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Derivation of the Energy-Time Uncertainty Relation

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Abstract

A new derivation from first principles is given of the energy-time uncertainty relation in quantum mechanics. A canonical transformation is made in classical mechanics to a new canonical momentum, which is energy E , and a new canonical coordinate T , which is called *tempus*, conjugate to the energy. *Tempus* T , the canonical coordinate conjugate to the energy, is conceptually different from the time t in which the system evolves. The Poisson bracket is a canonical invariant, so that energy and *tempus* satisfy the same Poisson bracket as do p and q . When the system is quantised, we find the energy-time uncertainty relation $\Delta E \Delta T \geq \hbar/2$. For a conservative system the average of the *tempus* operator \hat{T} is the time t plus a constant. For a free particle and a particle acted on by a constant force, the *tempus* operators are constructed explicitly, and the energy-time uncertainty relation is explicitly verified.

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1 Introduction

Ever since the discovery of quantum mechanics, the energy-time uncertainty relation, $\Delta E \Delta t \geq \hbar/2$, has had a different basis than the position-momentum uncertainty relation, $\Delta q \Delta p \geq \hbar/2$. In nonrelativistic quantum mechanics time t is a parameter, not an operator as is the coordinate q . Therefore the usual quantum mechanical average and uncertainty in the time cannot be calculated as in the case of the coordinate. For this reason, it has been suggested that the energy-time uncertainty relation simply be eliminated from quantum mechanics [1]. This radical suggestion overlooks the many different approaches to the physically useful energy-time uncertainty relation.

We briefly mention some of the derivations of the energy-time uncertainty relation, even though some are only heuristic [2, ?]. (1) The energy-time uncertainty relation can be obtained by showing that a wave packet of width Δt in time requires a spread in angular frequencies $\Delta \omega$ such that $\Delta \omega \Delta t \sim 1$. Together with the Planck relation $E = \hbar \omega$, we obtain $\Delta E \Delta t \sim \hbar$ [4]. (2) The energy-time uncertainty relation is derived by reducing it to the position-momentum uncertainty relation [3, ?]. If $\Delta E = (\partial E / \partial p) \Delta p = v \Delta p$ and $\Delta t = \Delta q / v$, where v is the group velocity, then $\Delta E \Delta t = \Delta p \Delta q \geq \hbar/2$. In this context, the meaning of ΔE and Δt is not precisely given. (3) Mandelstam and Tamm [6] showed that the energy-time uncertainty relation can be derived from the Heisenberg equation of motion for an arbitrary operator \hat{A} and the generalized uncertainty relation for \hat{A} and the Hamiltonian \hat{H} [3, ?]. If we define Δt to be the smallest $\Delta t_A \equiv \Delta A / (d \langle \hat{A} \rangle / dt)$ for all operators \hat{A} , and $\Delta E = \Delta H$, we also recover the energy-time uncertainty relation [8, ?]. (4) Wigner [10] has considered the expectation value of t^n in the state $\psi = \psi(q, t)$ to be

$$\langle t^n \rangle = \frac{\int_{-\infty}^{\infty} dt t^n |\psi(q, t)|^2}{\int_{-\infty}^{\infty} dt |\psi(q, t)|^2}. \quad (1)$$

Therefore he defines the square of the uncertainty $(\Delta t)^2 = \langle t^2 \rangle - \langle t \rangle^2$ at constant q and defines $(\Delta E)^2$ similarly. He shows that $\Delta E \Delta t \geq \hbar/2$ holds at constant q . Cook [11] uses this form of the uncertainty relation to explain a thought experiment of Einstein [12]. (5) Many authors [13]-[29] have tried to find a time operator \hat{t} , which satisfies the canonical commutation relation $[\hat{t}, \hat{H}] = i\hbar$. Then the energy-time uncertainty relation can be derived by the same procedure as for the position-momentum uncertainty relation [7]. These attempts have to contend with Pauli's argument [30] that a time

operator cannot exist in general because the spectrum of the time operator is continuous and unbounded, while the spectrum of the Hamiltonian may be discrete and is bounded from below. The role between the time t as a parameter and the time operator \hat{t} is often not clear in these attempts. A clear distinction is however made by Razavy [17, ?]. (6) In scattering theory it is the time of arrival which is important, and attempts have been made to find an operator which describes this quantity

