



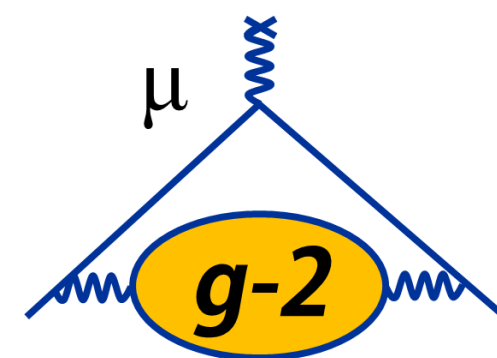
Measuring the Muon Anomalous Magnetic Moment to High Precision

David Flay

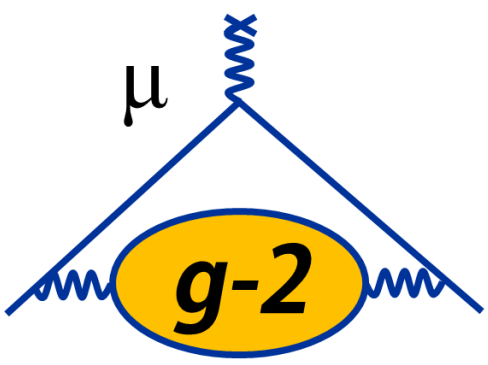
Electromagnetic Interactions with Nucleons and Nuclei, Paphos, Cyprus

31 October 2019

This document was prepared by [Muon G-2 Collaboration] using the resources of the Fermi National Accelerator Laboratory (Fermilab), a U.S. Department of Energy, Office of Science, HEP User Facility. Fermilab is managed by Fermi Research Alliance, LLC (FRA), acting under Contract No. DE-AC02-07CH11359.



Outline



Introduction

- The Magnetic Moment and the Anomaly
- Recent Theoretical Efforts

The Muon $g-2$ Experiment at Fermilab

- Experimental Technique
- Overview of Operations to Date
- Analysis Status

Summary

The Magnetic Moment and the Anomaly



The Magnetic Moment

$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$
A Feynman diagram showing a muon loop. Two blue lines representing muons form a triangle. A red wavy line representing a photon is attached to the top vertex, with an 'x' at its end. The label 'mu' is placed near the bottom-left vertex.

- Magnetic moment connected to spin via dimensionless g-factor
- Dirac: $g = 2$ for $s = 1/2$ particles (1928)
- Hyperfine structure experiments on hydrogen: $g \neq 2$ (Nafe, Nelson, Rabi 1947)
 - Anomalous contribution
 $a \equiv (g-2)/2 = \alpha/2\pi$ (Schwinger, QED, 1948)
 - Radiative corrections from virtual particles in loops

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The Muon Anomaly a_μ

$$a_\mu = \text{Schwinger } O(\alpha) \quad \text{Vacuum polarization } O(\alpha^2)$$

QED **EW** **QCD**

The Magnetic Moment and the Anomaly

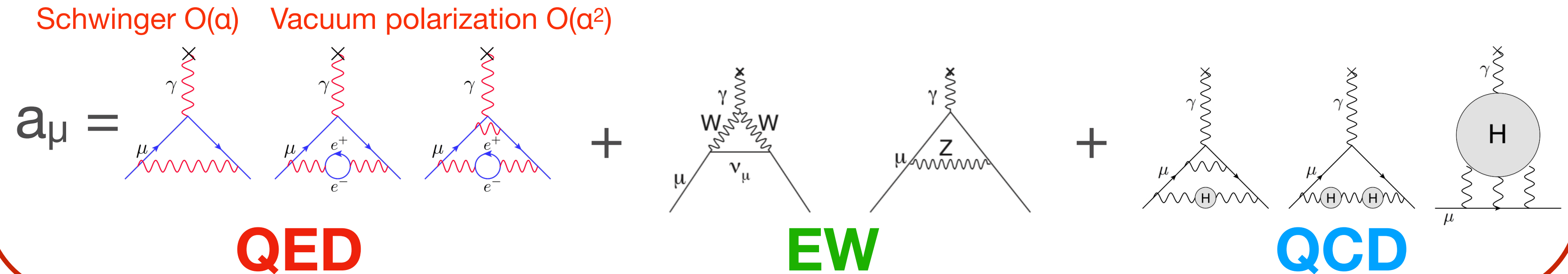


The Magnetic Moment

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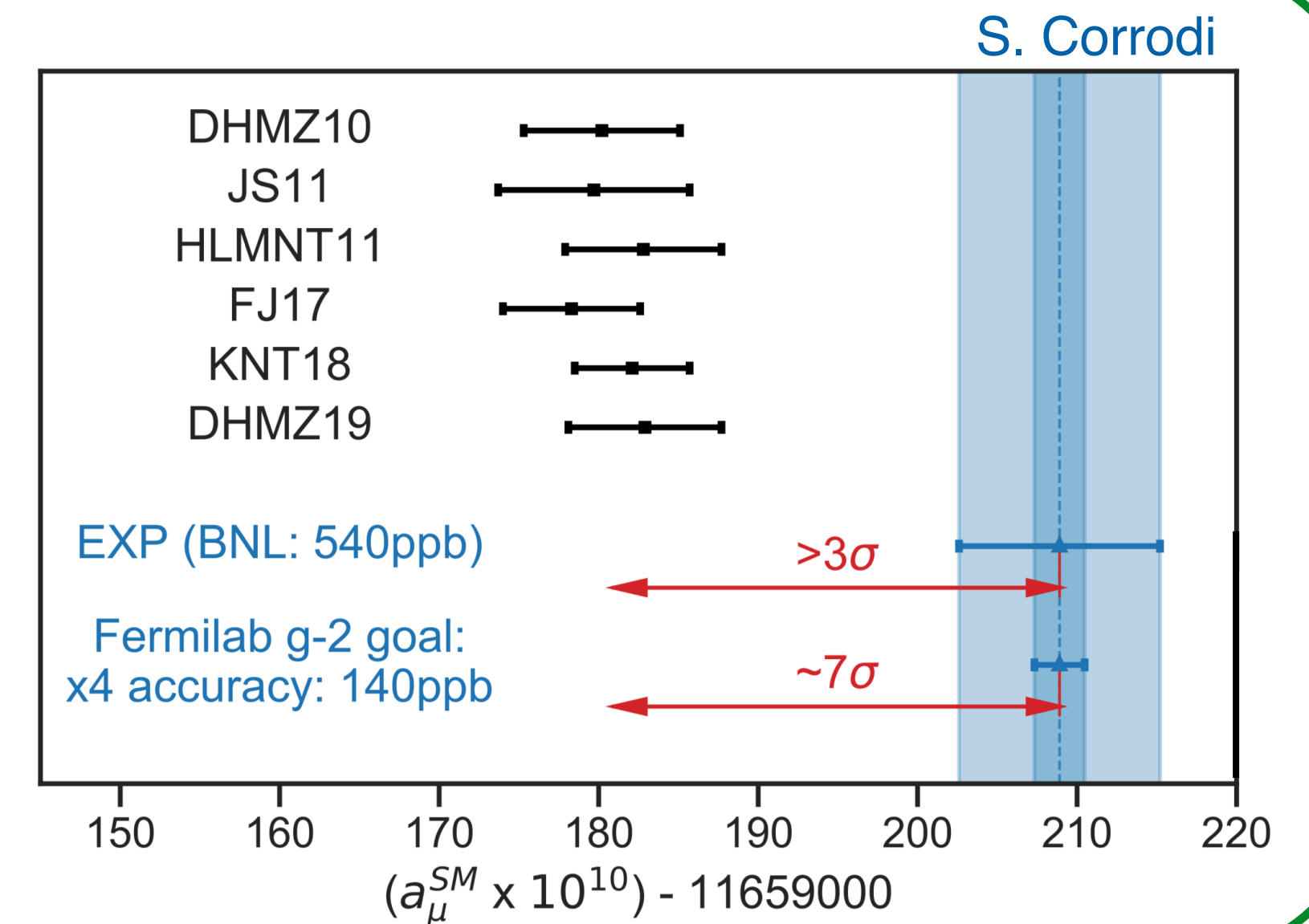
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The Muon Anomaly a_μ

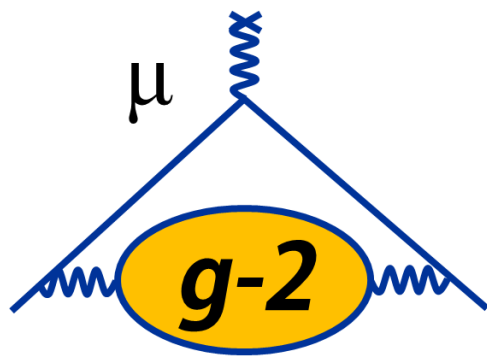


Current Status

- **Disagreement** between experiment and theory at $> 3\sigma$
- **Improvements**
 - **Experiment:** more statistics, reduced systematics
 - **Theory:** focus on QCD uncertainties



The Magnetic Moment and the Anomaly



Lattice groups making excellent progress (HVP LO, NLO, HLbL)

Calculation of the hadronic vacuum polarization contribution to the muon anomalous magnetic moment

T. Blum, P.A. Boyle, V. Gülpers, T. Izubuchi, L. Jin, C. Jung, A. Jüttner, C. Lehner, A. Portelli, J.T. Tsang
(Submitted on 22 Jan 2018)

We present a first-principles lattice QCD+QED calculation at physical pion mass of the leading-order hadronic vacuum polarization contribution to the muon anomalous magnetic moment. The total contribution of up, down, strange, and charm quarks including QED and strong isospin breaking effects is found to be $a_{\mu}^{\text{HVP LO}} = 715.4(16.3)(9.2) \times 10^{-10}$, where the first error is statistical and the second is systematic. By supplementing lattice data for very small q^2 we significantly improve the precision of our calculation. We also calculate the leading-order hadronic vacuum polarization contribution to the muon $g-2$ from lattice QCD.

Comments: 12 pages, 11 figures
Subjects: High Energy Physics – Lattice (hep-lat); High Energy Physics – Phenomenology (hep-ph)
Cite as: [arXiv:1801.07224 \[hep-lat\]](https://arxiv.org/abs/1801.07224)
(or [arXiv:1801.07224v1 \[hep-lat\]](https://arxiv.org/abs/1801.07224v1) for this version)

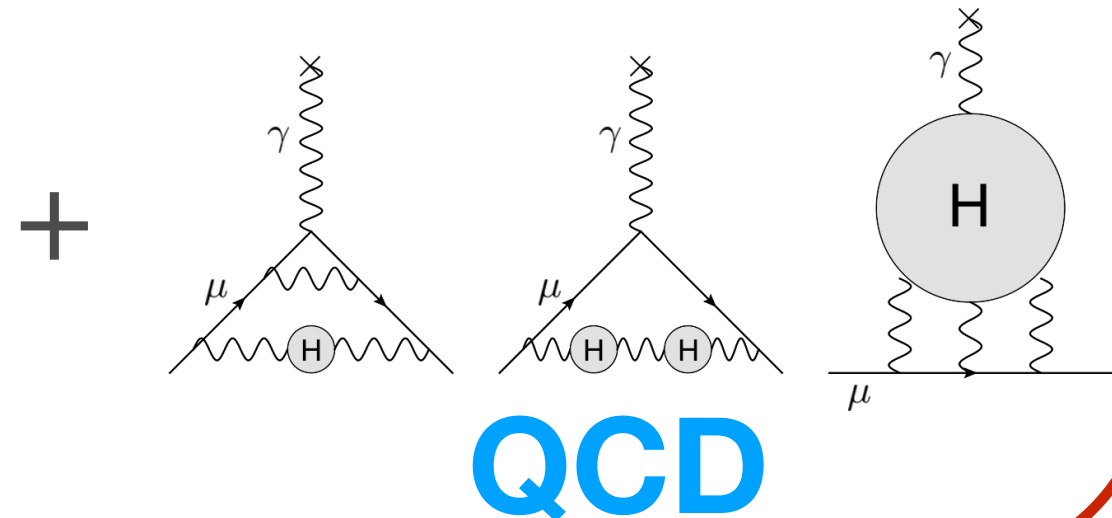
B. Chakraborty, C. T. H. Davies, J. Koponen, G. P. Lepage, and R. S. Van de Water (Fermilab Lattice, HPQCD, and MILC Collaborations)
Phys. Rev. D **98**, 094503 – Published 9 November 2018

Article References Citing Articles (1) PDF HTML Export Citation

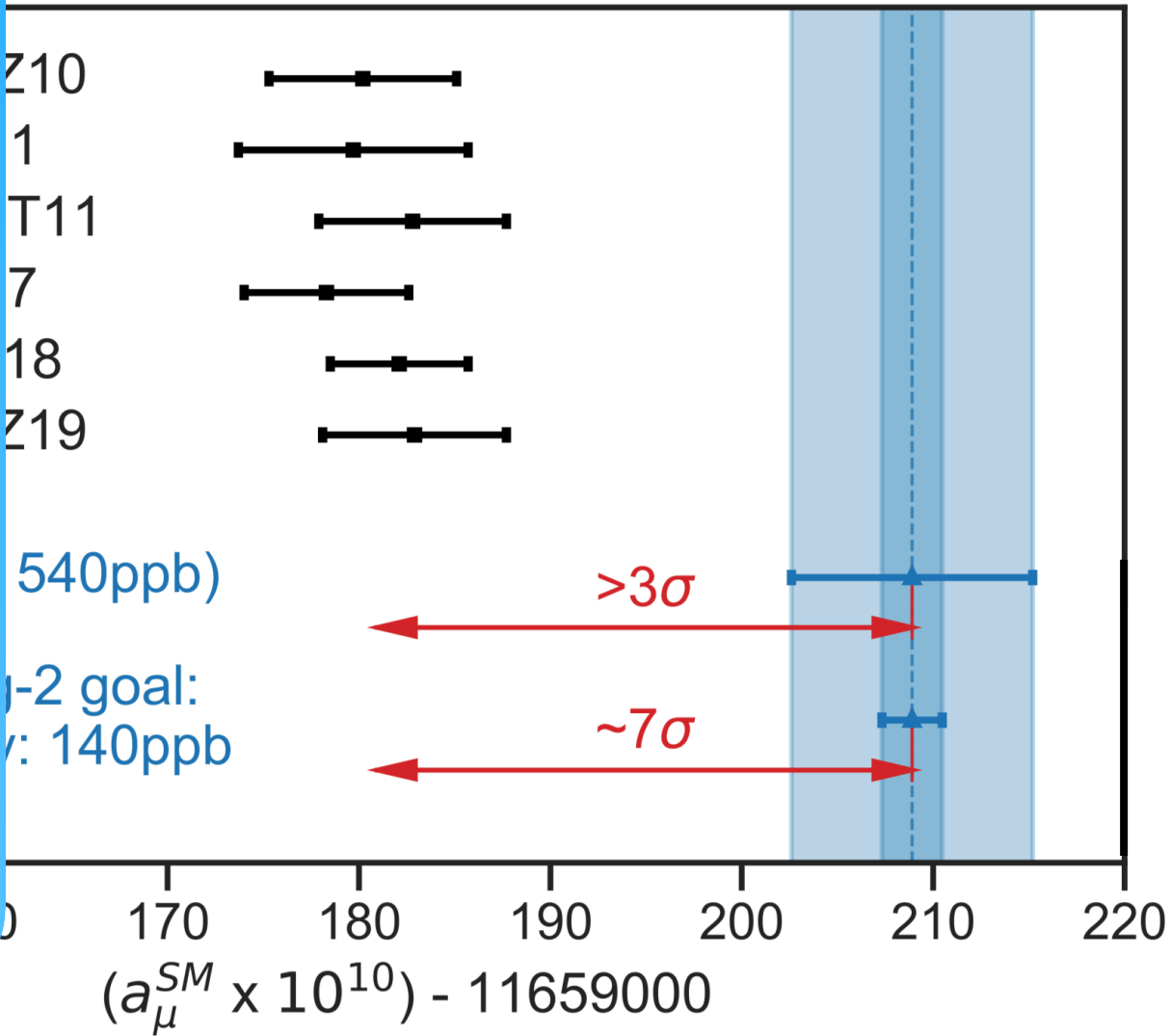
ABSTRACT

We introduce a new method for calculating the $O(\alpha^3)$ hadronic-vacuum-polarization contribution to the muon anomalous magnetic moment from *ab initio* lattice QCD. We first derive expressions suitable for computing the higher-order contributions either from the renormalized vacuum polarization function $\hat{\Pi}(q^2)$ or directly from the lattice vector-current correlator in Euclidean space. We then demonstrate the approach using previously published results for the Taylor coefficients of $\hat{\Pi}(q^2)$ that were obtained on four-flavor QCD gauge-field configurations with physical light-quark masses. We

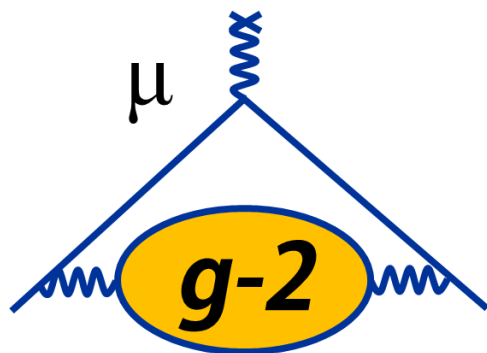
a_{μ}



S. Corrodi



The Magnetic Moment and the Anomaly



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We present a first-principles lattice QCD+QED calculation at physical pion mass of the hadronic vacuum polarization contribution to the muon anomalous magnetic moment. The total contribution to the muon anomalous magnetic moment, including strong isospin breaking effects is found to be $a_\mu^{\text{HVP LO}} = 715.4(16.3)(9.2) \times 10^{-10}$. By supplementing lattice data for very small pion masses we significantly improve the precision of our calculation. We also present a calculation of the light-quark QED correction at physical pion mass, which is found to be $a_\mu^{\text{QED}} = 10.5(0.4) \times 10^{-10}$. The total contribution to the muon anomalous magnetic moment is $a_\mu = 116.59(6) \times 10^{-10}$. This result is in excellent agreement with the experimental value $a_\mu^{\text{exp}} = 116.592(6) \times 10^{-10}$.

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Subjects: High Energy Physics – Lattice (hep-lat); High Energy Physics – Phenomenology (hep-ph)
Cite as: arXiv:1801.07224 [hep-lat]
(or arXiv:1801.07224v1 [hep-lat] for this version)

Higher-order muon $g-2$ from lattice QCD

B. Chakraborty, C. T. H. Davies, M. Lüscher, S. Meinel, M. Papinenko, R. Ruffing, J. Shigemitsu, S. Strey, A. V. Manohar, J. Zanotti
Phys. Rev. D **98**, 094501 (2018)

Article References



2nd g-2 Theory Initiative Meeting in June 2018

JOHANNES GUTENBERG UNIVERSITÄT MAINZ

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Theory Group

Institute for Nuclear Physics

Second Plenary Workshop of the Muon g-2 Theory Initiative

18 June 2018 - 22 June 2018

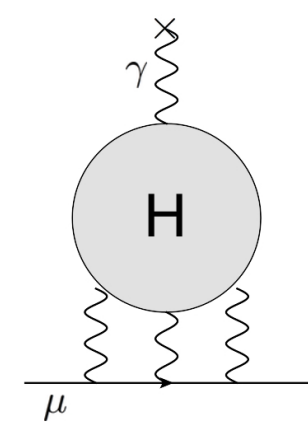
In the coming years, experiments at Fermilab and at J-PARC plan to reduce the uncertainties on the already very precisely measured anomalous magnetic moment of the muon by a factor of four. The goal is to resolve the current tantalizing tension between theory and experiment of three to four standard deviations. On the theory side the hadronic corrections to the anomalous magnetic moment are the dominant sources of uncertainty. They must be determined with better precision in order to unambiguously discover whether or not new physics effects contribute to this quantity.

There are a number of complementary theoretical efforts underway to better understand and quantify the hadronic corrections, including dispersive methods, lattice QCD, effective field theories, and QCD models. The Muon (g-2) Theory Initiative was formed in order to facilitate interactions between the different groups through organizing a series of workshops. The goal of this workshop is to bring together theorists from the different communities to discuss, assess, and compare the status of the various efforts, and to map out strategies for obtaining the best theoretical predictions for these hadronic corrections in advance of the experimental results.

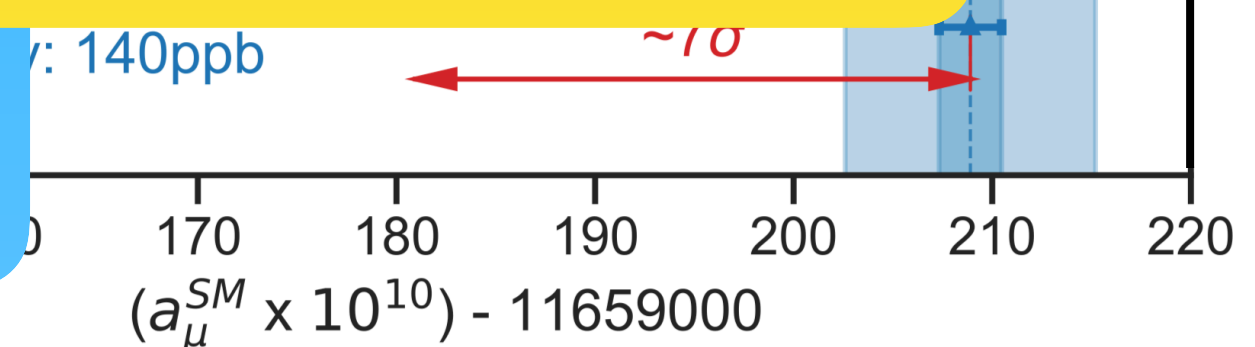
Dates	June 18, 2018 - June 22, 2018
Timezone	GMT+2
Location	Helmholtz-Institut Mainz Staudinger Weg 18, 55128 Mainz, Ground Floor

Second Workshop of the Muon g-2 Theory Initiative
Helmholtz-Institut Mainz
Staudinger-Weg 18
55128 Mainz
18 - 22 June 2018

First Circular
Second Circular
Indico website



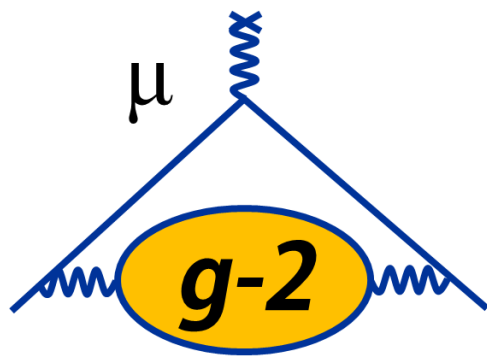
Corrodi



A Feynman diagram for the muon $g-2$ experiment. It shows a muon (μ) at the top vertex, with a wavy line (photon) extending downwards. This photon forms a loop with a fermion (solid line). The loop is labeled $g-2$.

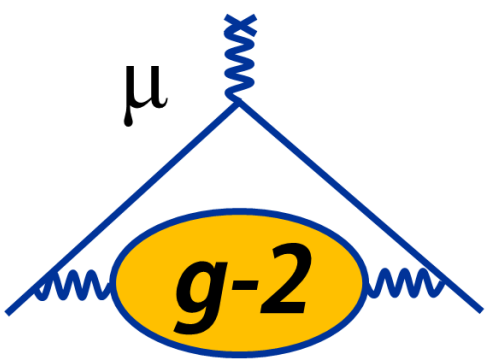
The image displays two Feynman diagrams for the muon g-2 calculation. The left diagram shows a muon line with a photon loop and a Z boson exchange. The right diagram shows a muon line with a photon loop and a Z boson exchange, with a muon line continuation.

a_μ Theoretical Status



Contribution	Value (x 10 ⁻¹¹)	Reference
QED	116 584 718.95 ± 0.08	PRL 109 111808 (2012)
EW	153.6 ± 1.0	PRD 88 053005 (2013)

a_μ Theoretical Status

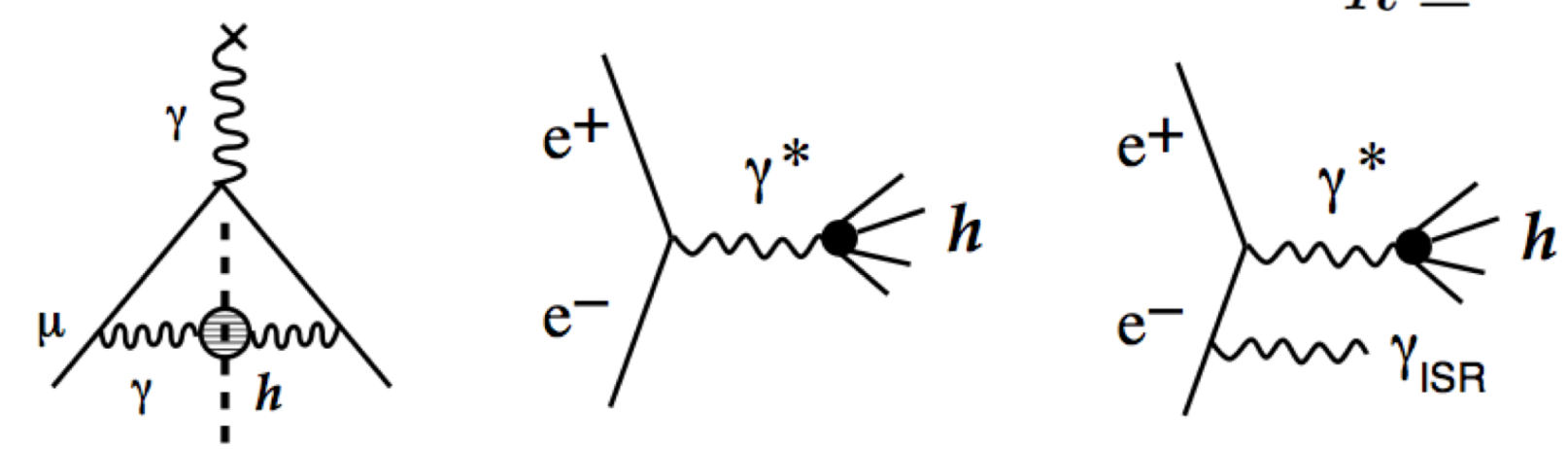


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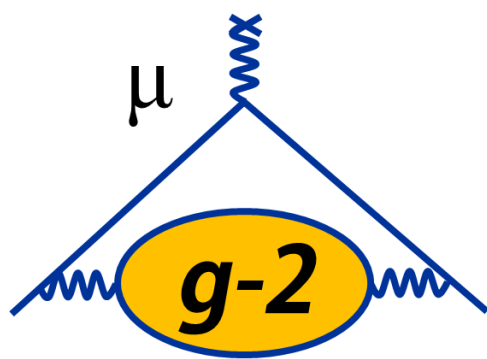
HVP (LO): Lowest-Order Hadronic Vacuum Polarization

- Critical input** from e^+e^- colliders (data from SND, CMD3, BaBar, KLOE, Belle, BESIII), extensive physics program running to reduce $\delta a_\mu^{\text{HVP}}$ to $\sim 0.3\%$ in coming years
- Progress on the lattice**: Calculations at physical π mass; approaching goal of $\delta a_\mu^{\text{HVP}} \sim 1\%$ (cross-check with e^+e^- data)

$$a_\mu^{\text{had;LO}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2 \int_{m_\pi^2}^\infty \frac{ds}{s^2} K(s) R(s)$$
$$R \equiv \frac{\sigma_{\text{tot}}(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

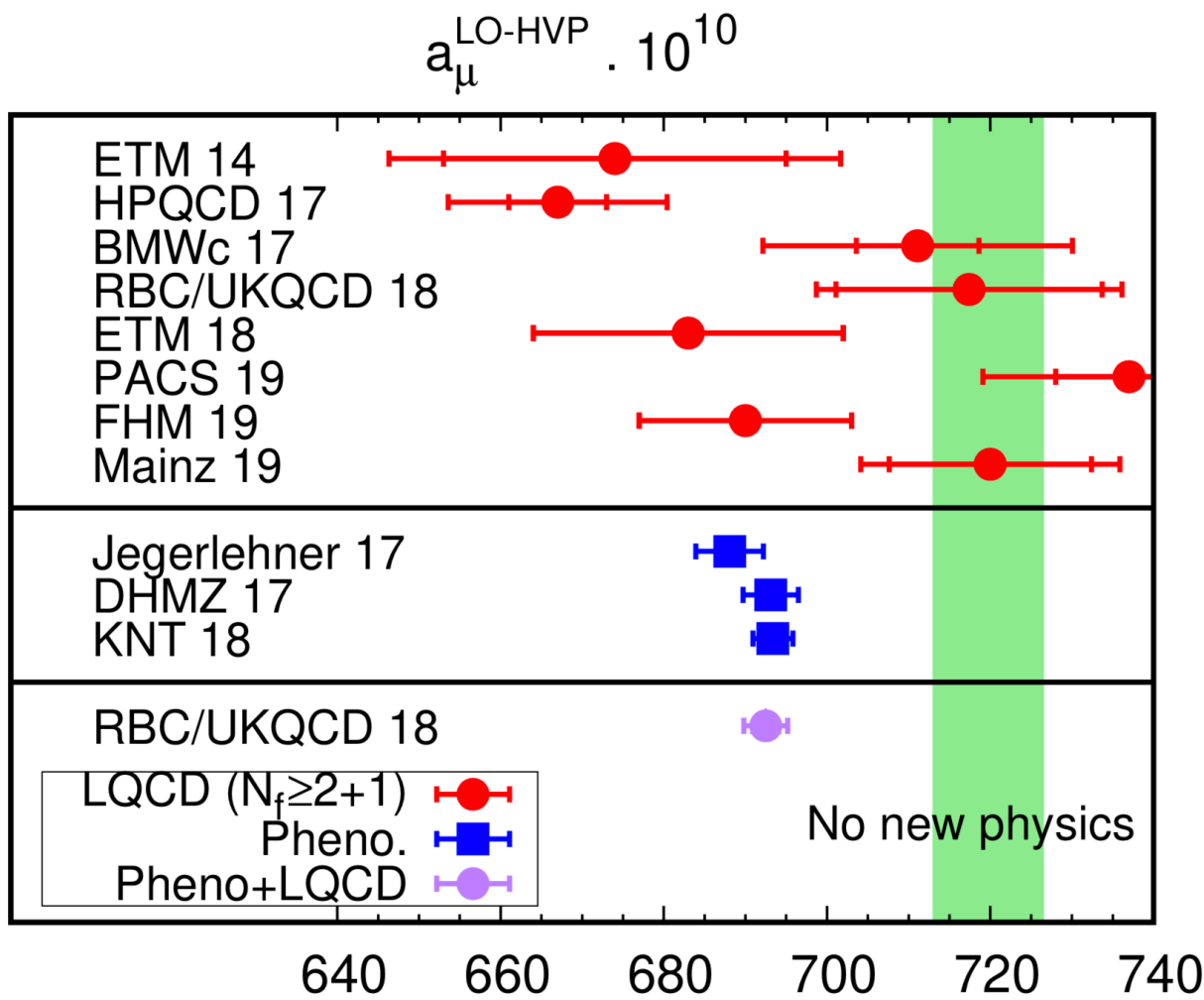


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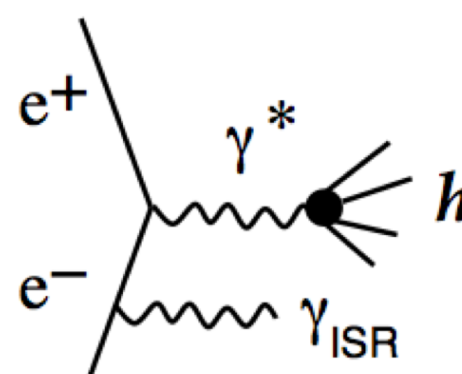
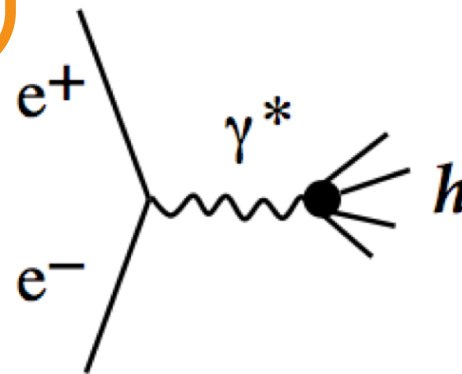
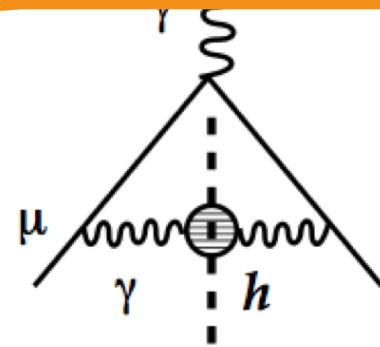
[prepared by K. Miura for WP]



