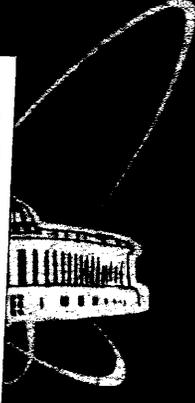
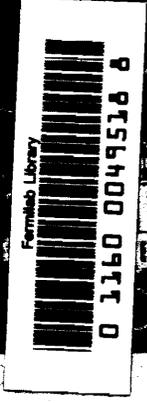


JINR E2-95-226



СООБЩЕНИЯ
ОБЪЕДИНЕННОГО
ИНСТИТУТА
ЯДЕРНЫХ
ИССЛЕДОВАНИЙ

Дубна

JINR E2-95-226

A.I.Titov, Shin Nan Yang*, Yongseok Oh*

POLARIZATION OBSERVABLES
IN ϕ -MESON PHOTOPRODUCTION

FERMILAB

NOV 20 1995

LIBRARY

*Department of Physics, National Taiwan University, Taipei,
Taiwan 10764, Republic of China

AMY

1995

Introduction

Recently, the strangeness content of ordinary nucleons and nuclei has attracted considerable interest [1-4]. This is connected with experimental evidence of nonzero strangeness matrix elements in the proton which comes from the measurement of the pion-nucleon σ -term at low energies [5] and deep-inelastic polarized muon scattering from the proton [6]. At the same time, it has been argued [7] that both measurements could be understood with a little or without strangeness content in the nucleon. Therefore, it is natural to consider other independent tests related to the strangeness content [1,3,8]. One of them is to use the leptonproduction and photoproduction of the ϕ meson from a proton, as it has been suggested by Henley *et al.* [1,2]. The idea is to determine the amount of the $s\bar{s}$ -admixture in the nucleon by isolating the direct knockout contribution to the measured cross section. Since the quark content of ϕ is purely $s\bar{s}$, other processes violating the OZI rule are suppressed. In Ref. [2], the calculation has been carried out in a non-relativistic quark model (NRQM) and the direct knockout contributions have been found to be comparable with the prediction of the vector-meson dominance model (VDM) [9] of diffractive production if a (10-20)% strange quark admixture is assumed.

In Ref. [10], this model has been improved including the Lorentz-contraction effect, where the form factors and other spatial integrals are calculated within the relativistic harmonic oscillator model [11]. The results show that even with a 1 ~ 2% admixture of strange quarks, the direct knockout mechanisms are comparable with the diffractive VDM of ϕ meson electroproduction at large momentum transfers.

In this paper we show that the polarization observables are much more sensitive to the intrinsic content of the strangeness in the proton. As an example, we analyze two of them which can be measured experimentally: double beam-target polarization asymmetries, when the target is polarized normal to the scattering plane and beam-target asymmetry when the target is polarized along the photon momentum.

We denote the four momenta of the photon, initial and final proton, and the produced ϕ meson as q , p , p' , and p_ϕ , respectively. In the laboratory frame, we write $q \equiv (\nu, \mathbf{q})$, $\nu = |\mathbf{q}|$, $p \equiv (E_p, \mathbf{p}) = (M, \mathbf{0})$, where M is the nucleon mass, $p' \equiv (E_{p'}, \mathbf{p}')$, and $p_\phi \equiv (E_\phi, \mathbf{p}_\phi)$. We also represent the four-momentum transfer squared to the proton as $t \equiv (p - p')^2$ and

the photon–proton c.m. energy as $W^2 = (p + q)^2$. The photoproduction cross section is related to the invariant amplitude squared by

$$\frac{d\sigma_{fi}}{dt} = \frac{1}{4\pi} \frac{M^2}{(W^2 - M^2)^2} |T_{fi}|^2. \quad (1)$$

VDM diffractive photoproduction

Diffractive VDM ϕ photoproduction mechanism assumes that the incoming photon mixes into the ϕ meson and then scatters diffractively from the proton through the exchange of a Pomeron [9]. The corresponding amplitude has the form

$$T_{fi}^{\text{VDM}} = T_0^{\text{VDM}} \delta_{m_f, m_i} \hat{\mathbf{e}}_{\lambda_\phi}^* \cdot \hat{\mathbf{e}}_{\lambda_\gamma}, \quad (2)$$

where $m_{i,f}$ are the spin projections of the incoming and outgoing proton, and $\mathbf{e}_{\lambda_\gamma, \lambda_\phi}$ are the polarization vectors of the photon and ϕ meson, respectively. For simplicity, we restrict our consideration to the circularly polarized photon with helicity $\lambda_\gamma = \pm 1$. The amplitude T_0^V does not depend on the spin variables and is related to the unpolarized diffractive VDM photoproduction σ_0^{VDM} by Eq. (1). Using the usual exponential-like form with the parameters of Refs. [2,9] we find

$$T_0^V = \frac{W^2 - M^2}{M} \sqrt{4\pi\sigma_\gamma(W)b_\phi} \exp(-\frac{b_\phi}{2}|t - t_{\text{max}}|), \quad (3)$$

where t_{max} is a possible maximum value of t . The exponential slope b_ϕ and the total photoproduction cross section $\sigma_\gamma(W)$ are determined from experiment [12] as $b_\phi=4.01 \text{ GeV}^2$ and $\sigma_\gamma(W) \simeq 0.2 \mu\text{b}$ at $E_\gamma = 2 \text{ GeV}$. In the VDM diffractive photoproduction, the polarizations of incoming photon and target proton are transferred to the outgoing ϕ meson and recoil proton, respectively.

Knockout photoproduction

The main assumption of the “knockout” photoproduction model is the constituent quark wave function of hadrons with the momentum p_h written in the Fock space of the form [2]:

$$\begin{aligned}
|\phi; p_h\rangle &= |[s\bar{s}]^1, p_h\rangle, \\
|p; p_h\rangle &= \\
&\left\{ A[uud]^{\frac{1}{2}} + B \left(a_0[uud]^{\frac{1}{2}} \otimes [s\bar{s}]^0 \right)^{\frac{1}{2}} + a_1[uud]^{\frac{1}{2}} \otimes [s\bar{s}]^1 \right\}^{\frac{1}{2}}, p_h \rangle, \quad (4)
\end{aligned}$$

where B^2 is the strangeness admixture of the proton and (a_0^2, a_1^2) are the spin 0 and spin 1 fraction of $s\bar{s}$, respectively. It is supposed that the quarks in clusters are in a relative 1s-state with respect to the cluster c.m.-coordinate and $s\bar{s}$ is in a relative p -wave about the uud -cluster, because of the parity. Symbol \otimes in Eq. (4) means the vector addition of the orbital momentum L , ($\ell = 1$) and the momenta of uud and $s\bar{s}$ clusters in a proton. It is assumed that each quark cluster in Eq. (4) is described by a color-singlet spin-flavor wave function. The spatial structure of the hadron is specified by the effective confining quark-quark interaction which in the following is approximated by the relativistic harmonic oscillator potential [11].

Note that the relativistic harmonic oscillator model (RHOM) first considered in Ref. [11] enables one to take into account the Lorentz-contraction of the composite particle wave function. This essential relativistic effect becomes important at large $-q^2$ and provides an explanation of the dipole-type q^2 dependence of the elastic nucleon form factor. Due to this advantage, RHOM has been widely used for describing of the hadron properties at large momentum transfers [11,13,14], despite some theoretical difficulties which are inherent in the model [13].

For the hadron electromagnetic current, there are two ways of its construction: the first is the adoption of the non-relativistic electromagnetic current and the other is the development of some relativistic generalization [11,13,15]. For practical convenience, we use the non-relativistic expression for the hadron electromagnetic current in this calculation remembering that the main relativistic effect is related to the Lorentz-contraction of the hadron intrinsic wave functions. Its possible relativistic modification of Refs. [11,13] does not change the spin structure of the photoproduction amplitudes which determines the qualitative behavior of the polarization observables. Some quantitative difference between the two approaches associated with the relativistic normalization of the hadron center of mass - motion wave functions is not so essential at the initial photon energy $\nu \sim 2$ GeV and is within the accuracy of the model.