[31]-[36]. (7) In the context of the density matrix ρ , Eberly and Singh [37] have shown that $\partial\rho/\partial t$ plays the role of an "inverse time operator" from which the energy-time uncertainty relation can be derived [38]. (8) The role of the uncertainty relations in relativistic quantum theory has been discussed by a number of authors [39]-[43]. (9) The interpretation of the energy-time uncertainty relation is a part of quantum measurement theory [44]-[51], and is distinct from its derivation. The energy-time uncertainty relation for decaying states has been improved by Gislason, Sabelli, and Wood [52].

In this paper we give a new derivation of the energy-time uncertainty relation, which is based on first principles, and does not resort to any heuristic or *ad hoc* arguments. Our derivation is firmly based on classical mechanics and its quantization, and puts the energy-time uncertainty relation on the same footing as the position-momentum uncertainty relation. In classical mechanics we make a canonical transformation from the old canonical variables (q, p) to new canonical variables (q', p') , where the new canonical momentum p' is the energy E of the particle and the new canonical coordinate q' is a quantity T conjugate to the energy, which we call *tempus*. This *tempus* T has dimension of time, but is conceptually different from the time t in which the system evolves. Since T and E are canonically conjugate variables, and the Poisson bracket is a canonical invariant [53], their Poisson bracket is unity. Upon quantization, the *tempus* operator \hat{T} and the energy operator \hat{E} satisfy the canonical commutation relation. Therefore, the energy and the *tempus* operators satisfy the usual energy-time uncertainty relation $\Delta E \Delta T \geq \hbar/2$, as long as the expectation values of \hat{T} and \hat{T}^2 exist. For a conservative system the expectation value of the *tempus* operator is equal to the time t plus an irrelevant constant. The distinction between the *tempus* operator \hat{T} conjugate to the energy operator and the time t of evolution is thus made clear.

This approach is applied to two examples: (1) a free particle and (2) a particle acted on by a constant force. In both cases the *tempus* operator \hat{T} is constructed explicitly by quantizing the corresponding classical *tempus* T conjugate to the energy. In the subspace of the Hilbert space in which

it exists, the average value $\langle \hat{T} \rangle = t$ plus a constant. The energy-time uncertainty relation is explicitly shown to be satisfied.

In Section 2, a canonical transformation is made in classical mechanics to energy and its canonical conjugate *tempus*. Using Poisson brackets, we quantize the system in Section 3 and derive the energy-time uncertainty relation. In Section 4, examples of a free particle and a particle acted on by a constant force are considered, and it is shown that the energy-time uncertainty relation is explicitly satisfied. Finally, the conclusion is given in Section 5.

2 Energy and *Tempus* as Canonical Variables

In the Hamiltonian formulation of classical mechanics, the canonical coordinate q and the canonical momentum p conjugate to it satisfy Hamilton's equations

$$\dot{q} = \partial H / \partial p, \quad \dot{p} = -\partial H / \partial q, \quad (2)$$

where $H = H(q, p, t)$ is the Hamiltonian for a system of one degree of freedom, and the overdot denotes the total time derivative. Using a generating function S , we can make a canonical transformation to a new set of canonical variables (q', p') , which also satisfy Hamilton's equations [54]

$$\dot{q}' = \partial H' / \partial p', \quad \dot{p}' = -\partial H' / \partial q', \quad (3)$$

where the new Hamiltonian is $H' = H + \partial S / \partial t$.

If the new canonical momentum p' is chosen to be the energy E , the new canonical coordinate q' conjugate to it is called *tempus* and denoted as T . This *tempus* canonical coordinate conjugate to the energy is conceptually different from the time t in which the system evolves, and is a function of q , E and t [55] with the dimension of time.

