

ORIGIN OF NEUTRINO MASSES AND MIXINGS



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MAJOR THEORETICAL ISSUES IN NEUTRINO PHYSICS

— X —

i) WHY $m_\nu \ll m_{e,u,d} \dots$?

IS $m_\nu \neq 0$ AN EVIDENCE FOR ν_R ?

ii) MAXIMAL MIXING;
ATMOSPH. (+ PERHAPS SOLAR ?)

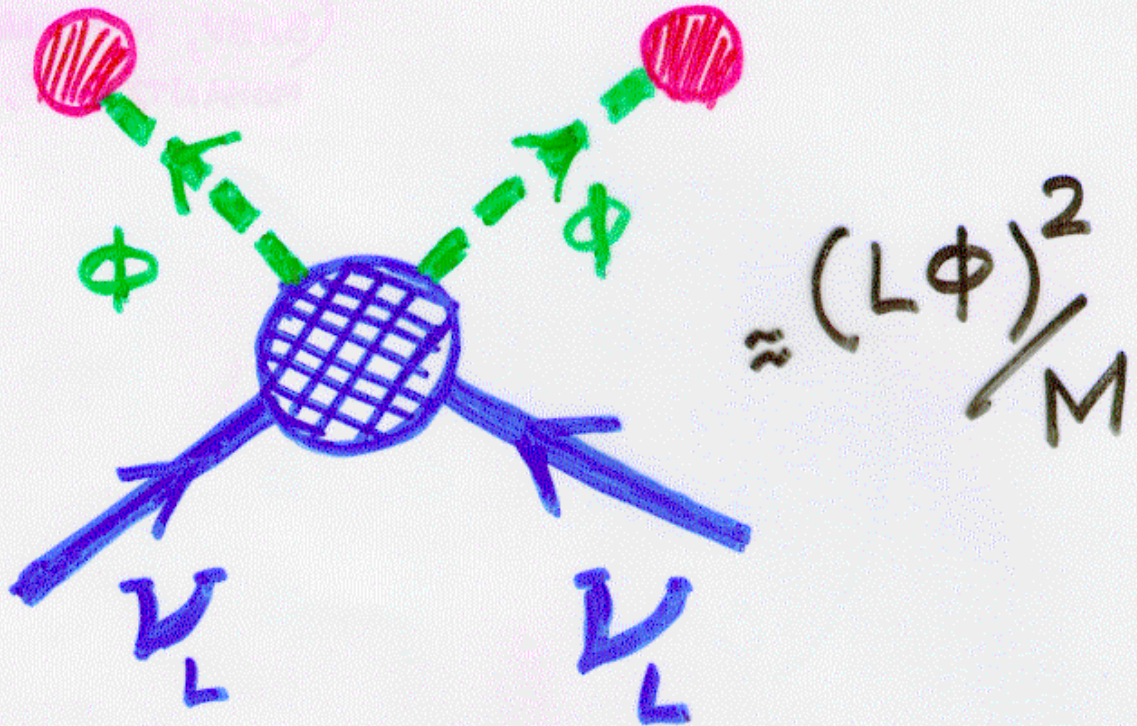
iii) ARE ν 'S DEGENERATE
IN MASS ?

LSND $\Rightarrow m_{\nu_\mu} \approx m_{\nu_\tau}$;

iv) ARE THERE STERILE ν 'S ?
IF SO, WHY ARE THEY LIGHT ?

LSND, SN γ -process etc.

m_ν : (GAUGE INV. \Rightarrow)



• MAJORANA, $m_\nu \sim \frac{f^2 V_{WK}}{M}$

SMALL $m_\nu \Rightarrow$

(i) $M \gg V_{WK}$, $f \sim 1$

(ii) $f \ll 1$, $M \sim V_{WK}$

(OR BOTH)

DOES $m_\nu \neq 0$ IMPLY
A ν_R ?