Considering interaction of the photon with the $s\bar{s}$ quarks in a proton, one can see that the amplitude does not vanish only when the initial

proton state has the antisymmetric $s\bar{s}$ component with spin $j_{s\bar{s}} = 0$ [2,10]. When the initial photon interacts with the u or d quark in the proton, the $s\bar{s}$ cluster is a spectator so that only the component with $j_{s\bar{s}} = 1$ in the initial wave function contributes. The relevant knockout amplitudes evaluated in the laboratory frame reads

$$T_{fi}^{s\bar{s}} = -T_0^{s\bar{s}} \{ \lambda_\gamma \langle \frac{1}{2}m_f 1 \rho | \frac{1}{2}m_i \rangle \sqrt{4\pi} Y_{1\rho}(\hat{\mathbf{p}}') \hat{\mathbf{e}}_{\lambda_\gamma} \cdot \hat{\mathbf{e}}_{\lambda_\gamma} \}, \quad (5)$$

$$T_{fi}^{uud} = T_0^{uud} \left\{ \sum_{i_c=\frac{1}{2},\frac{3}{2}} \left\langle \frac{1}{2}m' 1 \rho | i_c m_c \right\rangle \left\langle i_c m_c 1 m_\phi | \frac{1}{2}m_i \right\rangle \sqrt{4\pi} Y_{1\rho}^*(\hat{\mathbf{p}}_\phi) \right. \\ \left. \times [\lambda_\gamma \frac{1}{2} \sin \varepsilon \delta_{m',m_f} + \cos \varepsilon \delta_{m',m_f-\lambda_\gamma}] \right\}, \quad (6)$$

where $\cos^2 \varepsilon = (1 + 2(5p' \sin \theta_{p'q} / 3\mu\nu)^2)^{-1}$. Note that $\sin \varepsilon$ in (6) may be neglected because ε is smaller than 10^{-1} over the whole area of the kinematic variables of our interest and decreases fast in the region of $t \rightarrow t_{min}$ where the uud knockout dominates the total amplitude.

The amplitudes $T_0^{s\bar{s}}$ and T_0^{uud} are related to the corresponding single unpolarized knockout cross section as in Eq. (1) and are of the forms

$$T_0^{s\bar{s}} = \left(\frac{8\pi\alpha E_\phi E'_p}{M} \right)^{1/2} A^* B a_0 \frac{\mu_s \nu}{3M} F_{s\bar{s}}(\gamma_\phi, q_{\text{eff}}^{s\bar{s}}) F_{uud}(\gamma_{p'}, 0) V_{s\bar{s}}(\mathbf{p}'), \quad (7)$$

$$T_0^{uud} = \left(\frac{8\pi\alpha E_\phi E'_p}{M} \right)^{1/2} A^* B a_1 \frac{\mu\nu}{2M \cos \varepsilon} \\ \times F_{s\bar{s}}(\gamma_\phi, 0) F_{uud}(\gamma_{p'}, q_{\text{eff}}^{uud}) V_{uud}(\mathbf{p}_\phi), \quad (8)$$

where $\alpha = \frac{e^2}{4\pi}$, $\mu = M/M_{u,d} \simeq 3$, and $\mu_s = M/M_s \simeq 1.88$. Here $M_{u,d,s}$ are the mean energies or ‘‘constituent’’ masses of the corresponding quarks. When RHOM is used (see, for example, Ref. [11]) the spatial integrals are given as

$$F_{s\bar{s}}(\gamma_\phi, q_{\text{eff}}^{s\bar{s}}) = (\gamma_\phi)^{-1} \exp\left(-\frac{r_{s\bar{s}}^2 q_{\text{eff}}^{s\bar{s}2}}{6}\right), \\ F_{uud}(\gamma_{p'}, q_{\text{eff}}^{uud}) = (\gamma_{p'})^{-2} \exp\left(-\frac{r_{uud}^2 q_{\text{eff}}^{uud2}}{6}\right), \\ \frac{1}{(2\pi)^3} V_g(\mathbf{p}) = \frac{v_g(\mathbf{p})}{\int d\mathbf{p} v_g(\mathbf{p})}, \\ v_g(\mathbf{p}) = \mathbf{p}^2 \exp\left\{-\frac{1}{\omega} \left(\mathbf{p}^2 - x_g M \sqrt{\mathbf{p}^2 + M^2}\right)\right\}, \quad (9)$$

where

$$\begin{aligned}\gamma_\phi &= \frac{E_\phi}{M_\phi}, \quad q_{\text{eff}}^{s\bar{s}2} = 2\nu^2 - \frac{\nu}{E_\phi}(\nu^2 - \mathbf{p}'^2 + \mathbf{p}_\phi^2), \\ \gamma_{p'} &= \frac{E_{p'}}{M}, \quad q_{\text{eff}}^{uud2} = 2\nu^2 - \frac{\nu}{E_{p'}}(\nu^2 + \mathbf{p}'^2 - \mathbf{p}_\phi^2),\end{aligned}\quad (10)$$

and the subscript g denotes $s\bar{s}$ or uud with ($x_{s\bar{s}} = \frac{3}{5}$, $x_{uud} = \frac{2}{5}$).

Fixing parameters. The dimensional parameter r_{uud} is chosen so that the proton (uud) magnetic form factors are reproduced up to $Q^2 \simeq 10$ GeV². This gives $r_{uud} = r_p/c = 0.53$ fm, or when the proton radius is used $r_p = 0.83$ fm, the scale factor c is $c = 1.57$. The parameter $r_{s\bar{s}}$ is chosen to be $r_{s\bar{s}} = r_\phi/c$ and $r_\phi = 0.45$ fm, with the same scale factor of the uud cluster. And ω is determined so that the distribution $V_{s\bar{s}}(\mathbf{p})$ reduces to that of NRQM [2] as $|\mathbf{p}| \rightarrow 0$. We perform calculations with different strangeness probabilities: $B^2=0, 0.025, 1\%$, with assuming that $|a_0|=|a_1| = 1/\sqrt{2}$.

Displayed in Fig. 1 are unpolarized ϕ photoproduction cross sections of all the possible processes with $B^2=1\%$ and existing experimental data [12]. The main contribution comes from the diffractive VDM photoproduction. Only at large $|t|$ – near the kinematical limit, where the VDM photoproduction cross section is exponentially small, the uud knockout channel becomes dominant. But eliminating this extreme region we find, that the total unpolarized cross section is near VDM photoproduction and is not sensitive to a small admixture of the strangeness in a proton. The situation, however, changes dramatically for the polarization observables.