The generating function of the second type $S(q, E, t)$ which implements this canonical transformation satisfies [53]-[55]

$$p = \partial S(q, E, t) / \partial q \quad (4)$$

and the new canonical coordinate *tempus* is

$$T = \partial S(q, E, t) / \partial E. \quad (5)$$

By integrating Eq. (4), we obtain

$$S(q, E, t) = \int_{q_0}^q p(\bar{q}, E, t) d\bar{q} + S(q_0, E, t), \quad (6)$$

where q_0 is an arbitrary initial displacement. The arbitrary function $S(q_0, E, t)$ of E and t in Eq. (6) can often be chosen to be zero without loss of generality. To obtain the generating function in Eq. (6), it is necessary to find the canonical momentum p as function of q, E and t . The new generalized coordinate *tempus* T can be obtained as a function of q, E and t from Eq. (5). When this equation is solved for q as a function of T, E and t , we have the solution to the problem if we know T and E as functions of time t .

The new Hamiltonian H' in Eq. (3) for the new canonical variables $(q', p') = (T, E)$ is

$$H' = H + (\partial S / \partial t)_{q,E}. \quad (7)$$

The energy E may not be equal to the Hamiltonian H , so their difference is defined as

$$\Phi = H - E. \quad (8)$$

When Eqs. (7) and (8) are used in Hamilton's equations (3) for the new canonical variables (T, E) , we obtain [55]

$$\dot{T} = 1 + [\partial(\Phi + \partial S / \partial t) / \partial E]_{T,t} \quad (9)$$

and

$$\dot{E} = -[\partial(\Phi + \partial S / \partial t) / \partial T]_{E,t}. \quad (10)$$

In Eqs. (9) and (10) it is necessary to express $\Phi + (\partial S / \partial t)_{q,E}$ as a function of E, T and t before differentiation. Hamilton's equations (9) and (10) can be solved for $T(t)$ and $E(t)$. The solution to the original problem is $q = q(E(t), T(t), t) \equiv q(t)$.

For a conservative system, there is no explicit time dependence in S , so $\partial S / \partial t = 0$. In this case, the Hamiltonian H can be chosen to be the energy E , so $\Phi = H - E = 0$. Therefore, Eqs. (9) and (10) reduce to

$$\dot{T} = 1, \quad \dot{E} = 0. \quad (11)$$

The solution to Eq. (11) is $T = t - t_0$ and $E = E_0$, where t_0 and E_0 are constants. In this case *tempus* is equal to the time t minus a constant. The solution to the original problem is therefore $q = q(E, T) = q(E_0, t - t_0)$.

3 Quantization

In order to quantize a system, it is necessary to replace the Poisson bracket of two functions of the canonical variables by $(i\hbar)^{-1}$ times the commutator

of the corresponding operators. The Poisson bracket between $A = A(q, p, t)$ and $B = B(q, p, t)$ is defined as

$$\{A, B\}_{q,p} = \frac{\partial A}{\partial q} \frac{\partial B}{\partial p} - \frac{\partial B}{\partial q} \frac{\partial A}{\partial p}. \quad (12)$$

The Poisson bracket is a canonical invariant [53], so

$$\{A, B\}_{q,p} = \{A, B\}_{q',p'}, \quad (13)$$

where the new canonical variables (q', p') are related to the old canonical variables (q, p) by a canonical transformation. If $A = q$ and $B = p$, we have

$$1 = \{q, p\}_{q,p} = \{q', p'\}_{q',p'} = \{q', p'\}_{q,p} = \{T, E\}_{q,p}, \quad (14)$$

from Eqs. (12) and (13).

When the system is quantized, the operators \hat{q} and \hat{p} (\hat{T} and \hat{E}) corresponding to q and p (T and E) satisfy the commutation relations

$$[\hat{q}, \hat{p}] = i\hbar \quad (15)$$

and

$$[\hat{T}, \hat{E}] = i\hbar, \quad (16)$$

from Eq. (14).