(SABOHANI, R.N.M.,
1992)

STANDARD MODEL:

$$\begin{pmatrix} u_L \\ d_L \end{pmatrix}$$

$$u_R \quad d_R$$

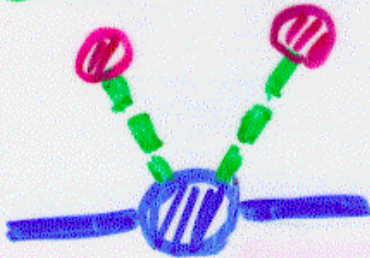
$$\begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$$

$$e_R$$

NO ν_R + B-L GLOBAL SYM

$$\Rightarrow m_\nu = 0$$

BLACK HOLES, WORM HOLES \Rightarrow ~~B-L~~



$$f \sim 1$$

$$M \sim M_{Pl}$$

$$m_\nu \lesssim \frac{V_{WH}^2}{M_{Pl}} \approx 10^{-5} \text{ eV}$$

BARBIERI, ELLIS, GAILLARD
AKHMEDOV, BEREZHIANI, SENJANKOVIĆ

(TOO SMALL !!)

STANDARD MODEL + TRIPLET HIGGS



MA, SARKAR '97

$$m_\nu \approx \frac{V_{WK}^2}{M_\Delta}$$

$$m_\nu \lesssim eV$$

$$\Rightarrow M_\Delta \gtrsim 10^{13} GeV$$

INCOMPATIBLE WITH SUPERSYMMETRY
+ BARYOGENESIS SINCE $T_R < 10^8 GeV$

(DELEPINE, SARKAR)

WEAK SCALE MODELS ($f \ll 1$)

(Zee '80
BABU '87)

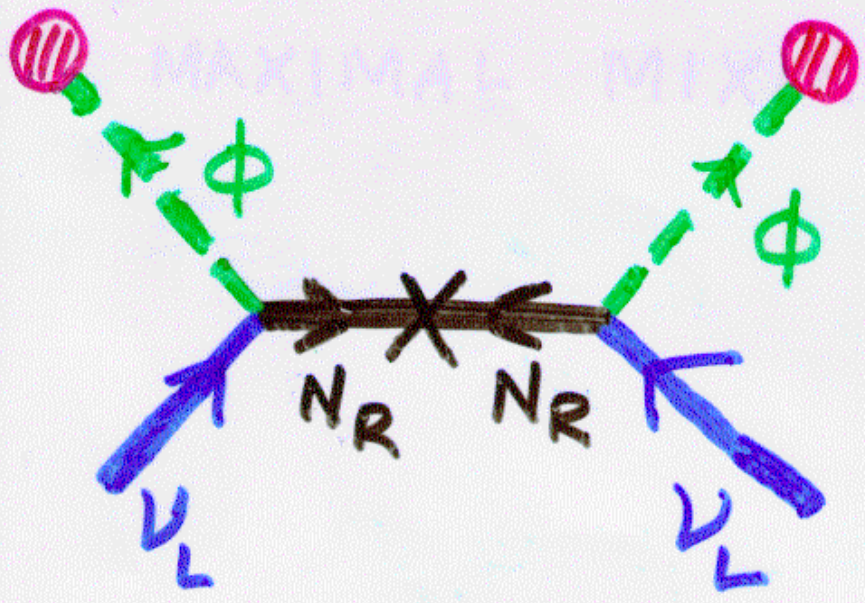


• SMALL DUE TO 2-LO (OR 1-LO)

$$\frac{\eta_B}{\eta_\gamma} = 0 \text{ (PROBLEM)}$$

• NO UNIFICATION.

MODEL WITH ν_R :



$$m_\nu \approx M_\nu^T \frac{1}{M_R} M_\nu \ll m_{l,q} \quad (\text{FOR } M_R \gg v)$$

SEE-SAW MECHANISM
(TYPE I).