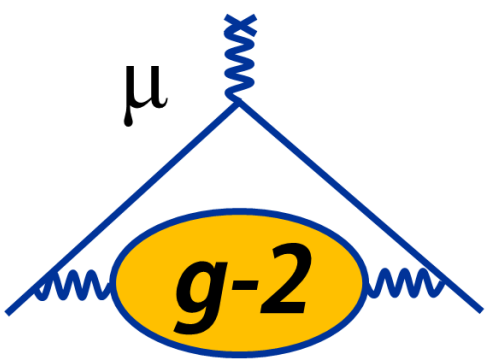
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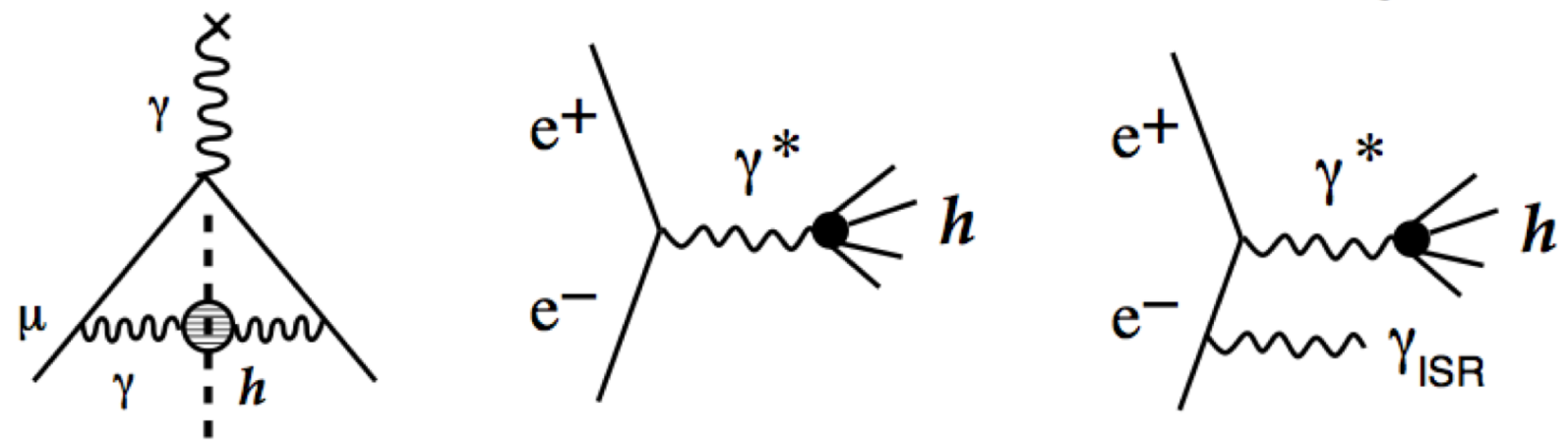


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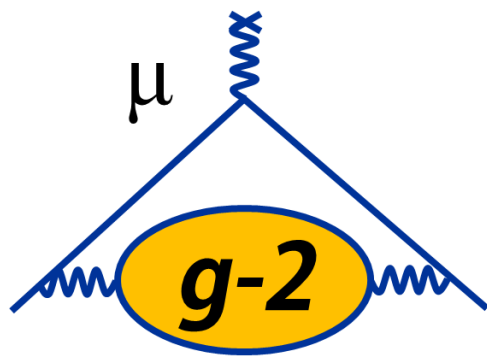
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a_μ Theoretical Status



New *ab initio* approaches [PRD 98 094503 (2018)] finding consistent result of $(-93 \pm 13) \times 10^{-11}$ — lattice making big strides

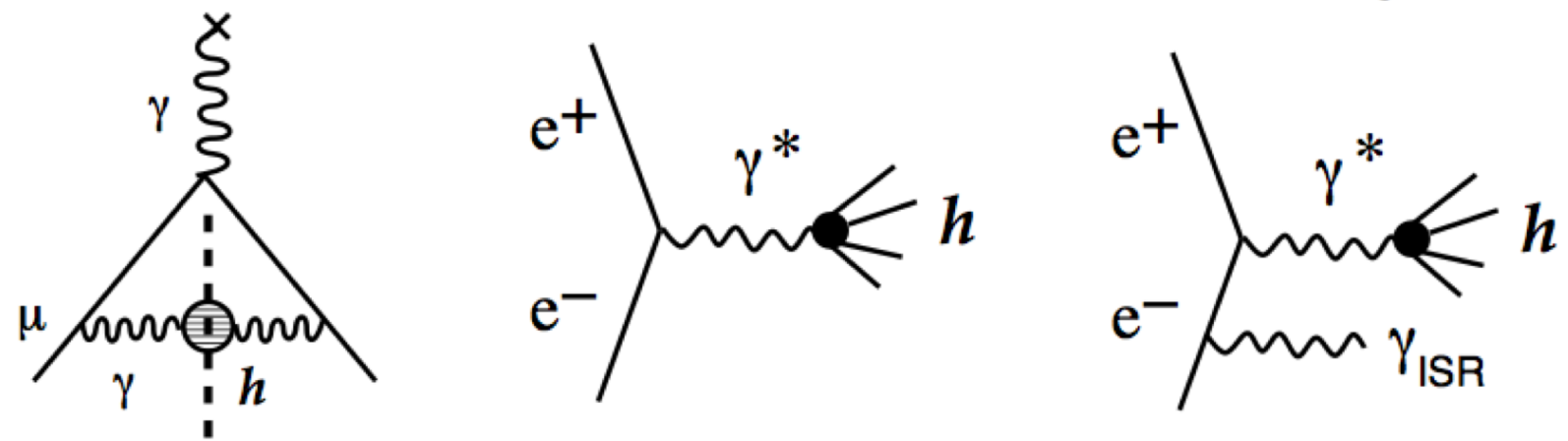
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HVP (LO): Lowest-Order Hadronic Vacuum Polarization

- Critical input** from e⁺e⁻ colliders (data from SND, CMD3, BaBar, KLOE, Belle, BESIII), extensive physics program running to reduce δa_μ^{HVP} to ~ 0.3% in coming years
- Progress on the lattice**: Calculations at physical π mass; approaching goal of δa_μ^{HVP} ~ 1% (cross-check with e⁺e⁻ data)

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HLbL (LO + NLO)	101 ± 26	PLB 735 90 (2014), EPJ Web Conf 118 01016 (2016)
Total SM	$116\,591\,818 \pm 43$ (368 ppb) $116\,591\,821 \pm 36$ (309 ppb)	

HLbL: Hadronic Light-by-Light

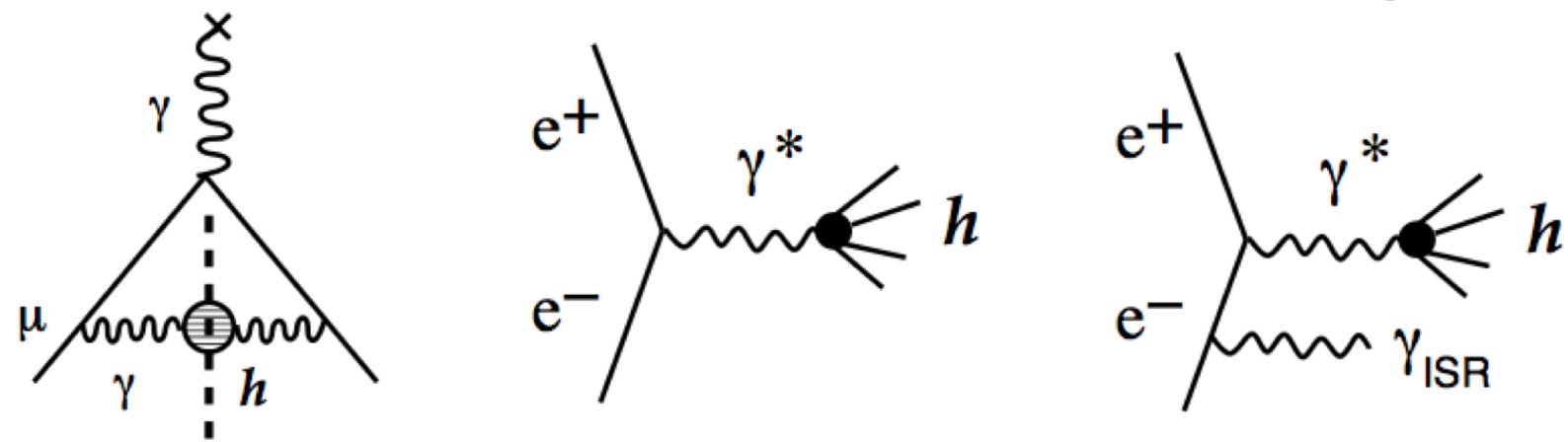
- Model dependent: based on χ PT + short-distance constraints (operator product expansion)
- Difficult to relate to data like HVP (LO); γ^* physics, π^0 data (BESIII, KLOE) important for constraining models
- Theory Progress:** New dispersive calculation approach; extend the lattice (finite volume, disconnected diagrams); Blum et al. making excellent progress

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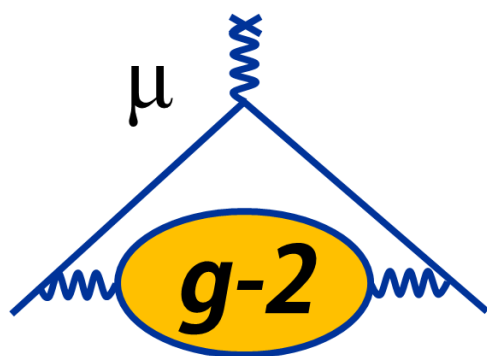
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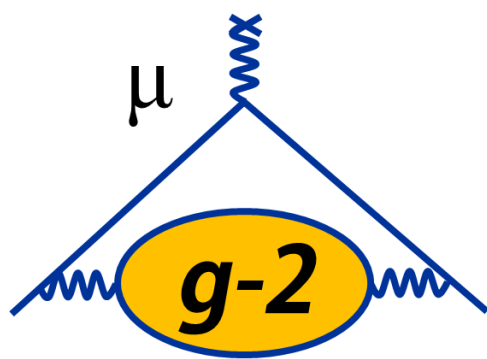
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Recent lattice & data-driven estimate [PRD 100 034520 (2019)] for $a_\mu^{\pi^0\text{-pole}}$ is consistent with lowest-meson dominance, + vector phenomenological models [PRD 51 4939 (2005), PRL 83 5230 (1999), EJC 21 659 (2001), PRD 65 073034 (2002), PRD 94 053006 (2016), EJC 75 586 (2015)]

$$\frac{\sigma_{\text{tot}}(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

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Builds confidence in HLbL term

See T. Blum's talk for updates on the lattice

HLbL: Hadronic Light-by-Light

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$$\frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

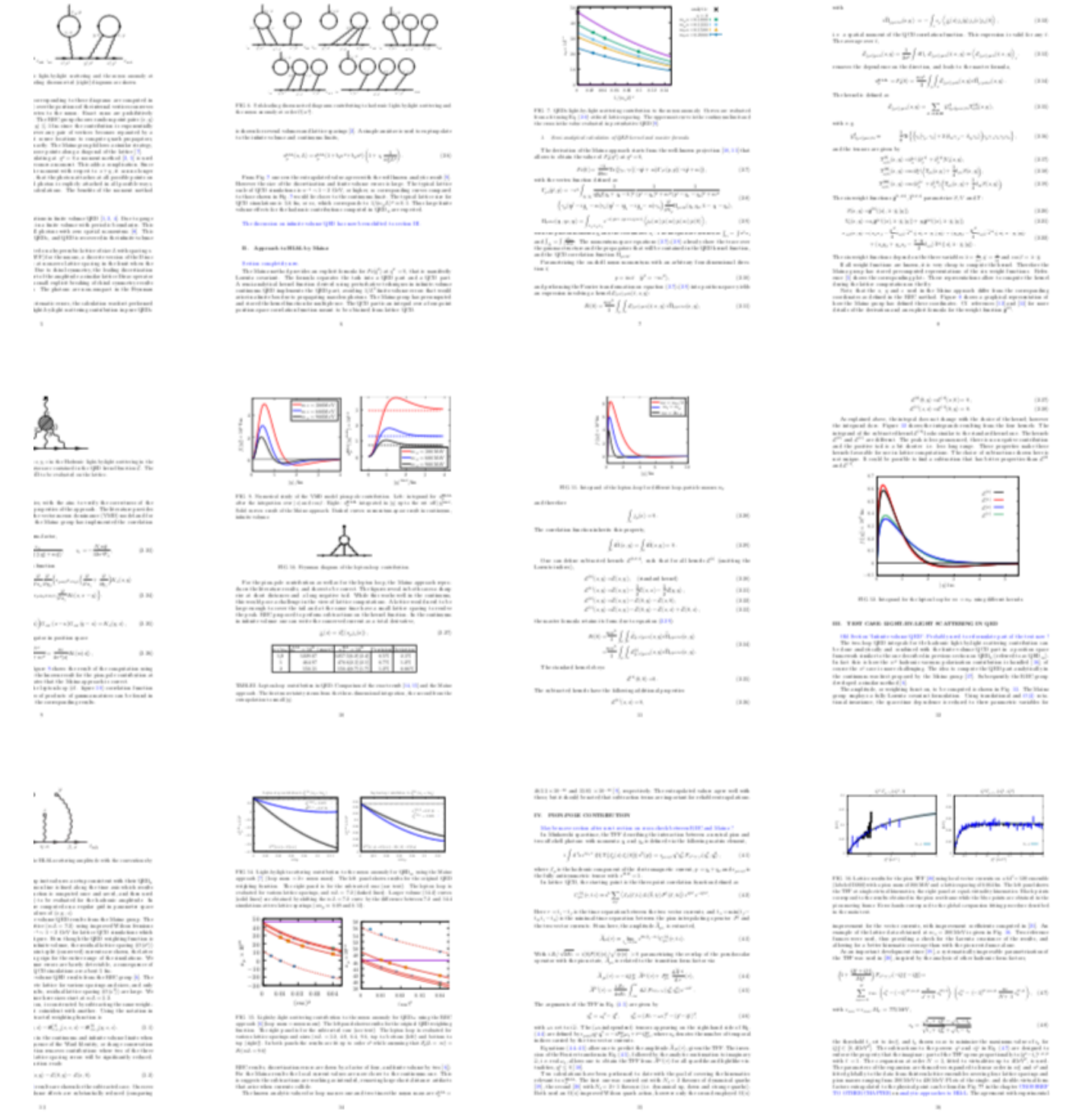
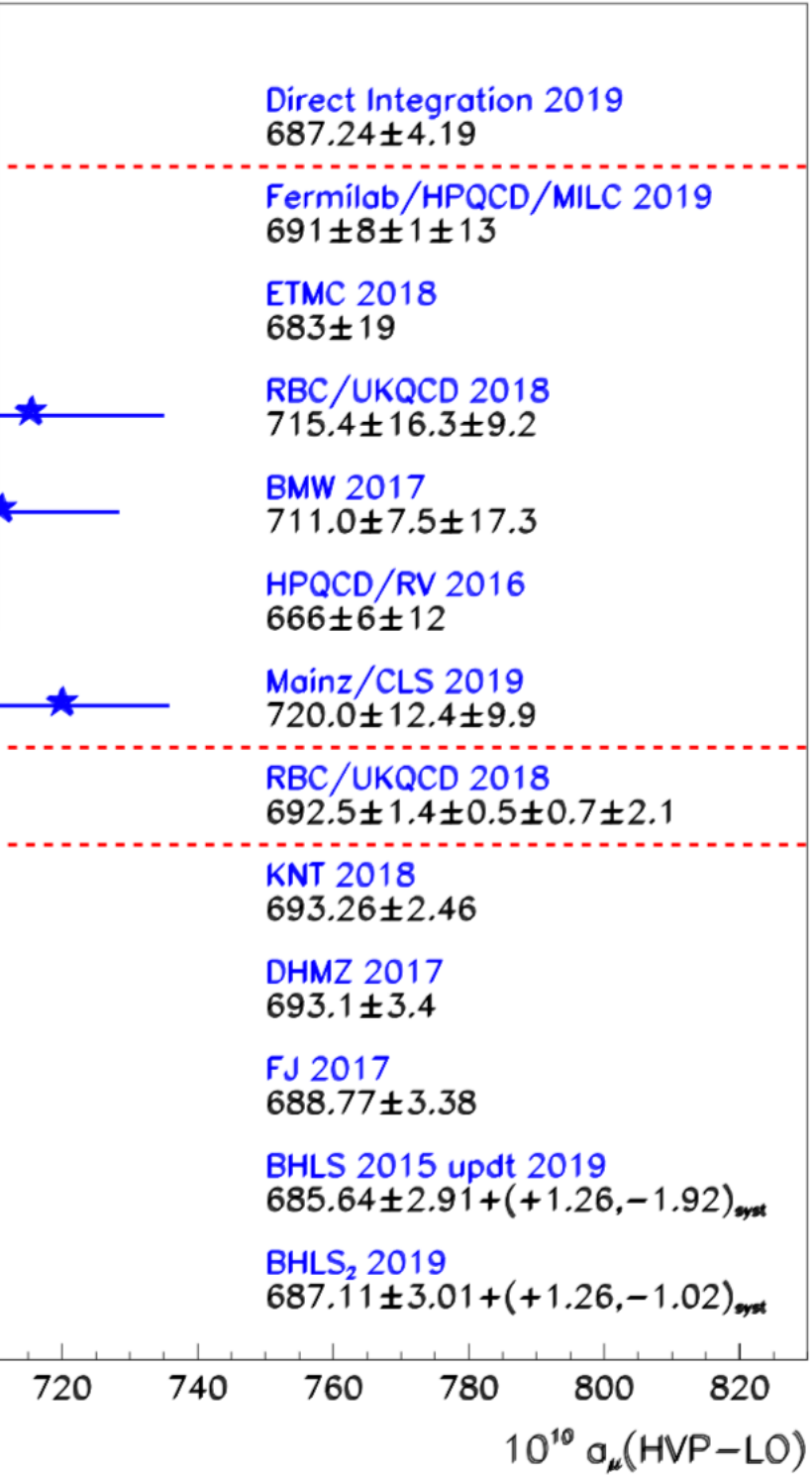
a_μ Theoretical Status

New *ab initio* approaches [PRD 98 094503 (2018)]
finding consistent result of $(-93 \pm 13) \times 10^{-11}$ —



- Theory groups are making steady progress to achieve competitive uncertainties on the same time scale as the FNAL experiment result
- White paper discussing all terms forthcoming

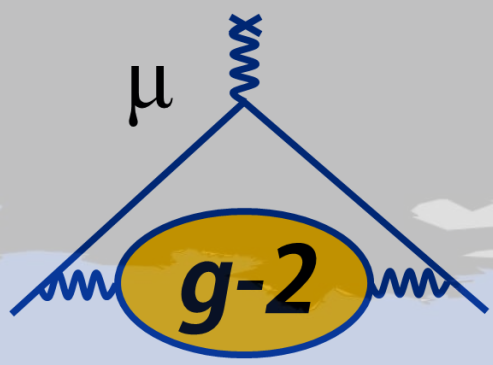
Progress on the lattice: Ca 5230 (1999), E 94 053006 (2018)





The Muon $g-2$ Experiment at Fermilab

Muon g-2: 33 Institutions, 7 countries, 203 Members

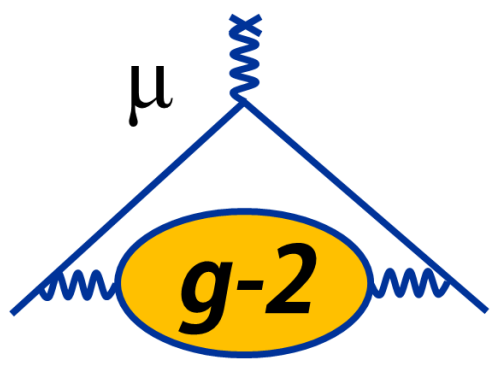
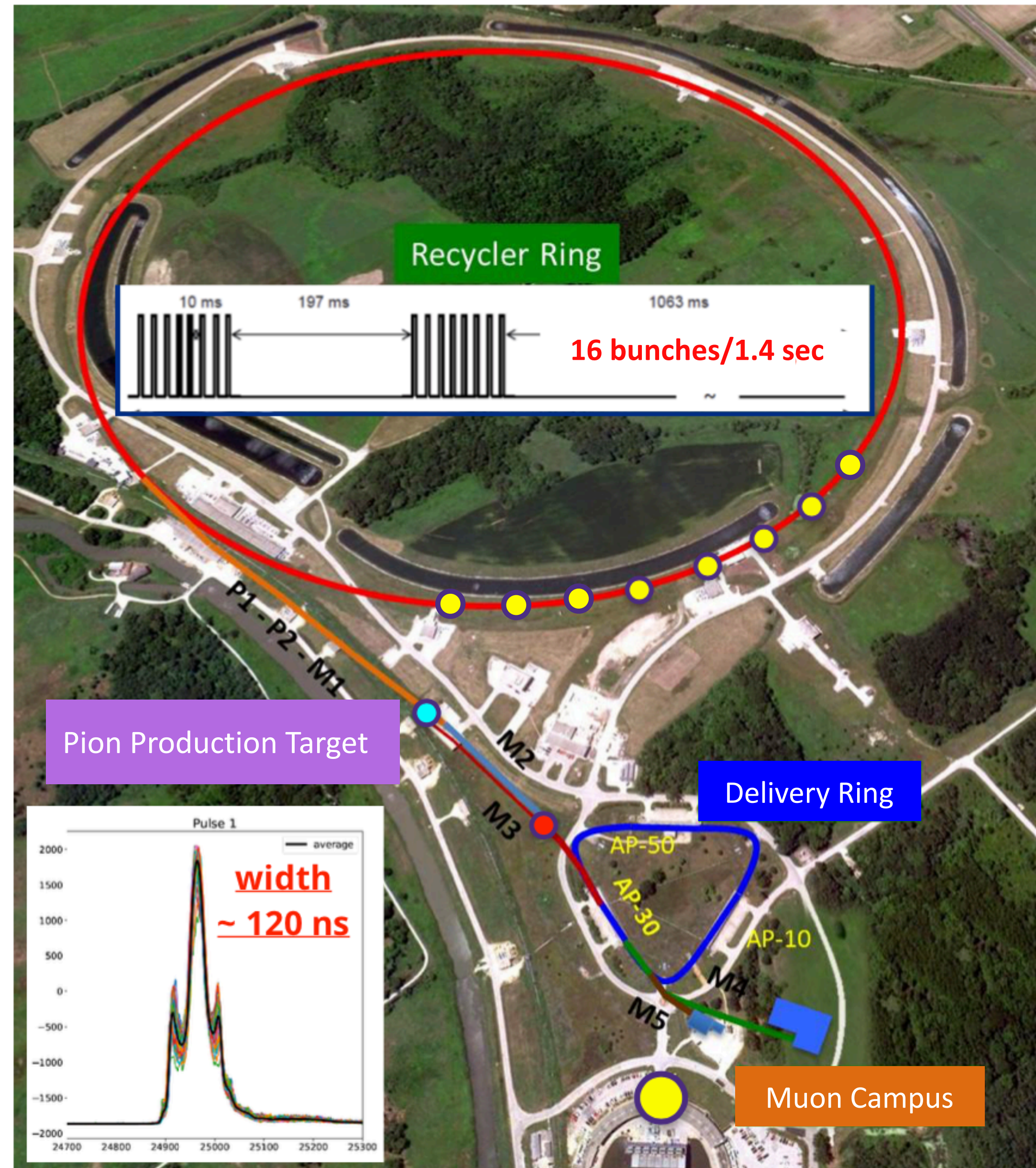


Why Fermilab?

- BNL limited by statistics
(540 ppb on 9×10^9 detected e^+)
- E989 goal: Factor of 21 more statistics
(2×10^{11} detected e^+)

Fermilab advantages

- ✓ Long beam line to collect $\pi^+ \rightarrow \mu^+$
- ✓ Much reduced amount of p, π in ring
- ✓ 4x higher fill frequency than BNL



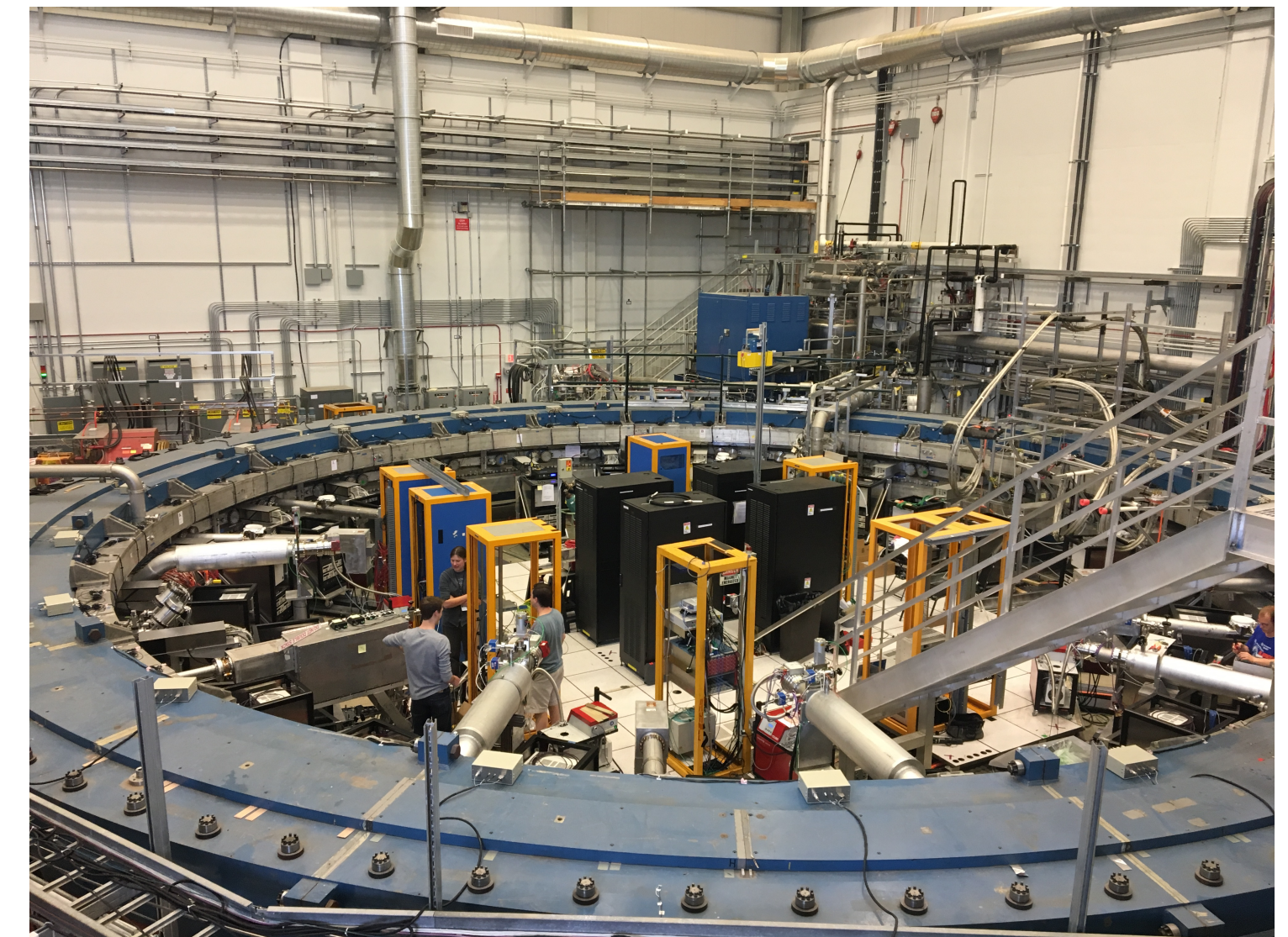
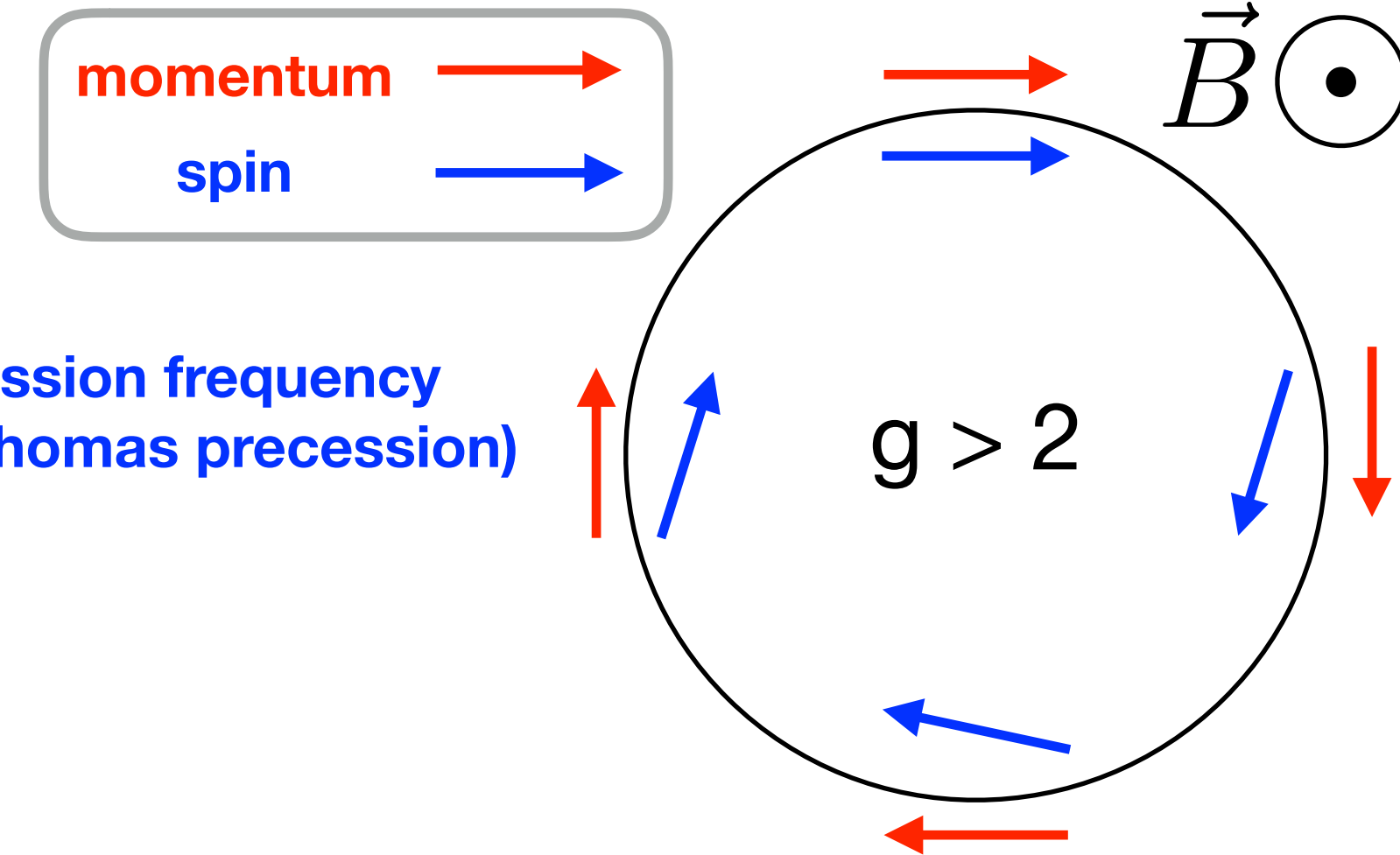
Measuring the Muon Anomaly

