

Measurements of the azimuthal anisotropy of prompt and nonprompt charmonia in PbPb collisions at

$$\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$$

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Abstract

The second-order (v_2) and third-order (v_3) Fourier coefficients describing the azimuthal anisotropy of prompt and nonprompt (from b-hadron decays) J/ψ , as well as prompt $\psi(2S)$ mesons are measured in lead-lead collisions at a center-of-mass energy per nucleon pair of $\sqrt{s_{\text{NN}}} = 5.02$ TeV. The analysis uses a data set corresponding to an integrated luminosity of 1.61 nb^{-1} recorded with the CMS detector. The J/ψ and $\psi(2S)$ mesons are reconstructed using their dimuon decay channel. The v_2 and v_3 coefficients are extracted using the scalar product method and studied as functions of meson transverse momentum and collision centrality. The measured v_2 values for prompt J/ψ mesons are found to be larger than those for nonprompt J/ψ mesons. The prompt J/ψ v_2 values at high p_T are found to be underpredicted by a model incorporating only parton energy loss effects in a quark-gluon plasma medium. Prompt and nonprompt J/ψ meson v_3 and prompt $\psi(2S)$ v_2 and v_3 values are also reported for the first time, providing new information about heavy quark interactions in the hot and dense medium created in heavy ion collisions.

Submitted to the Journal of High Energy Physics

1 Introduction

In high energy heavy ion collisions, a very strongly-interacting medium of deconfined quarks and gluons, known as the quark-gluon plasma (QGP), is created [1]. Quarkonia, bound states of a heavy quark and its antiquark, e.g., charmonia (J/ψ , $\psi(2S)$), and bottomonia ($Y(1S, 2S, 3S)$), are useful probes to study the properties of the QGP. Since heavy quarks are predominantly produced at the early stages of collisions by hard partonic scattering processes, they experience the whole space-time evolution of the medium. Models including static color screening [2, 3] and gluo-dissociation effects [4–6] inside a QGP medium predict a suppression of quarkonium yields in nucleus-nucleus (AA) collisions compared to proton-proton (pp) collisions [7]. On the other hand, quarkonia may also be created by recombination of quarks and antiquarks which were initially uncorrelated, thereby resulting in an enhancement of the measured yields [8].

Experimentally, the modification of particle yields in heavy ion collisions is quantified by the nuclear modification factor, defined as the ratio of yields in AA collisions to those in pp collisions scaled by the estimated number of binary nucleon-nucleon (NN) collisions. Results from the CERN SPS, BNL RHIC, and CERN LHC experiments show significant suppression of J/ψ mesons in AA collisions [9–19]. At the LHC, the ALICE Collaboration [17, 18] observed a weaker suppression of J/ψ mesons in the low transverse momentum (p_T) region below 3 GeV/ c than was seen by experiments at the SPS and RHIC. This result has been interpreted as a sign of enhanced recombination processes at LHC energies because of the larger production cross section of charm quarks [17, 18, 20, 21]. These measurements have provided information to constrain theoretical models that incorporate different in-medium effects on charmonium states in order to describe thermal characteristics of the QGP.

Another effective way to probe the dynamics of the QGP is the study of the azimuthal (ϕ) distribution of particles produced in heavy ion collisions. Azimuthal anisotropies of the particles can be characterized by Fourier coefficients (v_n) of the ϕ distribution with respect to the event plane, which corresponds to the direction of maximal particle density [22]. The presence of non-zero Fourier coefficients is referred to as collective flow by analogy to hydrodynamic models. In particular, the second-order (v_2) and the third-order (v_3) components are predominantly sensitive to initial collision geometry and event-by-event fluctuations, respectively [23]. Because of the initial anisotropic shape of the QGP in the transverse plane, quarkonia produced by initial hard-scattering processes experience different path lengths while traveling through the medium. Larger energy loss, and therefore larger suppression, can occur for quarkonia moving in directions corresponding to larger average path lengths. On the other hand, quarkonia created via recombination can have an asymmetric ϕ distribution as a result of asymmetries in quark densities [24, 25].

In LHC experiments, sizable v_2 values have been observed for J/ψ mesons in lead-lead (PbPb) collisions at center-of-mass energies per nucleon pair $\sqrt{s_{NN}} = 2.76$ and 5.02 TeV [26–29], indicating collective flow behavior of charm quarks. However, these results have limitations that prevent drawing firm conclusions on the origin of the v_2 for J/ψ mesons, given the large measurement uncertainties and the lack of separation between prompt and nonprompt J/ψ mesons. Measurements of nonprompt J/ψ mesons, which originate from bottom hadron decays, provide useful information on the propagation of bottom quarks in the QGP.

The CMS Collaboration has shown that in PbPb collisions the suppression of prompt J/ψ (of p_T above 6.5 GeV/ c) within jets depends upon the fragmentation [30]. Theoretical calculations suggest that jet quenching may drive the modulation of the ϕ distribution of J/ψ at high p_T [31]. These studies motivate the measurement of the azimuthal anisotropy for prompt J/ψ mesons at higher p_T to investigate the path length dependence of jet quenching [32, 33].

Azimuthal correlations for $\psi(2S)$ mesons are also a subject of interest. Prompt $\psi(2S)$ mesons have been found to be significantly more suppressed than prompt J/ψ mesons in PbPb collisions [19, 34]. Measurements of v_2 and v_3 values for prompt $\psi(2S)$ mesons are expected to provide constraints on in-medium effects for charmonium states with different binding energies in heavy ion collisions [21, 35].

This paper reports measurements of v_n values for prompt and nonprompt J/ψ mesons, as well as prompt $\psi(2S)$ meson, using PbPb data collected with the CMS detector at $\sqrt{s_{NN}} = 5.02$ TeV in 2018. The prompt and nonprompt J/ψ v_3 , and prompt $\psi(2S)$ v_n values are measured for the first time. The charmonium states are identified via their dimuon decay channel. The data set corresponds to an integrated luminosity of 1.61 nb^{-1} [36, 37], allowing the extraction of v_n values as functions of meson transverse momentum and PbPb event centrality (i.e., the degree of overlap of the two Pb ions). The numerical values of the results for this analysis are tabulated in the HEPDATA record [38].

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity (η) coverage provided by the barrel and endcap detectors. The forward hadron (HF) calorimeter uses steel as an absorber and quartz fibers as the sensitive material. The two halves of the HF are located 11.2 m from the interaction region, one on each end, and together they provide coverage in the range $3.0 < |\eta| < 5.2$. The HF calorimeters are composed of individual detector elements, or towers, having $\Delta\eta \times \Delta\phi = 0.175 \times 0.175$. Centrality is determined by the total transverse energy deposited in both of the HF calorimeters, and is defined as the fraction of the total hadronic inelastic nucleus-nucleus cross section, with 0% representing the largest overlap of the two colliding Pb nuclei [39].