Transverse Polarization

Consider first the asymmetry of the double polarized cross section when the target proton and the produced ϕ meson are quantized normal to the scattering plane, i.e., *transverse polarization*. In the ϕ -meson helicity basis, the VDM, $s\bar{s}$, uud knockout amplitudes are written as

$$\begin{aligned}T_{fi}^{\text{VDM}} &= T_0^V \delta_{m_f, m_i} d_{\lambda_\gamma, \lambda_\phi}^{(1)}(\theta_\phi), \\ T_{fi}^{s\bar{s}} &= -T_0^{s\bar{s}} \sqrt{\frac{3}{2}} \langle \frac{1}{2} m_f \ 1 \ \sigma \ | \ \frac{1}{2} m_i \rangle [\sqrt{2} \hat{p}'_\sigma] \lambda_\gamma d_{\lambda_\gamma, \lambda_\phi}^{(1)}(\theta_\phi),\end{aligned}\quad (11)$$

$$T_{fi}^{uud} = T_0^{uud} \sqrt{\frac{3}{2}} \sum_{i_c, m_c, m_\phi} \langle \frac{1}{2} m' 1 \sigma | i_c m_c \rangle \langle i_c m_c 1 m_\phi | \frac{1}{2} m_i \rangle [\sqrt{2} \hat{p}_{\phi, \sigma}^*] \\ \times \delta_{m', m_f - \lambda_\gamma} d_{m_\phi, \lambda_\phi}^{(1)}(\theta_\phi^*). \quad (12)$$

where

$$\hat{p}'_\sigma = -\frac{\sigma}{\sqrt{2}} \exp(-i\sigma\theta'_p), \quad \hat{p}_{\phi, \sigma} = -\frac{\sigma}{\sqrt{2}} \exp(-i\sigma\theta_\phi), \quad \theta_\phi^* = \frac{\pi}{2}, \quad (13)$$

and $\sigma = \pm 1, 0$. To express the polarized amplitudes, we use the following notation:

$$T_{b,v}^{t,r}, \quad (14)$$

where t, r, b , and v represent the polarization of the target, recoil, beam, and vector meson, respectively. After some exercises, we find

$$\begin{aligned} T_{+\lambda_\phi}^{++} &= T_0^V d_{1,\lambda_\phi}^1(\theta_\phi) + 0 && + T_0^{uud} \sqrt{2} d_1 e^{i\theta_\phi} d_{0,\lambda_\phi}^1(\theta_\phi), \\ T_{-\lambda_\phi}^{++} &= T_0^V d_{-1,\lambda_\phi}^1(\theta_\phi) + 0 && + 0, \\ T_{+\lambda_\phi}^{+-} &= 0 && - T_0^{s\bar{s}} e^{-i\theta'_p} d_{1,\lambda_\phi}^1(\theta_\phi) + 0, \\ T_{-\lambda_\phi}^{+-} &= 0 && + T_0^{s\bar{s}} e^{-i\theta'_p} d_{-1,\lambda_\phi}^1(\theta_\phi) - T_0^{uud} f(\theta_\phi), \\ T_{+\lambda_\phi}^{-+} &= 0 && - T_0^{s\bar{s}} e^{i\theta'_p} d_{1,\lambda_\phi}^1(\theta_\phi) + T_0^{uud} f^*(\theta_\phi), \\ T_{-\lambda_\phi}^{-+} &= 0 && + T_0^{s\bar{s}} e^{i\theta'_p} d_{-1,\lambda_\phi}^1(\theta_\phi) + 0, \\ T_{+\lambda_\phi}^{--} &= T_0^V d_{1,\lambda_\phi}^1(\theta_\phi) + 0 && + 0, \\ T_{-\lambda_\phi}^{--} &= T_0^V d_{-1,\lambda_\phi}^1(\theta_\phi) + 0 && - T_0^{uud} \sqrt{2} d_1 e^{-i\theta_\phi} d_{0,\lambda_\phi}^1(\theta_\phi), \end{aligned} \quad (15)$$

where $d_{1,2,3,4}$ are defined as

$$d_1 = \frac{1 + \sqrt{2}}{\sqrt{12}}, \quad d_2 = \frac{\sqrt{3}}{2}, \quad d_3 = \frac{2\sqrt{2} - 1}{\sqrt{12}}, \quad d_4 = \frac{\sqrt{2} - 1}{\sqrt{6}}. \quad (16)$$

and the function f is

$$f(\theta_\phi) = [d_2 e^{i\theta_\phi} d_{-1,\lambda_\phi}^1(\theta_\phi) + d_3 e^{-i\theta_\phi} d_{1,\lambda_\phi}^1(\theta_\phi)]. \quad (17)$$

Here we use the rotation matrix elements $d_{m', m}^1$ as in textbook [16].

$m'(\downarrow)m(\rightarrow)$	+1	0	-1	(18)
+1	$\frac{1}{2}(1 + \cos \theta)$	$\frac{1}{\sqrt{2}} \sin \theta$	$\frac{1}{2}(1 - \cos \theta)$	
0	$-\frac{1}{\sqrt{2}} \sin \theta$	$\cos \theta$	$\frac{1}{\sqrt{2}} \sin \theta$	
-1	$\frac{1}{2}(1 - \cos \theta)$	$-\frac{1}{\sqrt{2}} \sin \theta$	$\frac{1}{2}(1 + \cos \theta)$	

Also for $d_{m'm}^{1/2}$, we have

$$\begin{array}{c|cc} m'(\downarrow)m(\rightarrow) & +1/2 & -1/2 \\ \hline +1/2 & \cos \frac{\theta}{2} & \sin \frac{\theta}{2} \\ -1/2 & -\sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{array} \quad (19)$$

Beam-Vector-meson Asymmetry \mathcal{T}_{BV}

The beam-vector-meson asymmetry \mathcal{T}_{BV} is defined as

$$\mathcal{T}_{BV} \equiv \frac{|T_{++}^{uu}|^2 - |T_{+-}^{uu}|^2}{|T_{++}^{uu}|^2 + |T_{+-}^{uu}|^2}, \quad (20)$$

where

$$|T_{\pm\pm}^{uu}|^2 \equiv |T_{\pm\pm}^{++}|^2 + |T_{\pm\pm}^{+-}|^2 + |T_{\pm\pm}^{-+}|^2 + |T_{\pm\pm}^{--}|^2. \quad (21)$$

Using the expressions for the amplitudes, we have

$$\begin{aligned} |T_{++}^{uu}|^2 &= \frac{1}{2} \{ (T_0^V)^2 + (T_0^{s\bar{s}})^2 \} (1 + \cos \theta_\phi)^2 \\ &+ \frac{1}{2} (T_0^{uud})^2 \{ 1 + d_1^2 + (3d_1^2 - 1) \cos 2\theta_\phi \} \\ &- T_0^V T_0^{uud} d_1 (1 + \cos \theta_\phi) \cos \theta_\phi \\ &- \frac{1}{2} T_0^{s\bar{s}} T_0^{uud} (1 + \cos \theta_\phi) \{ d_2 \cos(\theta_{p'} + \theta_\phi) + d_3 \cos(\theta_{p'} - \theta_\phi) \}, \quad (22) \end{aligned}$$

and

$$\begin{aligned} |T_{+-}^{uu}|^2 &= \frac{1}{2} \{ (T_0^V)^2 + (T_0^{s\bar{s}})^2 \} (1 - \cos \theta_\phi)^2 \\ &+ \frac{1}{2} (T_0^{uud})^2 \{ 1 + d_1^2 + (3d_1^2 - 1) \cos 2\theta_\phi \} \\ &+ T_0^V T_0^{uud} d_1 (1 - \cos \theta_\phi) \cos \theta_\phi \\ &- \frac{1}{2} T_0^{s\bar{s}} T_0^{uud} (1 - \cos \theta_\phi) \{ d_2 \cos(\theta_{p'} + \theta_\phi) + d_3 \cos(\theta_{p'} - \theta_\phi) \}. \quad (23) \end{aligned}$$

Beam-Target Asymmetry \mathcal{T}_{BT}

The beam-target asymmetry \mathcal{T}_{BT} is defined as

$$\mathcal{T}_{BT} \equiv \frac{|T_{+u}^{+u}|^2 - |T_{-u}^{+u}|^2}{|T_{+u}^{+u}|^2 + |T_{-u}^{+u}|^2}, \quad (24)$$

where

$$|T_{\pm u}^{+u}|^2 \equiv |T_{\pm\mp}^{++}|^2 + |T_{\pm 0}^{++}|^2 + |T_{\pm\pm}^{+-}|^2 + |T_{\pm\pm}^{+-}|^2 + |T_{\pm 0}^{+-}|^2 + |T_{\pm\mp}^{+-}|^2. \quad (25)$$