- Inject polarized muon beam into magnetic storage ring
- Measure **difference** between spin precession and cyclotron frequencies
- If $g = 2$, $\omega_a = 0$
- $g \neq 2$, $\omega_a \equiv (e/m_\mu)a_\mu B$

$$\vec{\omega}_C = -\frac{e}{\gamma m} \vec{B} \quad \text{cyclotron frequency}$$

$$\vec{\omega}_S = -\frac{e}{\gamma m} \vec{B} (1 + \gamma a_\mu) \quad \text{spin precession frequency (Larmor, Thomas precession)}$$

$$\vec{\omega}_a \equiv \vec{\omega}_S - \vec{\omega}_C$$



Measuring the Muon Anomaly

- Inject polarized muon beam into magnetic storage ring
- Measure **difference** between spin precession and cyclotron frequencies
- If $g = 2$, $\omega_a = 0$
- $g \neq 2$, $\omega_a \approx (e/m_\mu)a_\mu B$
- Using $\hbar\omega_p = 2\mu_p|\vec{B}|$:

$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

↑
↑
↑

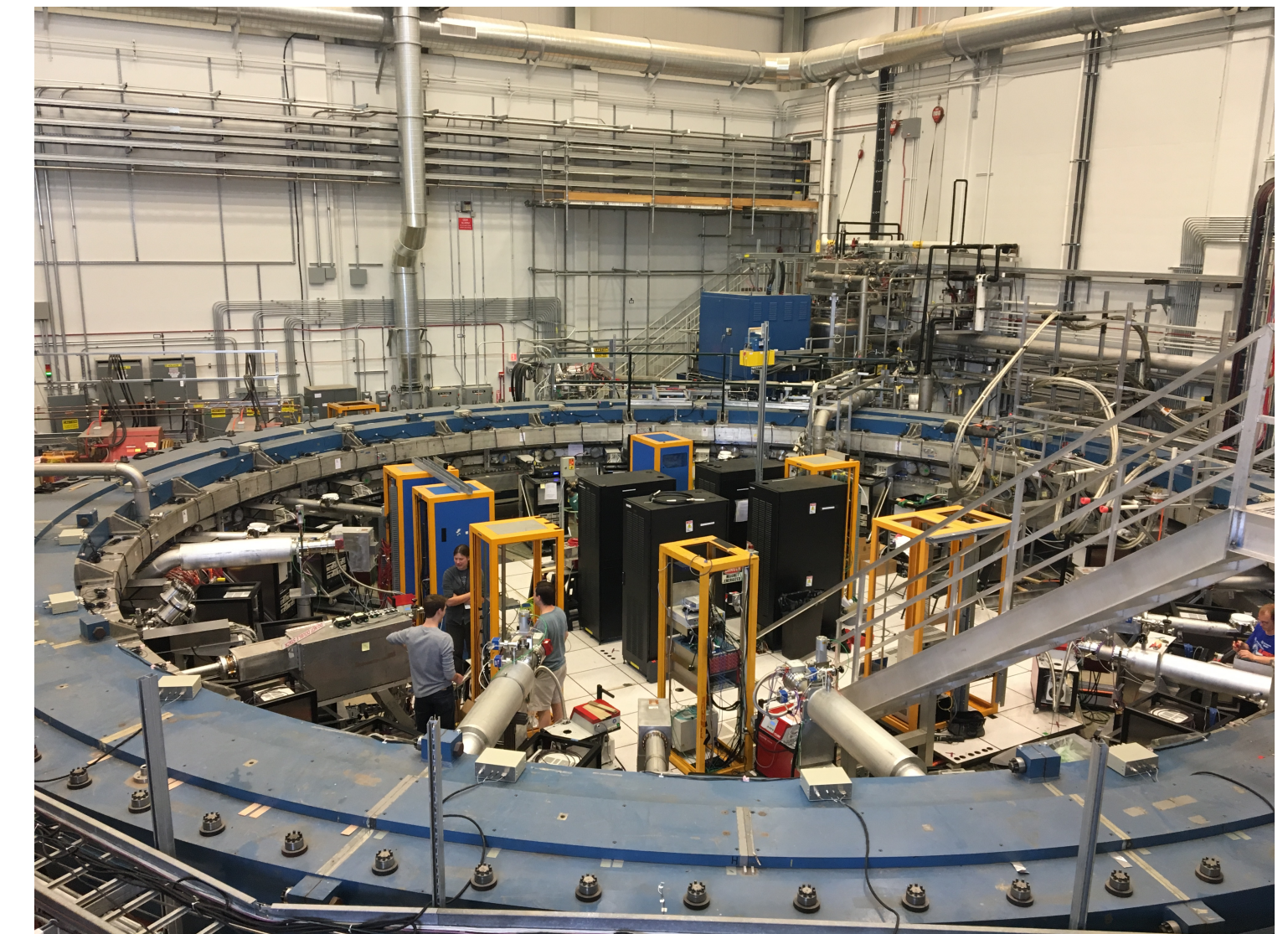
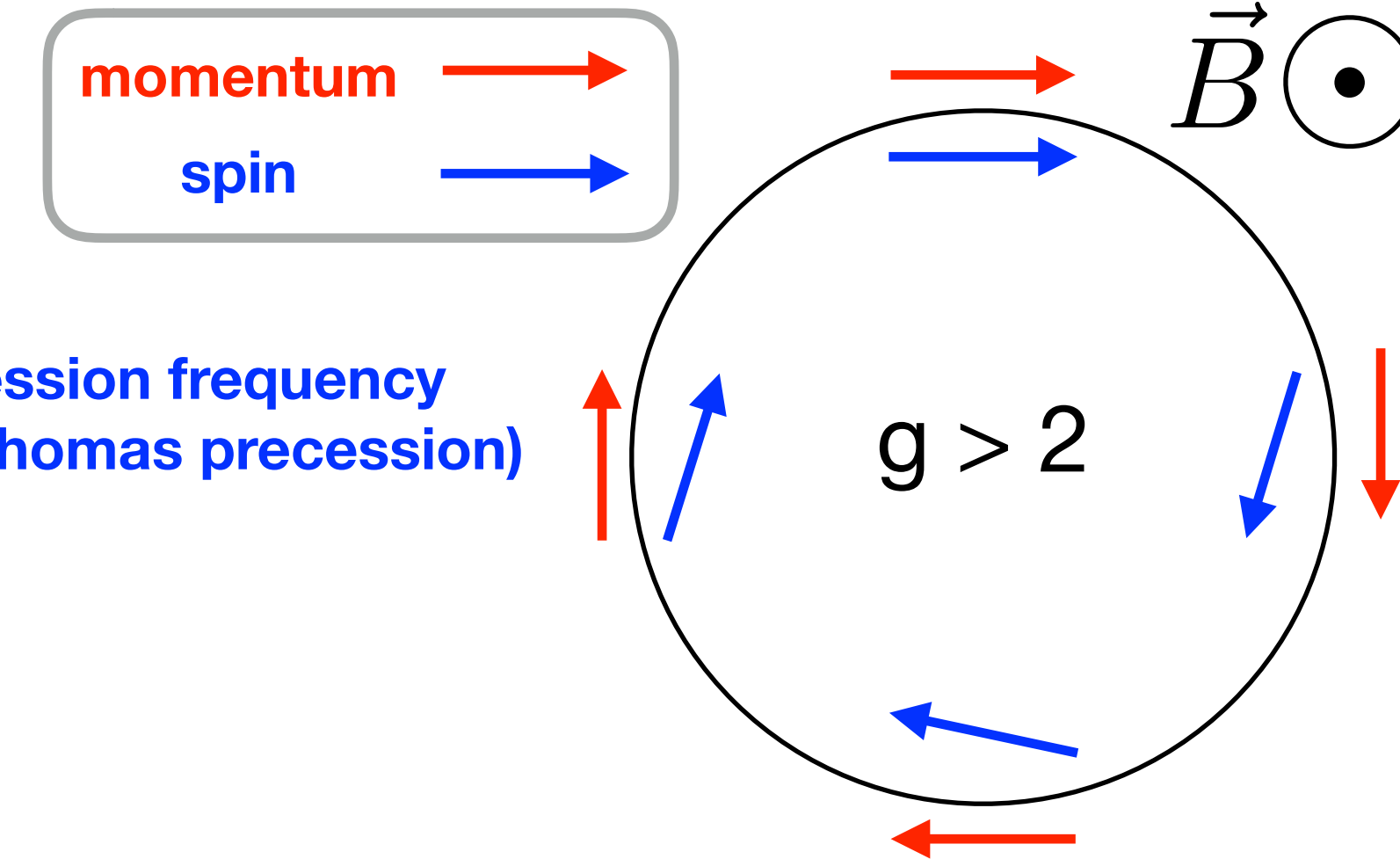
3 ppb
22 ppb
0.3 ppt

Rev. Mod. Phys. **88**, 035009 (2016)

$$\vec{\omega}_C = -\frac{e}{\gamma m} \vec{B} \quad \text{cyclotron frequency}$$

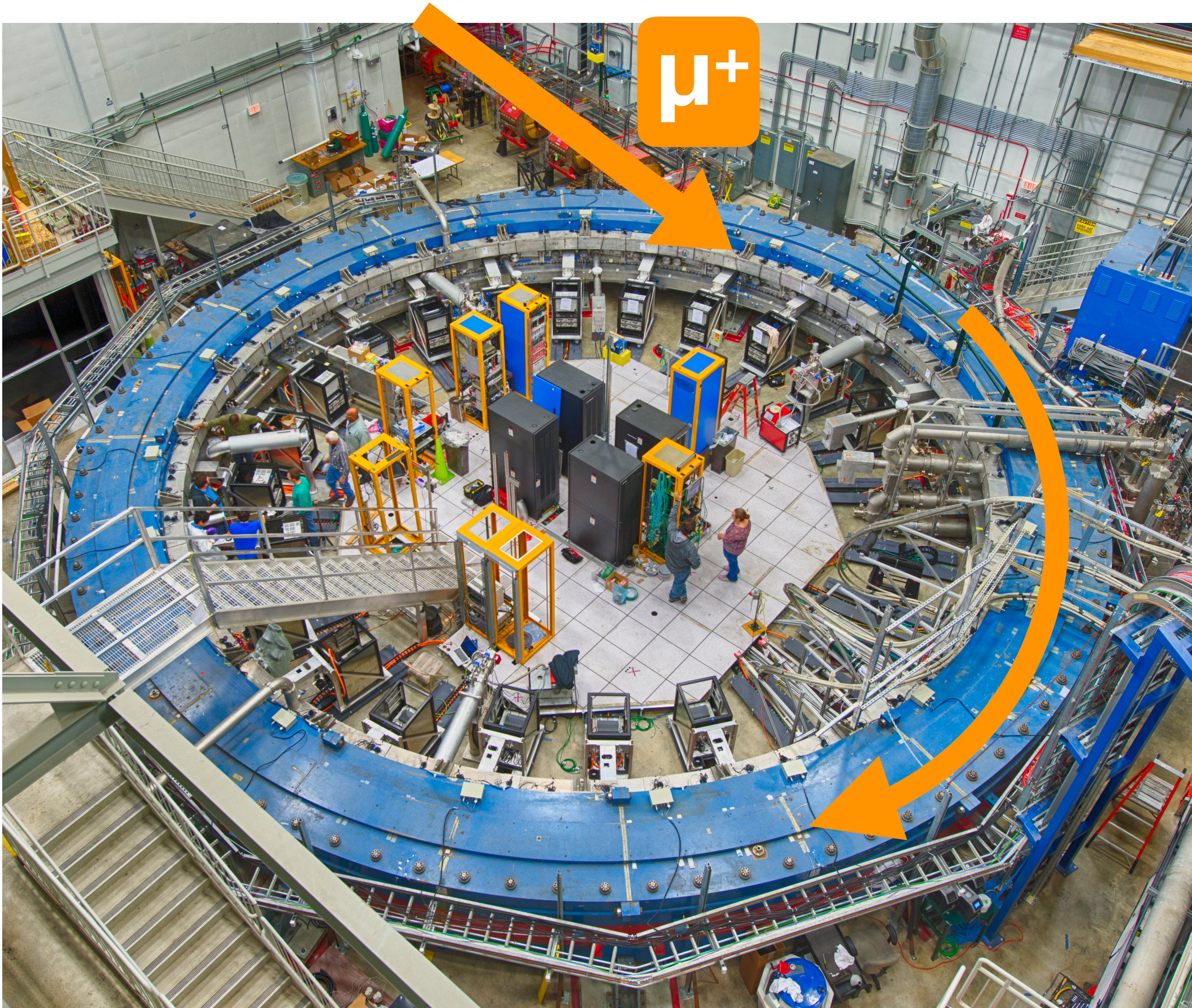
$$\vec{\omega}_S = -\frac{e}{\gamma m} \vec{B} (1 + \gamma a_\mu) \quad \text{spin precession frequency (Larmor, Thomas precession)}$$

$$\vec{\omega}_a \equiv \vec{\omega}_S - \vec{\omega}_C$$

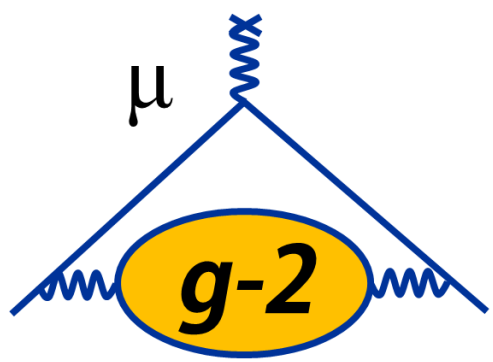


- We measure ω_a and ω_p separately
- Aiming for 70 ppb precision on each (systematic)
- Target: $\delta a_\mu = 140$ ppb; factor of 4 improvement over BNL**

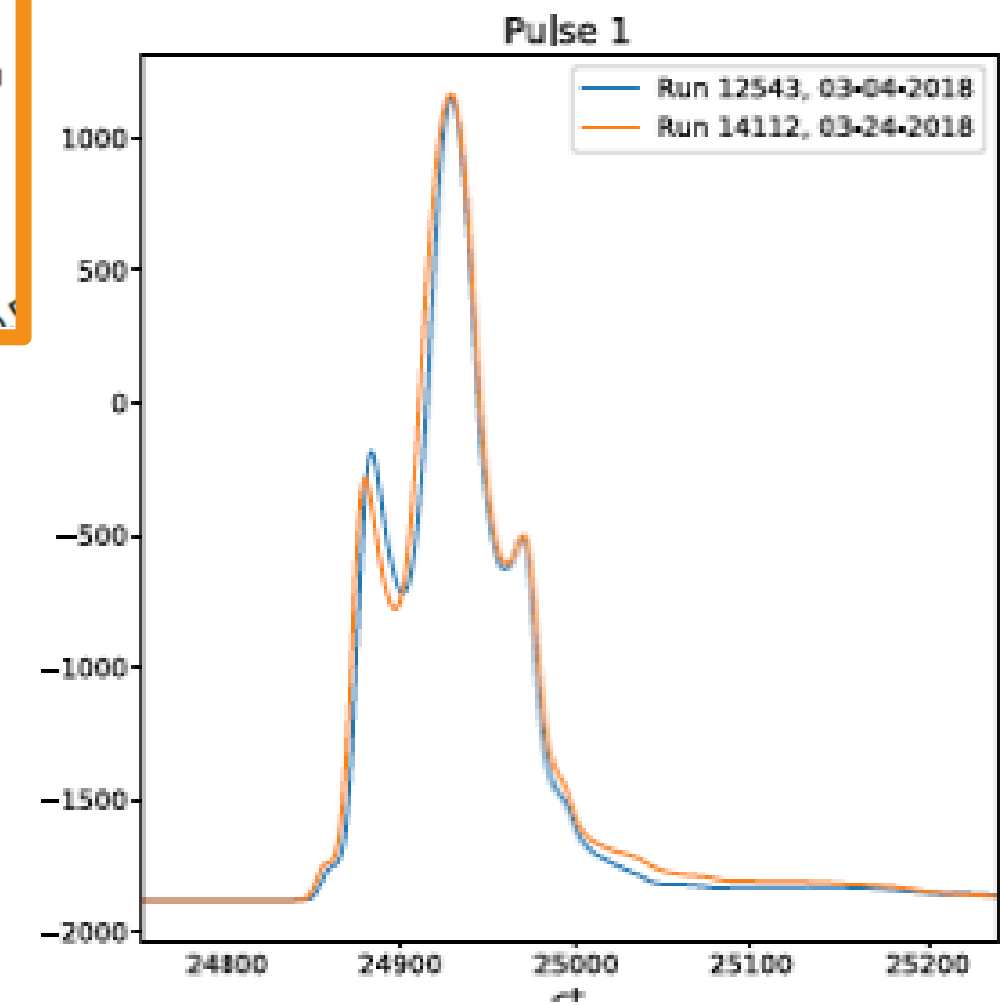
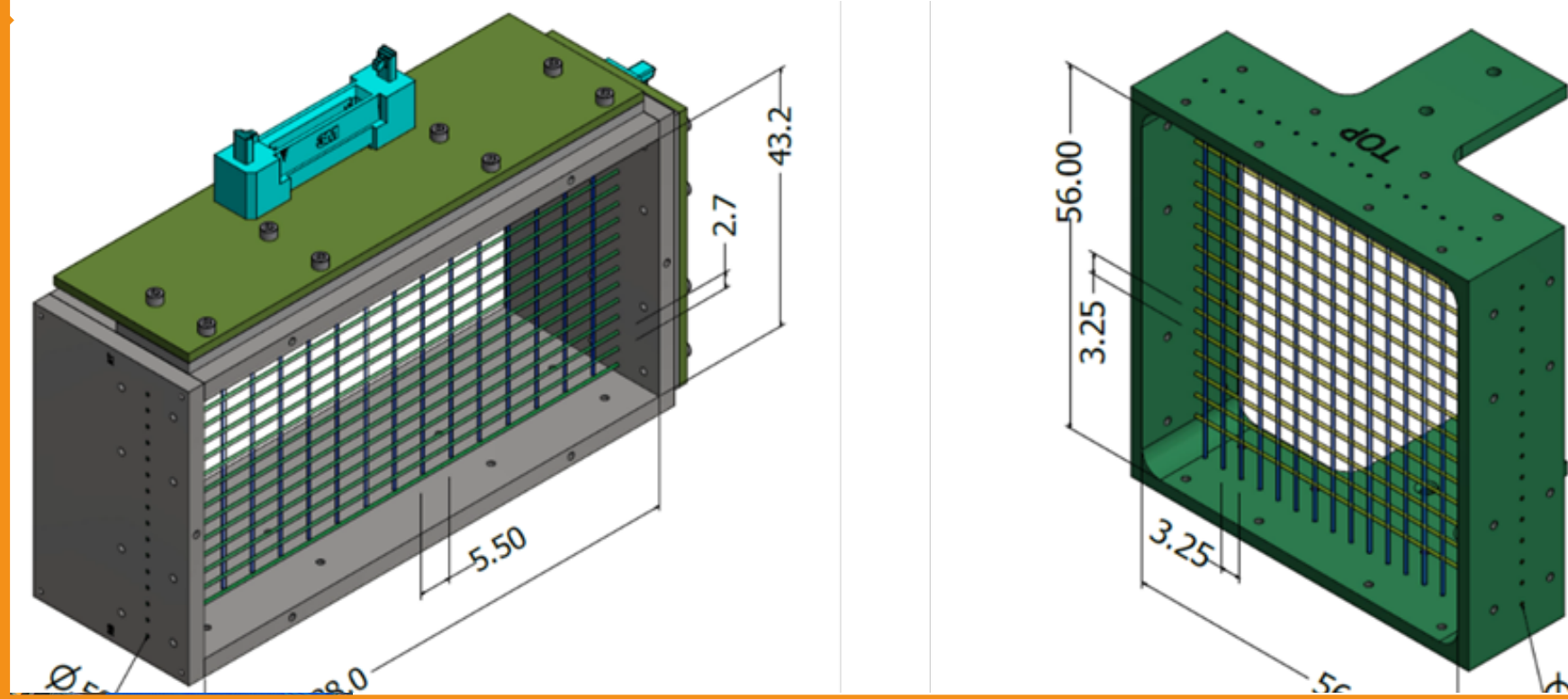
Muon Beam Injection



Muon Beam Injection



μ^+



Monitoring the Incoming Muon Beam

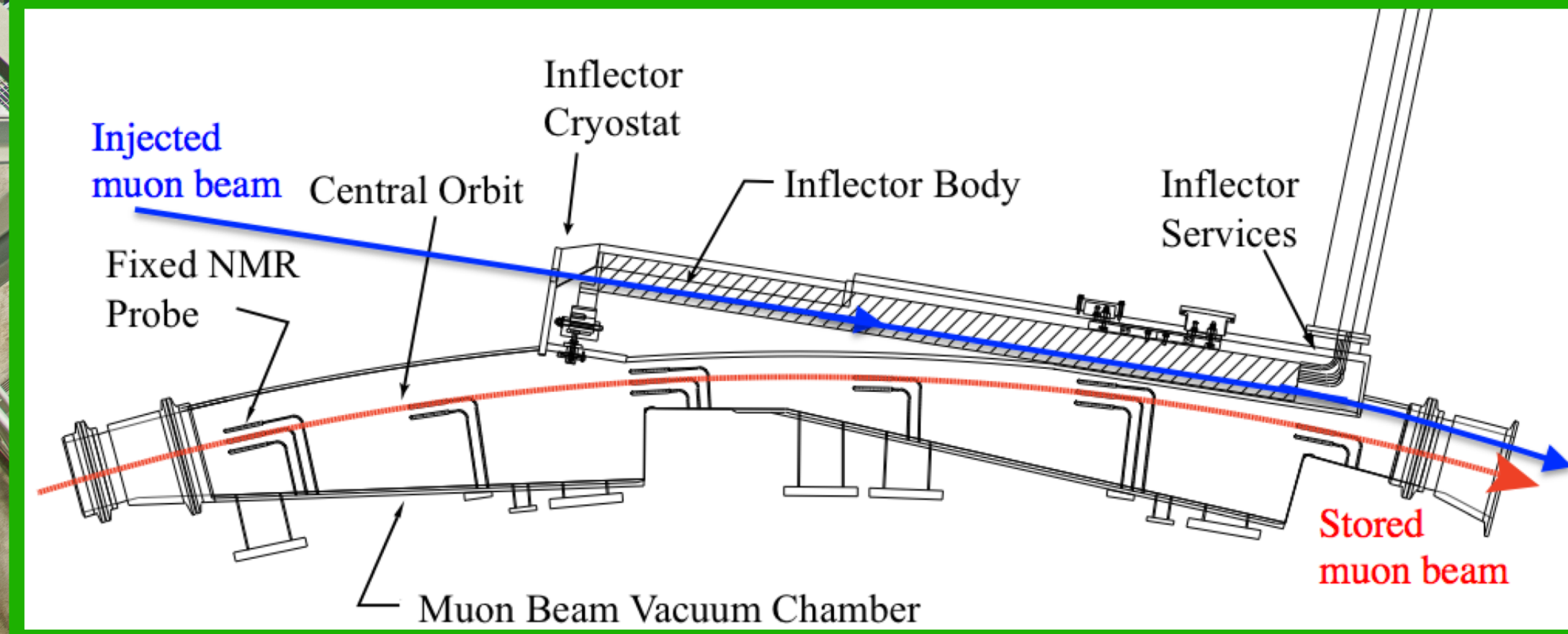
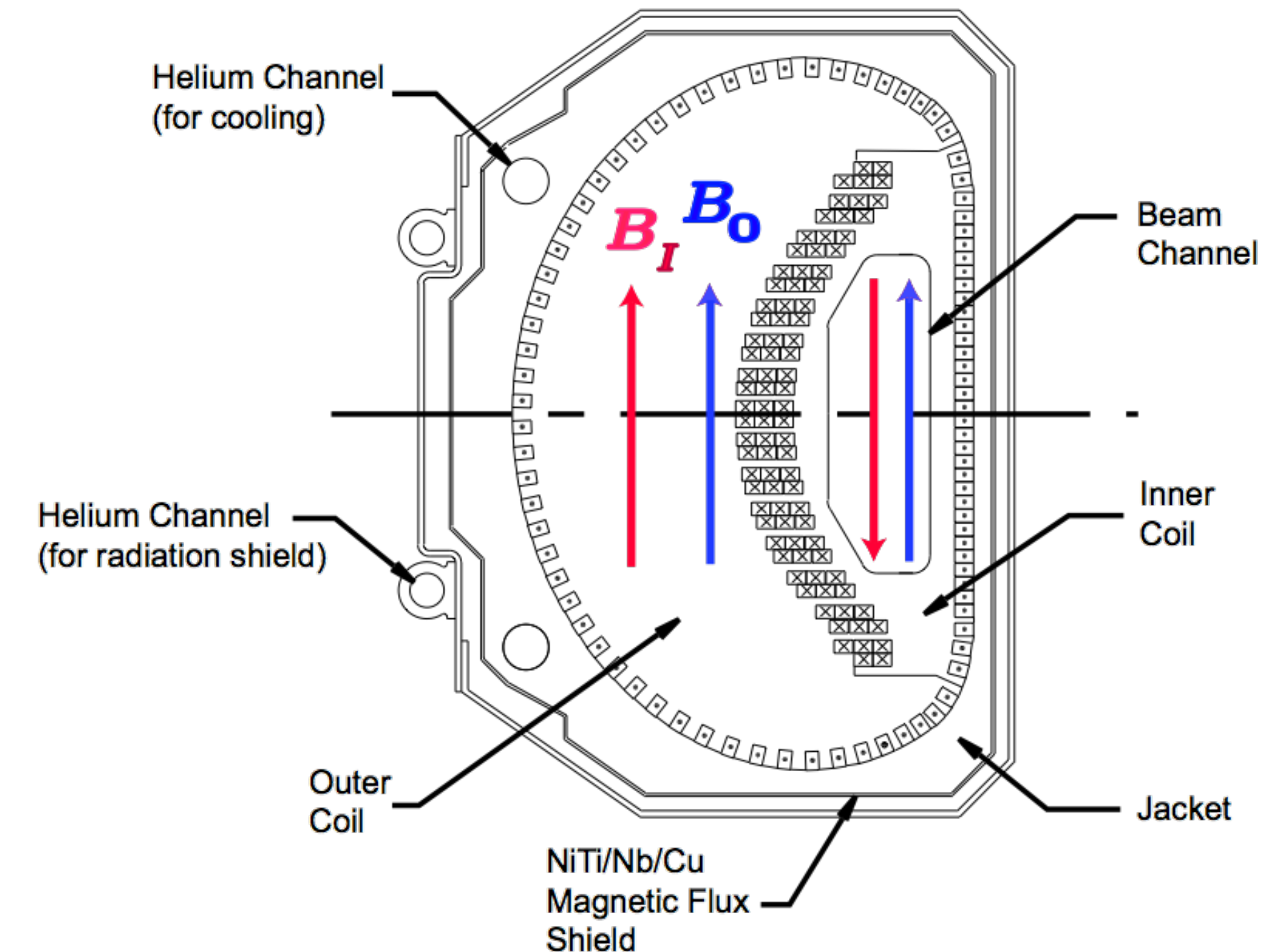
- **Scintillating Paddles:** Monitoring temporal distribution
- **Scintillating Fibers:** Map of transverse profile, guides μ tuning into the ring