In general, the uncertainty relation for the operators \hat{A} and \hat{B} in the state ψ is [56]

$$\Delta A \Delta B \geq \frac{1}{2} | \langle \hat{C} \rangle |, \quad (17)$$

where $i\hat{C} = [\hat{A}, \hat{B}]$ and $\langle \hat{C} \rangle \equiv \langle \psi | \hat{C} \psi \rangle$. The square of the uncertainty in \hat{A} is defined as

$$(\Delta A)^2 = \langle \psi | (\hat{A} - \langle \hat{A} \rangle)^2 \psi \rangle. \quad (18)$$

For the energy and *tempus* operators in Eq. (16) we have the uncertainty relation

$$\Delta E \Delta T \geq \hbar/2, \quad (19)$$

where the state ψ must be in the domain of \hat{E} , \hat{E}^2 , \hat{T} , and \hat{T}^2 . The choice of the new canonical momentum $p' = E$ and the new canonical coordinate $q' = T$ puts the energy-time uncertainty relation in Eq. (19) on the same solid basis as the momentum-position uncertainty relation $\Delta p \Delta q \geq \hbar/2$

[56]. In contrast to t , the new canonical *tempus* operator \hat{T} is conjugate to the energy operator \hat{E} . The uncertainty in the *tempus* operator has a well defined meaning for states that are in its domain.

For a conservative system the average of the *tempus* operator \hat{T} is the time t plus an irrelevant constant. The Schroedinger equation is

$$\hat{H}(q, \hat{p})\psi(q, t) = i\hbar\partial\psi(q, t)/\partial t, \quad (20)$$

for a system with a time-independent Hamiltonian $\hat{H} = H(q, \hat{p})$.

The time-dependent wave function $\psi(q, t)$ can be written as

$$\psi(q, t) = U(t)\psi(q, 0), \quad (21)$$

where $\psi(q, 0)$ is the wave function at time zero. The time evolution operator is

$$U(t) = \exp[-i\hat{H}t/\hbar], \quad (22)$$

where \hat{H} is the time-independent Hamiltonian operator.

From the commutation relation in Eq. (16) and $\hat{E} = \hat{H}$ we can show that the expectation value of the *tempus* operator \hat{T} is

$$\langle \psi(t) | \hat{T} \psi(t) \rangle = t + \langle \psi(0) | \hat{T} \psi(0) \rangle, \quad (23)$$

where $\langle \psi(0) | \hat{T} \psi(0) \rangle$ is a constant depending on the initial wave function. Therefore, Eq. (23) shows that the expectation value of the *tempus* operator is directly related to the time t of evolution of the system.

4 Examples

In this section, we explicitly construct the *tempus* operators for (1) a free particle and (2) a particle acted on by a constant force, and show that the energy-time uncertainty relation is satisfied.

4.1 Free particle

The *tempus* operator \hat{T} for a free particle is constructed from the corresponding classical expression, and the energy-time uncertainty relation is obtained.

For a free particle of mass m , the energy E in terms of the canonical momentum p is

$$E = p^2/2m. \quad (24)$$

Equation (24) can be solved for the canonical momentum $p = (2mE)^{1/2} \text{sgn}(p_0)$, where the sign of the initial momentum p_0 is $\text{sgn}(p_0)$. The generating function in Eq. (6) is

$$S(q, E, t) = (2mE)^{1/2}(q - q_0) \text{sgn}(p_0), \quad (25)$$

where we take the arbitrary function $S(q_0, E, t) = 0$. The new canonical coordinate *tempus* is obtained from Eq. (5) and is

$$T = (m/2E)^{1/2}(q - q_0) \text{sgn}(p_0). \quad (26)$$

In this way we have derived the canonical coordinate *tempus* conjugate to the energy, which is conceptually distinct from the time t in which the system evolves.

The solution to the classical problem is obtained by solving Eq. (26) for q and using the solutions of Eq. (11). We therefore obtain

$$q(t) = q_0 + v_0(t - t_0), \quad (27)$$

where the initial velocity $v_0 = p_0/m = (2E_0/m)^{1/2} \text{sgn}(p_0)$ and $T = t - t_0$.