Muons are measured in the range $|\eta| < 2.4$, with detection planes made using three technologies: drift tubes, cathode strip chambers, and resistive plate chambers. Matching muons to tracks measured in the silicon tracker results in a relative transverse momentum resolution, for muons with p_T up to 100 GeV/c, of 1% in the barrel and 3% in the endcaps. The p_T resolution in the barrel is better than 7% for muons with p_T up to 1 TeV/c [40]. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [41].

3 Event selection

Events of interest are selected using a two-tiered trigger system. The first level (Level-1), composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of about $4 \mu\text{s}$ [42]. The second level, known as the high-level trigger (HLT), consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage [43]. The trigger for this analysis was designed to have several stages. First, at least two muon candidates should be reconstructed at Level-1. At the HLT, the events are required to contain at least one Level-2 (L2) muon and another Level-3 (L3) muon, with requirements on their invariant mass being within $1\text{--}5 \text{ GeV}/c^2$. L2 muons are identified by matching the tracks to the hits in the outer muon spectrometer. For L3

muons, full tracks from the L2 muon tracks and the inner tracker information are used in the reconstruction, in addition requiring at least ten high-quality hits in the inner tracker [43].

To reject beam-related background processes, events are required to have at least one reconstructed primary vertex, and the clusters in the silicon pixel detector are required to be compatible with the vertex position. There must be at least two towers that contain energy deposits above 4 GeV in each of the HF detectors. The primary vertex, which is reconstructed from two or more tracks, is required to be located within 15 cm of the central point of the detector along the beam axis. In addition, the reconstructed muons are selected using a set of offline muon identification criteria [30]. To ensure the reconstruction efficiency of single muons is larger than 10%, muons within the following kinematic domains are selected:

$$\begin{aligned}
 p_{\text{T}}^{\mu} &> 3.5 \text{ GeV}/c && \text{for } |\eta^{\mu}| < 1.2, \\
 p_{\text{T}}^{\mu} &> (5.47 - 1.89|\eta^{\mu}|) \text{ GeV}/c && \text{for } 1.2 < |\eta^{\mu}| < 2.1, \text{ and} \\
 p_{\text{T}}^{\mu} &> 1.5 \text{ GeV}/c && \text{for } 2.1 < |\eta^{\mu}| < 2.4.
 \end{aligned} \tag{1}$$

4 Acceptance and efficiency corrections

Correction factors are applied to all results to account for the detector acceptance, as well as for the trigger, reconstruction, and selection efficiencies for $\mu^+\mu^-$ pairs. The correction factors are derived from simulated prompt and nonprompt J/ψ and $\psi(2S)$ meson samples which are then embedded in simulated PbPb collisions. Corrections are evaluated in the same bins of p_{T} , centrality, and rapidity (y) used in the v_n analyses. The meson samples are generated with PYTHIA 8.212 [44] using the CP5 underlying event tune [45], while the decay of bottom hadrons is simulated with EVTGEN 1.3.0 [46], which provides a better description of the kinematic distribution of nonprompt charmonia. These meson events are then embedded into simulated PbPb collision events generated using HYDJET 1.9 [47], and propagated through the CMS detector with the GEANT4 package [48]. The p_{T} distributions for simulated J/ψ and $\psi(2S)$ mesons are compared to those in data with fine p_{T} intervals, and the ratios of data over the simulation are used to reweight the simulation to better describe the data as a function of p_{T} . This weighting procedure accounts for possible mismodeling of the J/ψ and $\psi(2S)$ meson kinematics. Moreover, the reconstructed charmonia in data include feed-down contributions from heavier quarkonium state decays. These contributions are not explicitly considered in the simulation, but their effect on the kinematic distribution of the simulated charmonia is taken into account by the p_{T} reweighting procedure.

The acceptance in a given analysis bin is defined as the fraction of generated J/ψ and $\psi(2S)$ mesons in that bin that decay into two muons passing the kinematic requirements defined in Eq. (1), and reflects the geometrical coverage of the CMS detector. The value of the acceptance increases with dimuon p_{T} and varies from 0.05–0.76.

The efficiency in a given analysis bin is defined as the ratio of the number of reconstructed J/ψ and $\psi(2S)$ mesons from which both muons pass the selections described in Section 3 and the number of generated J/ψ mesons from which both muons pass the kinematic requirements defined in Eq. (1). Individual components of the efficiency (trigger, track reconstruction, and muon identification) are measured using single muons from J/ψ meson decays in both collision data and simulated samples, using a “tag-and-probe” (T&P) technique [49, 50]. The ratio in the efficiency between the data and simulated samples is used as a correction factor for the efficiency extracted from the simulation. The reciprocals of the acceptance and efficiency correction factors are used as event-by-event weights when the distributions of invariant mass and average v_n are computed.

5 Signal extraction

The signal candidates for J/ψ ($\psi(2S)$) mesons are extracted in the kinematic range of $6.5 < p_T < 50$ GeV/ c , $|y| < 2.4$, and centrality 0–90% (0–60%). The results are reported as functions of p_T and $\langle N_{\text{part}} \rangle$, where $\langle N_{\text{part}} \rangle$ is the average number of participating nucleons in a given centrality interval [39]. The centrality range for p_T -dependent results is limited to 10–60% to ensure a large anisotropy of the QGP. The most central collisions (0–10%) are excluded due to the small eccentricity in the initial geometry of the medium, while the most peripheral collisions (60–90%) are excluded because of the poor resolution in extracting the v_n coefficients. In addition, the lower p_T region is extended down to 3 GeV/ c (4 GeV/ c) for J/ψ ($\psi(2S)$) in the forward rapidity region. Details of the kinematic selection and $\langle N_{\text{part}} \rangle$ values for each centrality interval can be found in Tables 2 and 3 in Appendix A. The mass ranges of 2.6–3.5 and 3.3–4.1 GeV/ c^2 are used to study the J/ψ and $\psi(2S)$ mesons, respectively.

The separation of the prompt and nonprompt charmonium production components relies on the measurement of a secondary $\mu^+\mu^-$ vertex displaced from the primary collision vertex. The pseudo-proper decay length [51] is defined as $\ell_{J/\psi} = L_{xyz} m_{J/\psi} c / |p_{\mu\mu}|$, where L_{xyz} is the distance between the primary and dimuon vertices, $m_{J/\psi}$ is the world average value [52] of the J/ψ meson mass (assumed for all dimuon candidates), and $p_{\mu\mu}$ is the dimuon momentum. To extract prompt and nonprompt J/ψ meson yields, the invariant mass spectrum of $\mu^+\mu^-$ pairs and their $\ell_{J/\psi}$ distribution are fitted using a two-dimensional (2D) extended unbinned maximum likelihood fit. The parameters of different components of the 2D probability density function (pdf) are obtained through fits to the invariant mass and the $\ell_{J/\psi}$ distributions, which are further used in the final 2D fits. These fits are performed for each v_n interval in the measured kinematic region. The evaluation of v_n coefficients is detailed in Section 6. Figure 1 shows the dimuon invariant mass and $\ell_{J/\psi}$ projection of the 2D fit in the ranges $3.0 < p_T < 4.5$ GeV/ c and $0.0 < v_2 < 0.3$ for the 10–60% centrality interval.