Gell-Mann, Ramond, Slansky '79
YANAGIDA '79
R.N.M., SENJANOVIĆ '79

ADVANTAGES OF V_R :

- RESTORES QUARK-LEPTON SYMMETRY
- GAUGE SYM. BECOMES LEFT-RIGHT SYMMETRIC $C SO(10)$

$$SU(2)_{\vec{W}_L} \times SU(2)_{\vec{W}_R} \times U(1)_B^{B-L}$$

$$\begin{pmatrix} u_L \\ d_L \end{pmatrix} \xleftrightarrow{P} \begin{pmatrix} u_R \\ d_R \end{pmatrix}$$

$$\begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \xleftrightarrow{P} \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$$

$$\bullet \quad SU(2)_R \times U(1)_{B-L} \xrightarrow{V_R} U(1)_Y$$

$$Q = I_{3L} + I_{3R} + \frac{B-L}{2}$$

$$\Rightarrow \Delta I_{3R} = \frac{1}{2} \Delta L : M_{NR} \sim V_R$$

m_ν PROBES THE SCALE OF

W_R (PARITY BREAKING)

e.g.

SOLAR $\Rightarrow m_{\nu\mu} \sim 10^{-3} \text{ eV}$

ATMOS. $m_{\nu\tau} \sim 10^{-1} \text{ eV}$

(IF $m_{\nu e} \ll m_{\nu\mu} \ll m_{\nu\tau}$)

$M_{W_R} \gtrsim 10^{12} \text{ GeV}$

$\Rightarrow \frac{\eta_B}{\eta_\gamma} \Rightarrow$ WORKS VERY WELL FOR HIGH SCALES.

$\bullet N_R \rightarrow L\phi \Rightarrow \frac{\eta_B}{\eta_\gamma} \neq 0 \Rightarrow$ SPHALERONS

(FUKUGITA, YANAGIDA '84)

\Downarrow
 $\frac{\eta_B}{\eta_\gamma}$

\Rightarrow "IMPROVES" UNIFICATION OF COUPLINGS \Rightarrow SO(10) GUTS.

MSSM $\Rightarrow \alpha_s(M_Z) \sim 0.13$
EXPT. $\sim 0.119 \pm 0.004$

(CASAS et al. 00 + ...)

TWO TYPES OF SEESAW:

TYPE I:

$$M_\nu = -M_{\nu D}^T M_R^{-1} M_{\nu D}$$

$$M_R \sim f V_R, \quad M_{\nu D} \sim M_{e,u}$$

$$\Rightarrow m_{\nu_1} \ll m_{\nu_2} \ll m_{\nu_3}$$

TYPE II (TRUE IN MOST MODELS)

$$M_\nu = f \frac{V_{wk}^2}{V_R} - M_{\nu D}^T M_R^{-1} M_{\nu D}$$

$$\Rightarrow m_{\nu_1} \sim m_{\nu_2} \sim m_{\nu_3} \quad \text{POSSIBLE !!}$$

$$\text{OR } m_{\nu_2} \sim m_{\nu_3} \quad \text{WITH SAME CP;}$$

NEUTRINO MIXINGS

(MAKI-NAKAGAWA-SAKATA MATRIX)

DEF.

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_\nu \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}; \quad \begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix} = U_e \begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix}$$

WEAK

MASS

WEAK

MASS

$$U_{MNS} = U_e^\dagger U_\nu$$

(i) SMA MSW:

$$U_{MNS} =$$

$$\begin{pmatrix} 1 & \epsilon & 0 \\ \frac{\epsilon}{\sqrt{2}} & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ \frac{\epsilon}{\sqrt{2}} & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \text{ OR } \begin{pmatrix} 1 & \epsilon & 0 \\ -\epsilon & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$



(ii) BIMAXIMAL

$$U_{MNS} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ \frac{1}{2} & \frac{1}{2} & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

VISSANI; 97
 BARGER, PAKVASA, WEILER, WHISNIANT '98
 JEZABEK, SUMNO '98
 BALTA, GOLDHABER, ... '98
 ALTARELLI, FERUGLIO '98
 (FRITESCH, XING)
 (KOIDE, FUSAKA)



(MSW LARGE ANGLE OR ν_0 FOR SOLAR)

KEY QUESTION

• WHY θ_{MNS} SO DIFFERENT

FROM θ_{CKM} ?