In order to construct a self-adjoint *tempus* operator, Eq. (26) can be expressed in terms of p by substituting the energy in Eq. (24) into it. Then we obtain the canonical *tempus* coordinate for $q_0 = 0$ to be

$$T = mq/p. \quad (28)$$

Equation (28) may be quantized to obtain the *tempus* operator \hat{T} by replacing q and p by their corresponding operators \hat{q} and \hat{p} , and writing the operator in a symmetric way [14]-[16] [19, ?]

$$\hat{T} = \frac{1}{2}m(\hat{q}\hat{p}^{-1} + \hat{p}^{-1}\hat{q}). \quad (29)$$

The domain of \hat{T} is not the whole Hilbert space because of the presence of the operator \hat{p}^{-1} . Nevertheless, for the states in its domain, the operator \hat{T} has useful properties. The argument of Pauli [30] applies only to Hermitian operators defined on the whole Hilbert space. From the commutation relations in Eq. (15) for \hat{q} and \hat{p} , we can derive for \hat{T} in Eq. (29) and the operator \hat{E} corresponding to Eq. (24) the commutation relation in Eq. (16).

The uncertainty relation in Eq. (19) for energy and time can be obtained if we choose a wave function which is in the domain of \hat{E} , \hat{E}^2 , \hat{T} , and \hat{T}^2 . Since the *tempus* operator in Eq. (29) involves \hat{p}^{-1} , it is more convenient to work in momentum space where the realization of the commutation relation in Eq. (15) is $\hat{p} = p$ and $\hat{q} = i\hbar\partial/\partial p$. We choose a wave function in momentum space at time $t = 0$ to be

$$\phi(p, 0) = Np^n \exp(-a_0 p^2), \quad (30)$$

where a_0 is a real positive constant, n is an arbitrary integer ≥ 2 , and N is the normalization constant. Since the Schroedinger equation in momentum space is

$$\hat{H}(\hat{q}, p, t)\phi(p, t) = i\hbar\partial\phi(p, t)/\partial t, \quad (31)$$

the solution at time $t \neq 0$ with the Hamiltonian \hat{H} equal to the energy E in Eq. (24) is

$$\phi(p, t) = Np^n \exp(-ap^2), \quad (32)$$

where $a = a_0 + it/2m\hbar$ is a complex function of time. With the wave function in Eq. (32), the average of the *tempus* operator \hat{T} in Eq. (29) is

$$\langle \phi | \hat{T} \phi \rangle = \langle \hat{T} \rangle = t, \quad (33)$$

the time t in which the system evolves, as expected from Eq. (23).

The square of the uncertainty in the *tempus* operator \hat{T} at time t is

$$\begin{aligned} (\Delta T)^2 &\equiv \langle \phi | (\hat{T} - \langle \hat{T} \rangle)^2 \phi \rangle \\ &= \langle \hat{T}^2 \rangle - \langle \hat{T} \rangle^2 \\ &= (2a_0 m \hbar)^2 / (n - 3/2). \end{aligned} \quad (34)$$

The square of the uncertainty in the energy operator \hat{E} at time t is

$$\begin{aligned} (\Delta E)^2 &\equiv \langle \phi | (\hat{E} - \langle \hat{E} \rangle)^2 \phi \rangle \\ &= \langle \hat{E}^2 \rangle - \langle \hat{E} \rangle^2 \\ &= (n + 1/2)(4ma_0)^{-2}. \end{aligned} \quad (35)$$

From Eqs. (34) and (35) the product of ΔE and ΔT is

$$\Delta E \Delta T = [1 + 2/(n - 3/2)]^{1/2} \hbar/2 \geq \hbar/2 \quad \text{for } n \geq 2, \quad (36)$$

so the uncertainty relation is established.