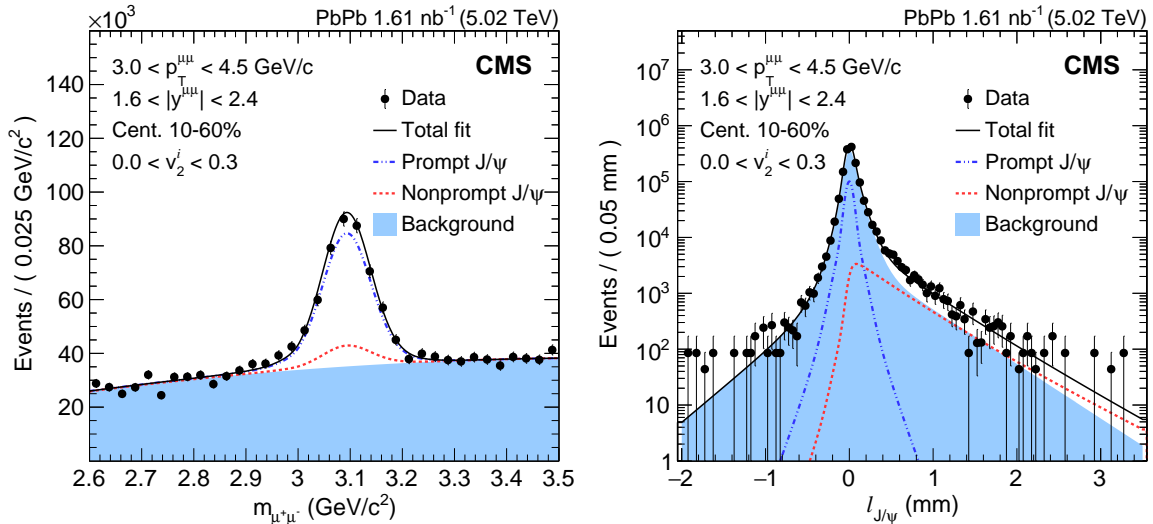


Figure 1: Invariant mass (left) and $\ell_{J/\psi}$ (right) distributions for the bin of $3.0 < p_T < 4.5$ GeV/ c and $0.0 < v_2 < 0.3$ in centrality 10–60%. The solid lines represent the total fit, while the dashed and dash-dotted lines represent the prompt and nonprompt components, respectively. The background contributions are shown by the filled blue histograms.

Sums of two Crystal Ball (CB) functions [53], with different widths but common mean and tail parameters, are used to extract the signal yield values from the invariant mass distribution. Two separate sums of CB functions are used to describe the prompt and nonprompt contri-

butions to the J/ψ signal. The tail parameters, as well as relative contributions of the two CB functions and the ratio of their widths, are fixed to the values obtained from simulation. The background component for the invariant mass distribution is described by a Chebyshev polynomial function of order N , where N is determined by performing the log-likelihood ratio (LLR) test [54]. This procedure is done for each analysis bin, while the tail and width ratio parameters are kept fixed. The polynomial order is chosen as the minimum so that increasing it does not significantly improve the quality of the fit. In most analysis bins, the first-order Chebyshev polynomial function is used as the background function.

Prompt and nonprompt signals and background components of $\ell_{J/\psi}$ distributions are parameterized using collision data and simulated samples. The *sPlot* technique [55] is used to obtain $\ell_{J/\psi}$ distributions of J/ψ signal and background from data, which have different invariant mass distributions. The correlation between the mass and $\ell_{J/\psi}$ is found to be very small, such that its impact on the final results is negligible. This justifies the use of the dimuon mass distribution to separate signal and background distributions of $\ell_{J/\psi}$ using the *sPlot* [30] technique. The $\ell_{J/\psi}$ resolution is obtained by fitting the negative tail of the J/ψ signal $\ell_{J/\psi}$ distribution in data using a sum of two Gaussian functions. Note that $\ell_{J/\psi}$ can have negative values as a result of the finite detector resolution. The shape of the $\ell_{J/\psi}$ distributions of prompt J/ψ signals is described by the $\ell_{J/\psi}$ resolution. The $\ell_{J/\psi}$ distribution of nonprompt J/ψ signals is parameterized by the sum of three exponential functions convoluted with the resolution function. The parameters of the three exponential functions are determined by a fit of the $\ell_{J/\psi}$ distribution of simulated nonprompt J/ψ mesons at the generator level. The background $\ell_{J/\psi}$ distribution in data is fit to a sum of exponential functions. To extract the yields of prompt and nonprompt J/ψ mesons, a fit to the $m_{\mu^+\mu^-}$ and $\ell_{J/\psi}$ distributions is performed, where the free parameters include the J/ψ meson yield, nonprompt J/ψ meson fraction, and background yield. A more detailed description of the fitting procedure can be found in Ref. [19].

For $\psi(2S)$ mesons, the nonprompt component is reduced by placing tight constraints on the $\ell_{J/\psi}$ distribution. The constraint, studied in simulated samples, rejects the nonprompt $\psi(2S)$ contamination while maintaining a prompt $\psi(2S)$ purity of at least 90%. The remaining contribution of nonprompt $\psi(2S)$ mesons ranges from 5 to 10% and is assigned as a systematic uncertainty for the nonprompt $\psi(2S)$ contamination.

6 Extraction of v_n

The v_n ($n = 2$ and 3) values of J/ψ and $\psi(2S)$ candidates are determined using the scalar product (SP) method [56]. The Q-vectors are defined in the complex plane $Q_n = \sum_{k=1}^M \omega_k e^{in\phi_k}$, obtained using the tracker or calorimeters. Here, M is the multiplicity of particles in the tracker or the number of towers for HF; ϕ is the azimuthal angle of the particle or the tower; ω_k is the weighting factor, that is the p_T of a particle for the tracker or the transverse energy deposited in an HF tower. In this analysis, Q-vectors from three subevents are calculated using the tracker at the mid-pseudorapidity ($|\eta| < 0.75$) and similarly the two HF calorimeters covering the forward ($3 < \eta < 5$, HF+) and backward ($-5 < \eta < -3$, HF-) regions. Two different procedures are applied in extracting v_2 of J/ψ and $\psi(2S)$. The procedure for extracting $\psi(2S)$ v_2 has been established in previous publications [56], while a new method is applied to the extraction of J/ψ v_2 with the advantage of avoiding making assumptions of background candidate v_2 . The

v_n coefficient for J/ψ or $\psi(2S)$ mesons is obtained as follows:

$$v_n \{SP\} \equiv \frac{\langle Q_n^{J/\psi, \psi(2S)} Q_{nA}^* \rangle}{\sqrt{\frac{\langle Q_{nA} Q_{nB}^* \rangle \langle Q_{nA} Q_{nC}^* \rangle}{\langle Q_{nB} Q_{nC}^* \rangle}}}. \quad (2)$$