• CAN THE SEESAW FORMULA PROVIDE THE ANSWER?

$$U_{MNS} = U_\ell^+ U_\nu$$

- (i) $U_\ell \approx 1$; $U_\nu \Rightarrow$ LARGE θ_ν
- (ii) $U_\nu \approx 1$; $U_\ell \Rightarrow$ LARGE θ_ℓ
- (iii) $U_\nu \approx 1$; $U_\ell \approx 1$; DYNAMICS

"NATURALNESS"

"MY MODEL IS MORE NATURAL
THAN YOURS"

A MORE OBJECTIVE CRITERION:

"ANY RELATION BETWEEN
PARAMETERS IN A THEORY
SHOULD BE GUARANTEED
BY A SYMMETRY"

CASE A:

$$M_\nu = -M_{\nu D}^T M_R^{-1} M_{\nu D}$$

SO(10)
R/LR

$M_L = \text{DIAGONAL}$

$$M_{\nu D} \sim M_e \text{ OR } M_u$$

• GENERIC $M_R \Rightarrow$ SMALL Θ_ν

$\Rightarrow M_R$ MUST HAVE TEXTURE

↑ RESTRICT M_R BY SYMMETRIES

(HORIZONTAL OR $U(1)$ OR ...)

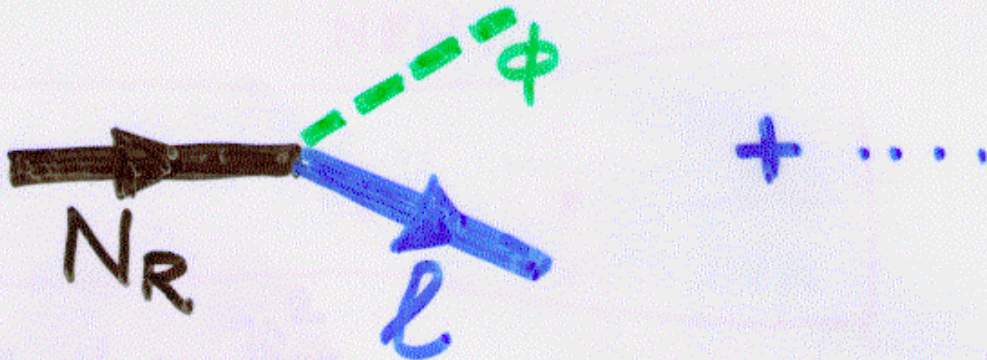
\Rightarrow POSSIBLE TO GET

SMALL U_{e2} + LARGE $U_{\mu 3}$

+ SMALL U_{e3} (CHOOZ, PALO VERDE)

(ARE WE FREE TO CHOOSE ANY M_R ?)

$\frac{n_B}{n_\gamma}$ AND M_R



SAKHAROV CONDITION:

$$\Gamma_{N_R} < H$$

$$f_{ij} = \frac{M_{R,ij}}{v_R} \quad \text{CONSTRAINED!!}$$

e.g.
USUALLY

$$M_{N_1} \ll M_{N_2} \ll M_{N_3} \quad (M_{N_i} < T_{RH})$$

\uparrow N_1 -DECAY $\Rightarrow n_L/n_\gamma$

$$\frac{M_{R,13}}{M_{R,33}} < 10^{-3}; \quad \frac{M_{R,12}}{M_{R,22}} < 10^{-2}$$

3) LOPSIDED UNIFICATION

ALBRIGHT, BABU, BARR '98
 IRGES, ELWOOD, LAVIGNAC, RAMOND
 SATO, YANAGIDA '98
 OKAMURA, TANIMOTO '00
 ALBRIGHT, BARR; BARR, DORSNER

ASSUME : $M_{\nu D}$, M_R GENERIC

BUT M_l PECULIAR !

SU(5) : $\begin{pmatrix} \bar{d}_{R,i} \\ \nu_L \\ e_L \end{pmatrix} \Rightarrow M_e = M_d^T$

$\Rightarrow \theta_l$ RELATED TO $\theta_{q,R}$

↓
(NOT OBSERVED)

\Rightarrow ASYM. $M_{l,q}$

C) DYNAMICALLY GENERATED LARGE MIXINGS : (BALASI, DIGHE, PARIDA, R.N.M. '00)

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AT SEE SAW SCALE,

$\mu = V_R$: BOTH θ_ν & θ_e SMALL.