As $n \rightarrow \infty$, Eq. (36) shows that $\Delta E \Delta T = \hbar/2$, i.e., equality holds in the uncertainty relation Eq. (36). The reason equality holds is that in the limit as $n \rightarrow \infty$, there is no distinction between $n+4, n+3, \dots, n-2$. Then we have

$$(i) \text{ as } n \rightarrow \infty \quad (\hat{T} - \langle \hat{T} \rangle) \phi = \gamma (\hat{E} - \langle \hat{E} \rangle) \phi, \quad (37)$$

where γ is a constant, and

(ii)

$$\begin{aligned} &\langle [(\hat{T} - \langle \hat{T} \rangle)(\hat{E} - \langle \hat{E} \rangle) + (\hat{E} - \langle \hat{E} \rangle)(\hat{T} - \langle \hat{T} \rangle)] \rangle \\ &= 2\text{Re} \langle (\hat{T} - \langle \hat{T} \rangle) \phi | (\hat{E} - \langle \hat{E} \rangle) \phi \rangle = 0, \end{aligned} \quad (38)$$

which are the two conditions for equality to hold in the uncertainty relation [56].

4.2 Particle acted on by a constant force

For a particle acted on by a constant force F_0 , the Hamiltonian H is

$$H = p^2/2m - F_0 q = E, \quad (39)$$

which is also the energy E . Equation (39) can be solved for the canonical momentum p ,

$$p = (2m)^{1/2}(E + F_0 q)^{1/2}. \quad (40)$$

When Eq. (40) is substituted into Eq. (6) for the generating function S , and the integral is performed, we obtain

$$S(q, E, t) = (2/3F_0)(2m)^{1/2} [(E + F_0 q)^{3/2} - (E + F_0 q_0)^{3/2}] + S(q_0, E, t). \quad (41)$$

The canonical coordinate *tempus* T conjugate to the energy is obtained from Eq. (5), which gives

$$T(q, E, t) = F_0^{-1}(2m)^{1/2}(E + F_0 q)^{1/2} + T_0, \quad (42)$$

where T_0 is a constant. Since the Hamiltonian H in Eq. (39) is conservative, Hamilton's equations (11) have the solution $T = t - t_0$ and $E = E_0$, where t_0 and E_0 are constants. If Eq. (42) is solved for q , we obtain

$$q(t) = q_0 + v_0 t + (1/2)at^2, \quad (43)$$

where $q_0 = q(0)$ is the initial position, $v_0 = \dot{q}(0)$ is the initial velocity, and $a = \ddot{q}(0) = F_0/m$ is the constant acceleration. Equation (43) is the solution to the classical problem.

In order to obtain a Hermitian *tempus* operator \hat{T} , it is necessary to express *tempus* in Eq. (42) in terms of q and p . When Eq. (39) for the energy is substituted into Eq. (42) for the canonical coordinate *tempus* we obtain $T = p/F_0$. When this expression is quantized, we obtain [23]

$$\hat{T} = \hat{p}/F_0, \quad (44)$$

where \hat{p} is the canonical momentum operator. Because the potential energy in Eq. (39) is not bounded from below, Pauli's argument [30] for the nonexistence of a Hermitian time operator does not apply to this example.

By a direct calculation we can show that the *tempus* operator \hat{T} in Eq. (44) and the energy operator \hat{E} corresponding to Eq. (39) satisfy the canonical commutation relation in Eq. (16). Therefore, from Eq. (17) the energy-time uncertainty relation in

Eq. (19) is also satisfied.

5 Conclusion

In this paper a new derivation of the energy-time uncertainty relation is given based on first principles. It is not necessary to make use of any *ad hoc* or heuristic assumptions. A classical canonical transformation is made from the old canonical variables (q, p) to new canonical variables (q', p') , where the new canonical momentum p' is chosen to be the energy E and the new canonical coordinate q' conjugate to the energy is called *tempus* T . This *tempus* canonical coordinate has the dimension of time, but is conceptually different from the time t in which the system evolves. In general, *tempus* T is a function of q, E and t . The energy E and *tempus* T have the same Poisson bracket as do p and q , since the Poisson bracket is a canonical invariant. When the system is quantized, the operators \hat{E} and \hat{T} satisfy the same commutation relations as do the operators \hat{p} and \hat{q} . Therefore, \hat{E} and \hat{T} satisfy the same uncertainty relation as do \hat{p} and \hat{q} . The energy-time uncertainty relation is therefore put on the same basis as the momentum-position uncertainty relation. The approach used here to derive the energy-time uncertainty relation has firm basis in classical mechanics, which makes a clear distinction between the time of evolution t and the canonical coordinate