Muon Beam Injection

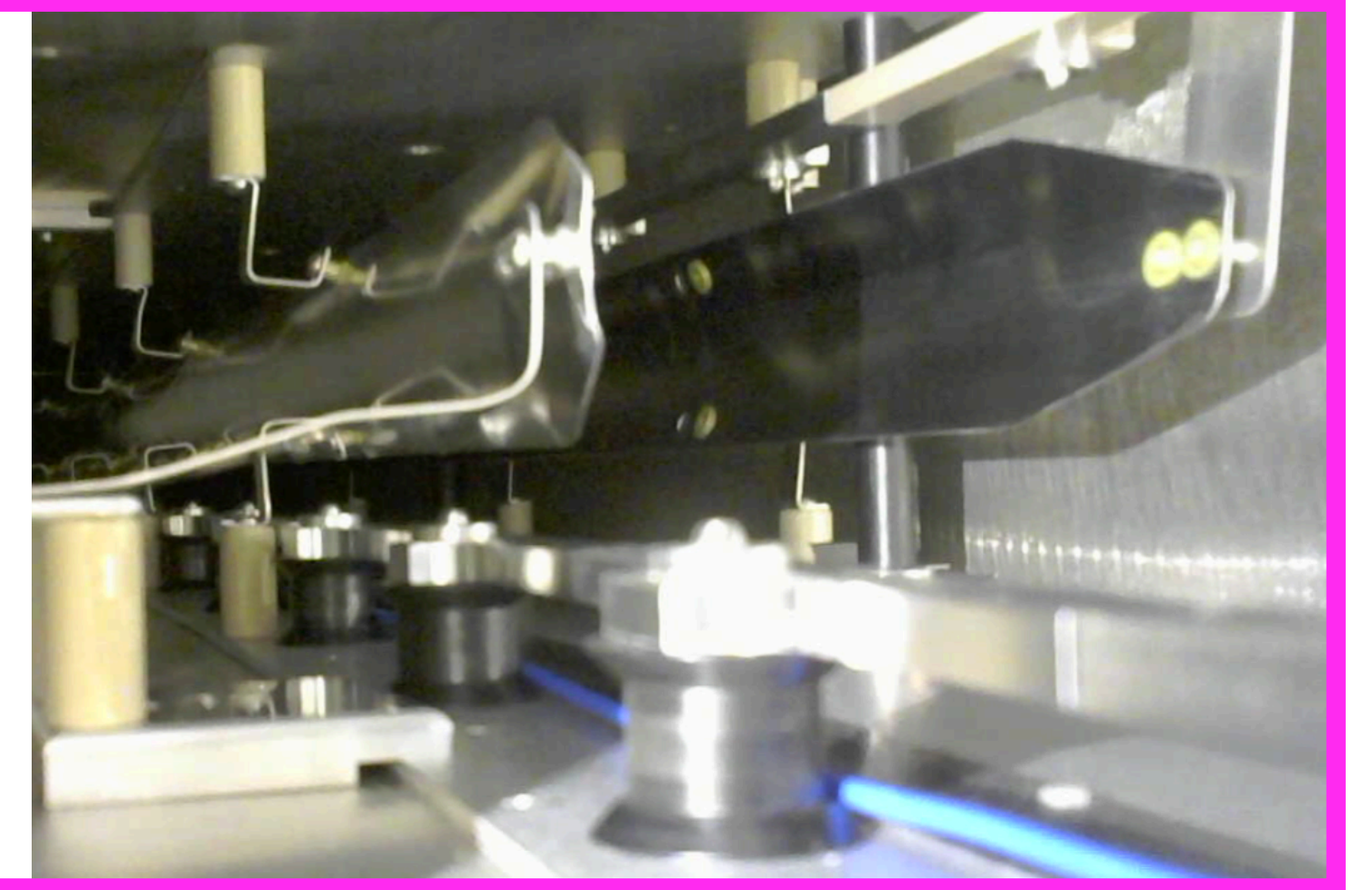
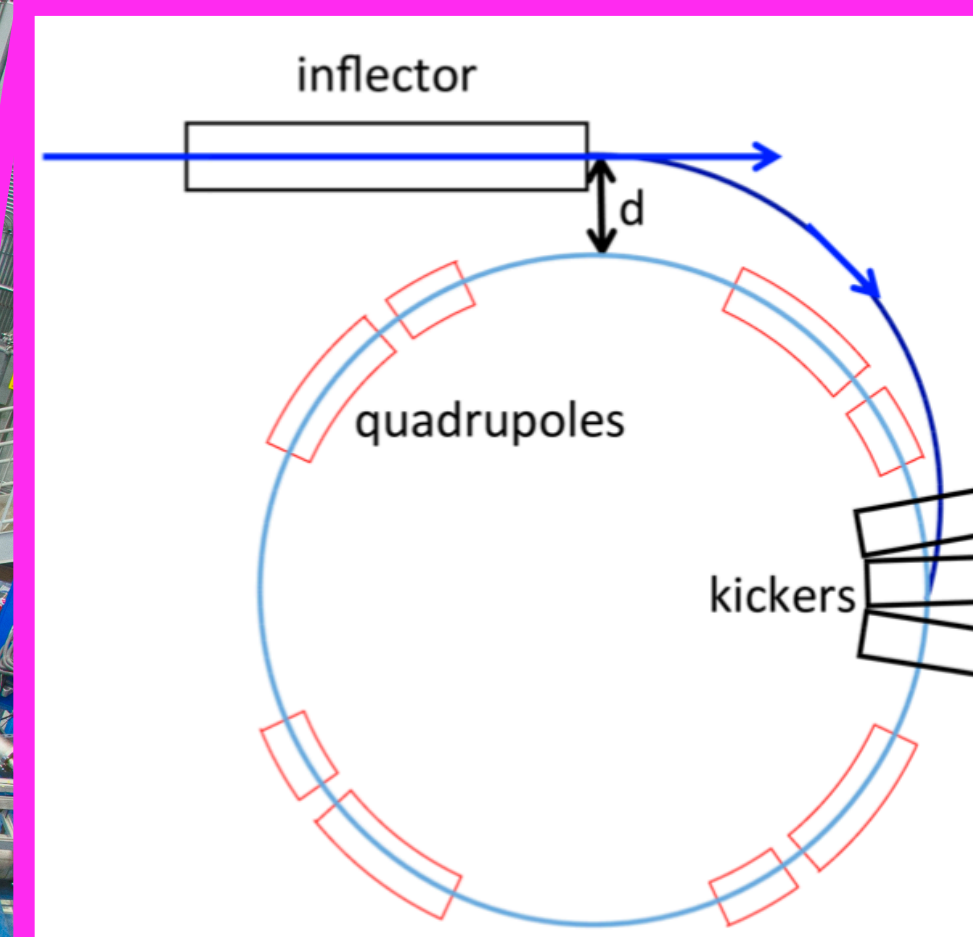
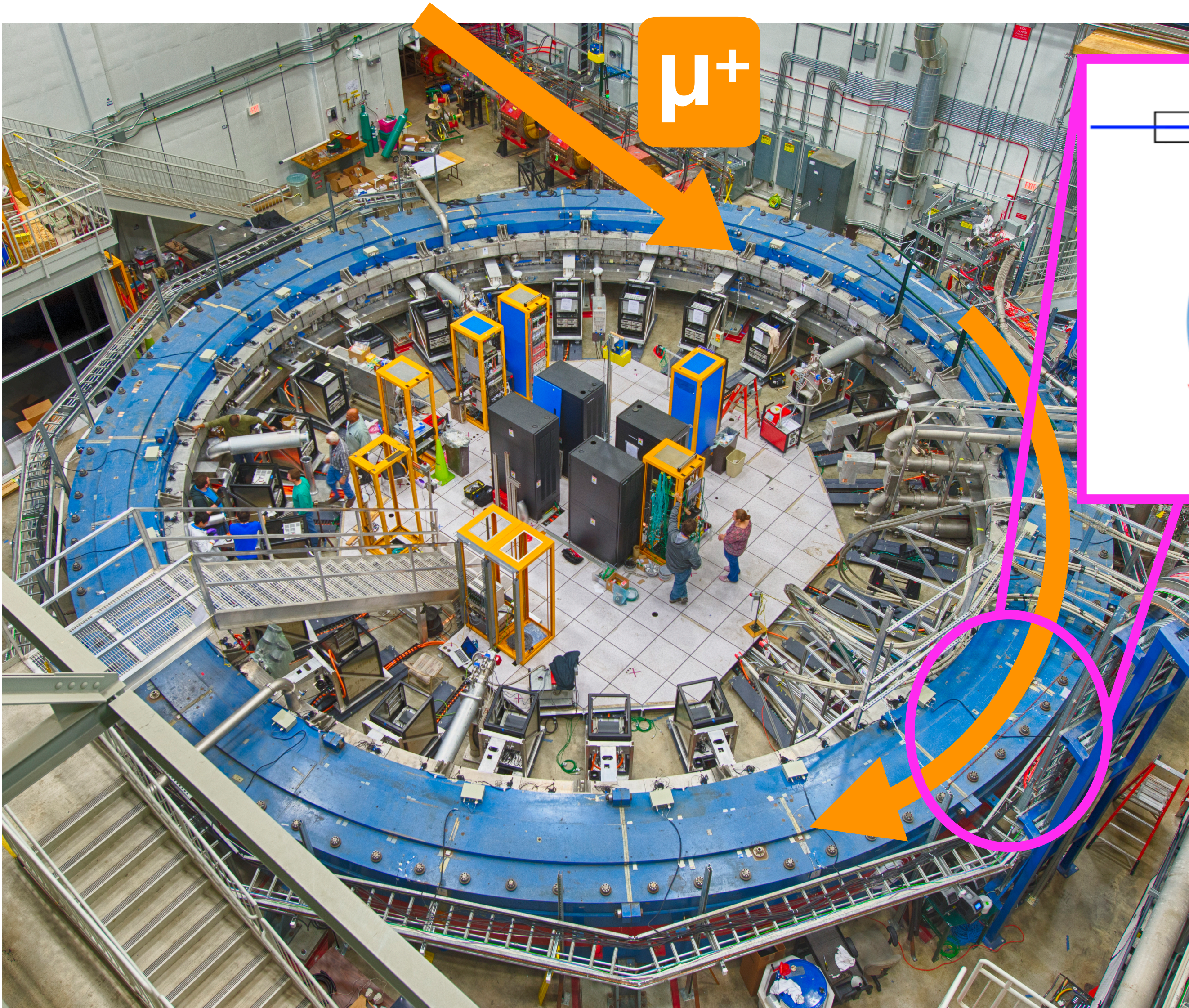


Inflector Magnet

- Need to cancel field in beam channel
- Prevents strong deflection of the beam
- Minimal perturbation to storage magnetic field



Muon Beam Storage and Focusing



3 Kicker Magnets

- After inflector, muons enter storage region at $r = 77$ mm outside central closed orbit
- Deliver pulse in < 149 ns to muon beam
- Steer muons onto stored orbit

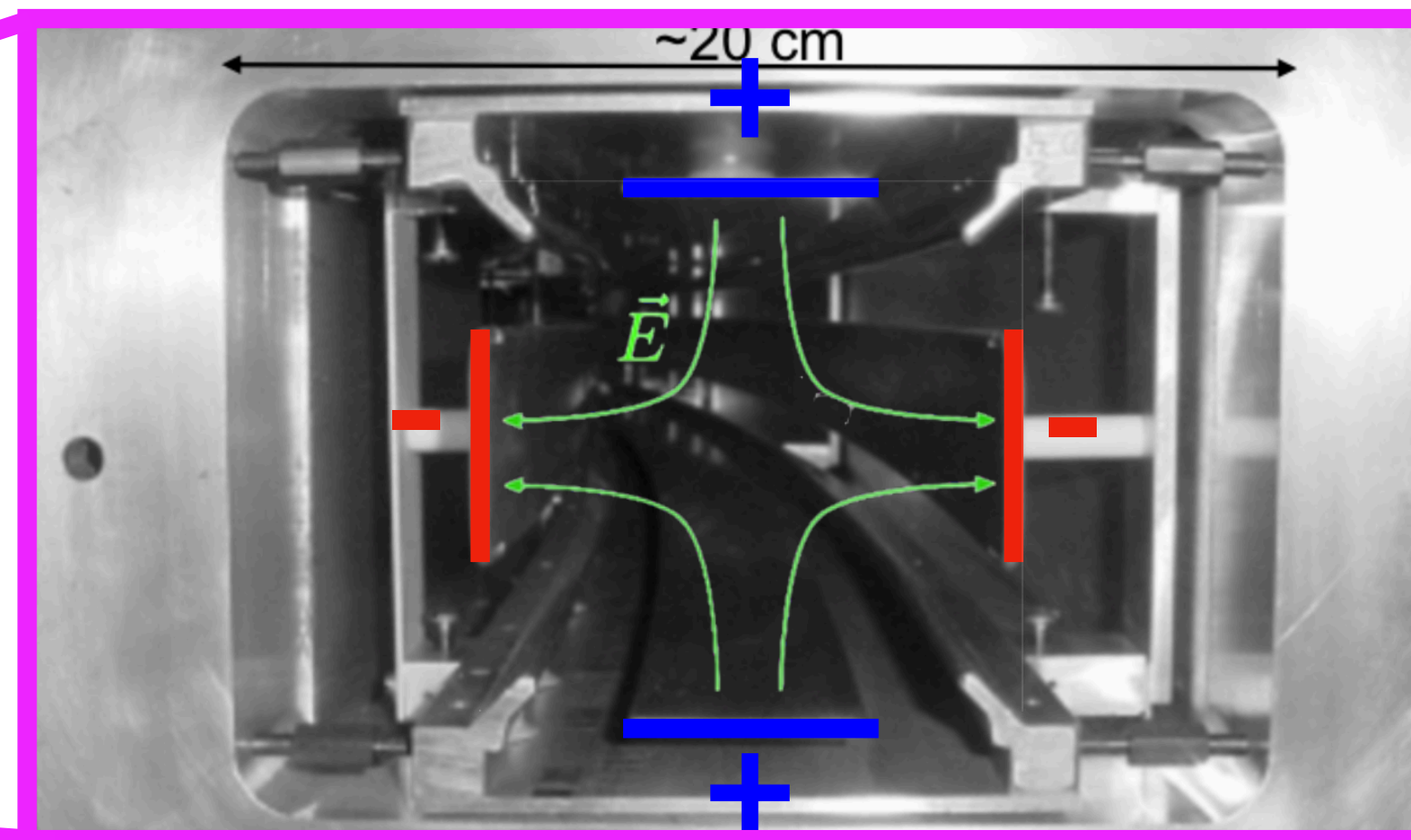
Muon Beam Storage and Focusing



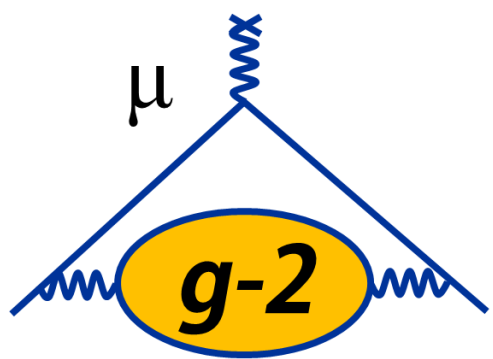
μ^+

Electrostatic Quadrupoles

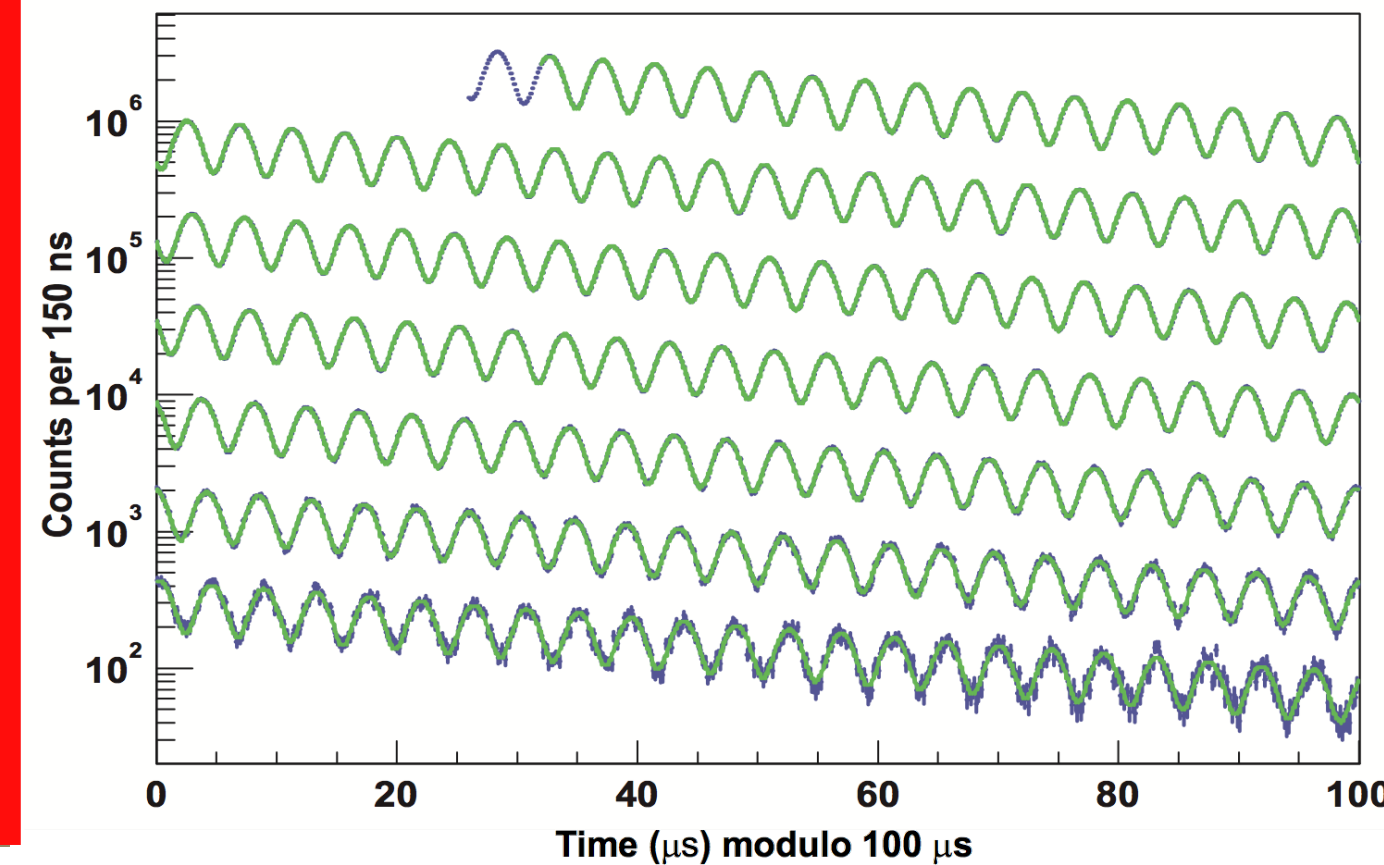
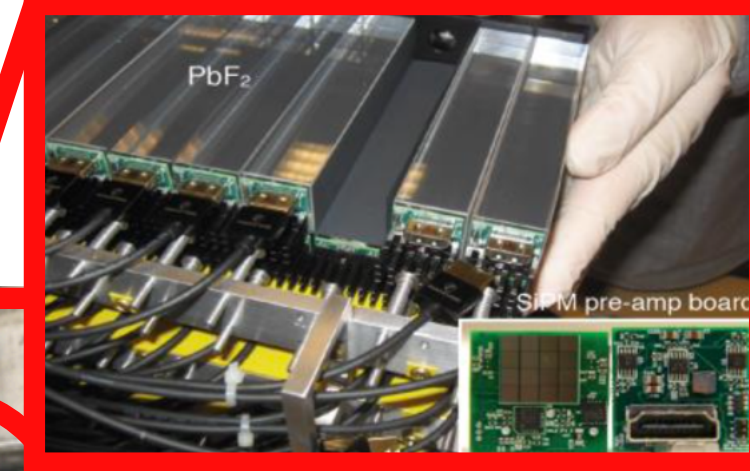
- Drives the muons towards the central part of storage region vertically
- Aluminum electrodes cover $\sim 43\%$ of total circumference



Measuring Muon Spin Precession (ω_a)



μ^+

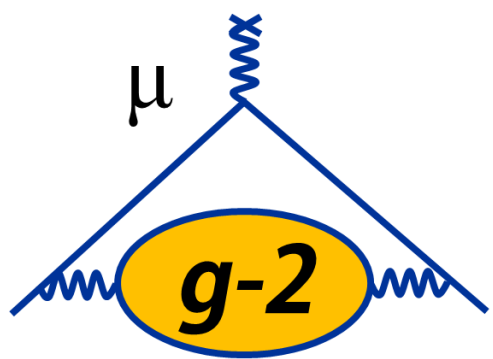


24 finely-segmented PbF₂ crystal calorimeters

- Self-analyzing decay: $\mu^+ \rightarrow e^+ \bar{\nu}_\mu \nu_e$
- Highest-energy e^+ emitted preferentially along muon spin
- Results in sinusoidally-oscillating arrival time of these e^+ in calorimeters

UMassAmherst

Measuring Muon Spin Precession (ω_a)

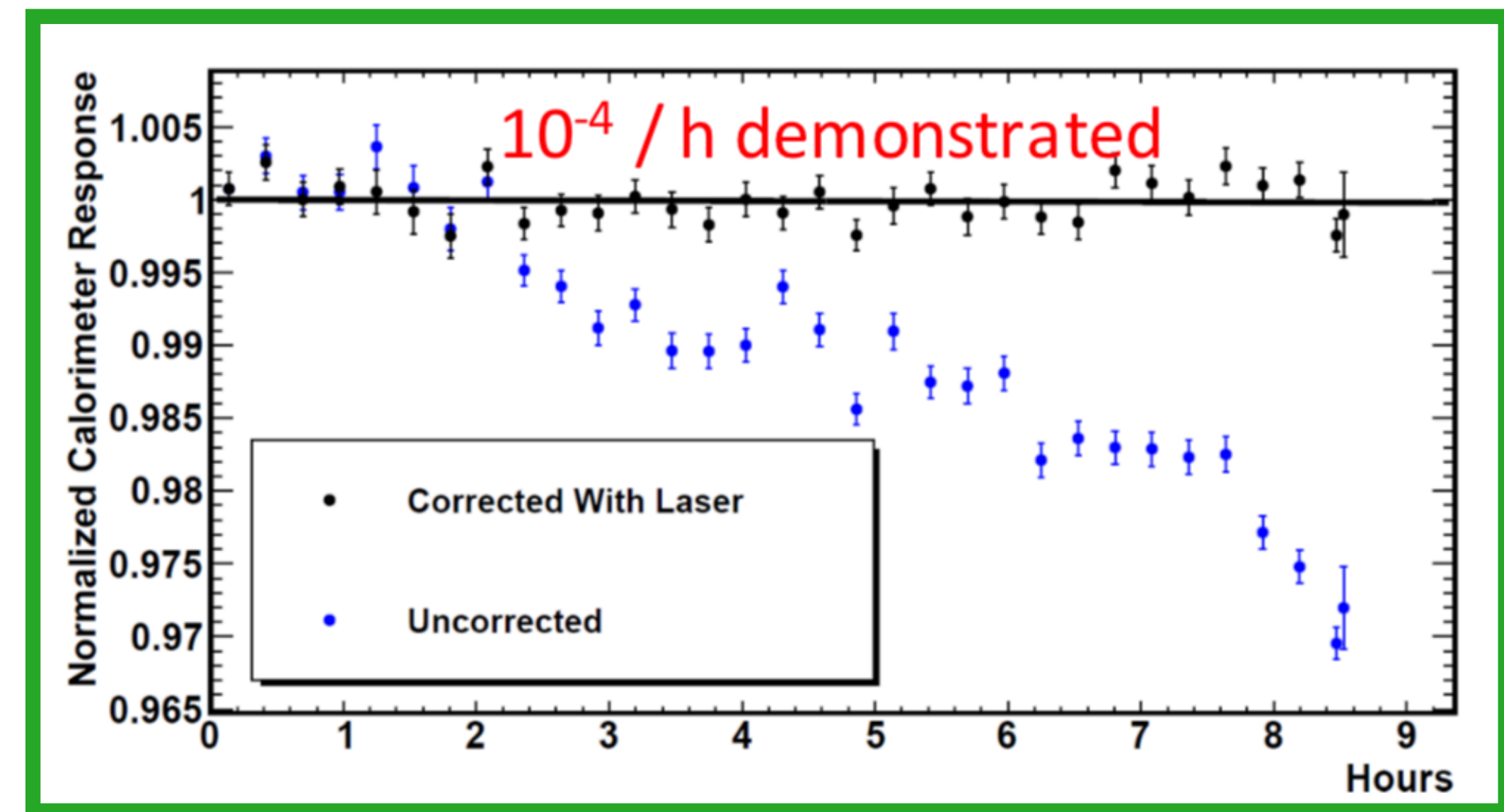


μ^+

Laser System

- Calibrate calorimeter gain response throughout data taking
- Demonstrated stability to $10^{-4}/\text{hr}$

Inside the laser hut



Measuring Muon Spin Precession (ω_a)

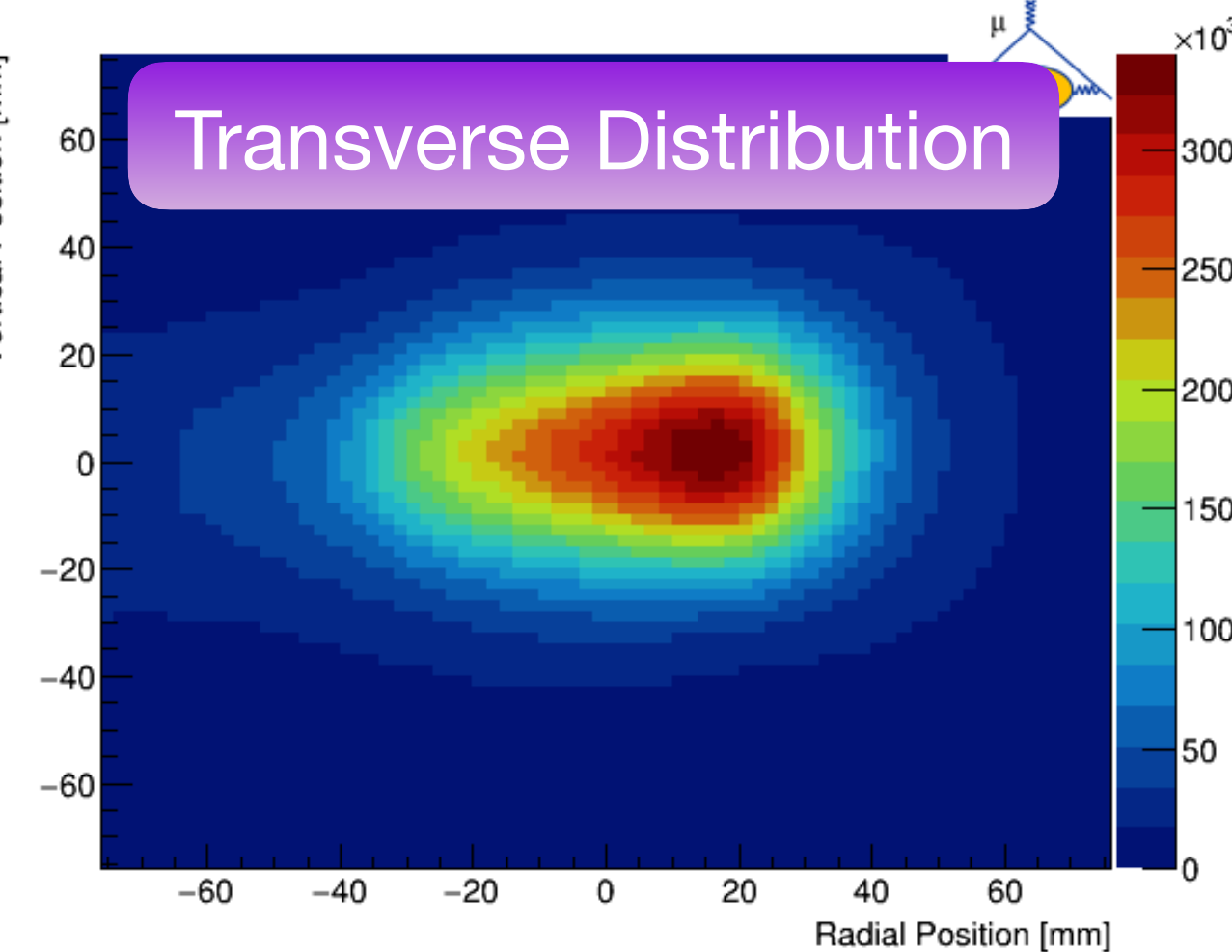


μ^+

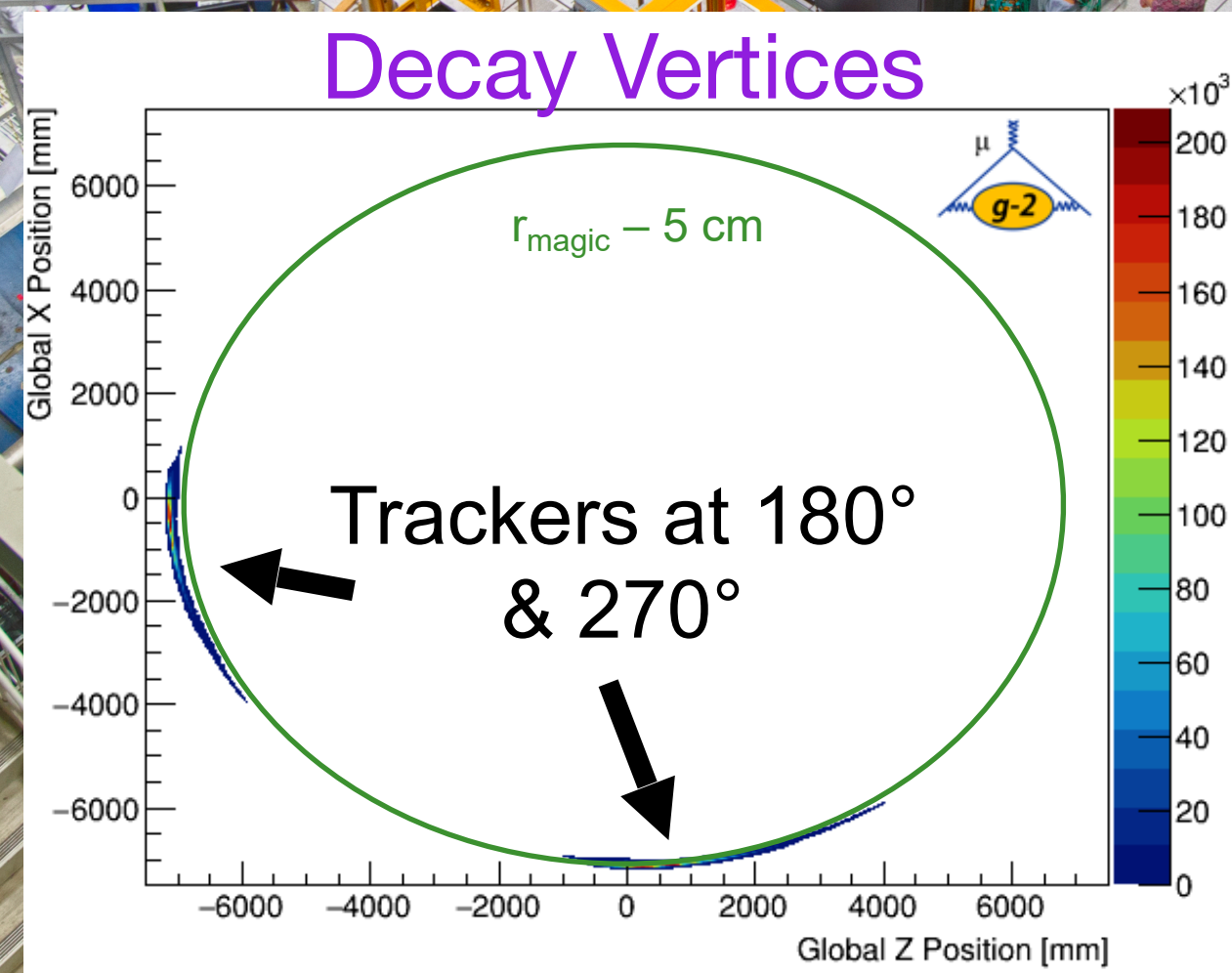
Storage Region

Muon's view of the storage region

Transverse Distribution



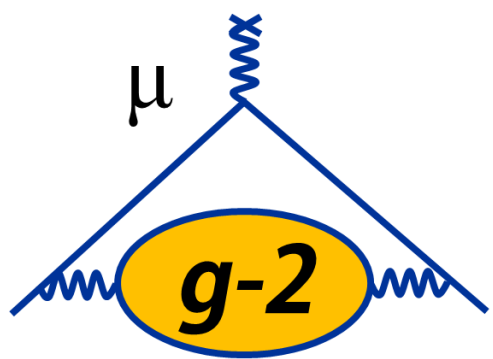
Decay Vertices



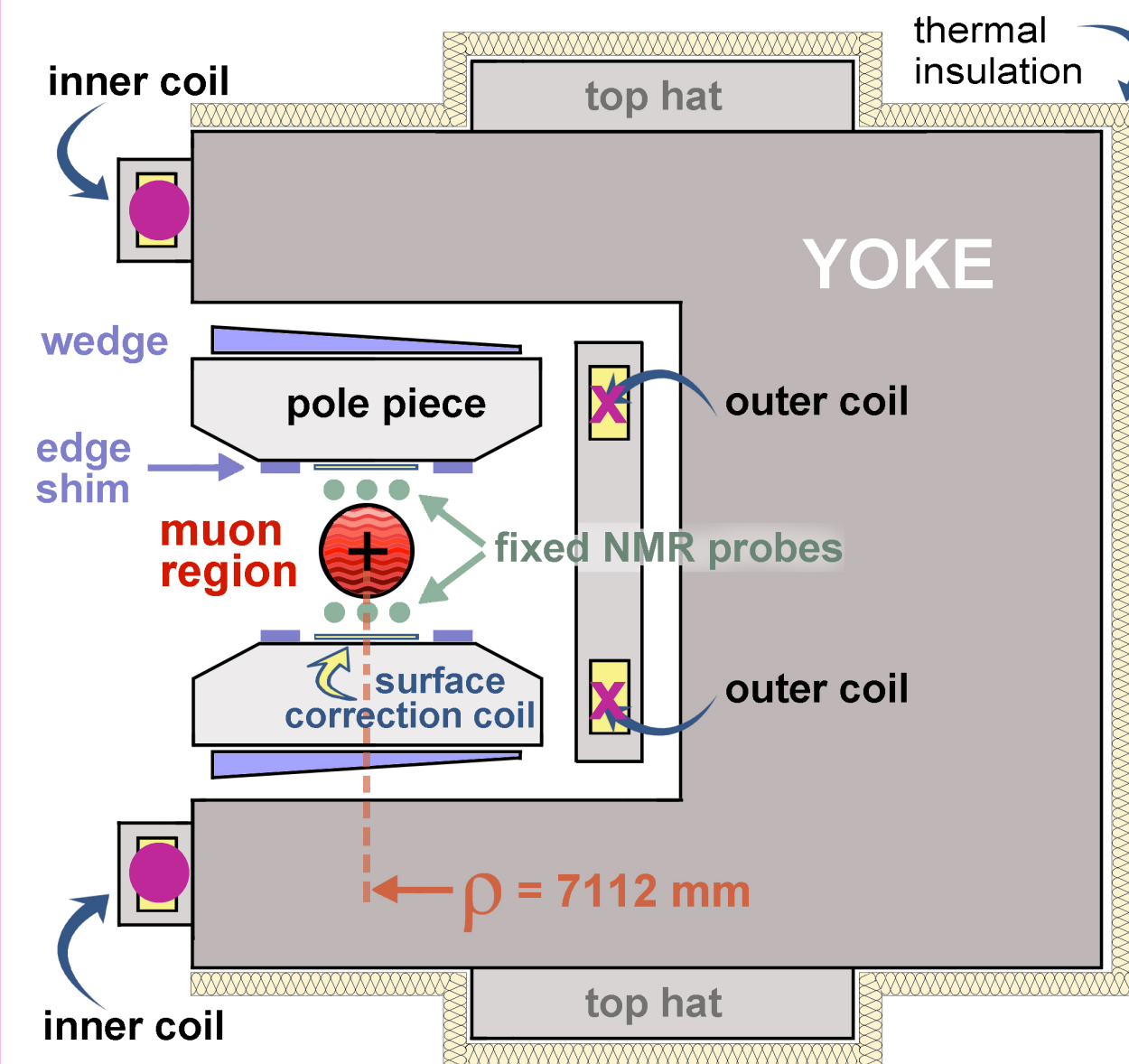
Two straw-tracker stations

- Reconstruct the muon beam distribution from e^+ hits
- Tracker module: 128 straws/module
- 8 modules per station