⇒ IF TWO ν 'S ARE NEARLY DEGENERATE AND SAME CP

⇒ RADIATIVE CORRECTIONS CAN MAGNIFY θ_ν AT M_Z .

WHAT ARE RADIATIVE CORRECTIONS?



BABU, LEUNG, PATALEONE
CHANKOWSKI, PLUCINIEK
ELLIS, LOILA; MA;
IBARA, NAVARO,
HABA et. al., TAKASUGI
et. al.

$$M_\nu(M_Z) = I_{RAD} M_\nu(V_R) I_{RAD}$$

$$I_\alpha = 1 - \epsilon_\alpha$$

CONSIDER 2-GEN. (ν_μ, ν_τ) :

AT $\mu = V_R$, SEESAW SCALE

$$M_{V_R} = \begin{pmatrix} m_0 & \delta \\ \delta & m_0 + \delta' \end{pmatrix}; \quad \delta \ll \epsilon \ll m_0$$

$$\tan 2\theta = \frac{2\delta}{\delta'} \ll 1$$

(AS IN QUARK SECTOR)

EXTRAPOLATE TO WEAK SCALE:

$$M(V_{wk}) \simeq \begin{pmatrix} m_0 & \delta \\ \delta & m_0 + \delta' - \epsilon_\tau m_0 \end{pmatrix}$$

$(\epsilon_\mu \ll \epsilon_\tau)$

$$\Rightarrow \tan 2\bar{\theta} = \frac{2\delta}{\delta' - \epsilon_\tau m_0}$$

IF $m_0 \epsilon_\tau \approx \delta'$, $\bar{\theta}$ MAXIMAL!!

• NO SYM. NEEDED;

• QUARK MIXINGS DO NOT MAGNIFY SINCE $m_b \gg m_s, m_t \gg m_c$

DOES IT REALLY WORK?

YES

MSSM:

$$\varepsilon_z = + \frac{h_z^2}{16\pi^2} \ln \frac{V_R}{V_{WK}} \approx 10^{-5} \tan^2 \beta$$

$$\tan \beta = 10, \quad \Delta m^2 \approx 10^{-3} \text{ eV}^2 \quad (m = 1 \text{ eV})$$

(ATMOSPHERIC VALUE),

• REQUIRES

$$\Delta m_{23}^2 < 0$$

- IMPLEMENTABLE IN GUT MODELS WITH Q-L UNIFICATION
- REQUIRES SUSY WITH LARGER $\tan \beta$: (TEST)
- REQUIRE SMA MSW (TEST)

HOW NATURAL IS IT TO GET "BIMAXIMAL" MIXING ?

$$U_{MNS}(\text{Bimax}) = U_{23}(\theta = \frac{\pi}{4}) U_{12}(\theta = \frac{\pi}{4})$$

$$M_\nu = \begin{pmatrix} \delta & F & F \\ F & A & -A \\ F & -A & A \end{pmatrix}$$

(R.N.M., NUSSINOV
ALTARELLI, FERUGI)

TWO CASES:

AKHMEDOV, BRANCO, REBEL
STECH '00
CHEN, MAHANTHAPPA

(i) $\delta \ll A \ll F$

$$\Delta m_{23}^2 > 0$$

(ii) $\delta \ll F \ll A$

$$\Delta m_{23}^2 < 0$$

GENERAL PROBLEM:

EITHER M_R FINE TUNED OR TOO LARGE

$$f_{13} \Rightarrow \eta_B / \eta_X = 0$$

BOTTOM LINE

i) SEVERAL SIMPLE MECHANISMS
TO UNDERSTAND CASE (i)
($\theta_{e\mu}$ SMALL, $\theta_{\mu\tau}$ LARGE)
EXIST.

ii) NOT EASY TO UNDERSTAND
CASE (ii) NATURALLY
IN SEE-SAW PICTURE
WITH SYM. MASS MATRICES.