tempus T conjugate to the energy. For a conservative classical system, *tempus* T is equal to the time t plus a constant. For a conservative quantum system, the average of the *tempus* operator \hat{T} is equal to the time t plus a constant.

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References

- [1] M. Bunge, *Can. J. Phys.* **48**, 1410 (1970).
- [2] V. V. Dodonov and V. I. Man'ko, "Generalization of the uncertainty relations in quantum mechanics" in *Invariants and the Evolution of Nonstationary Quantum Systems*, edited by M. A. Markov (Nova Science, Commack, N. Y., 1989), pp. 3-101.
- [3] H. C. Ohanian, *Principles of Quantum Mechanics* (Prentice-Hall, Englewood Cliffs, NJ, 1990), pp. 133-138.
- [4] D. B. Beard and G. B. Beard, *Quantum Mechanics with Applications* (Allyn and Bacon, Boston, 1970), pp. 19-21.
- [5] C. Cohen-Tannoudji, B. Diu and F. Laloë, *Quantum Mechanics*, Vol. 1 (John Wiley, New York, 1977), p. 252.
- [6] L. Mandelstam and I. G. Tamm, *J. Phys. U.S.S.R.* **9**, 249 (1945).
- [7] A. Messiah, *Quantum Mechanics*, Vol. I (North-Holland, Amsterdam, 1961), pp. 129-139, 319-320.
- [8] M. Grabowski, *Lett. Math. Phys.* **8**, 455 (1984).
- [9] K. Bhattacharyya, *J. Phys. A: Math. Gen.* **16**, 2993 (1983).
- [10] E. P. Wigner, "On the time-energy uncertainty relation", in *Aspects of Quantum Theory*, edited by A. Salam and E. P. Wigner (Cambridge, London, 1972), pp. 237-247.

- [11] L. F. Cook, *Am. J. Phys.* **48**, 142 (1980).
- [12] N. Bohr, "Discussion with Einstein on epistemological problems in atomic physics", in *Albert Einstein: Philosopher-Scientist*, edited by P. A. Schilpp (Open Court, LaSalle, Illinois, 1970), pp. 201-241.
- [13] F. Engelmann and E. Fick, *Nuovo Cim. Suppl.* **XII** (Sec. X) 63 (1959).
- [14] H. Paul, *Ann. Phys.* **9**, 252 (1962).
- [15] E. Fick and F. Engelmann, *Zeit. f. Phys.* **175**, 271 (1963)..
- [16] F. Engelmann and E. Fick, *Zeit. f. Phys.* **178**, 551 (1964).
- [17] M. Razavy, *Am. J. Phys.* **35**, 955 (1967).
- [18] M. Razavy, *Nuovo Cim.* **63B**, 271 (1969).
- [19] D. M. Rosenbaum, *J. Math. Phys.* **10**, 1127 (1969).
- [20] M. Bauer and P. A. Mello, *Ann. Phys. (N.Y.)* **111**, 38 (1978).
- [21] M. Bauer, *Ann. Phys. (N.Y.)* **150**, 1 (1983).
- [22] I. Fujiwara, *Prog. Theor. Phys.* **44**, 1701 (1970).
- [23] I. Fujiwara, *Prog. Theor. Phys.* **64**, 18 (1980).
- [24] I. Fujiwara, K. Wakita and H. Yoro, *Prog. Theor. Phys.* **64**, 363 (1980).
- [25] T. Goto, K. Yamaguchi and N. Sudo, *Prog. Theor. Phys.* **66**, 1525 (1981).
- [26] T. Goto, S. Naka and K. Yamaguchi, *Prog. Theor. Phys.* **66**, 1915 (1981).
- [27] E. Recami, "A time operator and the time-energy uncertainty relation", in *The Uncertainty Principle and Foundations of Quantum Mechanics*, edited by W. C. Price and S. S. Chissick (John Wiley, New York, 1977), pp. 21-28.
- [28] V. S. Olkhovskiy, E. Recami and A. J. Gerasimchuk, *Nuovo Cimento* **22A**, 263 (1974).
- [29] E. Papp, *Int. J. Theoret. Phys.* **10**, 385 (1974).