Magnet Anatomy



Current direction indicated by ● and x



g-2 Magnet in Cross Section

$B = 1.45 \text{ T}$ (~5200 A)

- Power supply with feedback to fine-tune field in real time

12 C-shaped yokes

- 3 upper and 3 lower poles per yoke
- 72 total poles

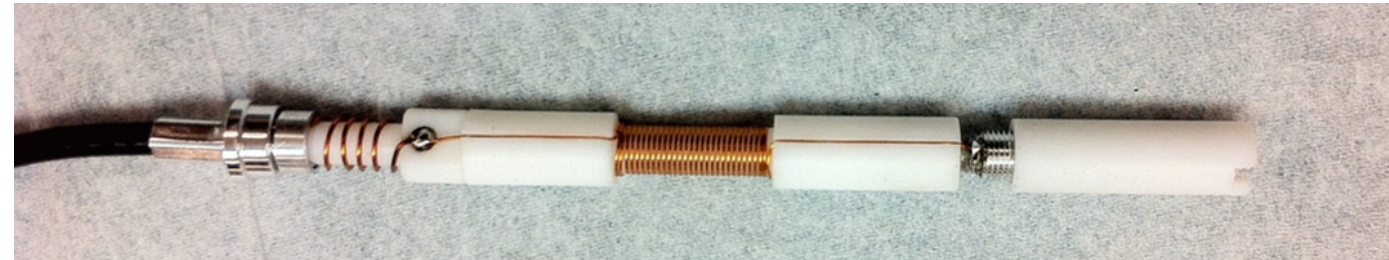
Field Shape

- Determined by positioning of pole pieces, wedge-shaped pieces of steel, programmable surface coils

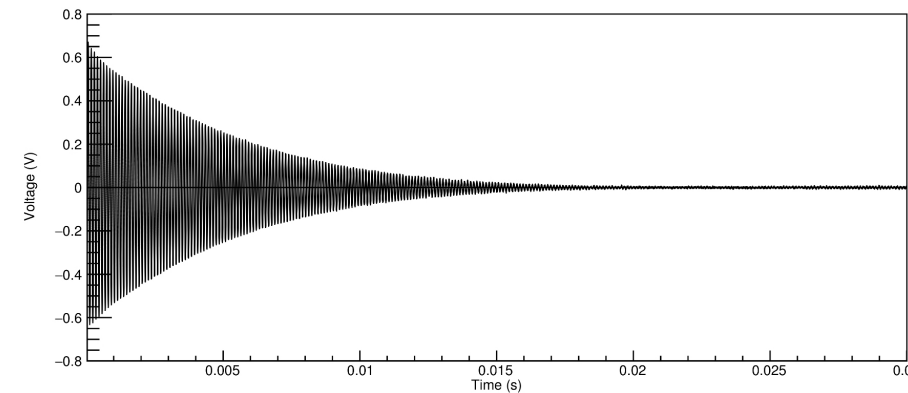
Monitoring and Mapping the Magnetic Field



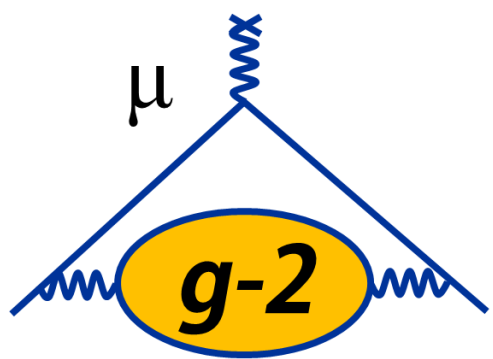
Pulsed NMR



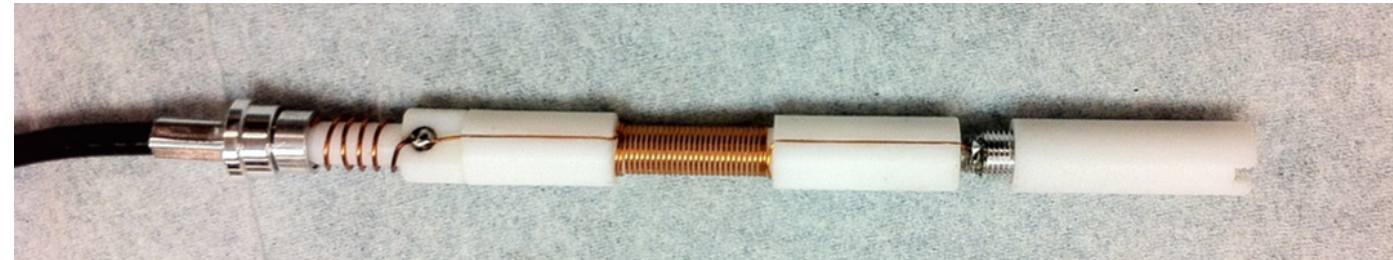
- Deliver $\pi/2$ pulse to probe, induce & record the free-induction decay (FID)
- Extracted frequency precision: 10 ppb/FID



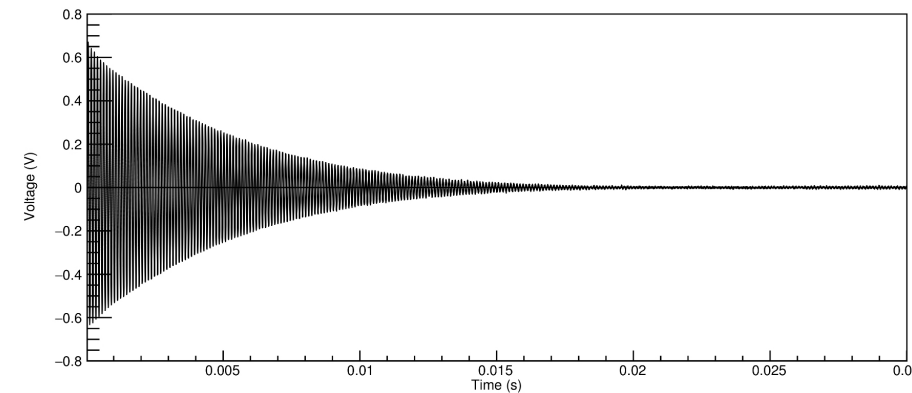
Monitoring and Mapping the Magnetic Field



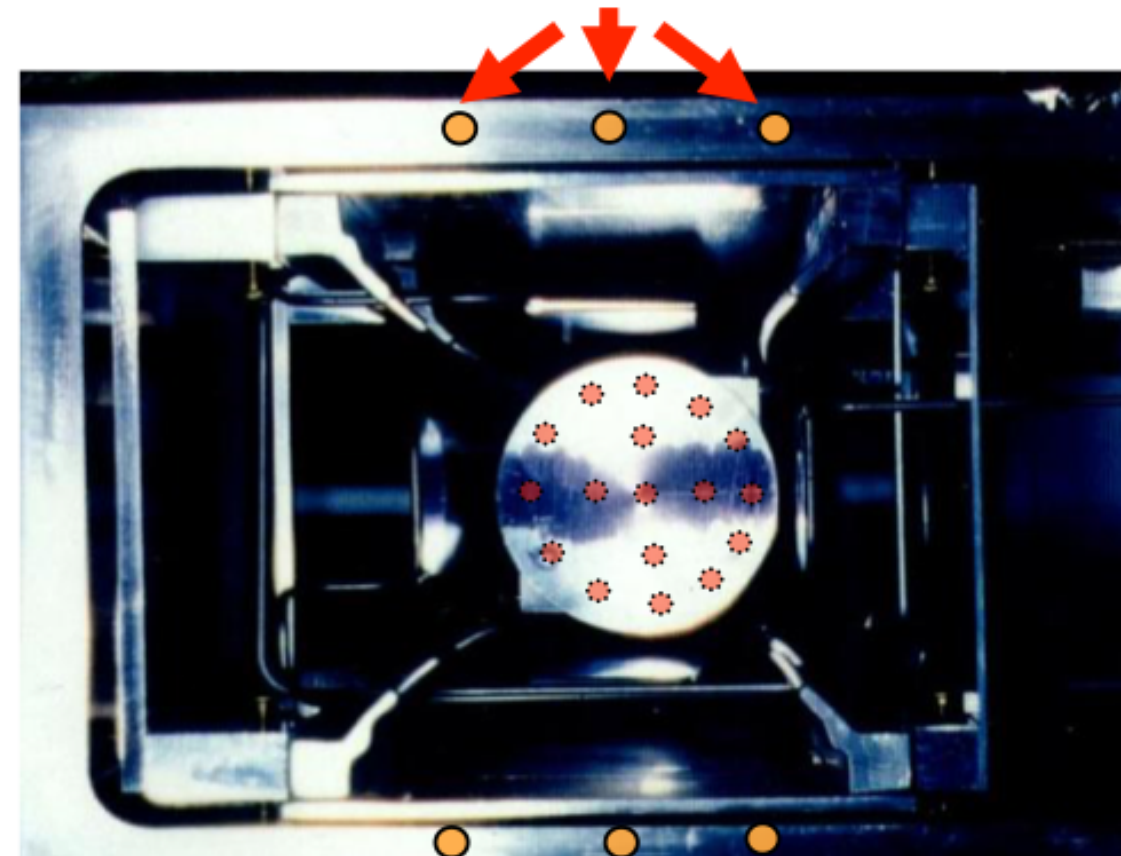
Pulsed NMR



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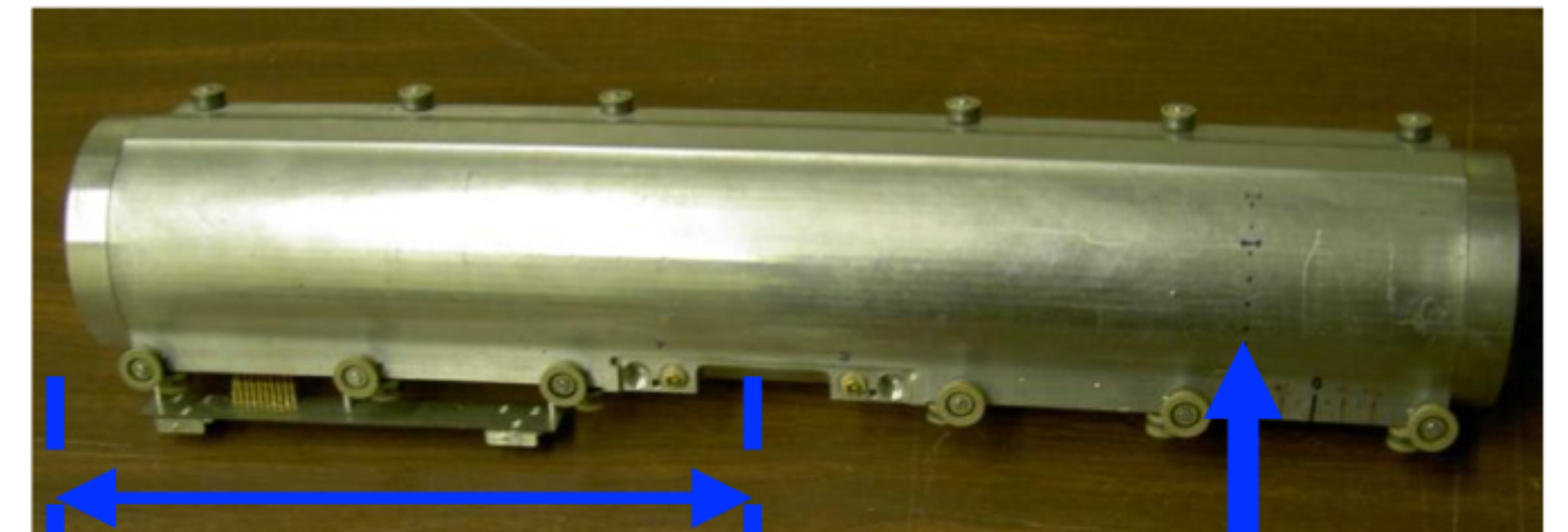


Fixed probes on vacuum chambers



- Measure field while muons are in ring
– 378 probes **outside** storage region

Trolley matrix of 17 NMR probes

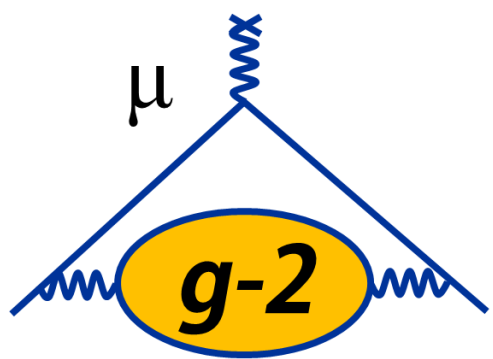


Electronics,
Microcontroller,
Communication

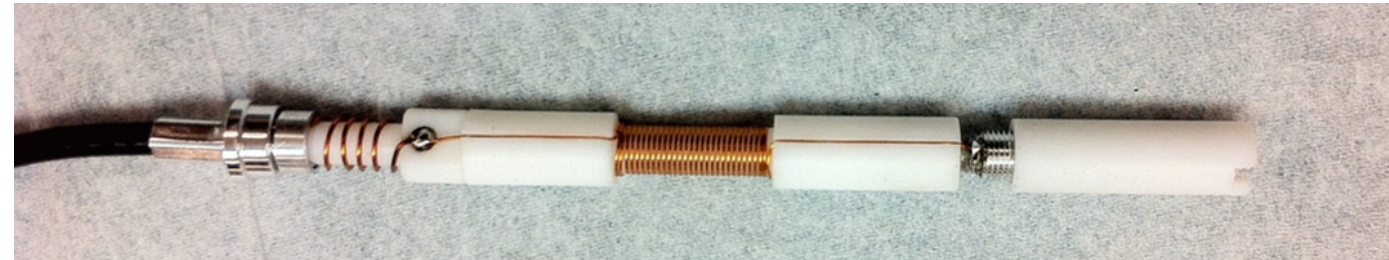
Positon of NMR probes

- Measure field in storage region during **specialized runs** when **muons are not being stored**

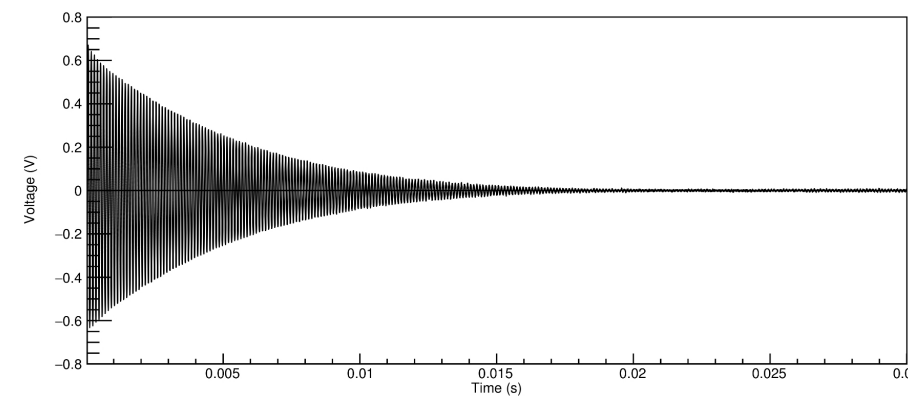
Monitoring and Mapping the Magnetic Field



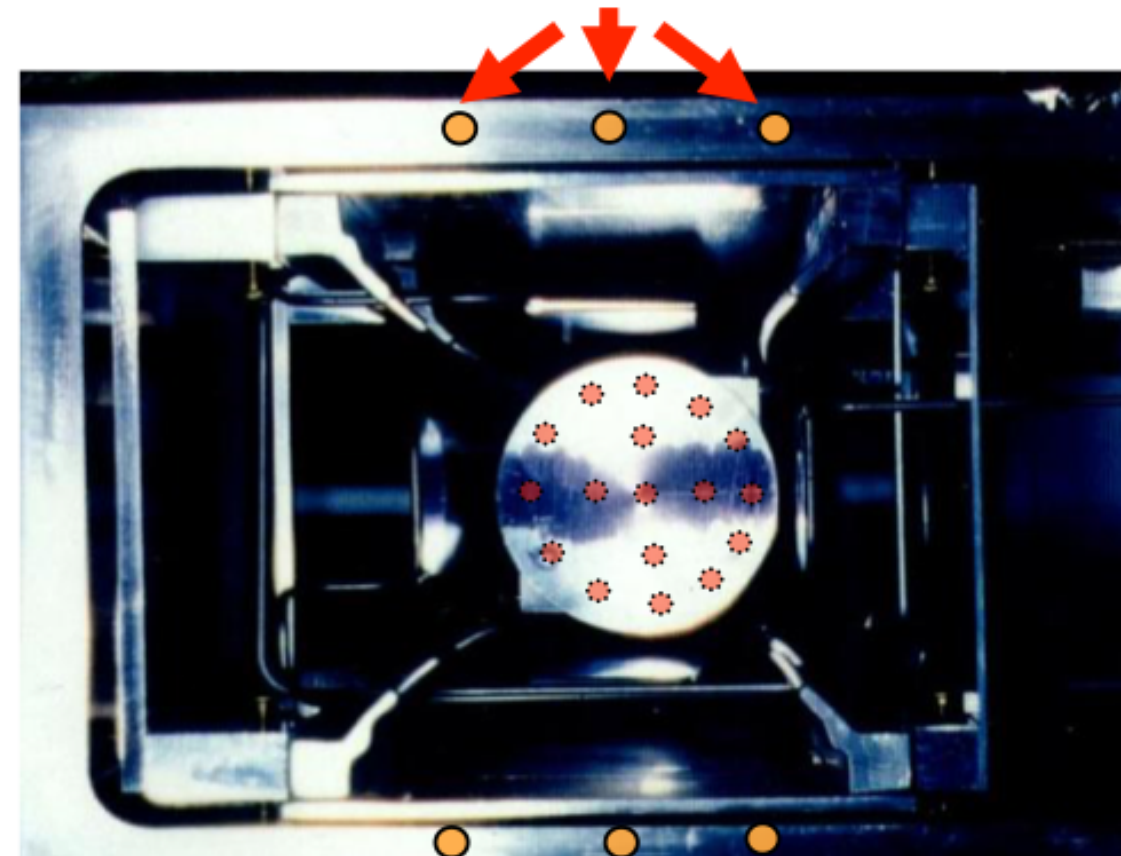
Pulsed NMR



- Deliver $\pi/2$ pulse to probe, induce & record the free-induction decay (FID)
- Extracted frequency precision: 10 ppb/FID

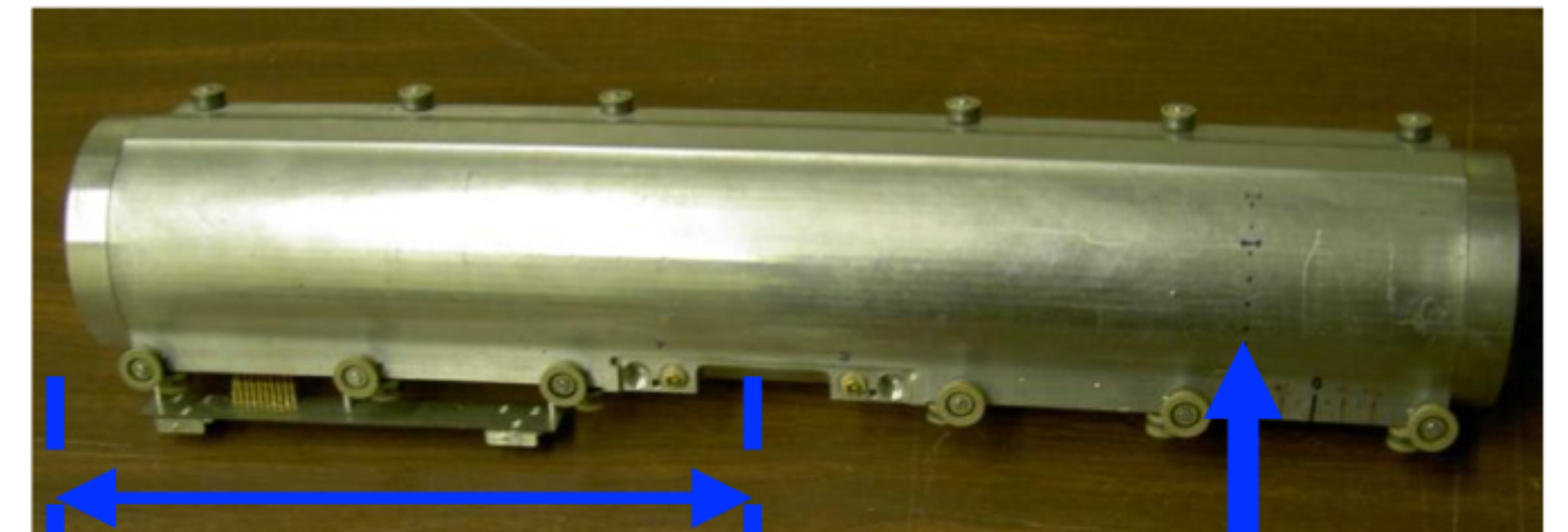


Fixed probes on vacuum chambers



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Trolley matrix of 17 NMR probes



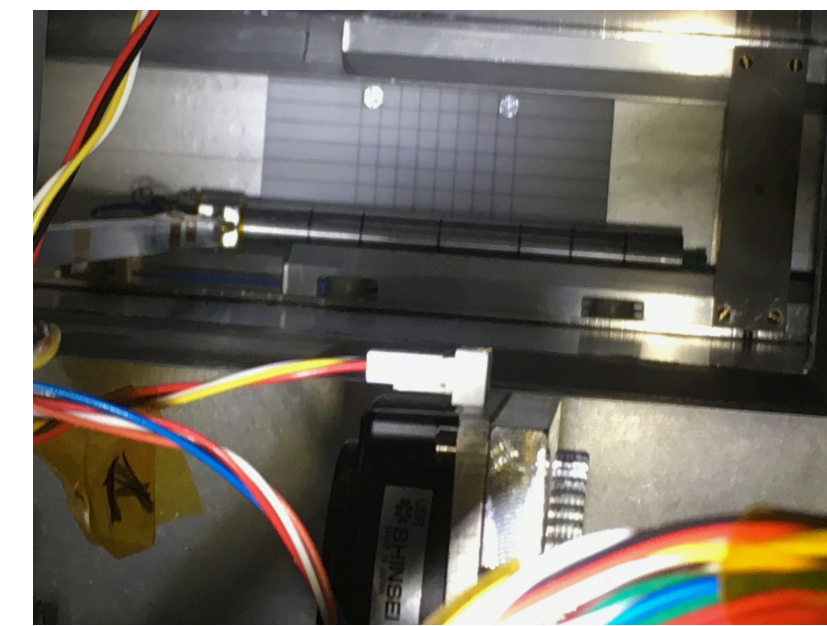
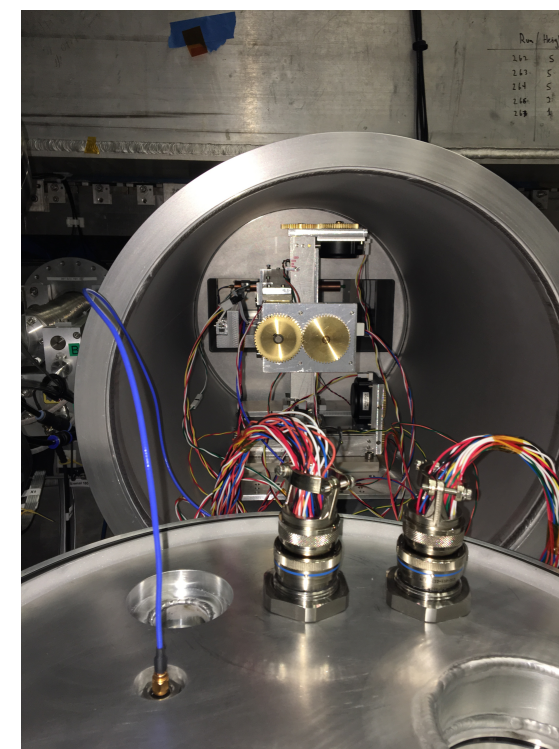
Electronics,
Microcontroller,
Communication

Position of NMR probes

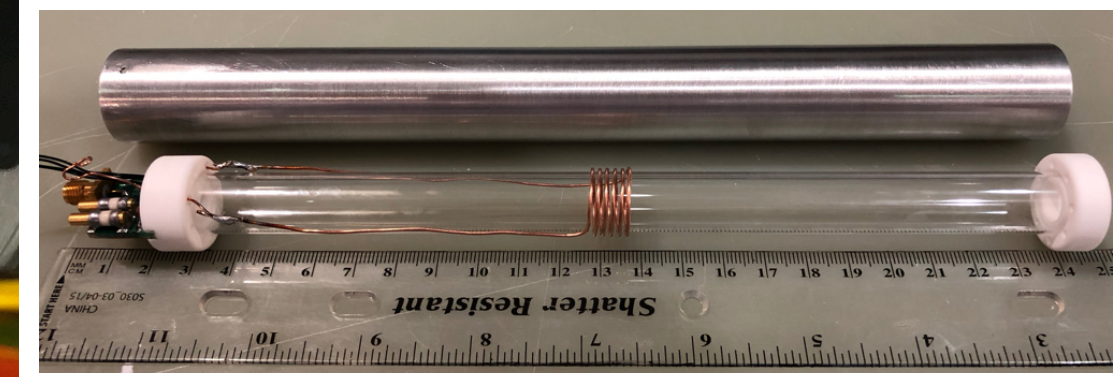
- Measure field in storage region during **specialized runs** when **muons are not being stored**

- **Trolley** probes **calibrated to free-proton Larmor frequency**

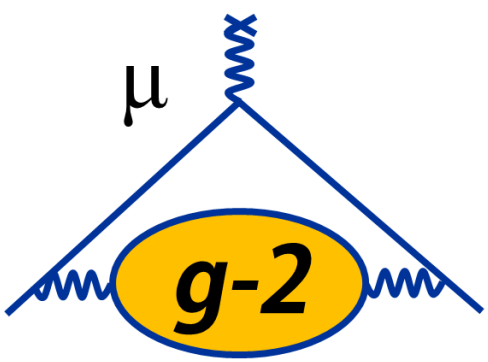
- Calibrate trolley probes using a special probe that uses a water sample
- Measurements in specially-shimmed region of ring



Plunging Probe



Systematic Uncertainty Comparison: E821 and E989



$$a_{\mu} = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_{\mu}}{m_e} \frac{g_e}{2}$$

- New hardware (calorimeters, trackers, NMR)
- Improved analysis techniques
- Reduce uncertainties by at least a factor of 2.5

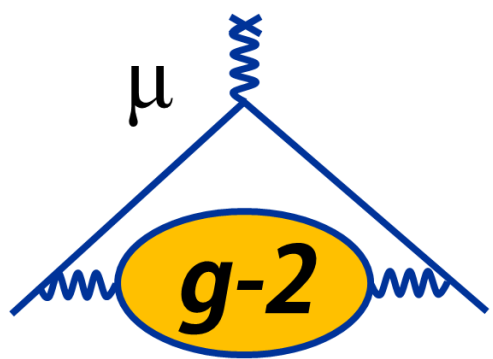
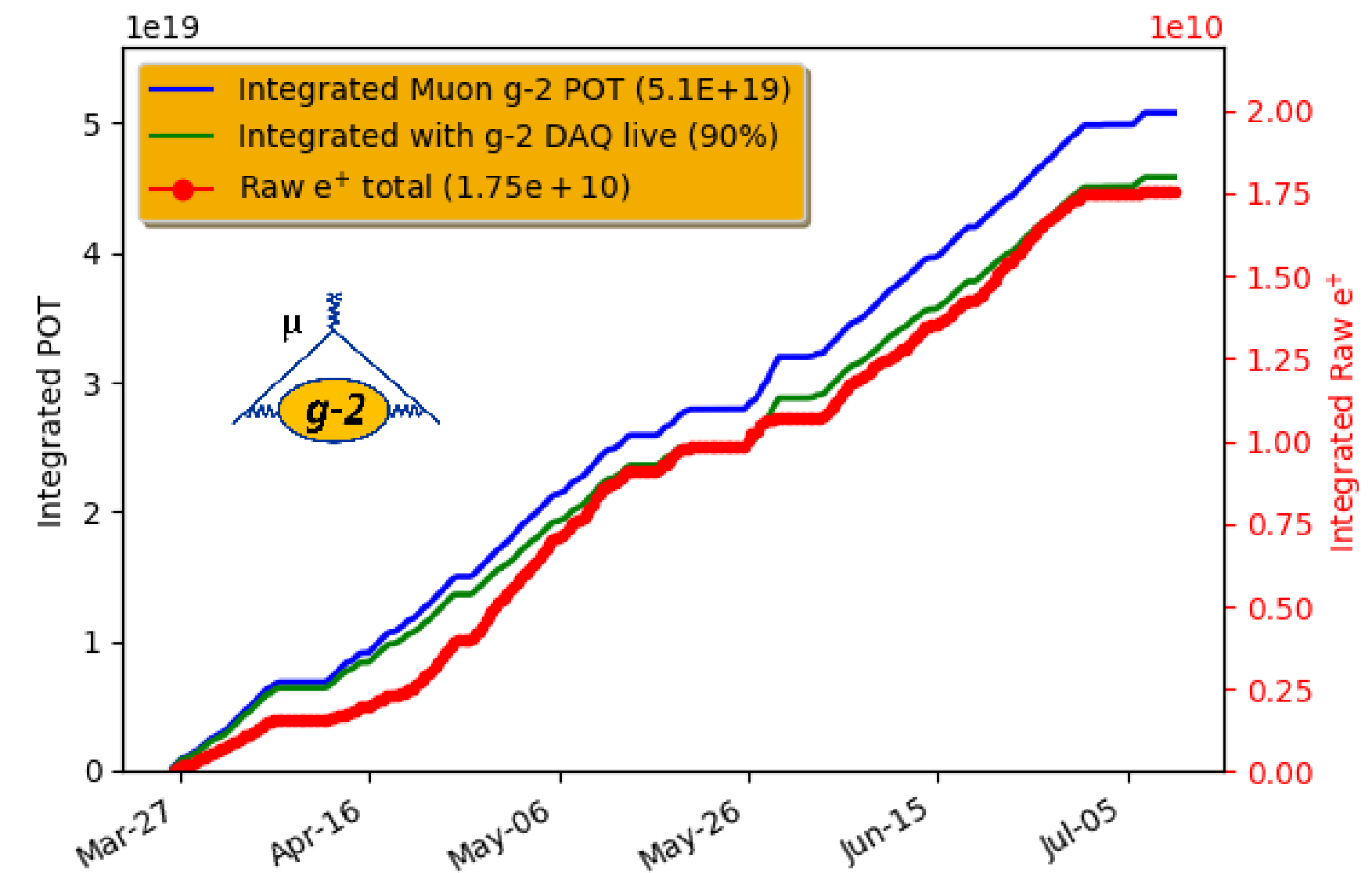
ω _a Goal: Factor of 3 Improvement		
Category	E821 (ppb)	E989 Goal (ppb)
Gain Changes	120	20
Lost Muons	90	20
Pileup	80	40
Horizontal CBO	70	< 30
E-field/pitch	110	30
Quadrature Sum	214	70