- $\Delta m_{23}^2 > 0$ OR < 0 (V-FACTORY)
CAN DISCRIMINATE BETWEEN MODELS

SO(10), m_ν AND τ_{PROTON}

SO(10) - PRIME CANDIDATE FOR
GUT WITH $m_\nu \neq 0$.

Q. CAN INCREASING $\tau(p \rightarrow \nu k)$
LIMIT BY A FACTOR OF TEN
RULE OUT SO(10) MODELS
FOR ν -MASS ?

ANS. **NO!!**

IT WILL RULE OUT ONE OR TWO
OUT OF MANY INTERESTING MODELS

SO(10), m_ν , τ_p :

$$M_\nu = - M_{\nu D}^T M_R^{-1} M_{\nu D}$$

$$\text{SO}(10) : M_{\nu D} \rightarrow M_u$$

$$M_R = ?$$

SEVERAL SOURCES: (ACHIMAN, MERTEK)

$$\mathcal{L}_Y = f_{126} \psi_m \psi_m \Delta_{126} + \frac{f_{16}^{(1)}}{M} \psi_m \bar{\psi}_H \psi_m \bar{\psi}_H$$

$$+ \frac{f_{16}^{(2)}}{M} \psi_m \psi_m \psi_H \psi_H$$

$$M_R \sim f_{126} V_R + \frac{f_{16}^{(1)}}{M} V_R^2 + \frac{f_{16}^{(2)}}{M} V_R^2$$

$$A(p \rightarrow \nu_n k) \sim \frac{f_{16}^{(2)}}{M} V_R$$

ONLY IF $f_{126} = 0$; $f_{16}^{(1)} = 0 \Rightarrow \tau_p, m_\nu$ RELATED
(BABU, PATI, WILCZEK)

ISSUE 4:

• SOLAR, ATMOS. + LSND
 REQUIRE 4th ν_4 ($m_{\nu_4} \lesssim 10^{-3} \text{ eV}$)

• LEP-SLC Z-WIDTH
 $\Rightarrow N_\nu = 3$

$\Rightarrow \nu_4 = \nu_s$

$SU(2)_L \times U(1)_Y$
 SINGLET
 (STERILE ν)



M_ν TOO LARGE



M_{ν_s} TOO LARGE

WHY IS THE STERILE ν_s ULTRALIGHT ?

TWO CLASSES OF MODELS:

A) MIRROR UNIVERSE MODEL

OUR WORLD

W^\pm, Z, γ, \dots

$u, d, e, \nu \dots$

MIRROR WORLD

W^\pm, Z, γ, \dots

u', d', e', ν', \dots

↑ GRAVITY ↑

↙ "STERILE UNIVERSE"

TWO VERSIONS:

FOOT & VOLKAS '95

SYMMETRIC

$$M_{W,Z} = M_{W',Z'}$$

BEREZHIANI, R.N.M. '95

ASYMMETRIC

$$M_{W',Z'} \approx (10-20) M_{W,Z}$$

$$M_{q',l'} \approx (10-20) M_{q,l}$$

BOTH VERSIONS PREDICT ULTRALIGHT STERILE ν 's:

ν_e, ν_μ, ν_τ

LIGHT DUE TO B-L

← GRAVITY →

$\nu_e', \nu_\mu', \nu_\tau'$

LIGHT DUE TO B'-L'

NU-MASSES AND MIXINGS

ARISE FROM PLANCK SCALE EFFECTS (AND ARE SMALL).

OPERATORS

$$\frac{(LH)^2}{M_{Pl}^2}, \quad \frac{LH \cdot L'H'}{M_{Pl}^2}, \quad \frac{(L'H')^2}{M_{Pl}^2} \Rightarrow$$

DEF.

$$\langle H' \rangle / \langle H \rangle \approx \xi$$

$$M_\nu = 10^{-5} \text{ eV.} \begin{pmatrix} \nu_e & \nu_s \\ 1 & \xi \\ \xi & \lambda \xi^2 \end{pmatrix}$$

● $\lambda \sim 1$; $\xi \sim 20 \Rightarrow$ MSW SMALL ANGLE

WHAT IS THE MIRROR UNIVERSE

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REALLY LIKE?