Again, we have

$$|T_{+u}^{+u}|^2 = (T_0^V)^2 + (T_0^{s\bar{s}})^2 + (T_0^{uud})^2 2d_1^2 - T_0^V T_0^{uud} 2d_1 \cos^2 \theta_\phi, \quad (26)$$

$$\begin{aligned} |T_{-u}^{+u}|^2 &= (T_0^V)^2 + (T_0^{s\bar{s}})^2 + (T_0^{uud})^2 (d_2^2 + d_3^2) \\ &\quad - T_0^{s\bar{s}} T_0^{uud} \{(1 + \sin \theta_\phi) d_2 \cos(\theta_{p'} + \theta_\phi) \\ &\quad + (1 - \sin \theta_\phi) d_3 \cos(\theta_{p'} - \theta_\phi)\}. \end{aligned} \quad (27)$$

Target-Vector-meson Asymmetry \mathcal{T}_{TV}

The target-vector-meson asymmetry \mathcal{T}_{TV} is defined as

$$\mathcal{T}_{TV} \equiv \frac{|T_{u+}^{+u}|^2 - |T_{u-}^{+u}|^2}{|T_{u+}^{+u}|^2 + |T_{u-}^{+u}|^2}, \quad (28)$$

where

$$|T_{u\pm}^{+u}|^2 \equiv |T_{+\pm}^{++}|^2 + |T_{+\pm}^{+-}|^2 + |T_{-\pm}^{++}|^2 + |T_{-\pm}^{+-}|^2. \quad (29)$$

And we get

$$\begin{aligned} |T_{u+}^{+u}|^2 &= \frac{1}{2} \{(T_0^V)^2 + (T_0^{s\bar{s}})^2\} (1 + \cos^2 \theta_\phi) \\ &\quad + \frac{1}{2} (T_0^{uud})^2 \{1 + d_1^2 + (3d_1^2 - 1) \cos 2\theta_\phi\} \\ &\quad - T_0^V T_0^{uud} d_1 (1 + \cos \theta_\phi) \cos \theta_\phi \\ &\quad - \frac{1}{2} T_0^{s\bar{s}} T_0^{uud} (1 - \cos \theta_\phi) \{d_2 \cos(\theta_{p'} + \theta_\phi) + d_3 \cos(\theta_{p'} - \theta_\phi)\}, \end{aligned} \quad (30)$$

and

$$\begin{aligned} |T_{u-}^{+u}|^2 &= \frac{1}{2} \{(T_0^V)^2 + (T_0^{s\bar{s}})^2\} (1 + \cos^2 \theta_\phi) \\ &\quad + \frac{1}{2} (T_0^{uud})^2 \{1 + d_1^2 + (3d_1^2 - 1) \cos 2\theta_\phi\} \\ &\quad + T_0^V T_0^{uud} d_1 (1 - \cos \theta_\phi) \cos \theta_\phi \\ &\quad - \frac{1}{2} T_0^{s\bar{s}} T_0^{uud} (1 + \cos \theta_\phi) \{d_2 \cos(\theta_{p'} + \theta_\phi) + d_3 \cos(\theta_{p'} - \theta_\phi)\}. \end{aligned} \quad (31)$$

Target-Recoil Asymmetry \mathcal{T}_{TR}

The target-recoil asymmetry \mathcal{T}_{TR} is defined as

$$\mathcal{T}_{TR} \equiv \frac{|T_{uu}^{++}|^2 - |T_{uu}^{+-}|^2}{|T_{uu}^{++}|^2 + |T_{uu}^{+-}|^2}, \quad (32)$$

where

$$|T_{uu}^{+\pm}|^2 \equiv |T_{++}^{+\pm}|^2 + |T_{+0}^{+\pm}|^2 + |T_{+-}^{+\pm}|^2 + |T_{-+}^{+\pm}|^2 + |T_{-0}^{+\pm}|^2 + |T_{--}^{+\pm}|^2. \quad (33)$$

Therefore, we get

$$|T_{uu}^{++}|^2 = 2(T_0^V)^2 + 2(T_0^{uud})^2 d_1^2 - 2T_0^V T_0^{uud} d_1 \cos^2 \theta_\phi, \quad (34)$$

and

$$\begin{aligned} |T_{uu}^{+-}|^2 = & 2(T_0^{s\bar{s}})^2 + 2(T_0^{uud})^2 (1 - d_1^2) \\ & - T_0^{s\bar{s}} T_0^{uud} \{ (1 + \sin \theta_\phi) d_2 \cos(\theta_{p'} + \theta_\phi) \\ & + (1 - \sin \theta_\phi) d_3 \cos(\theta_{p'} - \theta_\phi) \}. \end{aligned} \quad (35)$$

Inspection of the above expression shows that neither of the asymmetries contain interference between the VDM photoproduction and $s\bar{s}$ knockout which is proportional to $T_0^V T_0^{s\bar{s}}$. This means that the asymmetries are not sensitive to the contribution of the $s\bar{s}$ knockout mechanism or to the amplitude a_0 ; Whereas they are sensitive to the uud - knockout (or the amplitude a_1), dominant at large $|t|$ where, however, the cross sections are exponentially small. As an example, Fig.2 shows the target-recoil asymmetry. At small $|t|$ it is proportional to $1 - 2(T_0^{s\bar{s}}/T_0^V)^2$. We choose $B^2 = 0, 0.0025, 0.01$; $a_0^2 = a_1^2 = 0.5$. At large $|t|$ it depends on the combination of $(T_0^{uud})^2$ and $T_0^V T_0^{uud}$, where the term proportional to $(T_0^{uud})^2$ is dominant because of sharply decreasing T_0^V with increasing $|t|$. So, at large $|t|$ the result is not responsive to the phase of a_1 qualitatively.

Longitudinal Polarization

In this case, the photon and the produced ϕ -meson are quantized parallel to the scattering plane. We choose the z -axis to be parallel to the photon momentum. In the helicity basis, the amplitudes read

$$\begin{aligned}
T_{fi}^{VDM} &= T_0^V d_{\lambda_i, \lambda_f}^{1/2}(\theta'_p) d_{\lambda_\gamma, \lambda_\phi}^1(\theta_\phi), \\
T_{fi}^{s\bar{s}} &= -T_0^{s\bar{s}} \sqrt{\frac{3}{2}} \langle \frac{1}{2} m_f 1 \sigma | \frac{1}{2} \lambda_i \rangle \lambda_\gamma [\sqrt{2} \hat{p}'_\sigma] d_{m_f, \lambda_f}^{1/2}(\theta'_p) d_{\lambda_\gamma, \lambda_\phi}^1(\theta_\phi), \\
T_{fi}^{uud} &= T_0^{uud} \sqrt{\frac{3}{2}} \langle \frac{1}{2} m_f - \lambda_\gamma 1 \sigma | i_c m_c \rangle \langle i_c m_c 1 m_\phi | \frac{1}{2} \lambda_i \rangle [\sqrt{2} \hat{p}'_{\phi, \sigma}] \\
&\quad \times d_{m_f, \lambda_f}^{1/2}(\theta'_p) d_{m_\phi, \lambda_\phi}^1(\theta_\phi).
\end{aligned} \tag{36}$$

where

$$\begin{aligned}
\hat{p}'_0 &= \cos \theta_{p'}, & \hat{p}'_{\pm} &= \mp \frac{1}{\sqrt{2}} \sin \theta_{p'}, \\
\hat{p}_{\phi, 0} &= \cos \theta_{p_\phi}, & \hat{p}_{\phi, \pm} &= \mp \frac{1}{\sqrt{2}} \sin \theta_{p_\phi}.
\end{aligned} \tag{37}$$