ω _p Goal: Factor of 2.5 Improvement		
Category	E821 (ppb)	E989 Goal (ppb)
Field Calibration	50	35
Trolley Measurements	50	30
Fixed Probe Interpolation	70	30
Muon Convolution	30	10
Time-Dependent Fields	–	5
Others	100	50
Quadrature Sum	170	70

Run 1 Overview

- Data taking period: April—July 2018
- A number of changing conditions as we optimized hardware
- Accumulated $\sim 1.1 \times$ BNL statistics (after data quality cuts) — $\delta\omega_a(\text{stat}) \sim 410 \text{ ppb}$
- Field uniformity $\sim 2x$ better than BNL

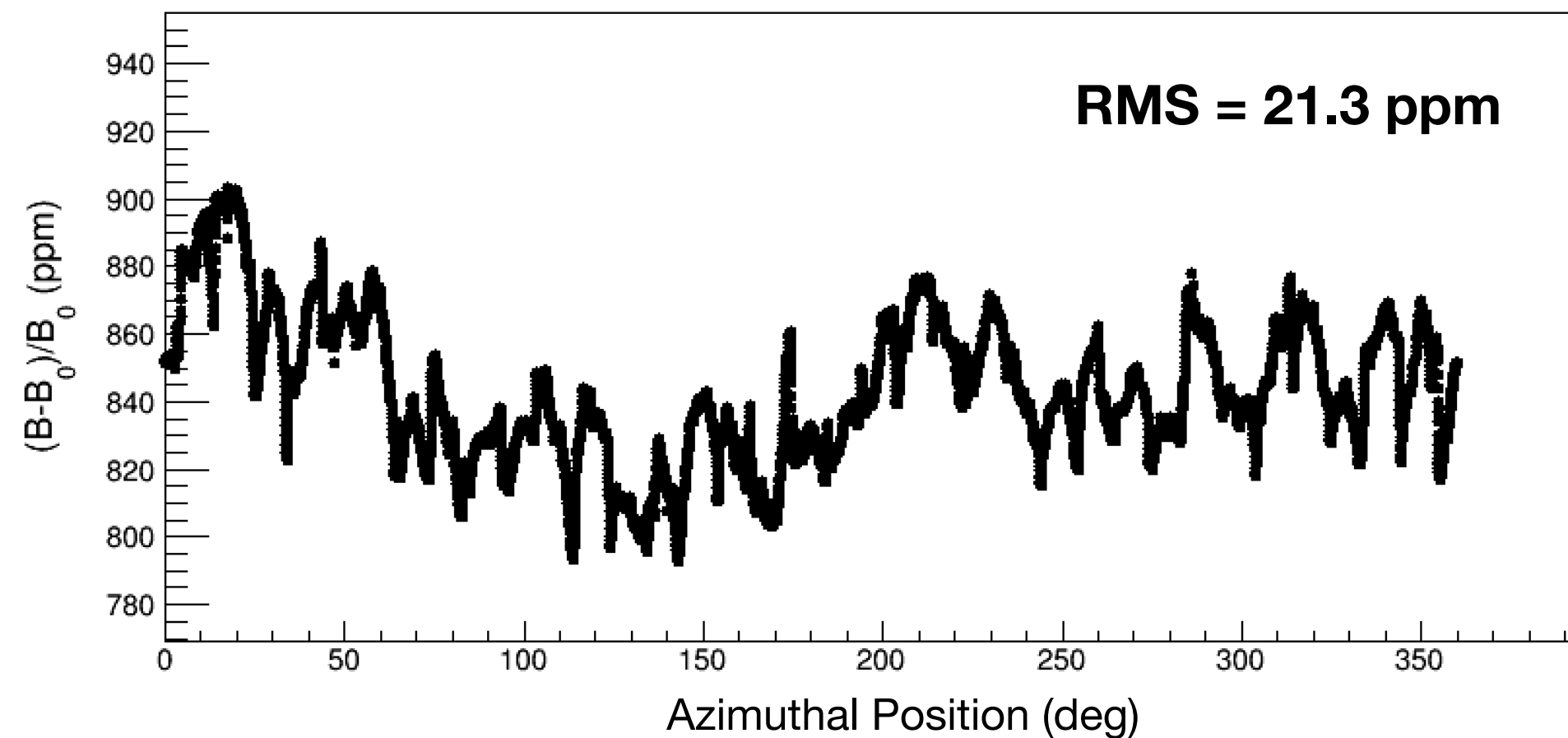
Accumulated statistics



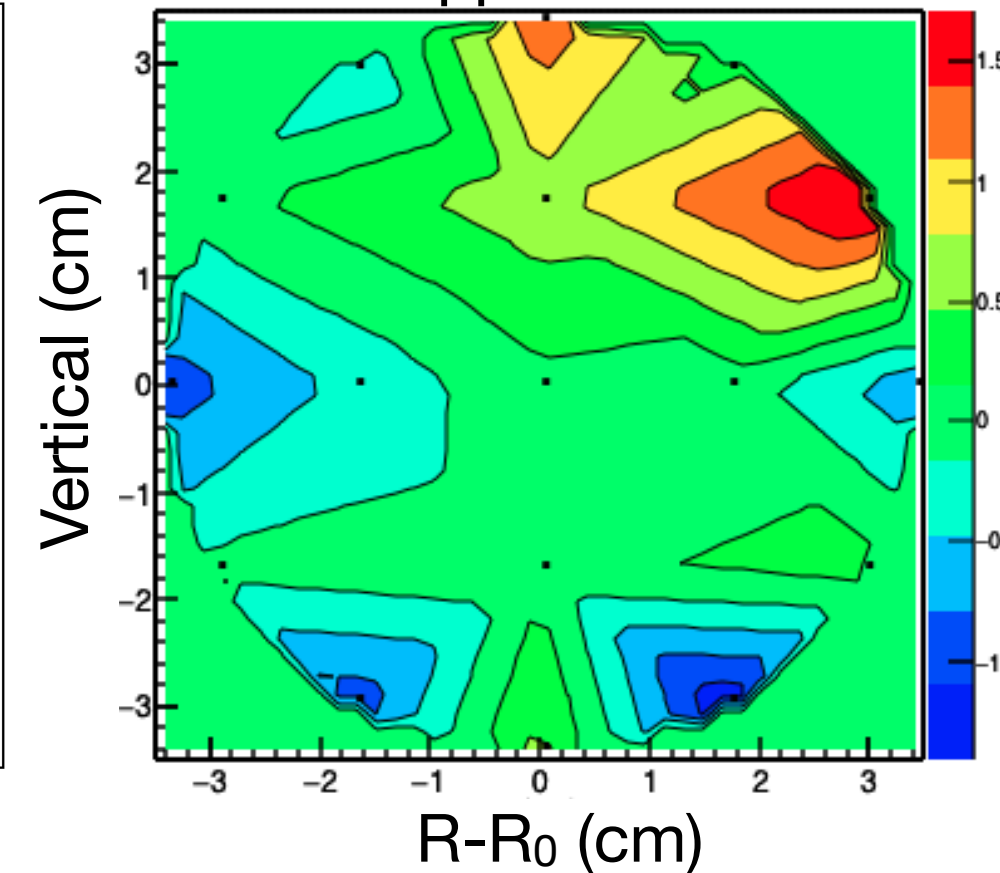
Typical Field Map

Dipole Moment

RMS = 21.3 ppm



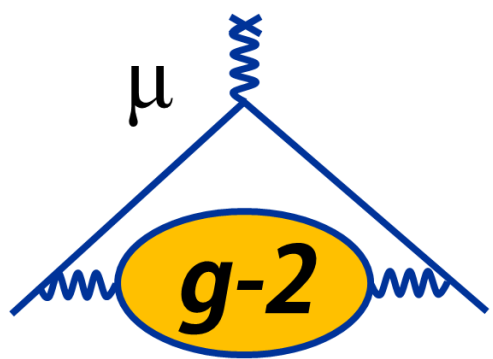
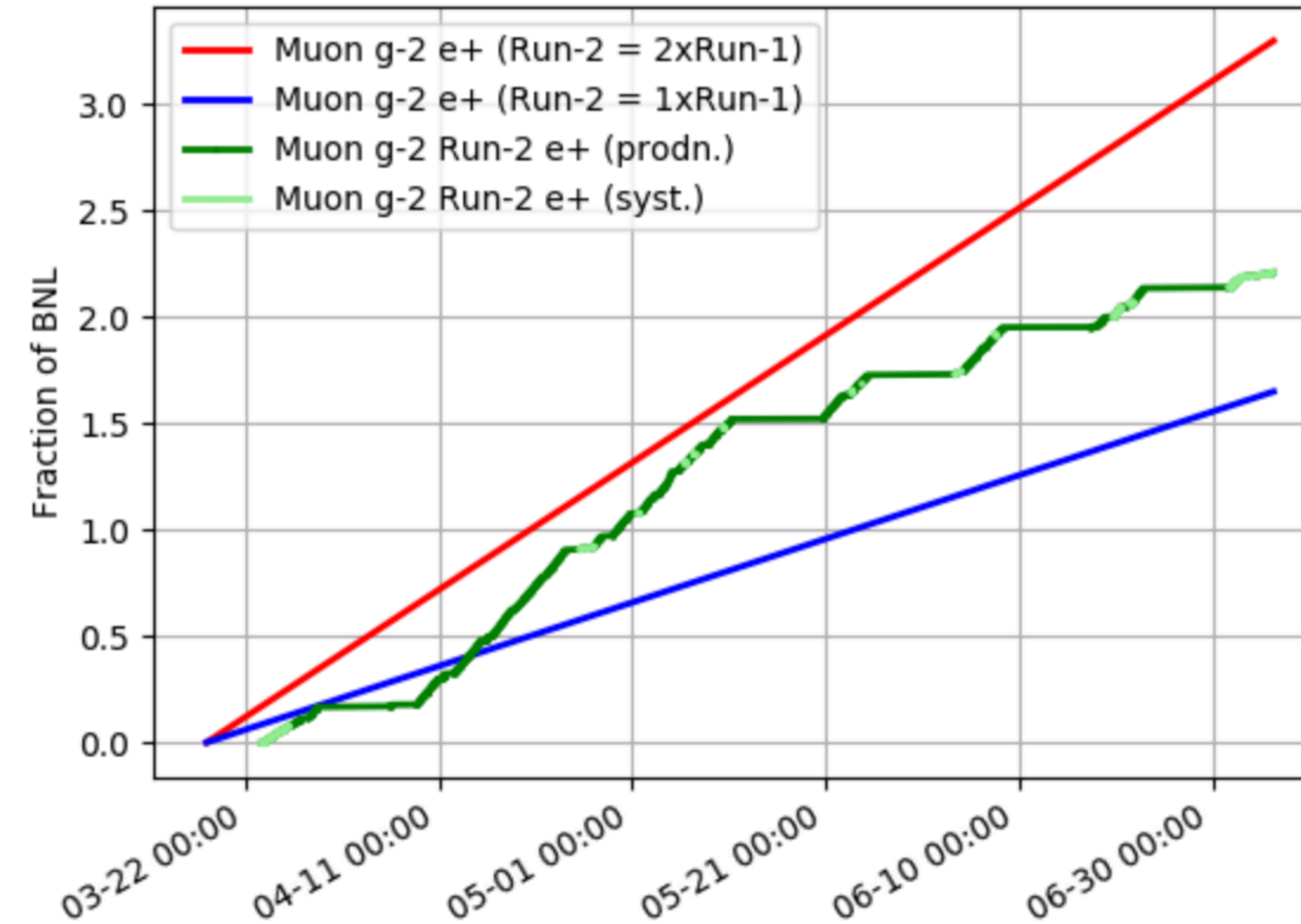
Azimuthal average
250-ppb contours



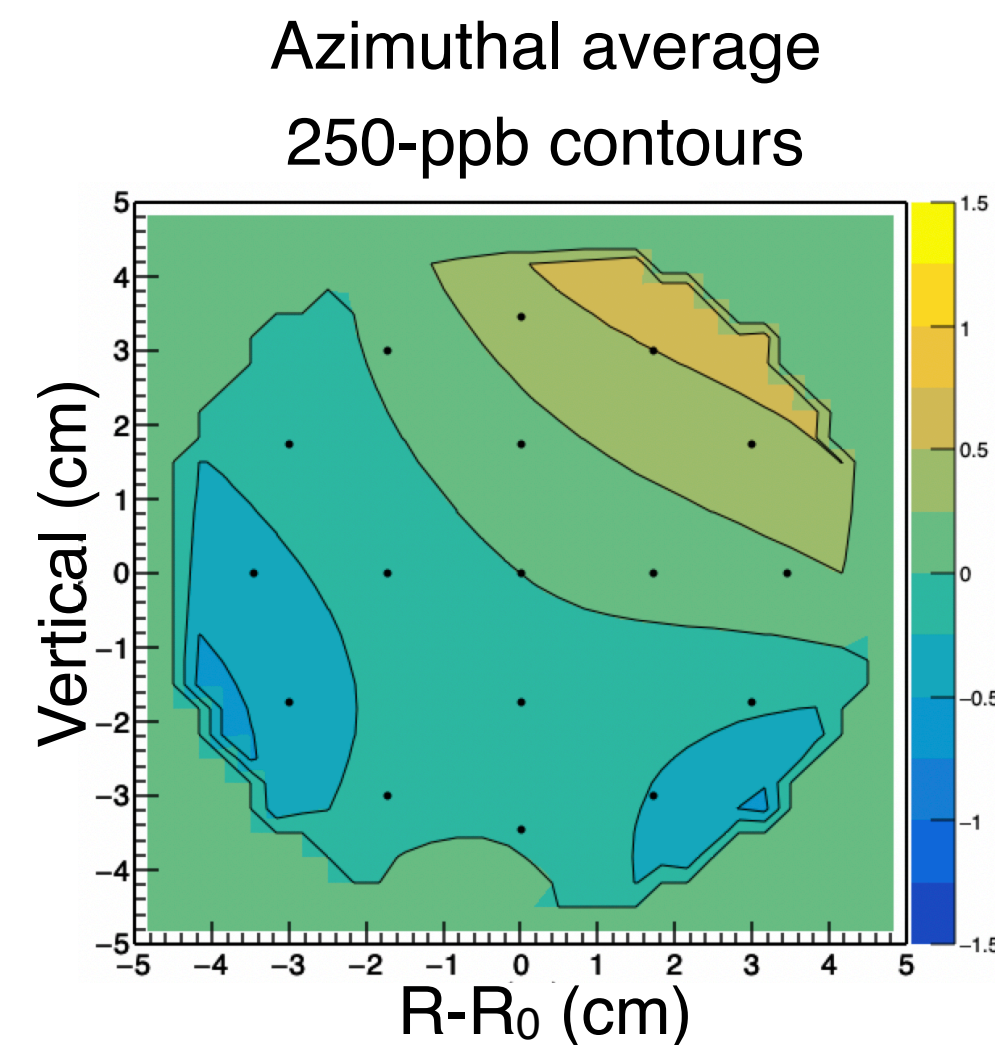
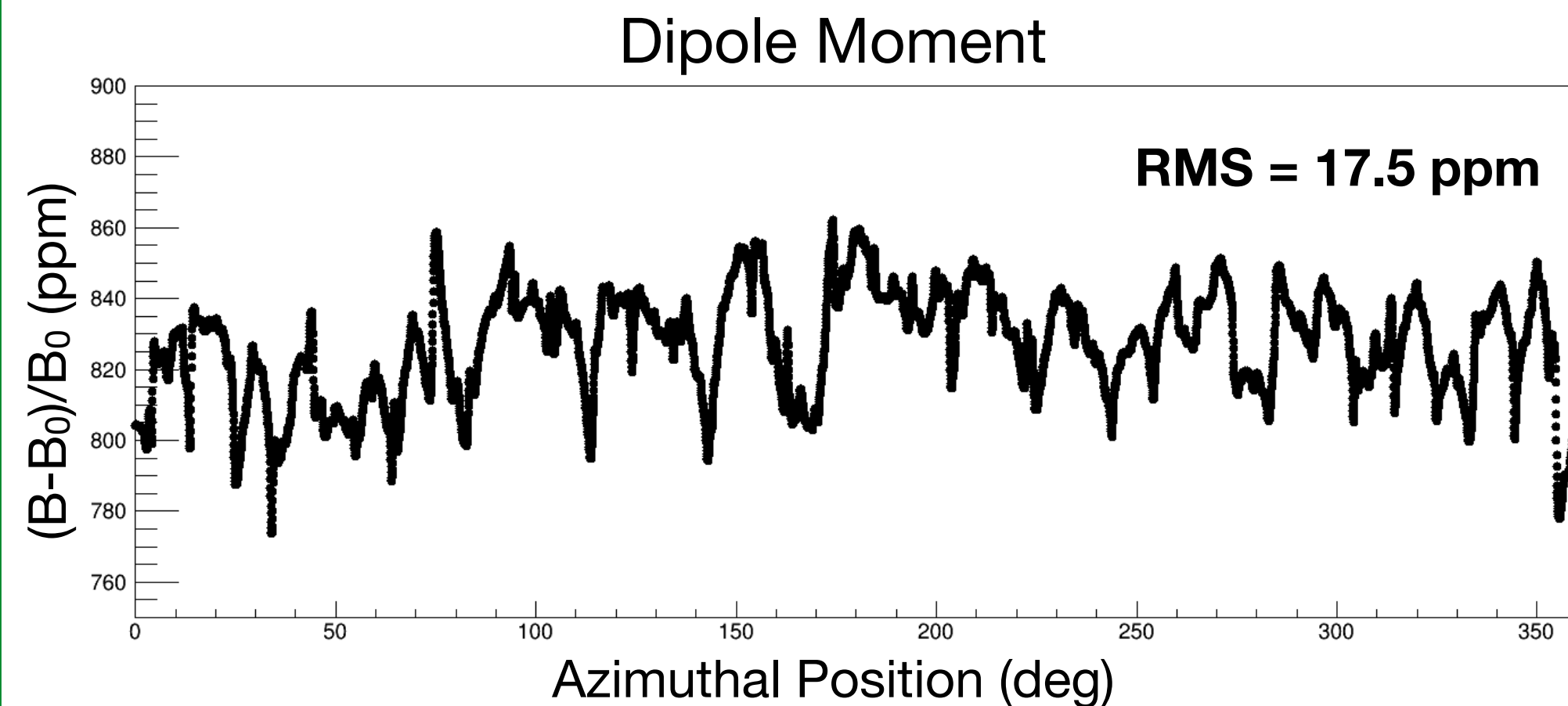
Run 2 Overview

- Data taking period: March—July 2019
- Contiguous data set
- Accumulated $\sim 1.9 \times$ BNL statistics (before data quality cuts)
- Field uniformity in very good condition

Accumulated statistics



Typical Field Map



Run 3

- To start in mid-November
- Aim to **triple** statistics accumulated to date
- Direct continuation of Run 2

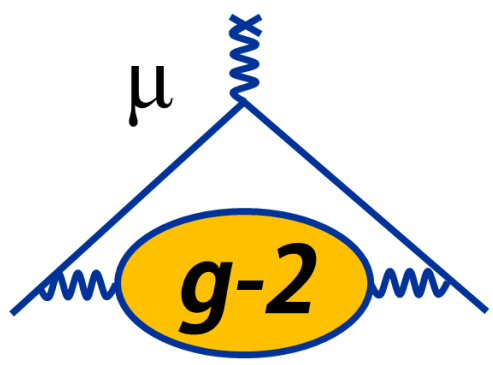
Run 1 Analysis Status — ω_a

$$a_\mu = \frac{\boxed{\omega_a}}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

Run 1 Analysis Status: ω_a

- Account for a number of effects that can affect the extraction of ω_a

$$N(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$$



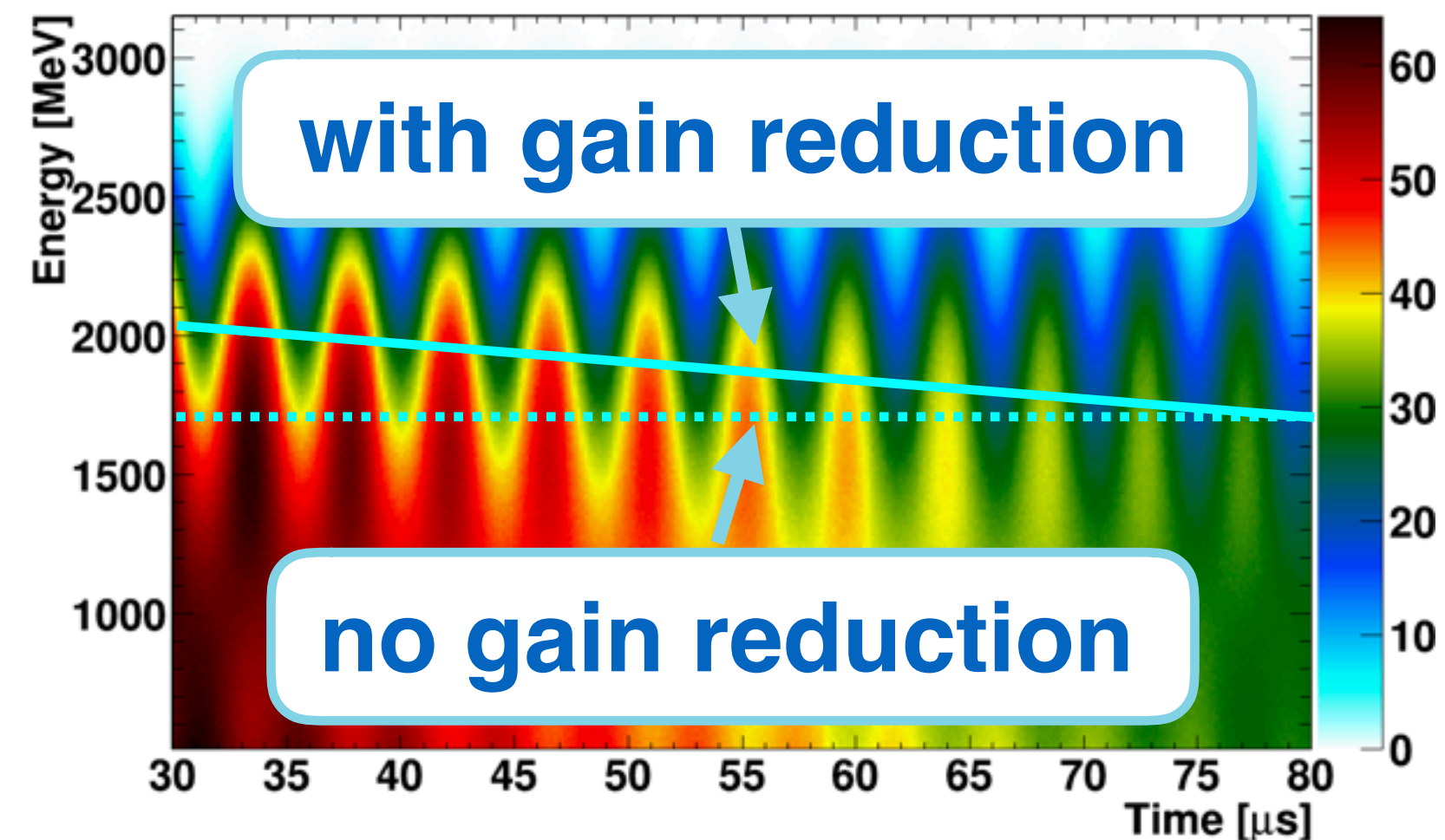
Run 1 Analysis Status: ω_a



- Account for a number of effects that can affect the extraction of ω_a

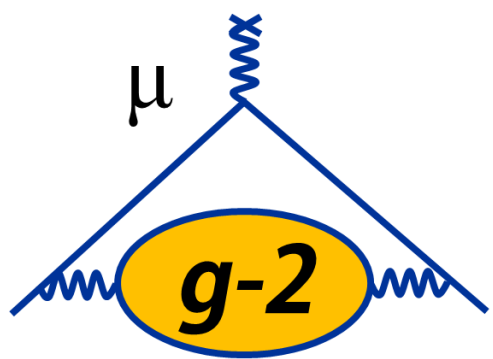
$$N(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$$

Detector effects



- Gain changes over time in calorimeters affects phase of signal: $N \rightarrow N(t)$, $A \rightarrow A(t)$, $\phi \rightarrow \phi(t)$
- Laser system provides corrections

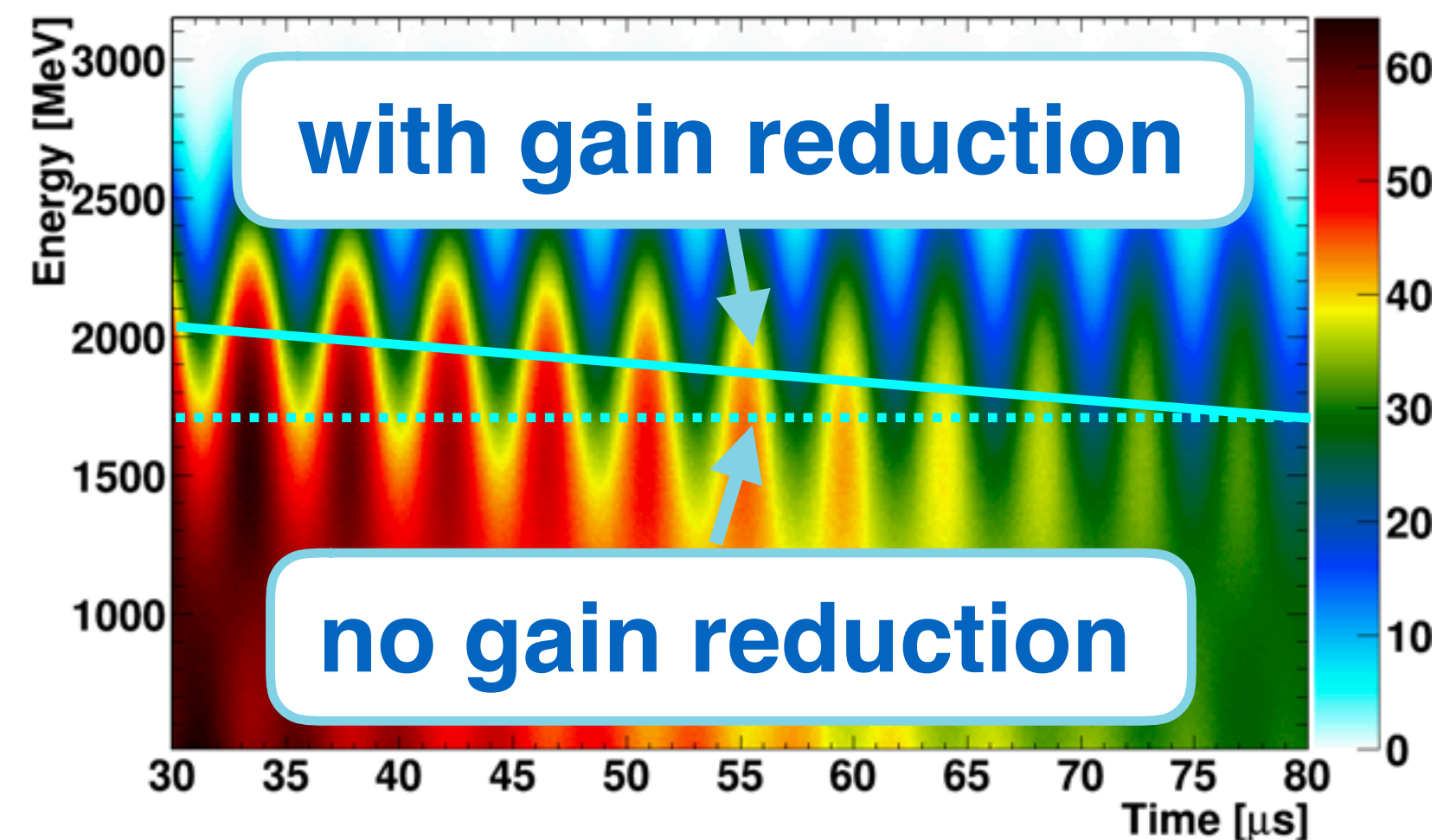
Run 1 Analysis Status: ω_a



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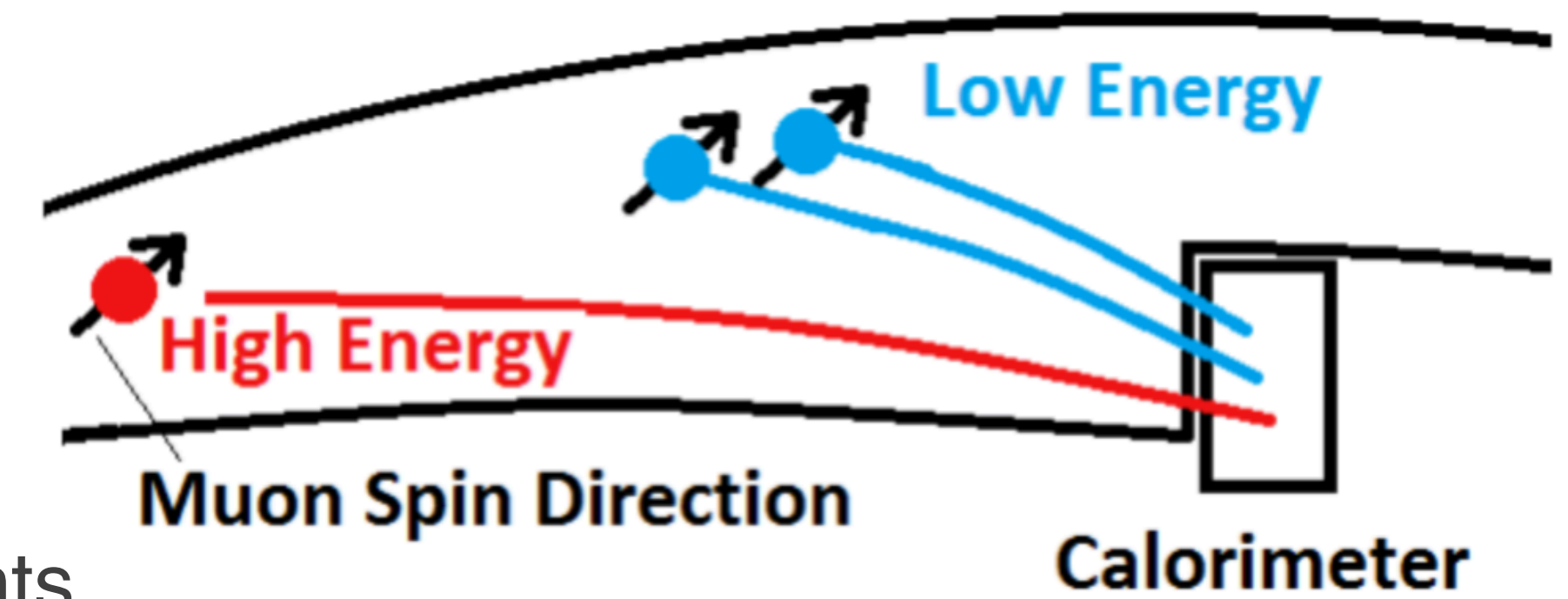
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Detector effects