The Q-vector of the J/ψ or $\psi(2S)$ candidate is defined as $Q_n^{J/\psi} = e^{in\phi}$, where ϕ is the azimuthal angle of the candidate. The subscripts A and B refer to either HF+ or HF−, depending on the rapidity of the J/ψ or $\psi(2S)$ candidate. Flattening and recentering procedures are applied to the Q-vectors relative to HF and the tracker for removing detector acceptance effects. To avoid short-range non-collective correlations, the η gap between the J/ψ or $\psi(2S)$ candidate and the detector used for the subevent determination is required to be at least three units of rapidity [56–58]. For this purpose, HF+ is selected for A (B) when J/ψ and $\psi(2S)$ candidates are produced at negative (positive) rapidity. The subscript C denotes the subevents taken from the tracker. The denominator in Eq. (2) is the correction factor to remove the finite resolution effect of the detectors and finite final-state multiplicity.

To extract the prompt and nonprompt J/ψ meson v_n values, the J/ψ candidates are classified into fine v_n intervals of their flow vector SPs (i.e., Eq. (2) but not averaged over all J/ψ candidates). Then, a simultaneous invariant mass and $\ell_{J/\psi}$ fit, as described in Section 5, is performed in each v_n interval to extract the corresponding inclusive J/ψ meson yield and the nonprompt J/ψ meson fraction, and to obtain the SP distribution of both prompt and nonprompt J/ψ signal. The prompt J/ψ meson v_n values can be obtained by:

$$v_n^{\text{prompt}} = \frac{\sum (1 - f_i) v_n^i Y^i}{\sum (1 - f_i) Y^i}, \quad (3)$$

while the nonprompt J/ψ meson v_n values can be obtained by:

$$v_n^{\text{nonprompt}} = \frac{\sum f_i v_n^i Y^i}{\sum f_i Y^i}. \quad (4)$$

Here v_n^i is the center of the i -th SP bin, while f_i and Y^i are the nonprompt J/ψ meson fraction and inclusive J/ψ meson yield in the same bin, respectively.

A different method is used to extract the prompt $\psi(2S)$ meson flow harmonics. The average Q-vectors of all prompt $\psi(2S)$ meson candidates are divided into bins of invariant mass, as shown in the lower panel of Fig. 2. For every invariant mass bin, the prompt $\psi(2S)$ meson v_n coefficient is calculated using Eq. (2). The $\psi(2S)$ candidate invariant mass spectrum and v_n distribution as a function of invariant mass are fit simultaneously using binned χ^2 fits. The invariant mass spectrum is fit with the sum of two CB functions and a Chebyshev polynomial function, as described in Section 5. The v_n distribution is fit to the following formula:

$$v_n^{\text{Sig+Bkg}}(m_{\text{inv}}) = \alpha(m_{\text{inv}}) v_n^{\psi(2S)} + [1 - \alpha(m_{\text{inv}})] v_n^{\text{Bkg}}(m_{\text{inv}}), \quad (5)$$

where $v_n^{\psi(2S)}$ is the prompt $\psi(2S)$ meson v_n , which is independent of m_{inv} by definition; $v_n^{\text{Bkg}}(m_{\text{inv}})$ is the background v_n modeled as a first-order polynomial function of invariant mass; and $\alpha(m_{\text{inv}})$ is the $\psi(2S)$ signal fraction as a function of invariant mass:

$$\alpha(m_{\text{inv}}) = \frac{\text{Sig}_{\psi(2S)}(m_{\text{inv}})}{[\text{Sig}_{\psi(2S)}(m_{\text{inv}}) + \text{Bkg}(m_{\text{inv}})]}. \quad (6)$$

An example of a simultaneous fit to the mass distribution (upper panel) and v_n profile (lower panel) for $\psi(2S)$ mesons is shown in Fig. 2 for a representative analysis bin.

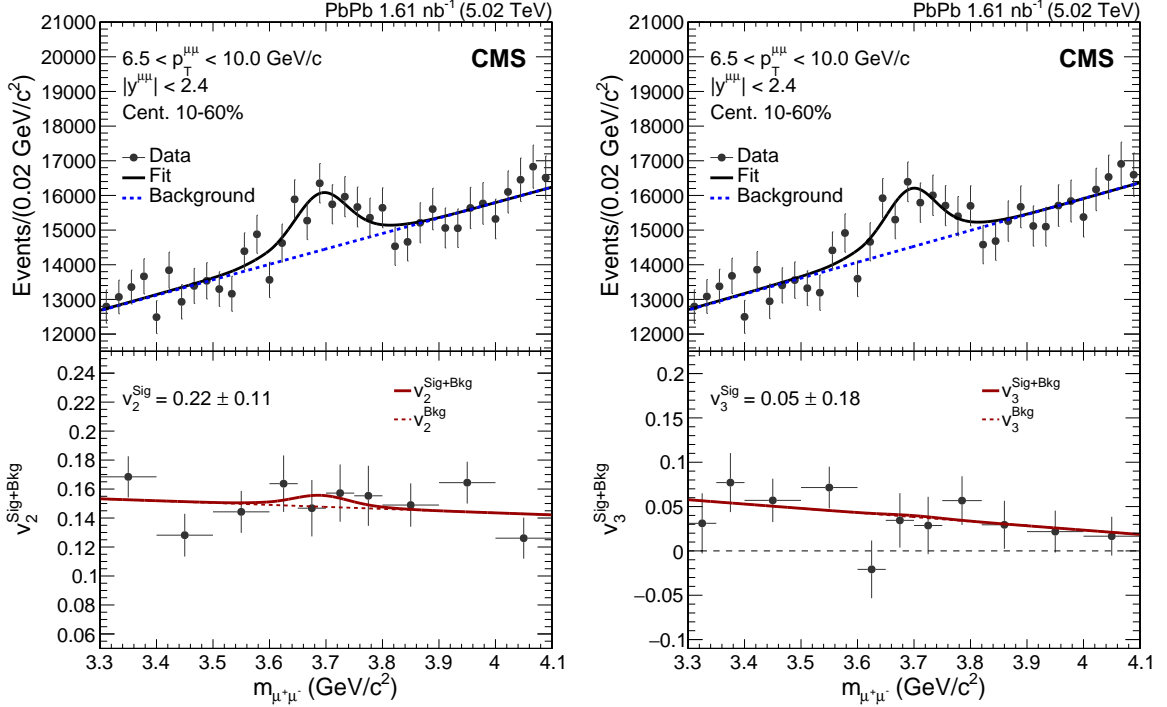


Figure 2: Simultaneous fits of the dimuon invariant mass spectrum and the $v_2^{\text{Sig+Bkg}}$ (left) and $v_3^{\text{Sig+Bkg}}$ (right) distributions for $\psi(2S)$ mesons, as defined by Eq. (5), for $6.5 < p_T < 10$ GeV/c and centrality 10–60%. In the upper panel, the solid black (signal + background) and dashed blue (background only) lines show the result of the mass fit. In the lower panel, the vertical lines of the v_n^{Sig} points represent the statistical uncertainty. The solid (v_n signal + background) and the dashed (v_n background) lines indicate the corresponding fit to the v_n distributions.