- [30] W. Pauli, "Die allgemeinen Prinzipien der Wellenmechanik", in *Encyclopedia of Physics*, Vol. V, Part 1, edited by S. Flügge (Springer-Verlag, Berlin, 1958), pp. 1-168. See especially footnote 1 on p. 60.
- [31] F. T. Smith, *Phys. Rev.* **118**, 349 (1960).
- [32] B. A. Lippmann, *Phys. Rev.* **151**, 1023 (1966).
- [33] G. R. Allcock, *Ann. Phys. (N.Y.)* **53**, 253 (1969).
- [34] G. R. Allcock, *Ann. Phys. (N.Y.)* **53**, 286 (1969).
- [35] G. R. Allcock, *Ann. Phys. (N.Y.)* **53**, 311 (1969).
- [36] M. Razavy, *Can. J. Phys.* **49**, 3075 (1971).
- [37] J. H. Eberly and L. P. S. Singh, *Phys. Rev. D* **7**, 359 (1973).
- [38] C. H. Blanchard, *Am. J. Phys.* **50**, 642 (1982).
- [39] E. Prugovecki, *Found. Phys.* **12**, 555 (1982).
- [40] R. Arshansky and L. P. Horwitz, *Found. Phys.* **15**, 701 (1985).
- [41] Y. S. Kim and M. E. Noz, *Found. Phys.* **9**, 375 (1979).
- [42] D. Han, M. E. Noz, Y. S. Kim and D. Son, *Phys. Rev. D* **27**, 3032 (1983).
- [43] P. E. Hussar, Y. S. Kim and M. E. Noz, *Am. J. Phys.* **53**, 142 (1985).
- [44] V. Fock and N. Krylov, *J. Phys., USSR*, **11**, 112 (1947).
- [45] Y. Aharonov and D. Bohm, *Phys. Rev.* **122**, 1649 (1961).
- [46] V. A. Fock, *Sov. Phys. - JETP* **15**, 784 (1962).
- [47] Y. Aharonov and D. Bohm, *Phys. Rev.* **134B**, 1417 (1964).
- [48] Y. Aharonov and J. L. Sajko, *Ann. Phys. (N.Y.)* **91**, 279 (1975).
- [49] Yu. I. Vorontsov, *Sov. Phys. - Uspekhi* **24**, 150 (1981).
- [50] M. Moshinsky, *Am. J. Phys.* **44**, 1037 (1976).

- [51] J. Rayski and J. M. Rayski, Jr., "On the meaning of the time-energy uncertainty relation", in *The Uncertainty Relation and Foundations of Quantum Mechanics*, edited by W. C. Price and S. S. Chissick (John Wiley, New York, 1977), pp. 13-20.
- [52] E. A. Gislason, N. H. Sabelli and J. W. Wood, *Phys. Rev. A* **31**, 2078 (1985).
- [53] H. Goldstein, *Classical Mechanics*, 2nd. ed. (Addison-Wesley, Reading, MA, 1980), pp. 378-401.
- [54] D. H. Kobe, *Am. J. Phys.* **56**, 252 (1988).
- [55] D. H. Kobe, "Canonical transformation to energy and 'Time' in classical mechanics", *Am. J. Phys.* (to be published).
- [56] E. Merzbacher, *Quantum Mechanics*, 2nd. ed. (John Wiley, New York, 1970), pp. 158-161.