The polarized amplitudes are obtained as:

$$\begin{aligned}
T_{+\lambda_\phi}^{++} &= \{T_0^V \cos \frac{\theta_{p'}}{2} - T_0^{s\bar{s}} \cos \frac{3\theta_{p'}}{2}\} d_{1, \lambda_\phi}^1(\theta_\phi) \\
&\quad + d_1 T_0^{uud} \cos \frac{\theta_{p'}}{2} [2 \cos \theta_\phi d_{1, \lambda_\phi}^1(\theta_\phi) + \sqrt{2} \sin \theta_\phi d_{0, \lambda_\phi}^1(\theta_\phi)]
\end{aligned} \tag{38}$$

$$\begin{aligned}
T_{-\lambda_\phi}^{++} &= \{T_0^V \cos \frac{\theta_{p'}}{2} + T_0^{s\bar{s}} \cos \frac{3\theta_{p'}}{2}\} d_{-1, \lambda_\phi}^1(\theta_\phi) \\
&\quad + T_0^{uud} \sin \frac{\theta_{p'}}{2} [d_3 \sin \theta_\phi d_{1, \lambda_\phi}^1(\theta_\phi) \\
&\quad + \sqrt{2} d_4 \cos \theta_\phi d_{0, \lambda_\phi}^1(\theta_\phi) + d_2 \sin \theta_\phi d_{-1, \lambda_\phi}^1(\theta_\phi)]
\end{aligned} \tag{39}$$

$$\begin{aligned}
T_{+\lambda_\phi}^{+-} &= \{T_0^V \sin \frac{\theta_{p'}}{2} - T_0^{s\bar{s}} \sin \frac{3\theta_{p'}}{2}\} d_{1, \lambda_\phi}^1(\theta_\phi) \\
&\quad + T_0^{uud} d_1 \sin \frac{\theta_{p'}}{2} [2 \cos \theta_\phi d_{1, \lambda_\phi}^1(\theta_\phi) + \sqrt{2} \sin \theta_\phi d_{0, \lambda_\phi}^1(\theta_\phi)]
\end{aligned} \tag{40}$$

$$\begin{aligned}
T_{-\lambda_\phi}^{+-} &= \{T_0^V \sin \frac{\theta_{p'}}{2} + T_0^{s\bar{s}} \sin \frac{3\theta_{p'}}{2}\} d_{-1, \lambda_\phi}^1(\theta_\phi) \\
&\quad - T_0^{uud} \cos \frac{\theta_{p'}}{2} [d_3 \sin \theta_\phi d_{1, \lambda_\phi}^1(\theta_\phi) + \sqrt{2} d_4 \cos \theta_\phi d_{0, \lambda_\phi}^1(\theta_\phi) \\
&\quad + d_2 \sin \theta_\phi d_{-1, \lambda_\phi}^1(\theta_\phi)]
\end{aligned} \tag{41}$$

$$\begin{aligned}
T_{+\lambda_\phi}^{-+} &= -\{T_0^V \sin \frac{\theta_{p'}}{2} + T_0^{s\bar{s}} \sin \frac{3\theta_{p'}}{2}\} d_{1, \lambda_\phi}^1(\theta_\phi) \\
&\quad + T_0^{uud} \cos \frac{\theta_{p'}}{2} [d_2 \sin \theta_\phi d_{1, \lambda_\phi}^1(\theta_\phi) - \sqrt{2} d_4 \cos \theta_\phi d_{0, \lambda_\phi}^1(\theta_\phi) \\
&\quad + d_3 \sin \theta_\phi d_{-1, \lambda_\phi}^1(\theta_\phi)]
\end{aligned} \tag{42}$$

$$T_{-\lambda_\phi}^{-+} = -\{T_0^V \sin \frac{\theta_{p'}}{2} - T_0^{s\bar{s}} \sin \frac{3\theta_{p'}}{2}\} d_{-1, \lambda_\phi}^1(\theta_\phi)$$

$$+T_0^{uud} d_1 \sin \frac{\theta_{p'}}{2} [\sqrt{2} \sin \theta_\phi d_{0,\lambda_\phi}^1(\theta_\phi) - 2 \cos \theta_\phi d_{-1,\lambda_\phi}^1(\theta_\phi)] \quad (43)$$

$$\begin{aligned} T_{+\lambda_\phi}^{--} &= \{T_0^V \cos \frac{\theta_{p'}}{2} + T_0^{s\bar{s}} \cos \frac{3\theta_{p'}}{2}\} d_{1,\lambda_\phi}^1(\theta_\phi) \\ &+ T_0^{uud} \sin \frac{\theta_{p'}}{2} [d_2 \sin \theta_\phi d_{1,\lambda_\phi}^1(\theta_\phi) - \sqrt{2} d_4 \cos \theta_\phi d_{0,\lambda_\phi}^1(\theta_\phi) \\ &+ d_3 \sin \theta_\phi d_{-1,\lambda_\phi}^1(\theta_\phi)] \end{aligned} \quad (44)$$

$$\begin{aligned} T_{-\lambda_\phi}^{--} &= \{T_0^V \cos \frac{\theta_{p'}}{2} - T_0^{s\bar{s}} \cos \frac{3\theta_{p'}}{2}\} d_{-1,\lambda_\phi}^1(\theta_\phi) \\ &- T_0^{uud} d_1 \cos \frac{\theta_{p'}}{2} [\sqrt{2} \sin \theta_\phi d_{0,\lambda_\phi}^1(\theta_\phi) - 2 \cos \theta_\phi d_{-1,\lambda_\phi}^1(\theta_\phi)] \end{aligned} \quad (45)$$

Beam-Vector-meson Asymmetry \mathcal{L}_{BV}

The beam-vector-meson asymmetry \mathcal{L}_{BV} is defined as

$$\mathcal{L}_{BV} \equiv \frac{|T_{++}^{uu}|^2 - |T_{+-}^{uu}|^2}{|T_{++}^{uu}|^2 + |T_{+-}^{uu}|^2}, \quad (46)$$

where

$$|T_{\pm\pm}^{uu}|^2 \equiv |T_{\pm\pm}^{++}|^2 + |T_{\pm\pm}^{+-}|^2 + |T_{\pm\pm}^{-+}|^2 + |T_{\pm\pm}^{--}|^2. \quad (47)$$

Using the expressions for the amplitudes, we have

$$\begin{aligned} |T_{++}^{uu}|^2 - |T_{+-}^{uu}|^2 &= 2\{(T_0^V)^2 + (T_0^{s\bar{s}})^2\} \cos \theta_\phi \\ &+ 4(T_0^{uud})^2 d_1 \{d_1 \cos \theta_\phi \cos 2\theta_\phi + d_4 \sin \theta_\phi \sin 2\theta_\phi\} \\ &+ 2(T_0^V)(T_0^{uud}) d_1 (\cos^2 \theta_\phi + \cos 2\theta_\phi) \\ &- 2(T_0^{s\bar{s}})(T_0^{uud}) \{\cos \theta_{p'} d_1 (\cos^2 \theta_\phi + \cos 2\theta_\phi) \\ &+ \sin \theta_{p'} (d_1 \sin \theta_\phi \cos \theta_\phi + d_4 \sin 2\theta_\phi)\}, \end{aligned} \quad (48)$$

$$\begin{aligned} |T_{++}^{uu}|^2 + |T_{+-}^{uu}|^2 &= \{(T_0^V)^2 + (T_0^{s\bar{s}})^2\} (1 + \cos^2 \theta_\phi) \\ &+ 2(T_0^{uud})^2 \{d_1^2 (1 + \cos^2 2\theta_\phi) + d_4^2 \sin^2 2\theta_\phi\} \\ &+ 2(T_0^V)(T_0^{uud}) d_1 \cos \theta_\phi (1 + \cos 2\theta_\phi) \\ &- 2(T_0^{s\bar{s}})(T_0^{uud}) \{\cos \theta_{p'} d_1 \cos \theta_\phi (1 + \cos 2\theta_\phi) \\ &+ \sin \theta_{p'} (d_1 \sin \theta_\phi + d_4 \cos \theta_\phi \sin 2\theta_\phi)\}. \end{aligned} \quad (49)$$