7 Systematic uncertainties

Systematic uncertainties in this analysis originate from various sources including the modeling of signal and background pdfs, and acceptance and efficiency corrections (trigger, track reconstruction, and muon identification). For uncertainties from the signal pdf, two factors are considered. The first factor is the choice of the signal pdf, and the second is the values of fixed parameters in the invariant mass fit. To evaluate uncertainty on the choice of two CB fit functions, an alternative pdf is formed by adding one CB function and one Gaussian function. The deviation of v_n to the nominal one is quoted as the uncertainty. Uncertainties from the fixed signal shape parameters are studied by releasing each parameter one at a time. The root-mean-square (RMS) of differences between the nominal and altered v_n results is assigned as the uncertainty. For the uncertainty in the background pdf, an exponential of a Chebyshev polynomial is considered as an alternative pdf for the invariant mass distribution. The variation in the v_n results is quoted as the systematic uncertainty.

The uncertainty due to the acceptance correction is evaluated by comparing results with and without p_T reweighting of the signal shape in simulated samples. The variation of efficiencies due to the T&P correction is considered. The change of results by varying the T&P correction within its uncertainty is quoted as an uncertainty in the final v_n results. In addition, the uncertainty due to hadronic event selection is considered. The collision event filter is varied, thus allowing for a migration of J/ψ and $\psi(2S)$ mesons across centrality boundaries. The effect on measured v_n values is taken as a systematic uncertainty. Then, the differences of resulting v_n values from the nominal ones are considered as the systematic uncertainty for the event

selection.

Different components of the $\ell_{J/\psi}$ pdf are varied to estimate the uncertainty in the nonprompt J/ψ meson fraction. The total (signal+background) distribution is used instead of individual distributions for the $\ell_{J/\psi}$ uncertainty distribution. The $\ell_{J/\psi}$ resolution obtained from prompt J/ψ meson simulation is used instead of that evaluated from data. The parameters used in the function to describe the bottom hadron decay length shape of the nonprompt J/ψ meson simulations are released in the final fit to data. Four exponential decay functions are also considered instead of three. The maximum deviation between each pair of two variations is taken as an uncertainty for bottom hadron decay length distribution of the nonprompt J/ψ . One additional exponential function on each side of the $\ell_{J/\psi}$ distribution of the background is used for the systematic uncertainty estimation. The RMS of these four sources is taken as the systematic uncertainty for the nonprompt J/ψ fraction.

For the $v_2^{\text{Bkg}}(m_{\text{inv}})$ modeling of $\psi(2S)$ mesons, a second-order polynomial is used as an alternative function. The deviation of v_n results from nominal ones is assigned as the systematic uncertainty. Uncertainties due to the nonprompt component in the $\psi(2S)$ meson v_n are derived from the deviation of v_n by using a looser $\ell_{J/\psi}$ requirement that has twice the residual efficiency for nonprompt $\psi(2S)$ mesons compared to the nominal case. The total uncertainty, calculated for each bin of p_T or $\langle N_{\text{part}} \rangle$, is determined by taking the quadratic sum of each individual source of uncertainty. In Table 1, the total systematic uncertainty is summarized along with its dominant sources. The ranges provided are illustrative and represent the absolute values of the minimum and maximum uncertainties averaged over all p_T and $\langle N_{\text{part}} \rangle$ bins.

Table 1: Summary of the absolute total systematic uncertainty and its dominant contributions, in units of 10^{-3} .

Uncertainty source	Prompt J/ψ		Nonprompt J/ψ		Prompt $\psi(2S)$	
	v_2	v_3	v_2	v_3	v_2	v_3
Total systematic	1.3–5.2	2.2–5.3	1.9–8.1	3.3–9.1	32.0–63.0	30.0–82.0
Signal parameters	0.4–3.0	0.1–3.0	0.9–5.0	0.1–3.0	3.0–10.0	10.0–30.0
Event selection	0.1–5.0	0.2–3.0	0.9–6.0	0.3–2.0	10.0–30.0	5.0–20.0
Nonprompt J/ψ fraction	0.3–4.0	0.8–4.0	0.4–6.0	2.0–9.0	—	—
Nonprompt $\psi(2S)$ contam.	—	—	—	—	20.0–50.0	10.0–40.0
Background v_n	—	—	—	—	1.0–20.0	0.7–50.0

8 Results

The v_2 and v_3 coefficients measured for prompt and nonprompt J/ψ mesons in PbPb collisions are shown in Fig. 3, as functions of p_T and $\langle N_{\text{part}} \rangle$. There is at most a weak p_T dependence observed for the prompt and nonprompt within their respective uncertainties. Results in Fig. 3 also show sizable v_2 values for prompt J/ψ mesons even at high p_T , which are consistent with the v_2 of charged hadrons in the high p_T region [32, 33].

The measured v_2 values for nonprompt J/ψ mesons are smaller than those for prompt J/ψ mesons in the studied region. Such a difference could not have been observed at $\sqrt{s_{\text{NN}}} = 2.76$ TeV [28] because of the limited amount of data. These results are compatible with measurements from the ALICE and ATLAS Collaborations [27, 29], and extend to a higher p_T region. A similar finding has been observed for prompt and nonprompt D^0 [59]. These observations can be contrasted to those found for bottomonium measurements, in which the v_2 values for Y mesons are found to be consistent with zero [56, 60].

In Fig. 3 (right), the v_2 values for prompt and nonprompt J/ψ mesons are found to be smallest in the most central PbPb collision events (0–10%), which corresponds to $\langle N_{\text{part}} \rangle > 356$, and the v_2 values increase toward mid-central collision events. This is expected from the eccentricity in the initial collision geometry of these collisions. The lower panel of the same figure shows measured v_3 values for prompt and nonprompt J/ψ components. No significant nonzero v_3 values are found in studied kinematic intervals.

In Fig. 4, prompt J/ψ meson v_2 values are compared to a theoretical prediction [31], where high- p_T prompt J/ψ are produced from the jet fragmentation. The medium response of jets in PbPb collisions is simulated using the linear Boltzmann transport framework and the v_2 values are calculated as a consequence of path length dependent energy loss of jets in the QGP. Note that the model is calculated with the centrality interval 10–30%, whereas the data is for the centrality interval 10–60%. A finite v_2 value is predicted by the model calculation, suggesting that path length dependent energy loss of jets in this p_T range needs to be taken into account to describe the prompt J/ψ meson v_2 . However, the magnitude of v_2 values in the data is found to be underpredicted by the model. As no strong centrality dependence in mid-central collision events is observed in data, this difference between the observed v_2 and theory calculation can not be explained by the different centrality intervals, indicating that additional contributions are needed to describe the prompt J/ψ meson v_2 at high p_T .