- Gain changes over time in calorimeters affects phase of signal: $N \rightarrow N(t)$, $A \rightarrow A(t)$, $\phi \rightarrow \phi(t)$
- Laser system provides corrections

Event pileup



- Low-energy events can mimic high-energy events in calorimeter
- Spin precession phase varies with energy — apparent high-energy decay carries phase of low-energy decays

Run 1 Analysis Status: ω_a

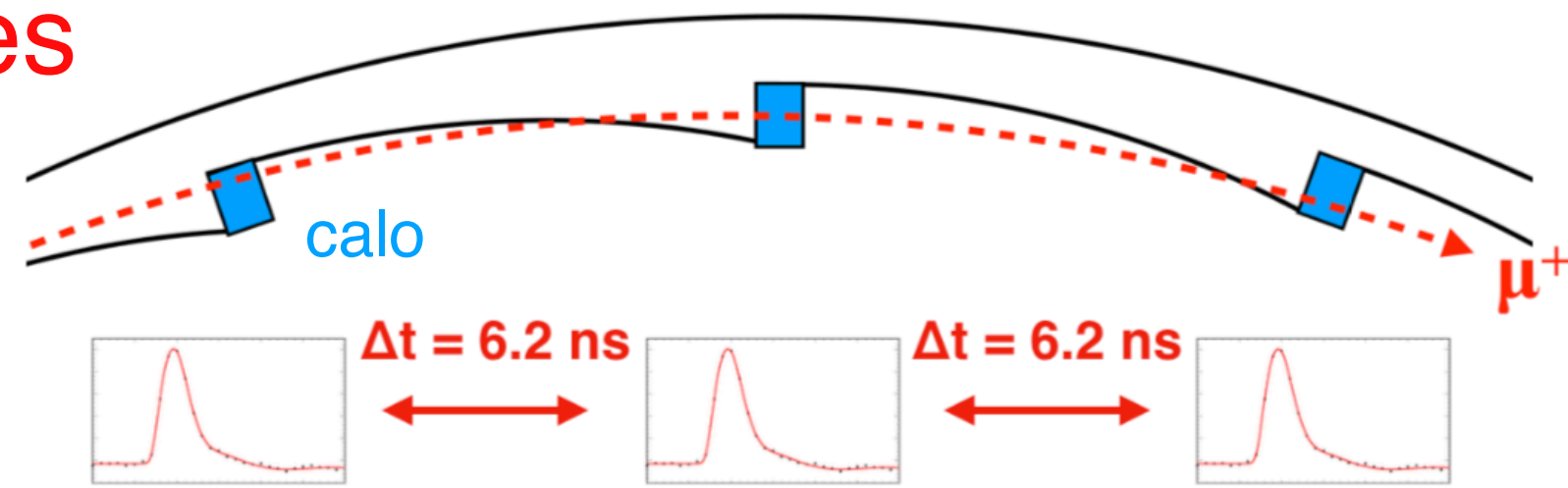


- Account for a number of effects that can affect the extraction of ω_a

$$N(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$$

Beam dynamics

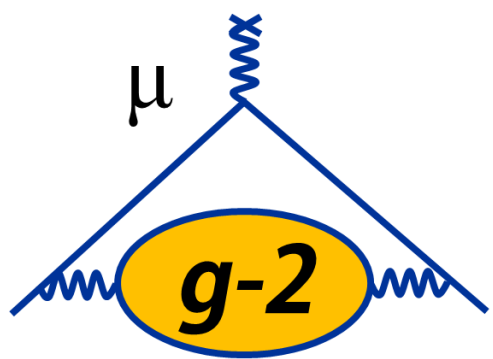
Muon losses



- Muons can leave storage ring by decaying or escaping
- Exhibit specific signature in multiple calorimeters
- Amplitude N_0 scaled by:

$$\Lambda(t) = 1 - K_{\text{loss}} \int_0^t e^{t'/\tau} L(t') dt'$$

Run 1 Analysis Status: ω_a

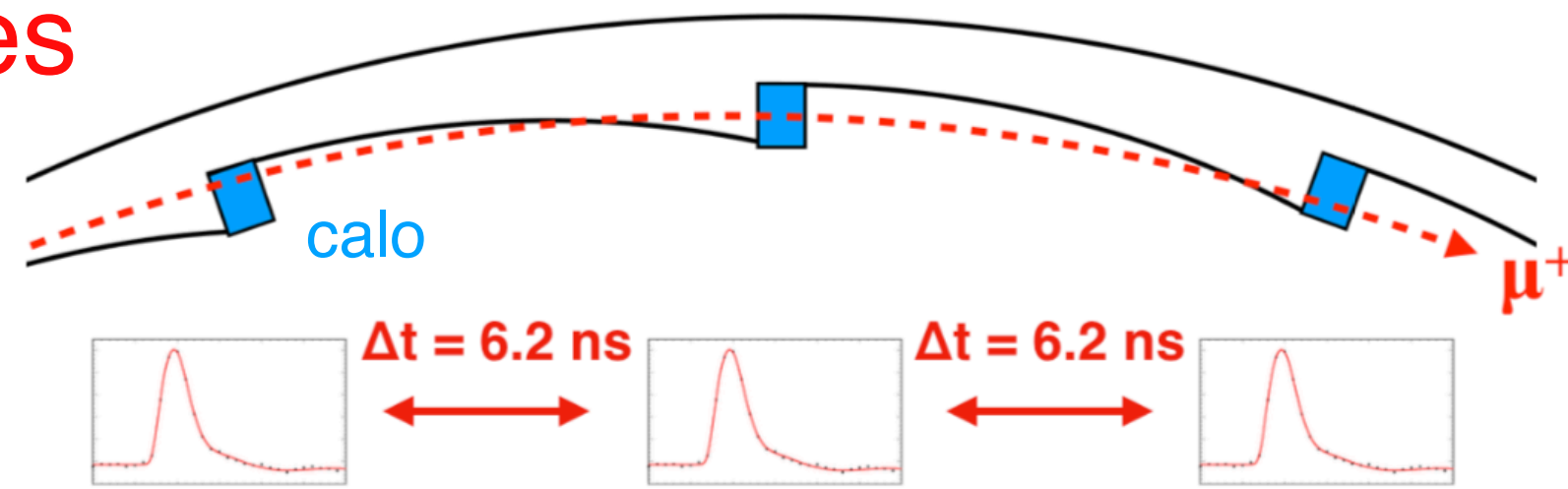


- Account for a number of effects that can affect the extraction of ω_a

$$N(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$$

Beam dynamics

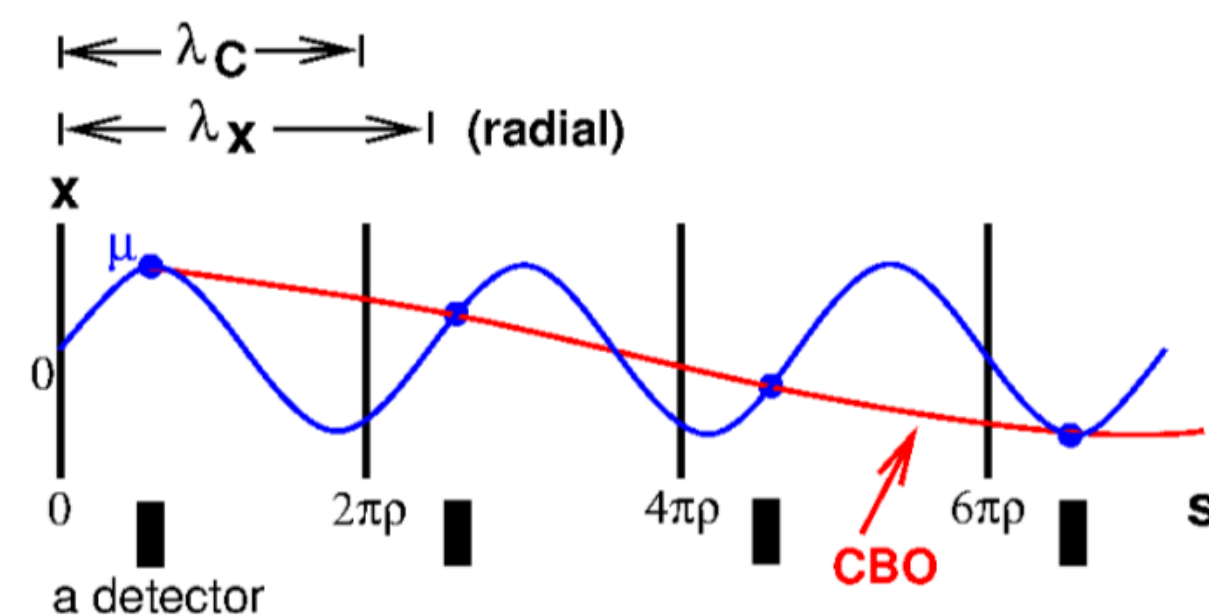
Muon losses



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- Exhibit specific signature in multiple calorimeters
- Amplitude N_0 scaled by:

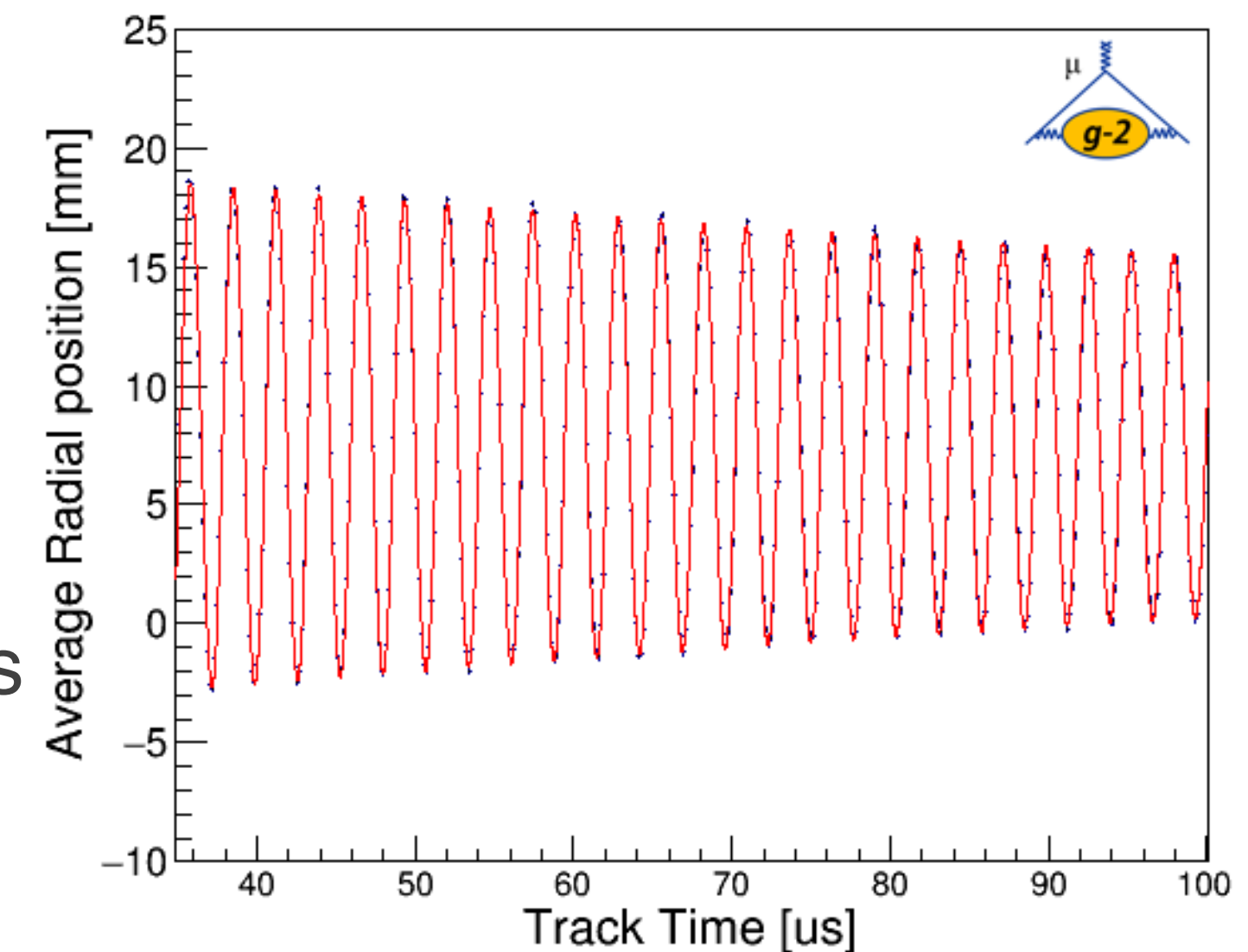
$$\Lambda(t) = 1 - K_{\text{loss}} \int_0^t e^{t'/\tau} L(t') dt'$$

Coherent betatron oscillations (CBO)

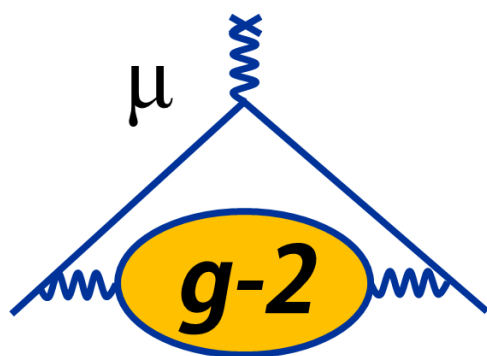


- Acceptance of calorimeters affected by coherent radial beam motion
- Amplitude N_0 scaled by:

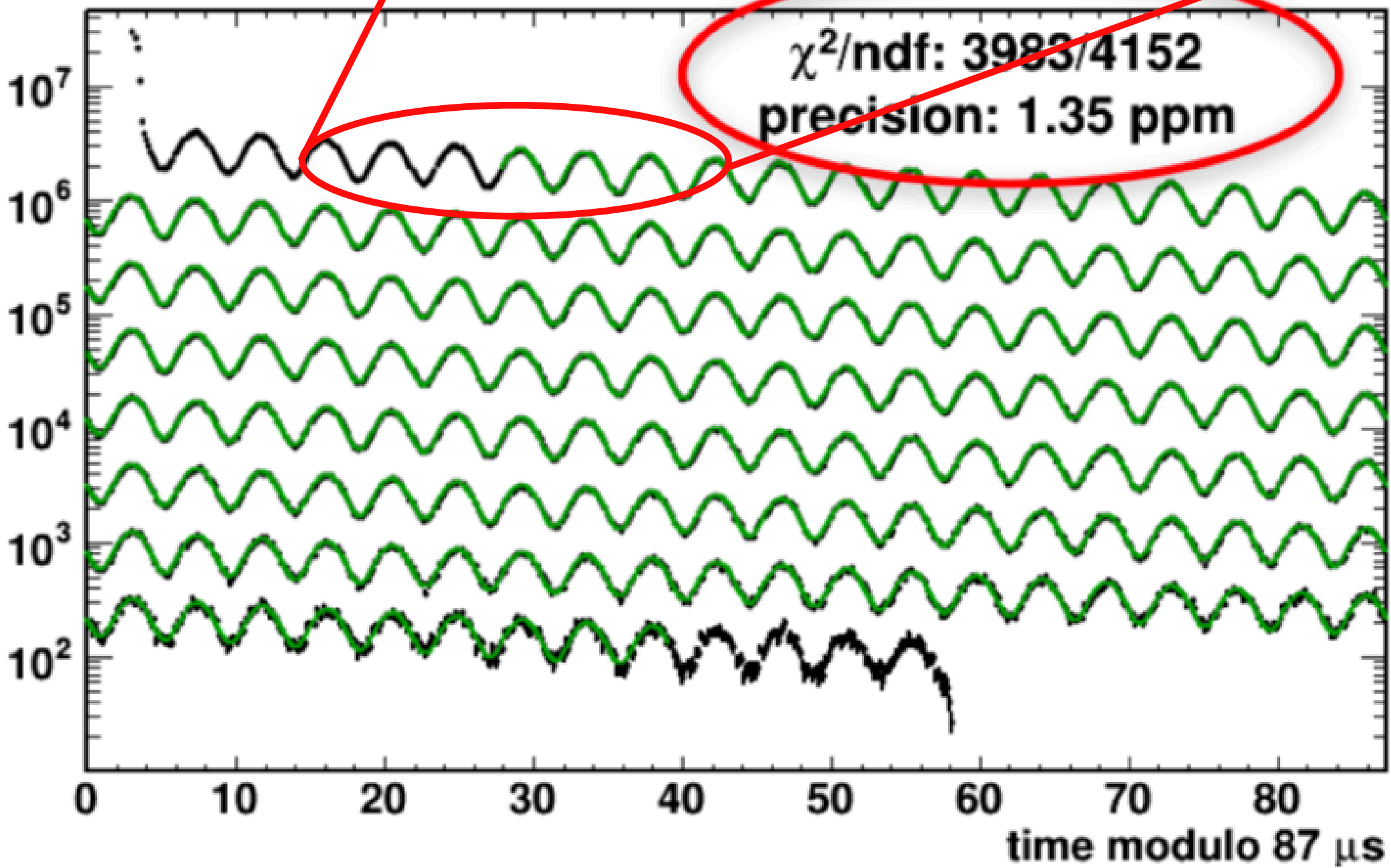
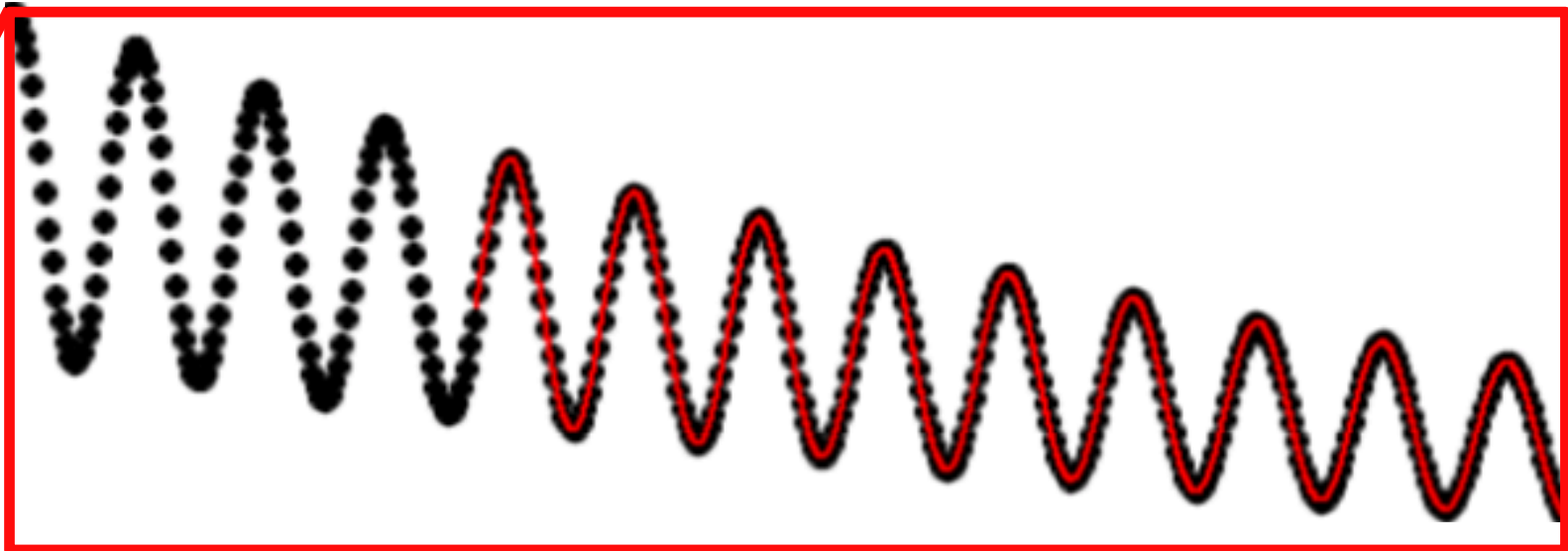
$$C(t) = 1 - e^{-t/\tau_{\text{CBO}}} A_1 \cos(\omega_{\text{CBO}} t + \phi_1)$$



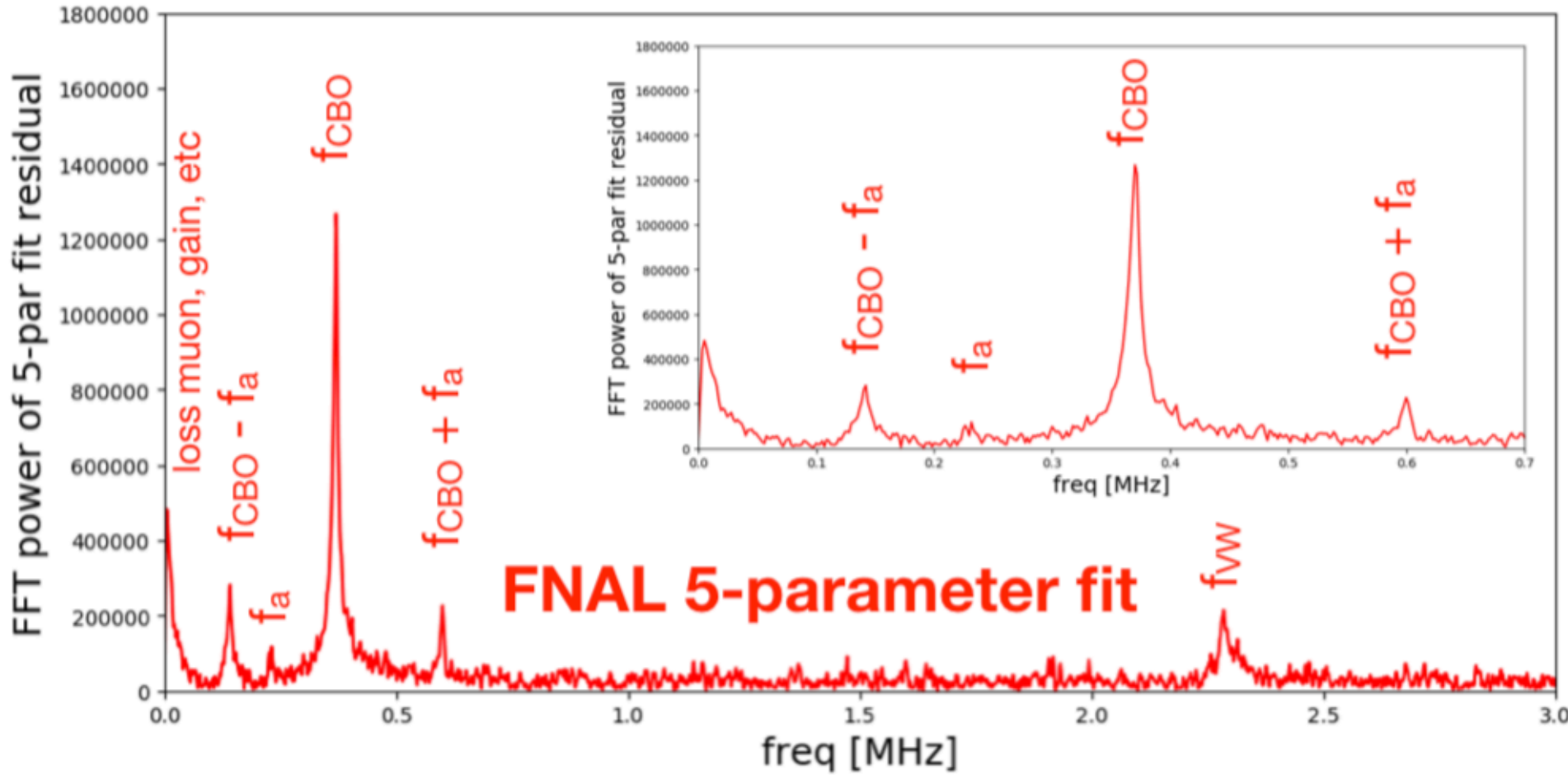
Run 1 Analysis Status: ω_a



Simple five-parameter fit

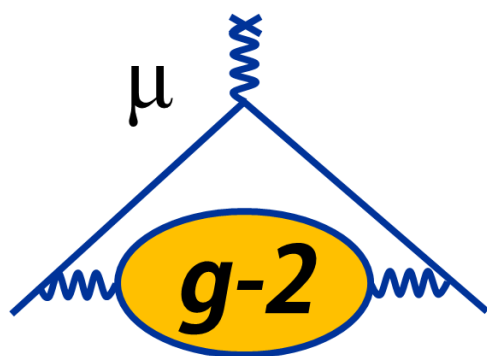


FFT of fit residuals



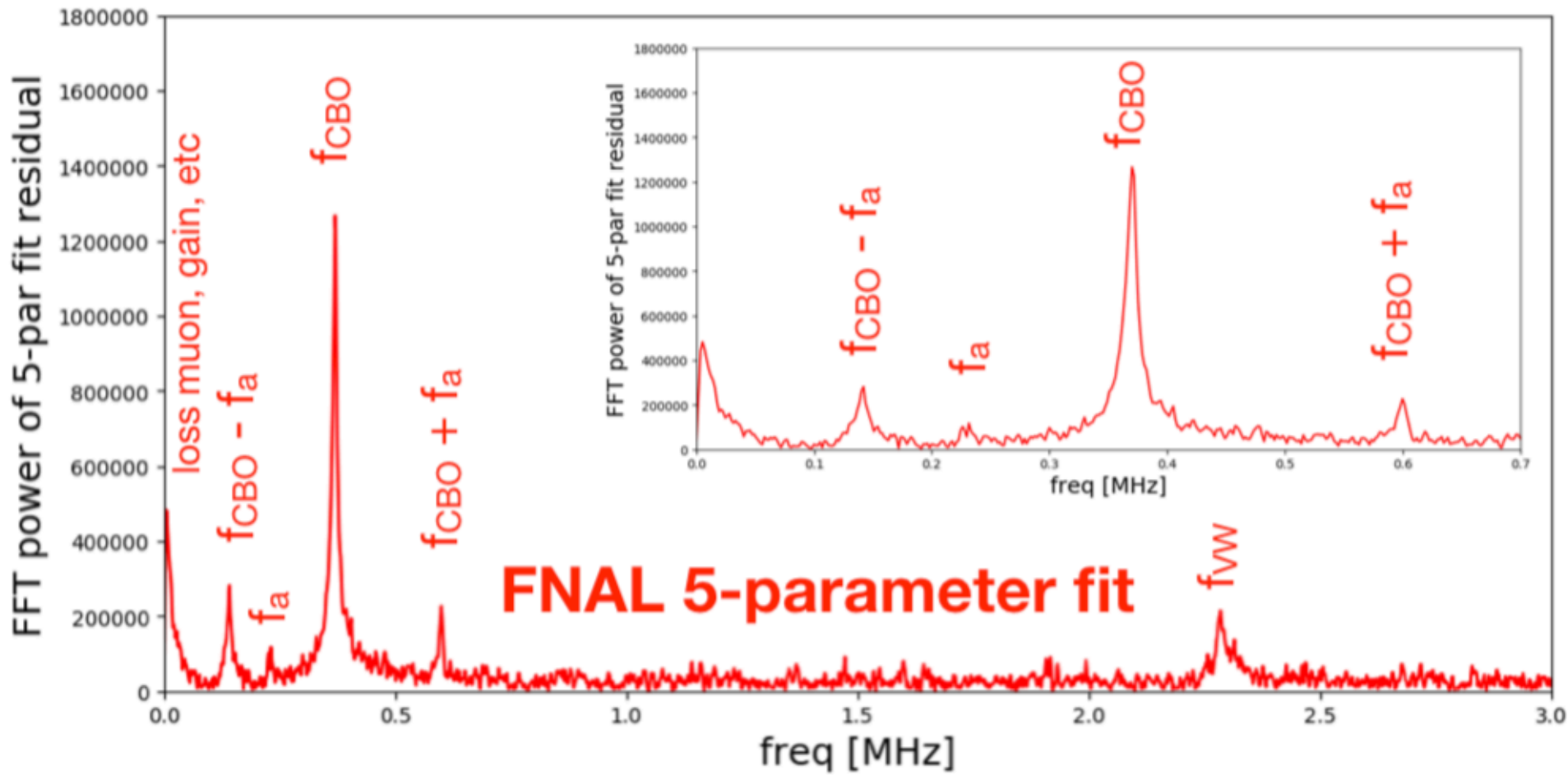
FNAL 5-parameter fit

Run 1 Analysis Status: ω_a



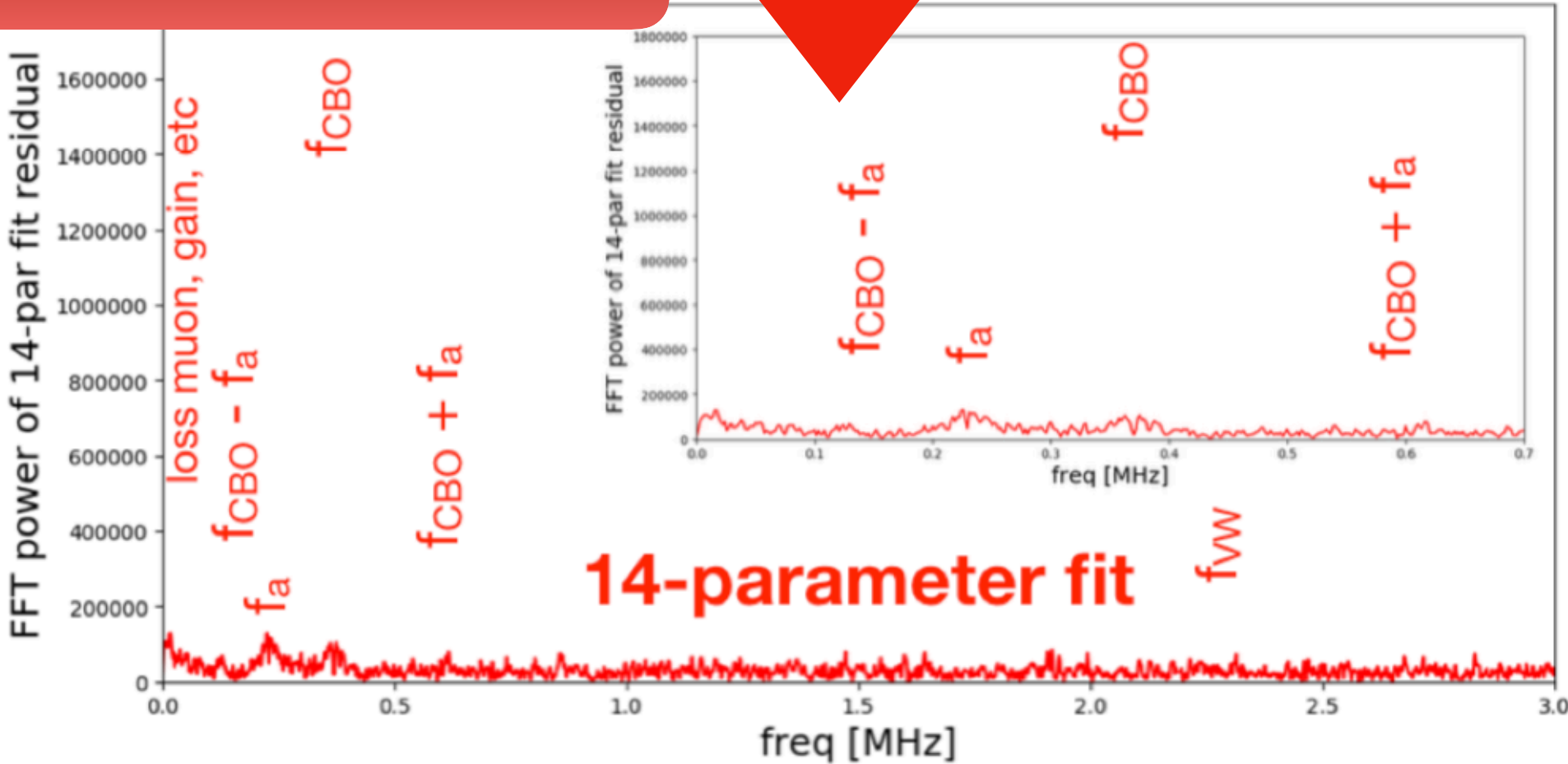
Simple five-parameter fit

FFT of fit residuals

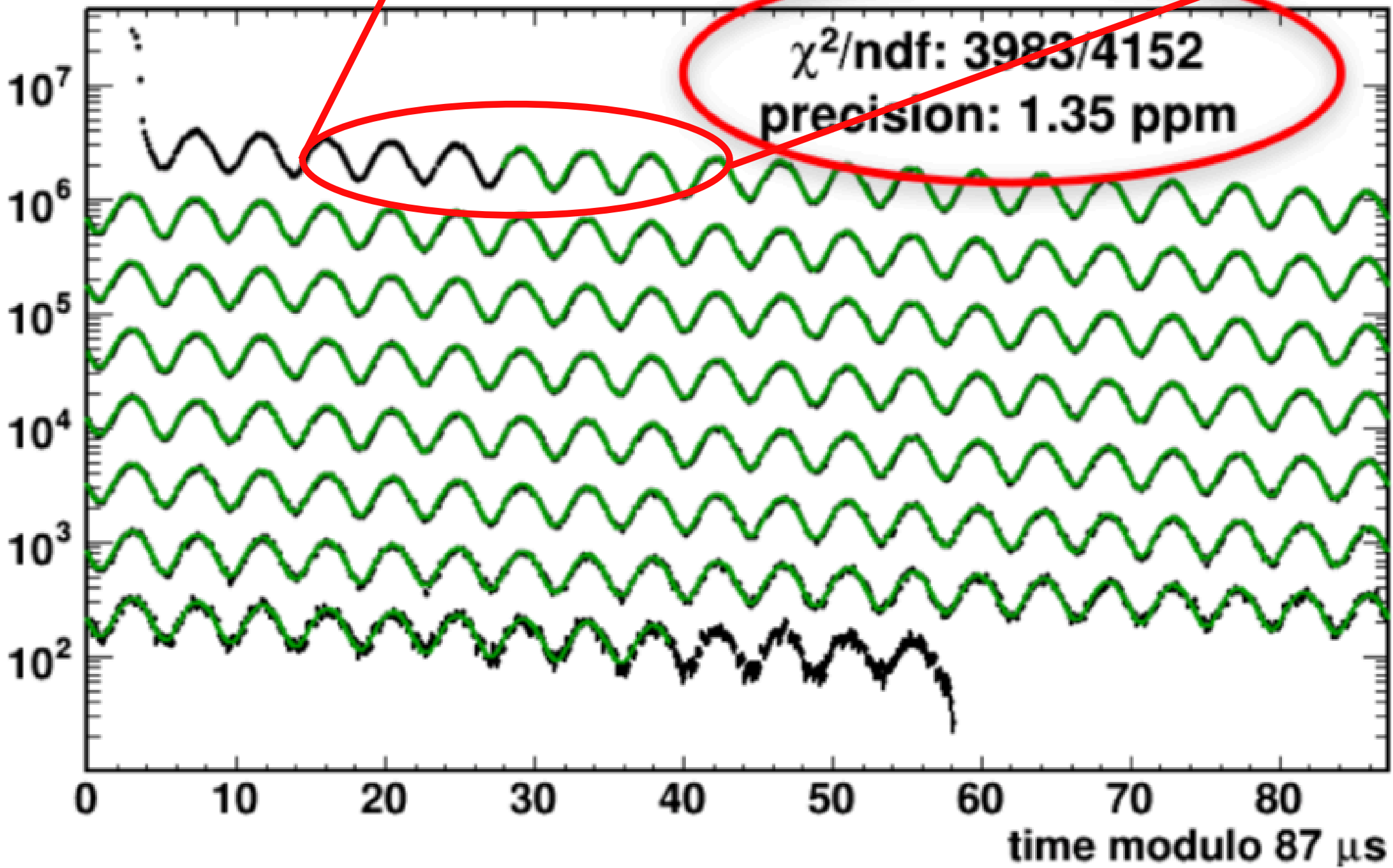


FNAL 5-parameter fit

Big improvements when accounting for CBO, lost muons,...