Beam-Target Asymmetry \mathcal{L}_{BT}

The beam-target asymmetry \mathcal{L}_{BT} is defined as

$$\mathcal{L}_{BT} \equiv \frac{|T_{+u}^{+u}|^2 - |T_{-u}^{+u}|^2}{|T_{+u}^{+u}|^2 + |T_{-u}^{+u}|^2}, \quad (50)$$

where

$$|T_{\pm u}^{+u}|^2 \equiv |T_{\pm+}^{++}|^2 + |T_{\pm 0}^{++}|^2 + |T_{\pm-}^{++}|^2 + |T_{\pm+}^{+-}|^2 + |T_{\pm 0}^{+-}|^2 + |T_{\pm-}^{+-}|^2. \quad (51)$$

Again, we have

$$\begin{aligned} |T_{+u}^{+u}|^2 - |T_{-u}^{+u}|^2 &= -2(T_0^{uud})^2 \{1 - 2d_1^2(1 + \cos^2 \theta_\phi)\} \\ &\quad -4(T_0^V)(T_0^{s\bar{s}}) \cos \theta_{p'} + 4(T_0^V)(T_0^{uud})d_1 \cos \theta_\phi \\ &\quad -2(T_0^{s\bar{s}})(T_0^{uud}) \{2d_1 \cos \theta_{p'} \cos \theta_\phi - d_2 \sin \theta_{p'} \sin \theta_\phi\}, \\ |T_{+u}^{+u}|^2 + |T_{-u}^{+u}|^2 &= 2\{(T_0^V)^2 + (T_0^{s\bar{s}})^2 + (T_0^{uud})^2\} + 4(T_0^V)(T_0^{uud})d_1 \cos \theta_\phi \\ &\quad -2(T_0^{s\bar{s}})(T_0^{uud}) \{2d_1 \cos \theta_{p'} \cos \theta_\phi + d_2 \sin \theta_{p'} \sin \theta_\phi\}. \end{aligned} \quad (52)$$

Target-Vector-meson Asymmetry \mathcal{L}_{TV}

The target-vector-meson asymmetry \mathcal{L}_{TV} is defined as

$$\mathcal{L}_{TV} \equiv \frac{|T_{u+}^{+u}|^2 - |T_{u-}^{+u}|^2}{|T_{u+}^{+u}|^2 + |T_{u-}^{+u}|^2}, \quad (53)$$

where

$$|T_{u\pm}^{+u}|^2 \equiv |T_{+\pm}^{++}|^2 + |T_{-\pm}^{++}|^2 + |T_{+\pm}^{+-}|^2 + |T_{-\pm}^{+-}|^2. \quad (54)$$

And we get

$$\begin{aligned} |T_{u+}^{+u}|^2 - |T_{u-}^{+u}|^2 &= 4(T_0^{uud})^2 d_1 (d_1 \cos \theta_\phi \cos 2\theta_\phi - d_4 \sin \theta_\phi \sin 2\theta_\phi) \\ &\quad -4(T_0^V)(T_0^{s\bar{s}}) \cos \theta_{p'} \cos \theta_\phi \\ &\quad +2(T_0^V)(T_0^{uud})(\cos^2 \theta_\phi + \cos 2\theta_\phi) \\ &\quad -2(T_0^{s\bar{s}})(T_0^{uud}) \{d_1 \cos \theta_{p'} (\cos^2 \theta_\phi + \cos 2\theta_\phi) \\ &\quad - \sin \theta_{p'} (d_1 \cos \theta_\phi \sin \theta_\phi + d_4 \sin 2\theta_\phi)\}, \end{aligned} \quad (55)$$

$$\begin{aligned} |T_{u+}^{+u}|^2 + |T_{u-}^{+u}|^2 &= \{(T_0^V)^2 + (T_0^{s\bar{s}})^2\} (1 + \cos^2 \theta_\phi) \\ &\quad +2(T_0^{uud})^2 \{d_1^2 (\cos^2 \theta_\phi + \cos^2 2\theta_\phi) + d_1^2 \sin^2 \theta_\phi + d_4^2 \sin^2 2\theta_\phi\} \\ &\quad +2(T_0^V)(T_0^{uud})d_1 \cos \theta_\phi (1 + \cos 2\theta_\phi) \\ &\quad -2(T_0^{s\bar{s}})(T_0^{uud}) \{d_1 \cos \theta_{p'} \cos \theta_\phi (1 + \cos 2\theta_\phi) \\ &\quad + \sin \theta_{p'} (d_1 \sin \theta_\phi + d_4 \cos \theta_\phi \sin 2\theta_\phi)\}. \end{aligned} \quad (56)$$

Target-Recoil Asymmetry \mathcal{L}_{TR}

The target-recoil asymmetry \mathcal{L}_{TR} is defined as

$$\mathcal{L}_{TR} \equiv \frac{|T_{uu}^{++}|^2 - |T_{uu}^{+-}|^2}{|T_{uu}^{++}|^2 + |T_{uu}^{+-}|^2}, \quad (57)$$

where

$$|T_{uu}^{+\pm}|^2 \equiv |T_{++}^{+\pm}|^2 + |T_{+0}^{+\pm}|^2 + |T_{+-}^{+\pm}|^2 + |T_{-+}^{+\pm}|^2 + |T_{-0}^{+\pm}|^2 + |T_{--}^{+\pm}|^2. \quad (58)$$

Therefore, we get

$$\begin{aligned} |T_{uu}^{++}|^2 - |T_{uu}^{+-}|^2 = & -2(T_0^{uud})^2 \cos \theta_{p'} \{ (1 - 2d_1^2) - 2d_1^2 \cos^2 \theta_\phi \} \\ & + (T_0^V)(T_0^{uud}) 2 \{ 2d_1 \cos \theta_{p'} \cos \theta_\phi + d_2 \sin \theta_{p'} \sin \theta_\phi \} \\ & - 2(T_0^{s\bar{s}})(T_0^{uud}) \{ 2d_1 \cos 2\theta_{p'} \cos \theta_\phi - d_2 \sin 2\theta_{p'} \sin \theta_\phi \}, \end{aligned} \quad (59)$$

$$\begin{aligned} |T_{uu}^{++}|^2 + |T_{uu}^{+-}|^2 = & 2 \{ (T_0^V)^2 + (T_0^{s\bar{s}})^2 + (T_0^{uud})^2 \} + 4(T_0^V)(T_0^{uud}) d_1 \cos \theta_\phi \\ & - 2(T_0^{s\bar{s}})(T_0^{uud}) \{ 2d_1 \cos \theta_{p'} \cos \theta_\phi - d_2 \sin \theta_{p'} \sin \theta_\phi \}. \end{aligned} \quad (60)$$

The beam-target and target-vector meson asymmetries are of the most interest because at small $|t|$ they are proportional to the interference $T_0^V T_0^{s\bar{s}}$ and can bring information on the amplitude a_0 . Figs.3,4 show the result of calculation of the beam-target asymmetry at different signs of a : $a_0 = a_1 < 0$ in Fig.3 and $a_0 = a_1 > 0$ in Fig.4. We choose $B^2 = 0, 0.0025, 0.01$. The result for all the other possible combinations a_0, a_1 can be reproduced straightforwardly by making use of the above expressions. Our calculation shows that at small $|t| \sim |t_{max}|$ the asymmetry may be as much as 30% when the proposed strangeness probability in the proton is only 0.25%.