Figure 5 shows the first v_n measurements of the azimuthal anisotropy for prompt $\psi(2S)$ mesons in heavy ion collisions together with the results for prompt J/ψ mesons. The v_2 values are found to be slightly larger for the prompt $\psi(2S)$ than for the prompt J/ψ mesons, especially at higher p_T and in peripheral PbPb collisions, but the statistical significance of this difference is too low to draw any conclusion, with p-values of at best 10%. The v_3 values are found to be consistent with zero in the measured bins.

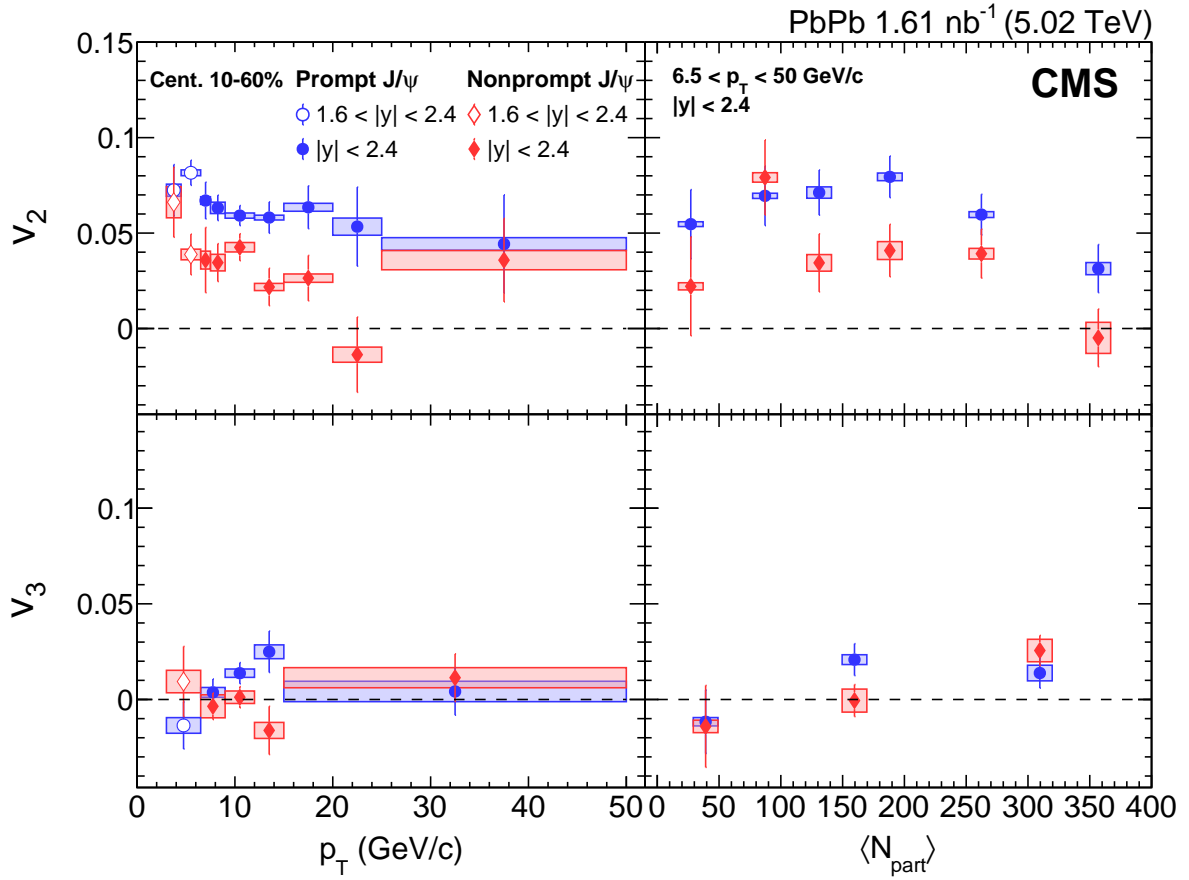


Figure 3: The v_2 (upper) and v_3 (lower) values, as functions of p_T (left) and $\langle N_{\text{part}} \rangle$ (right) for prompt and nonprompt J/ψ mesons. The results for $3 < p_T < 6.5$ and $6.5 < p_T < 50$ GeV/ c are studied in the rapidity range of $1.6 < |y| < 2.4$ and $|y| < 2.4$, respectively (left). The kinematic range for the right panel is $6.5 < p_T < 50$ GeV/ c and $|y| < 2.4$. The data points are positioned in the middle of each p_T bin. The vertical bars denote the statistical uncertainties and the rectangular bands show the systematic uncertainties.

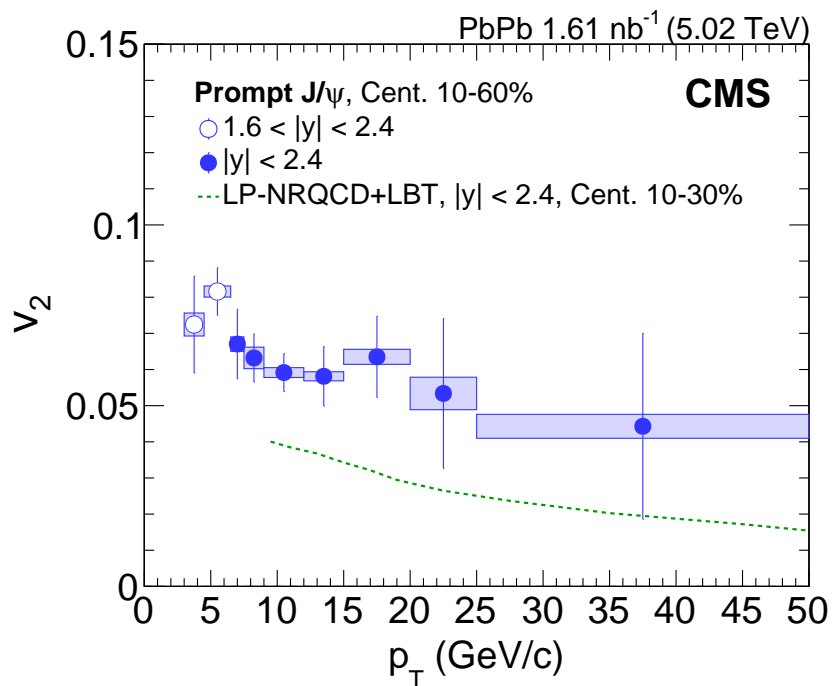


Figure 4: The v_2 values of prompt J/ψ mesons (blue circles) as a function of p_T in the 10–60% centrality range compared with a model calculation in Ref. [31] (green dashed line). The results for $3 < p_T < 6.5$ and $6.5 < p_T < 50$ GeV/ c are studied in the rapidity range of $1.6 < |y| < 2.4$ and $|y| < 2.4$, respectively. The vertical bars denote the statistical uncertainties and the rectangular bands show the systematic uncertainties.

