q ² Uncertainty	16 ppb
Al background	15 ppb
Rescattering Background	4 ppb
Total	49 ppb

2004

False Asymmetries	43 ppb
Polarization	23 ppb
Linearity	15 ppb
Radiative Corrections	7 ppb
Q ² Uncertainty	12 ppb
Al background	16 ppb
Rescattering Background	32 ppb
Total	63 ppb

Determining Q^2

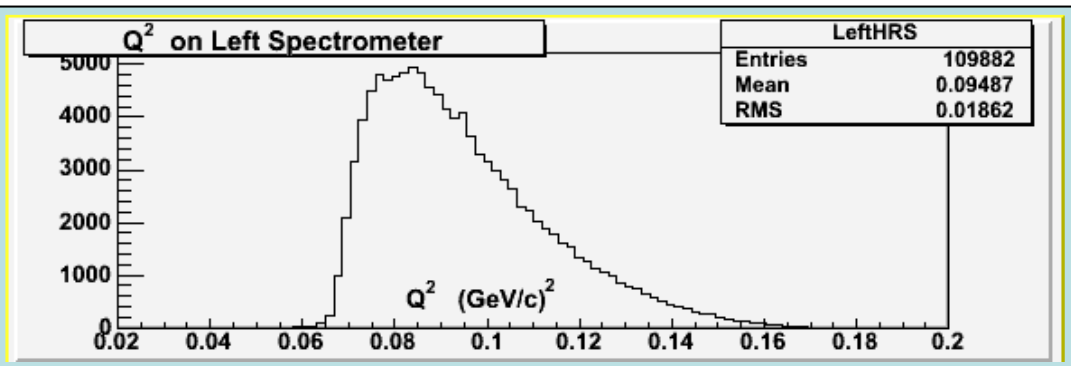
Asymmetry explicitly depends on Q^2 :

$$A_{PV} = \frac{-G_F Q^2}{4\pi\alpha\sqrt{2}} \left\{ \left(1 - 4\sin^2\theta_W\right) - \frac{\varepsilon G_E^p (G_E^n + G_E^s) + \tau G_M^p (G_M^n + G_M^s)}{\varepsilon (G_E^p)^2 + \tau (G_M^p)^2} \right\}$$

$$Q^2 = 2EE'(1 - \cos\theta)$$

Goal: $\delta_{Q^2} < 1\%$

Q^2 measured using standard HRS tracking package, with reduced beam current



- Central scattering angle must be measured to $\delta\theta < 0.5\%$
- Asymmetry distribution must be averaged over finite acceptance