As a summary, we analyzed polarization observables based on the RHOM that takes into account the Lorentz contraction effects of the composite particle wave function. We find that even with a less than 1% admixture of strange quarks, the deviations of the asymmetries from those of the pure diffractive VDM photoproduction are as much as 10 ~ 50% depending on the value of t when we take the knockout amplitudes into account. The strong t -dependence of the asymmetries could be crucial for testing experimentally the $s\bar{s}$ - uud cluster model for the strangeness content of the proton. Current experimental data [17] on the asymmetries are not sufficiently accurate to extract the knockout contribution

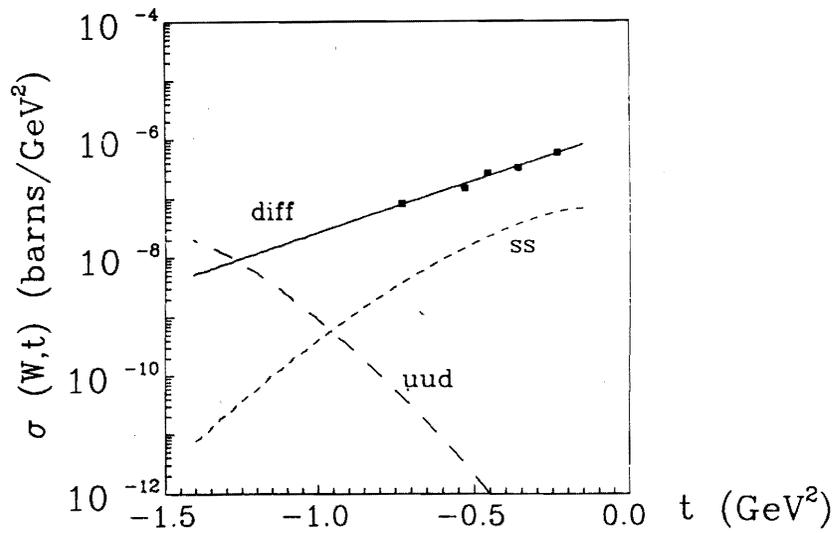


FIG. 1. The t -dependence of the unpolarized photoproduction cross section $d\sigma/dt$ with $W=2.1$ GeV and $Q^2 = 0.02$ GeV². The solid line corresponds to the diffractive cross section, the dashed line to the $s\bar{s}$ knockout cross section, and the dashed-dotted line to the uud knockout cross section. Experimental data [12] are given by dark circles for comparison.

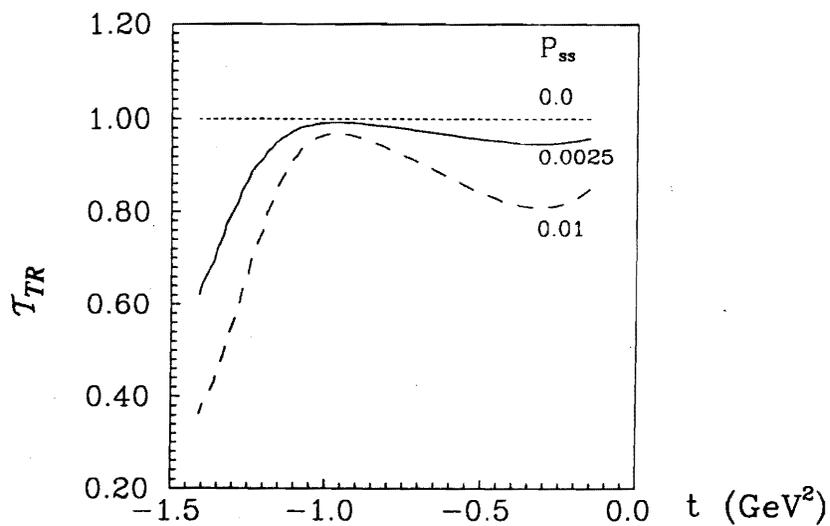


FIG. 2. The t -dependence of the transverse asymmetry T_{TR} with $B^2=0$ (solid line), 0.25% (dotted line), and 1% (dash-dotted line).

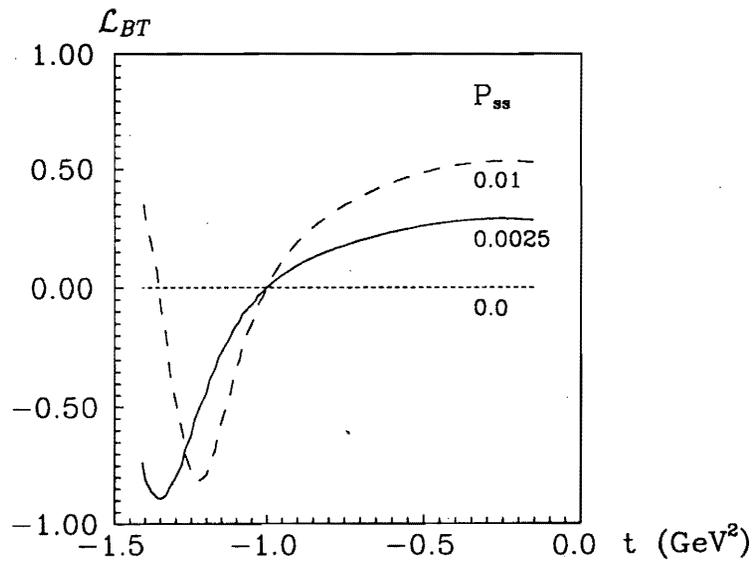


FIG. 3. The t -dependence of the longitudinal asymmetry \mathcal{L}_{BT} with $a_0 = a_1 < 0$ and $B^2=0$ (solid line), 0.25% (dotted line), and 1% (dash-dotted line).

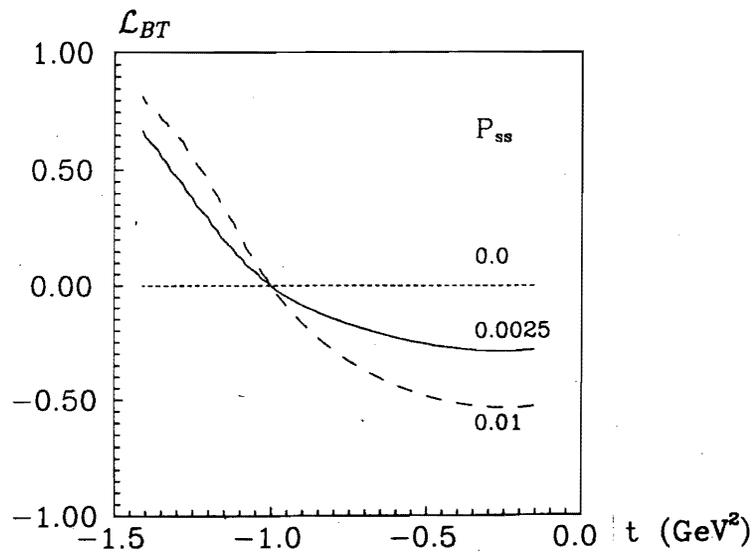


FIG. 4. The t -dependence of the longitudinal asymmetry \mathcal{L}_{BT} . Notation is the same as in Fig.3 but $a_0 = a_1 > 0$.

from the diffractive process. However, the similar effect should be seen in ϕ - electroproduction and it may be checked in future experiments at CEBAF.

One of us (A.I.T.) would like to acknowledge the warm hospitality of colleagues at the Department of Physics at National Taiwan University. This work was supported in part by the National Science Council of ROC under Grant No. NSC83-0208-M002-017 and in part by the Grant No. MP8300 from the International Science Foundation.