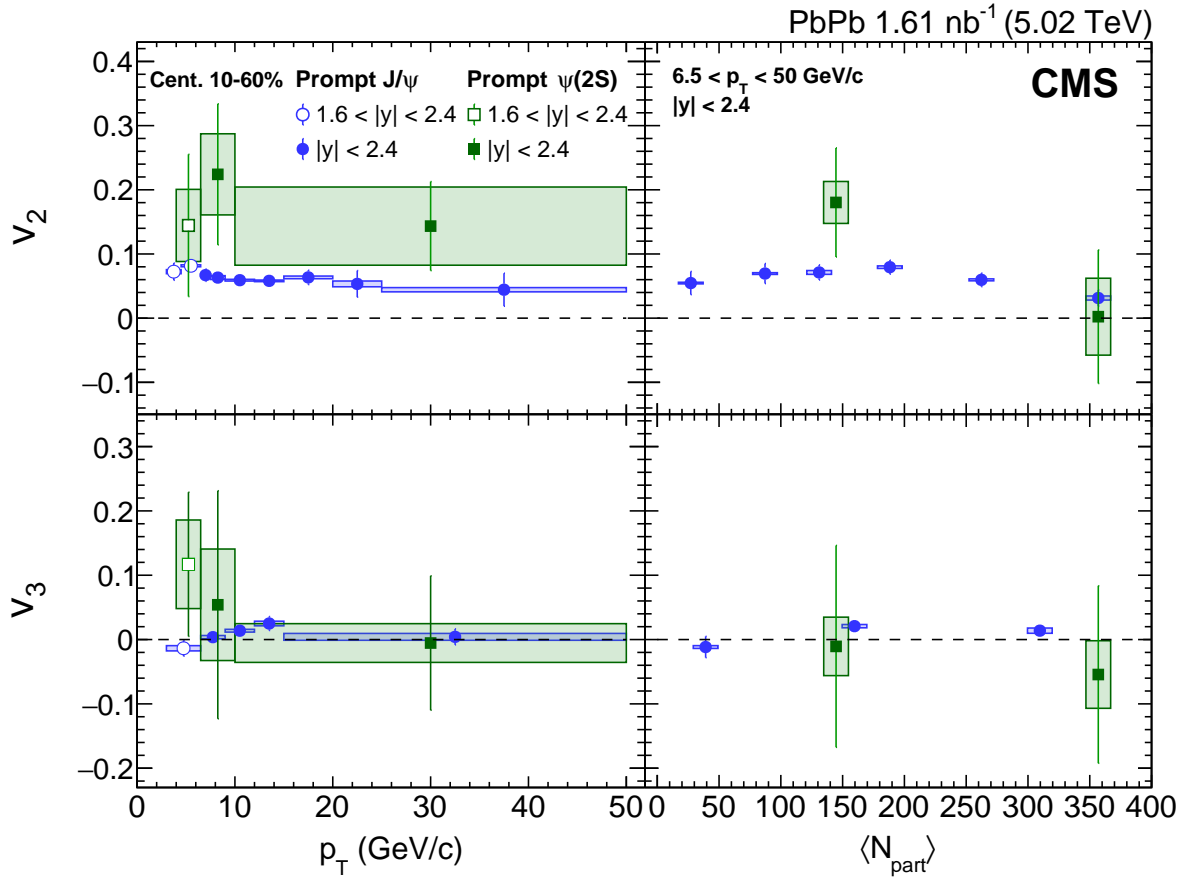


Figure 5: The v_2 (upper) and v_3 (lower) values, as functions of p_T (left) and $\langle N_{\text{part}} \rangle$ (right) for prompt J/ψ (blue circles) and prompt ψ(2S) (green squares) mesons. The results for $p_T < 6.5$ and $p_T > 6.5$ GeV/c are studied in the rapidity range of $1.6 < |y| < 2.4$ and $|y| < 2.4$, respectively (left). The kinematic range for the right panel is $6.5 < p_T < 50$ GeV/c and $|y| < 2.4$. The data points are positioned in the middle of each p_T bin. The vertical bars denote the statistical uncertainties and the rectangular bands show the systematic uncertainties.

9 Summary

The second-order (v_2) and third-order (v_3) Fourier coefficients of azimuthal distributions for prompt and nonprompt J/ψ , and prompt $\psi(2S)$ mesons are measured in PbPb collisions as functions of transverse momentum (p_T) and event centrality. The v_2 values for prompt and nonprompt J/ψ mesons both indicate a decreasing trend from mid-central towards central collision events. Within present experimental uncertainties, these v_2 values show no clear p_T dependence between 3 and 50 GeV/c. The prompt J/ψ meson v_2 values are found to be larger than those of nonprompt J/ψ mesons throughout the studied kinematic region, suggesting different in-medium effects for charm and bottom quarks. The prompt J/ψ meson v_2 values are larger than those from a model calculation which considered the prompt J/ψ production above 10 GeV/c from jet fragmentation. This discrepancy indicates additional effects are required to describe the observed sizable v_2 at high p_T for prompt J/ψ mesons. The v_3 values for prompt and nonprompt J/ψ mesons are reported for the first time, and found to be consistent with zero in the measured p_T and centrality intervals. The v_2 and v_3 values are also measured for prompt $\psi(2S)$ mesons and are compatible with, although slightly larger than, those of prompt J/ψ mesons. These J/ψ and $\psi(2S)$ meson measurements provide new information about heavy quark interactions in the hot and dense medium created in heavy ion collisions.

Acknowledgments

We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centers and personnel of the Worldwide LHC Computing Grid and other centers for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC, the CMS detector, and the supporting computing infrastructure provided by the following funding agencies: BMBWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, FAPERGS, and FAPESP (Brazil); MES and BNSF (Bulgaria); CERN; CAS, MoST, and NSFC (China); Minciencias (Colombia); MSES and CSF (Croatia); RIF (Cyprus); SENESCYT (Ecuador); MoER, ERC PUT and ERDF (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRI (Greece); NKFIH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); MSIP and NRF (Republic of Korea); MES (Latvia); LAS (Lithuania); MOE and UM (Malaysia); BUAP, CINVESTAV, CONACYT, LNS, SEP, and UASLP-FAI (Mexico); MOS (Montenegro); MBIE (New Zealand); PAEC (Pakistan); MES and NSC (Poland); FCT (Portugal); MESTD (Serbia); MCIN/AEI and PCTI (Spain); MOSTR (Sri Lanka); Swiss Funding Agencies (Switzerland); MST (Taipei); MHESI and NSTDA (Thailand); TUBITAK and TENMAK (Turkey); NASU (Ukraine); STFC (United Kingdom); DOE and NSF (USA).