(R.N.M., TEPLITZ
99-00)

ASSUME:

$$V_{\text{NR}}' \sim 10-20 V_{\text{NR}}$$

$$\Lambda_{\text{QCD}}' \sim 10-20 \Lambda_{\text{QCD}}$$

$$\Rightarrow m_{L',q'} \sim 10-20 m_{L,q}$$

$$m_{P',n'} \sim 10-20 m_{P,n} \text{ etc.}$$

MIRROR BARYON FRACTION IN Ω_m

•

$$\frac{\eta_{B'}}{\eta_B} \sim \frac{1}{6}$$

(ASYM. INFLATION)

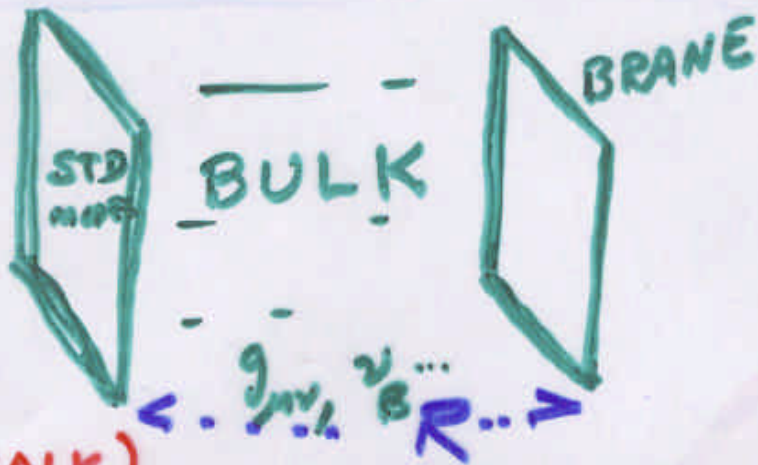
$$\rho_{B'} \approx 3 \rho_B$$

.24 .08

$$\Rightarrow \Omega_{B'} + \Omega_B \approx .32$$

- \Rightarrow MIRROR BARYONS CAN BE CDM.
- \Rightarrow CAN EXPLAIN MICROLENSING ANOMALY
- \Rightarrow RESOLVE HALO CORE DENSITY PUZZLE

ii) SMALL m_{ν_s} DUE TO
LARGE EXTRA DIMENSION



+ TeV
 STRING
 SCALE

(DIENES TALK)

STRING SYMMETRIES \oplus
 GLOBAL B-L OF STD MODE

\Rightarrow NO ARBITRARY ν_B -MASS!

$$\bar{\nu}_L \nu_{B,R} \Rightarrow m_{\nu} \approx \frac{M_{STR}}{M_{PL}} V_{WB} \approx 10^{-4} \text{ eV}$$

∞ -TOWER OF ν_s :

LOWEST STATE: $m_{\nu_s} \sim \frac{1}{R} \sim 10^{-3} \text{ eV}$

(HOW TO HANDLE HIGH m_{ν_s} (CONNECTED TO LARGE EXTRA DIM))

OVERALL APPRAISAL:

- HIGH SCALE SEE-SAW BY FAR THE MOST ATTRACTIVE WAY TO UNDERSTAND SMALL m_ν .

$M_R = \text{SCALE OF PARITY VIOLATION}$

- ATTRACTIVE SEESAW MODELS FOR SMALL $\theta_{e\mu}$ @ LARGE $\theta_{\mu\tau}$

- UNDERSTANDING "BI-MAXIMAL" IN SEESAW STILL A CHALLENGE !!

- LBL EXPTS CAN TEST MODELS WITH $\Delta m_{23}^2 > 0$ OR < 0 .

- STERILE $\nu_s \Rightarrow$ MIRROR UNIVERSE OR EXTRA LARGE DIMENSION ? (TRULY BEYOND CONVENTIONAL GUT'S)