REFERENCES

- [1] E. M. Henley, G. Krein, and A. G. Williams, *Bull. Am. Phys. Soc.* **34**, 1828 (1989); E. M. Henley, G. Krein, S. J. Pollock, and A. G. Williams, *Phys. Lett. B* **269**, 31 (1991).
- [2] E. M. Henley, G. Krein, and A. G. Williams, *Phys. Lett. B* **281**, 178 (1992).
- [3] D. Kaplan and A. V. Manohar, *Nucl. Phys. B* **310**, 527 (1988); R. D. McKeown, *Phys. Lett. B* **219**, 140 (1989); D. Kaplan, *ibid.* **275**, 137 (1992).
- [4] R. L. Jaffe, *Phys. Lett. B* **229**, 275 (1989).
- [5] T. P. Cheng and R. F. Dashen, *Phys. Rev. Lett.* **26**, 594 (1971); T. P. Cheng, *Phys. Rev. D* **13**, 2161 (1976); C. A. Dominguez and P. Langacker, *ibid.* **24**, 190 (1981); J. F. Donoghue and C. R. Nappi, *Phys. Lett. B* **168**, 105 (1986); J. Gasser, H. Leutwyler, and M. E. Sainio, *ibid.* **253**, 260 (1991).
- [6] EM Collaboration, J. Ashman *et al.*, *Phys. Lett. B* **206**, 364 (1988); *Nucl. Phys.* **B328**, 1 (1989).
- [7] M. Anselmino and M. D. Scadron, *Phys. Lett. B* **229**, 117 (1989); H. J. Lipkin, *ibid.* **256**, 284 (1991); R. L. Jaffe and H. J. Lipkin, *ibid.* **266**, 458 (1991).
- [8] J. Ellis, E. Gabathuler, and M. Karliner, *Phys. Lett. B* **217**, 173 (1989); D. Beck, *Phys. Rev. D* **39**, 3248 (1989).
- [9] D. W. G. S. Leith, in *Electromagnetic Interactions of Hadrons*, Vol. 1, edited by A. Donnachie and G. Shaw (Plenum, New York, 1978) and references therein; D. G. Cassel *et al.*, *Phys. Rev. D* **24**, 2787 (1981).
- [10] A. I. Titov, S. N. Yang, and Y. Oh, National Taiwan University Report (1994); A. I. Titov and S. N. Yang, in *Proceedings of the 13th International Conference on Particle and Nuclei*, Perugia, Italy, 1993, edited by A. Pascolini (World Scientific, Singapore, 1994).
- [11] K. Fujimura, T. Kobayashi, and M. Namiki, *Prog. Theor. Phys.*, **44**, 193 (1970); *Prog. Theor. Phys.*, **43**, 73 (1970).
- [12] H. J. Besch *et al.*, *Nucl. Phys.* **B70**, 257 (1974).
- [13] R. P. Feynman, M. Kislinger, and F. Ravndal, *Phys. Rev. D* **3**, 2107 (1971); R. G. Lipes, *ibid.* **5**, 2849 (1972); R. P. Feynman, *Photon-Hadron Interaction* (Benjamin, Reading, 1972).
- [14] Y. S. Kim and M. E. Noz, *Phys. Rev. D* **12**, 129 (1975); Y. Kizukuri, M. Namiki, and K. Okano, *Prog. Theor. Phys.* **61**, 559 (1979); V. V.

- Burov, S. M. Dorkin, V. K. Lukyanov, and A. I. Titov, *Z. Phys. A* **306**, 149 (1982).
- [15] S. J. Brodsky and J. R. Primack, *Ann. Phys. (N.Y.)* **52**, 315 (1969);
F. E. Close and Z. Li, *Phys. Rev. D* **42**, 2194 (1990).
- [16] A.R. Edmonds, "Angular Momentum in Quantum Mechanics", 2nd ed. (Princeton Univ. Press, Princeton, 1974)
- [17] H.-J. Behrend *et al.*, *Phys. Lett. B* **56**, 408 (1975); *Nucl. Phys. B* **144**, 22 (1978); H. J. Halpern *et al.*,
Phys. Rev. Lett. **29**, 1425 (1972).

Received by Publishing Department
on May 23, 1995.

The Publishing Department
of the Joint Institute for Nuclear Research
offers you to acquire the following books:

Index	Title
E7-93-274	Proceedings of the International School-Seminar on Heavy Ion Physics. Dubna, 1993 (2 volumes, 498 and 590 pp., stiffened cover, in English)
Д14-93-325	Proceedings of the Workshop «Activation Analysis in Environment Protection». Dubna, 1992 (in Russian and English)
Д1,9,13-93-459	Proceedings of the II Workshop on JINR Tau-Charm Factory. Dubna, 1993 (266 p. in Russian and English)
E7-94-19	Proceedings of the International Conference on Dynamical Aspects of Nuclear Fission. Smolenice, 1993 (318 p. in English)
94-55	Proceedings of the International Bogoliubov Memorial Meeting. Dubna, 1993 (216 p. in Russian and English)
E2-94-95	Proceedings of the International Symposium «Dubna Deuteron-93». Dubna, 1993 (390 p. in English)
Д14-94-269	Abstracts of the XXX Conference on Low Temperature Physics. Dubna, 1994 (2 volumes in Russian and English)
E7-94-270	Proceedings of the Workshop on Physical Experiments and First Results on Heavy Ion Storage and Cooler Rings. Smolenice, 1992 (324 p. in English)
E13-94-290	Abstracts of the 17th International Conference «Nuclear Tracks in Solids». Dubna, 1994 (296 p. in English)
E2-94-347	International Workshop «Symmetry Methods in Physics». In Memory of Professor Yu.A.Smorodinsky. Dubna, 1993 (2 volumes, 602 p. in English)
E4-94-370	Proceedings of the IV International Conference on Selected Topics in Nuclear Structure. Dubna, 1994 (412 p. in English)

Index	Title
E4-94-386	Proceedings of the VI Trilateral German-Russian-Ukrainian Seminar on High-Temperature Superconductivity. Dubna, 1994 (340 p. in English)

Please apply to the Publishing Department of the Joint Institute for Nuclear Research for extra information. Our address is:

Publishing Department
Joint Institute for Nuclear Research
Dubna, Moscow Region
141980 Russia
E-mail: adm@pubd.jinr.dubna.su.

Поляризационные характеристики в фоторождении ϕ -мезонов

Проведен анализ поляризационных экспериментов по образованию ϕ -мезонов на протонах на основе $uud-\bar{s}\bar{s}$ кластерной модели с целью возможного определения скрытой странности в протоне. Взаимодействие кварков описывается моделью релятивистского гармонического осциллятора, которая естественным образом учитывает лоренцевское сокращение продольного размера связанной системы. Показано, что предсказания поперечной асимметрии в дважды поляризационных экспериментах «мишень—протон отдачи» и продольной асимметрии в измерениях «пучок—мишень» очень сильно зависят от модели. Разность в предсказаниях $uud-\bar{s}\bar{s}$ кластерной модели и дифракционного рождения в модели векторной доминантности составляет 20–80% «странности» в протоне на уровне 0,2–0,3%.

Работа выполнена в Лаборатории теоретической физики им. Н.Н.Боголюбова ОИЯИ.

Сообщение Объединенного института ядерных исследований. Дубна, 1995

Polarization Observables in ϕ -Meson Photoproduction

We analyze polarization observables of the ϕ -meson photoproduction from a proton within the $uud-\bar{s}\bar{s}$ cluster model as a probe of strangeness in the proton. Our consideration is based on the relativistic quark model that takes into account the Lorentz-contraction effects of the intrinsic hadron wave functions. We show that transverse target-recoil and longitudinal beam-target double polarizations in the diffractive vector meson dominance model are different from those of the knockout mechanisms and the difference between them is as much as 20–80% if only 0.2–0.3% strangeness admixture in proton wave function is assumed. This strong effect may be studied experimentally.

The investigation has been performed at the Bogoliubov Laboratory of Theoretical Physics, JINR.

Communication of the Joint Institute for Nuclear Research. Dubna, 1995