Individuals have received support from the Marie-Curie program and the European Research Council and Horizon 2020 Grant, contract Nos. 675440, 724704, 752730, 758316, 765710, 824093, 884104, and COST Action CA16108 (European Union); the Leventis Foundation; the Alfred P. Sloan Foundation; the Alexander von Humboldt Foundation; the Belgian Federal Science Policy Office; the Fonds pour la Formation à la Recherche dans l'Industrie et dans l'Agriculture (FRIA-Belgium); the Agentschap voor Innovatie door Wetenschap en Technologie (IWT-Belgium); the F.R.S.-FNRS and FWO (Belgium) under the "Excellence of Science – EOS" – be.h project n. 30820817; the Beijing Municipal Science & Technology Commission, No. Z191100007219010; the Ministry of Education, Youth and Sports (MEYS) of the Czech

Republic; the Hellenic Foundation for Research and Innovation (HFRI), Project Number 2288 (Greece); the Deutsche Forschungsgemeinschaft (DFG), under Germany's Excellence Strategy – EXC 2121 “Quantum Universe” – 390833306, and under project number 400140256 - GRK2497; the Hungarian Academy of Sciences, the New National Excellence Program - ÚNKP, the NK-FIH research grants K 124845, K 124850, K 128713, K 128786, K 129058, K 131991, K 133046, K 138136, K 143460, K 143477, 2020-2.2.1-ED-2021-00181, and TKP2021-NKTA-64 (Hungary); the Council of Science and Industrial Research, India; the Latvian Council of Science; the Ministry of Education and Science, project no. 2022/WK/14, and the National Science Center, contracts Opus 2021/41/B/ST2/01369 and 2021/43/B/ST2/01552 (Poland); the Fundação para a Ciência e a Tecnologia, grant CEECIND/01334/2018 (Portugal); the National Priorities Research Program by Qatar National Research Fund; MCIN/AEI/10.13039/501100011033, ERDF “a way of making Europe”, and the Programa Estatal de Fomento de la Investigación Científica y Técnica de Excelencia María de Maeztu, grant MDM-2017-0765 and Programa Severo Ochoa del Principado de Asturias (Spain); the Chulalongkorn Academic into Its 2nd Century Project Advancement Project, and the National Science, Research and Innovation Fund via the Program Management Unit for Human Resources & Institutional Development, Research and Innovation, grant B05F650021 (Thailand); the Kavli Foundation; the Nvidia Corporation; the SuperMicro Corporation; the Welch Foundation, contract C-1845; and the Weston Havens Foundation (USA).

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A Kinematic regions

The kinematic ranges for J/ψ and $\psi(2S)$ mesons used in this study are summarized in Table 2. The centrality and $\langle N_{\text{part}} \rangle$ values used are calculated using a Glauber model simulation [39] and are summarized in Table 3.

Table 2: A Summary of the kinematic regions for J/ψ and $\psi(2S)$ in the centrality interval 10–60%.


	J/ψ	$\psi(2S)$
$ y < 2.4$	$6.5 < p_T < 50 \text{ GeV}/c$	$6.5 < p_T < 50 \text{ GeV}/c$
$1.6 < y < 2.4$	$3 < p_T < 6.5 \text{ GeV}/c$	$4 < p_T < 6.5 \text{ GeV}/c$

Table 3: The centrality classes and $\langle N_{\text{part}} \rangle$ values for PbPb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV used in this study.




















Centrality class	$\langle N_{\text{part}} \rangle$
0–10%	356.9 ± 0.9
10–20%	262.3 ± 1.3
20–30%	188.2 ± 1.4
30–40%	131.0 ± 1.4
40–50%	87.1 ± 1.3
50–90%	27.1 ± 0.6
0–20%	309.6 ± 1.0
20–40%	159.6 ± 1.4
40–90%	39.1 ± 0.7
10–60%	144.6 ± 1.2

B The CMS Collaboration




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





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

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D. Dobur , J. Knolle , L. Lambrecht , G. Mestdach, M. Niedziela , C. Rendón, C. Roskas , A. Samalan, K. Skovpen , M. Tytgat , N. Van Den Bossche , B. Vermassen, L. Wezenbeek 






















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


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



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

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






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










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



























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



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













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






















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
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












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
























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





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






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







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




























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


























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
















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G.M. Dallavalle^a , T. Diotallevi^{a,b} , F. Fabbri^a , A. Fanfani^{a,b} , P. Giacomelli^a ,
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S. Marcellini^a , G. Masetti^a , F.L. Navarria^{a,b} , A. Perrotta^a , F. Primavera^{a,b} ,
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P. Lenzi^{a,b} , M. Lizzo^{a,b} , M. Meschini^a , S. Paoletti^a , R. Seidita^{a,b} , G. Sguazzoni^a ,
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






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

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










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



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













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

































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



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



































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

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S. Lee , B.H. Oh , S.B. Oh , H. Seo , U.K. Yang , I. Yoon 







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


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





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




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







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














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
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








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






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









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





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





















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











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



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












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


















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M. Bachtis , R. Cousins , A. Datta , D. Hamilton , J. Hauser , M. Ignatenko , M.A. Iqbal , T. Lam , E. Manca , W.A. Nash , S. Regnard , D. Saltzberg , B. Stone , V. Valuev 























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R. Clare , J.W. Gary , M. Gordon , G. Hanson , G. Karapostoli , O.R. Long , N. Manganeli , W. Si , S. Wimpenny 















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









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N. Amin , C. Campagnari , M. Citron , G. Collura , A. Dorsett , V. Dutta , J. Incandela , M. Kilpatrick , J. Kim , A.J. Li , P. Masterson , H. Mei , M. Oshiro , M. Quinnan , J. Richman , U. Sarica , R. Schmitz , F. Setti , J. Sheplock , P. Siddireddy , D. Stuart , S. Wang 







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





























































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A. Perloff , C. Savard , N. Schonbeck , K. Stenson , K.A. Ulmer , S.R. Wagner ,
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H.W.K. Cheung , F. Chlebana , K.F. Di Petrillo , J. Dickinson , V.D. Elvira , Y. Feng ,
J. Freeman , A. Gandrakota , Z. Gecse , L. Gray , D. Green , S. Grünendahl ,
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L. Horyn , B. Jayatilaka , S. Jindariani , M. Johnson , U. Joshi , T. Klijsma ,
B. Klima , K.H.M. Kwok , S. Lammel , D. Lincoln , R. Lipton , T. Liu , C. Madrid ,
K. Maeshima , C. Mantilla , D. Mason , P. McBride , P. Merkel , S. Mrenna ,
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



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P. Avery , D. Bourilkov , L. Cadamuro , V. Cherepanov , R.D. Field , M. Kim ,
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
















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








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M.M. Baarmand , S. Butalla , T. Elkafrawy¹⁶ , M. Hohlmann , R. Kumar Verma ,
M. Rahmani , F. Yumiceva 












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C.E. Gerber , D.J. Hofman , D. S. Lemos , A.H. Merrit , C. Mills , S. Nanda , G. Oh ,
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















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





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


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







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
















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