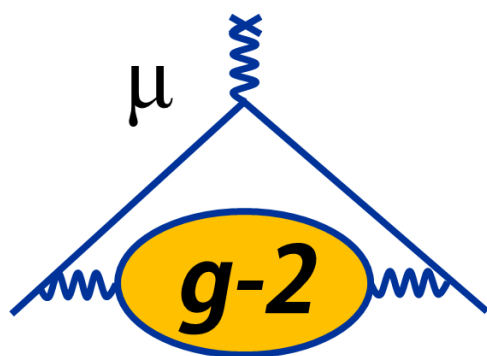


14-parameter fit



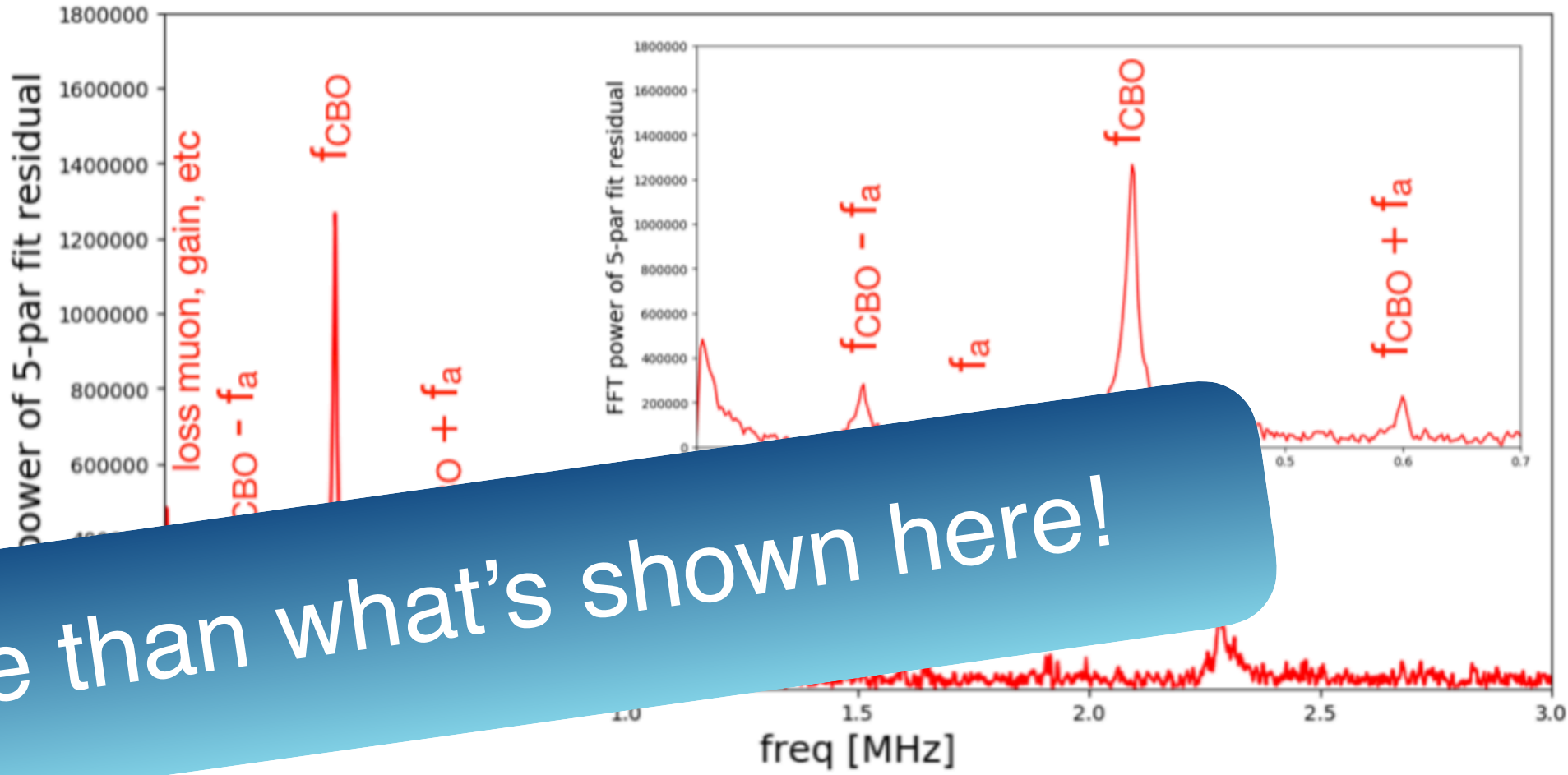
χ^2/ndf : 3983/4152
precision: 1.35 ppm

Run 1 Analysis Status: ω_a



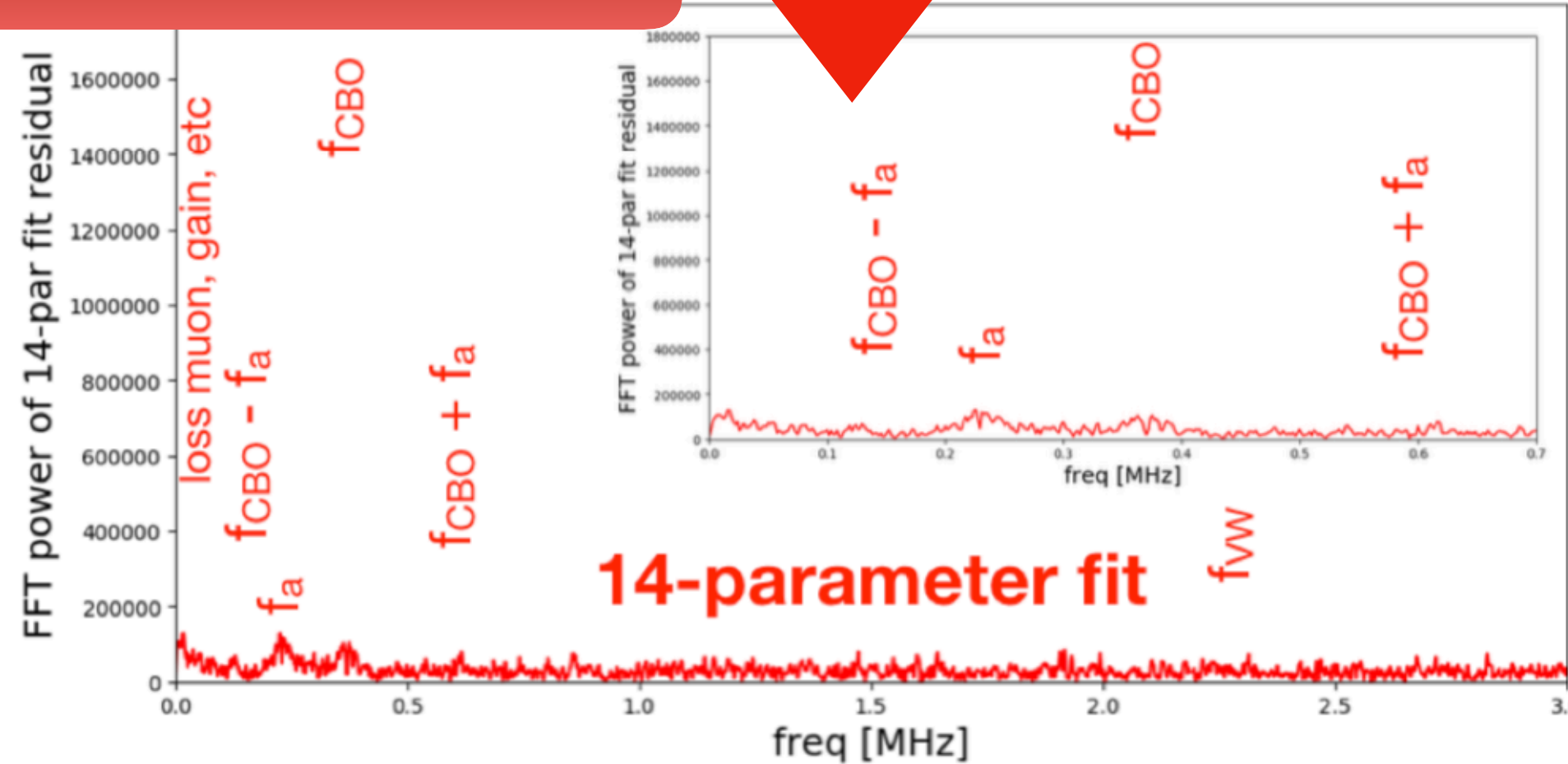
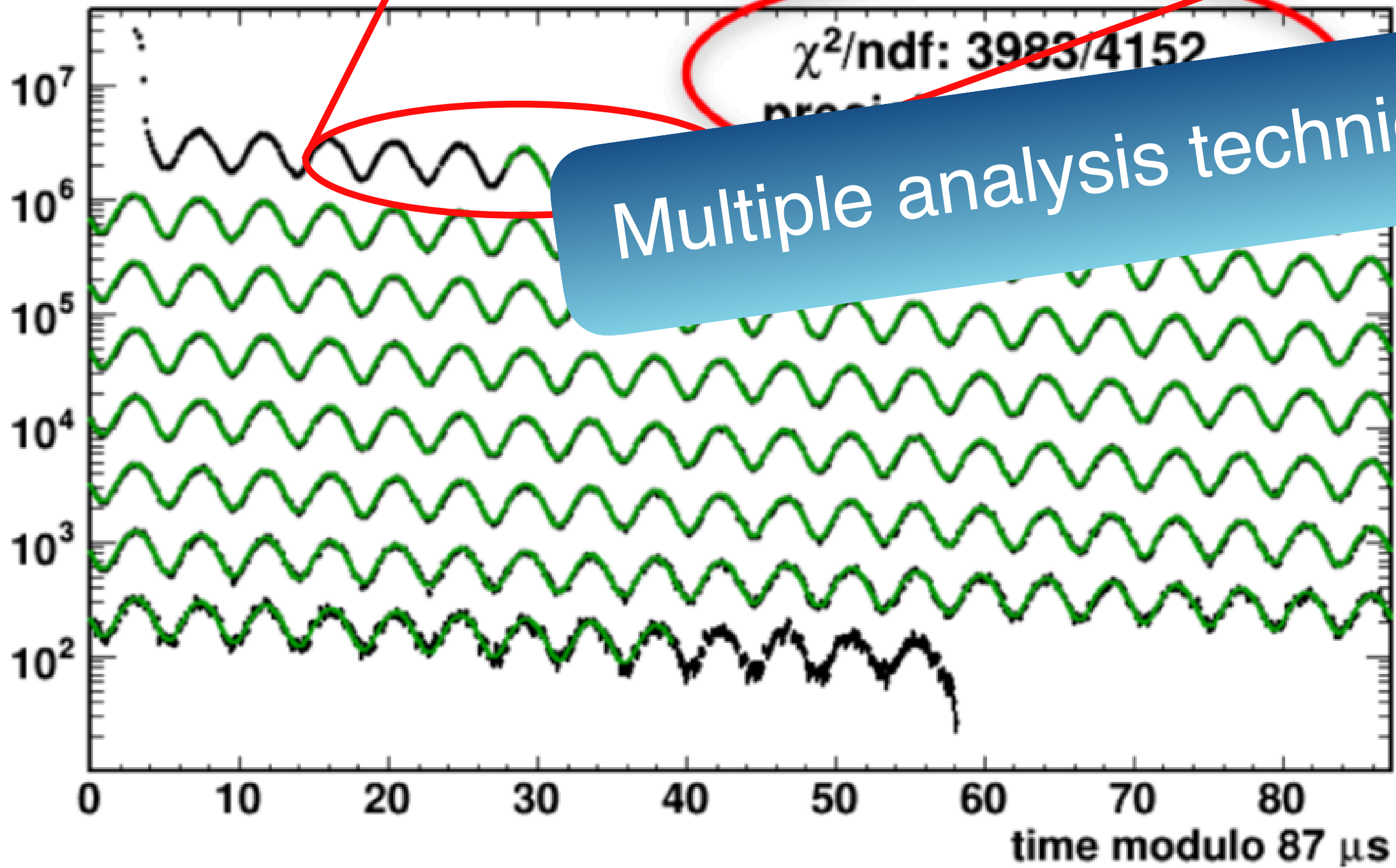
Simple five-parameter fit

FFT of fit residuals



Multiple analysis techniques — more than what's shown here!

Big improvements when accounting for CBO, lost muons,...

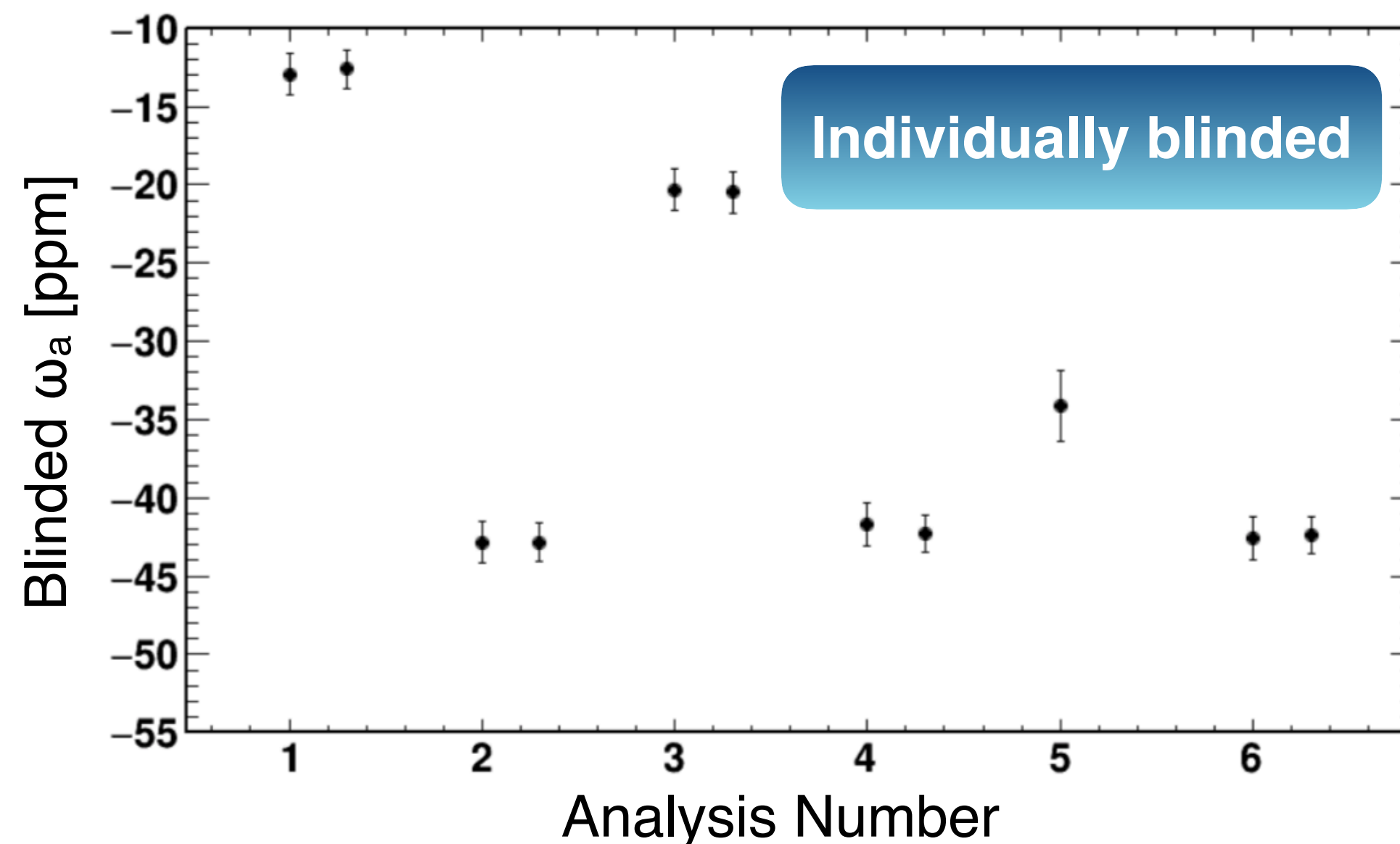


14-parameter fit

Run 1 Analysis Status: Relative Unblinding for ω_a



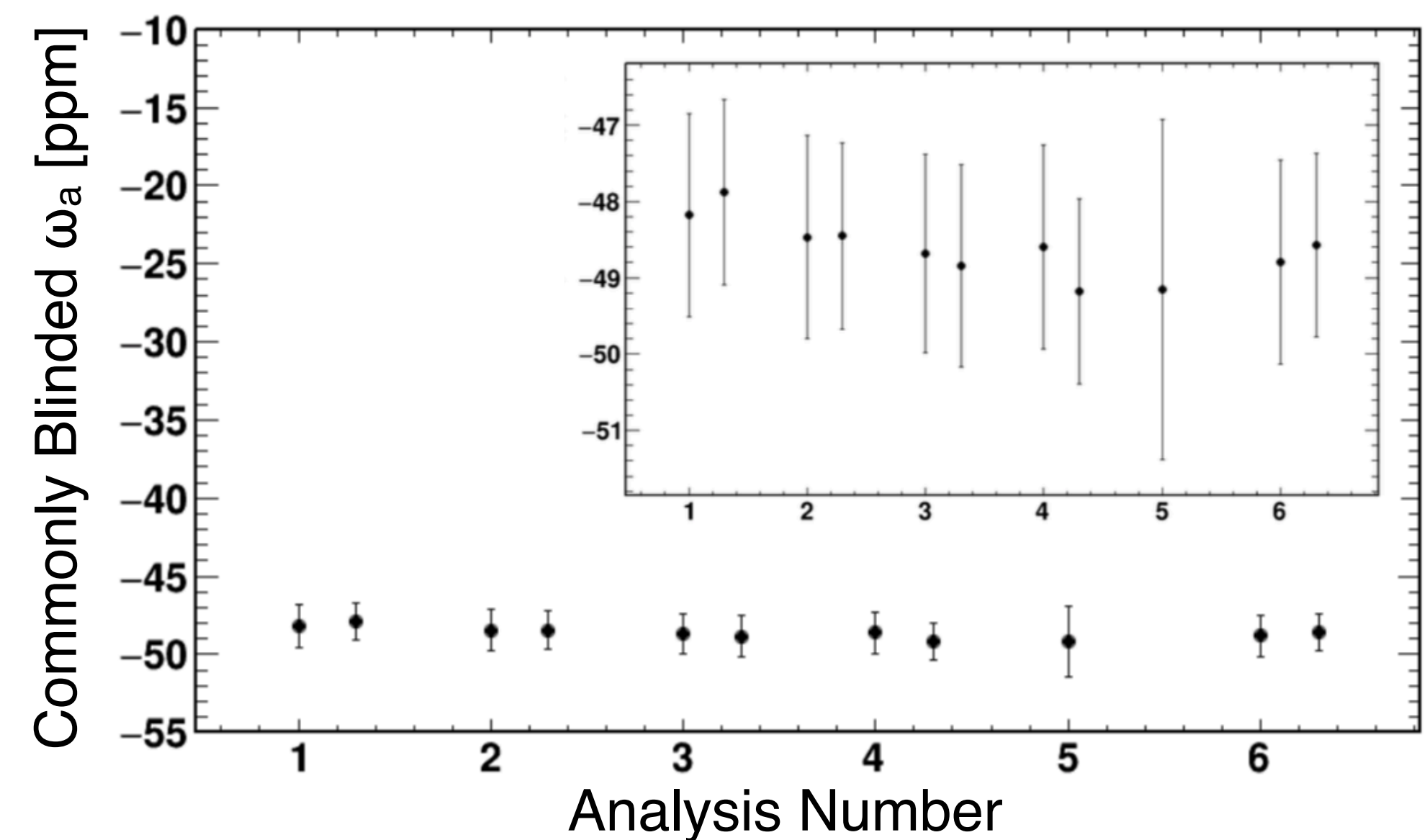
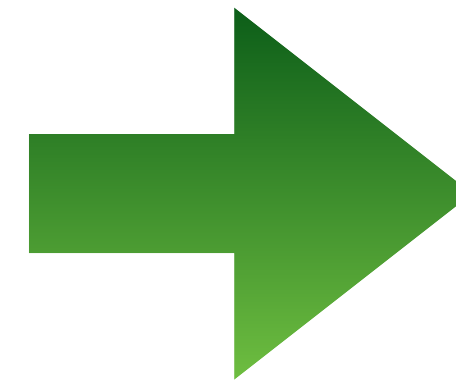
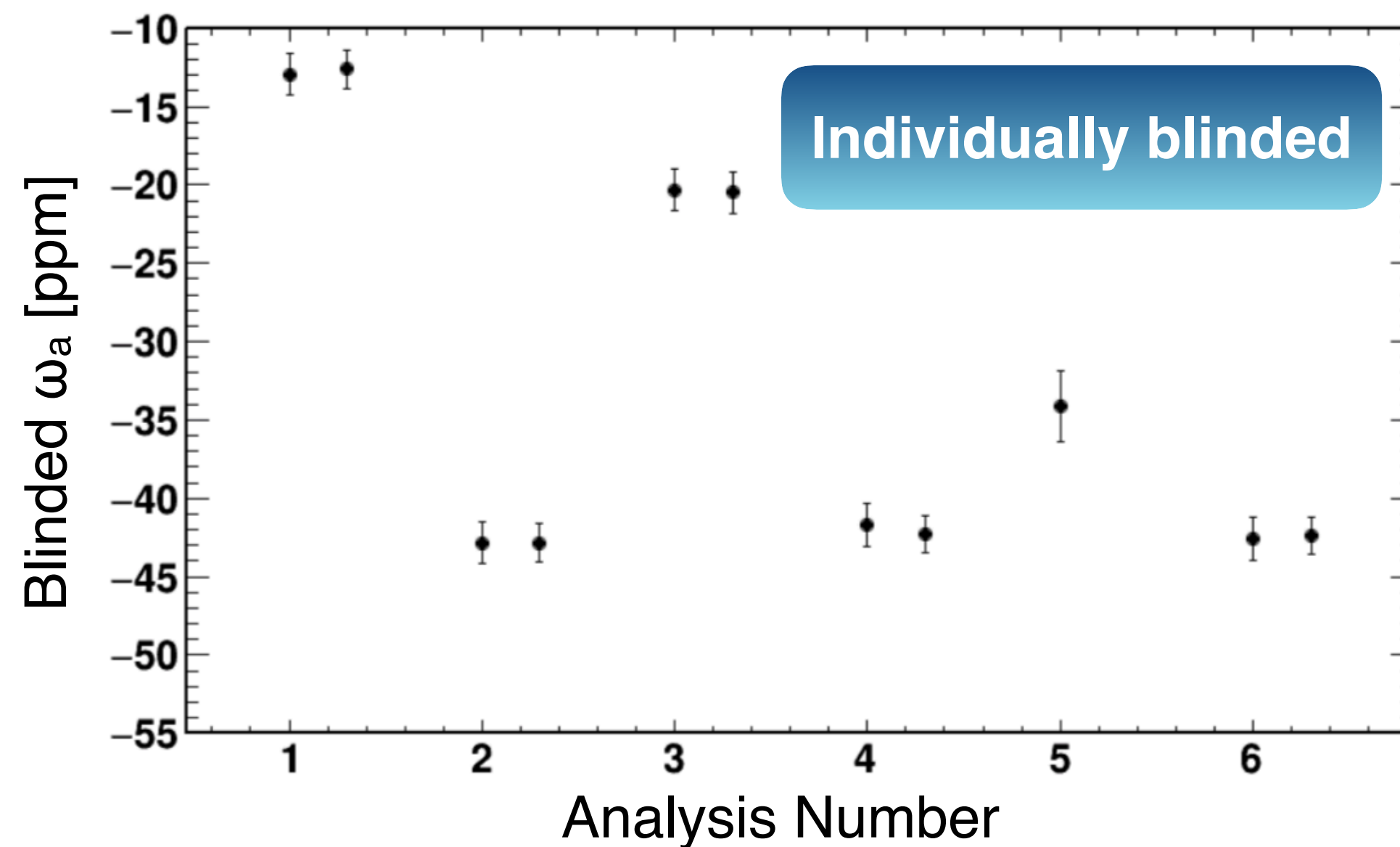
- Doubly-blinded in ω_a measurement: Clock tuned to $40 \text{ MHz} \pm 25 \text{ ppm}$
- Analyzers' results come with random frequency offset $\omega_a \rightarrow \omega_a \pm 25 \text{ ppm}$
- Recently compared results on **subset of data** at a **common blinded value**



Run 1 Analysis Status: Relative Unblinding for ω_a

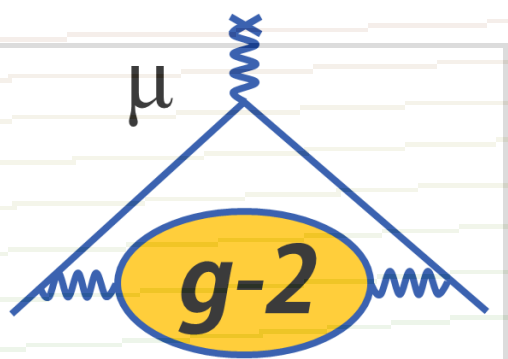


- Doubly-blinded in ω_a measurement: Clock tuned to $40 \text{ MHz} \pm 25 \text{ ppm}$
- Analyzers' results come with random frequency offset $\omega_a \rightarrow \omega_a \pm 25 \text{ ppm}$
- Recently compared results on **subset of data** at a **common blinded value**



- Consistent results at **common blinded value** builds confidence in our analyses

Paramagnetism

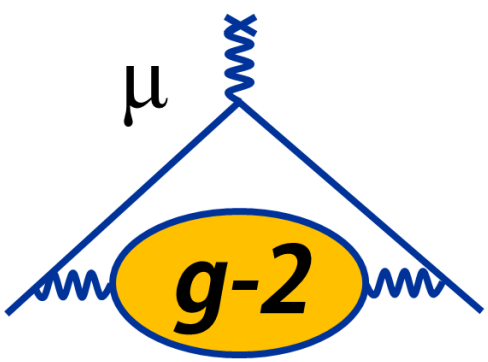
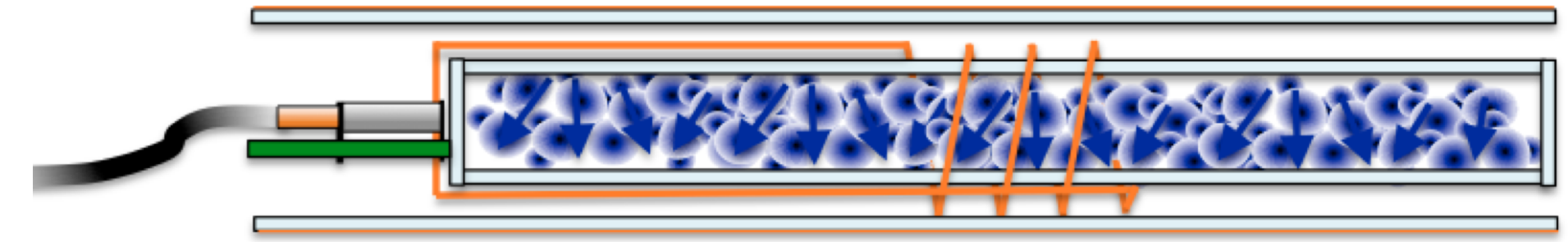
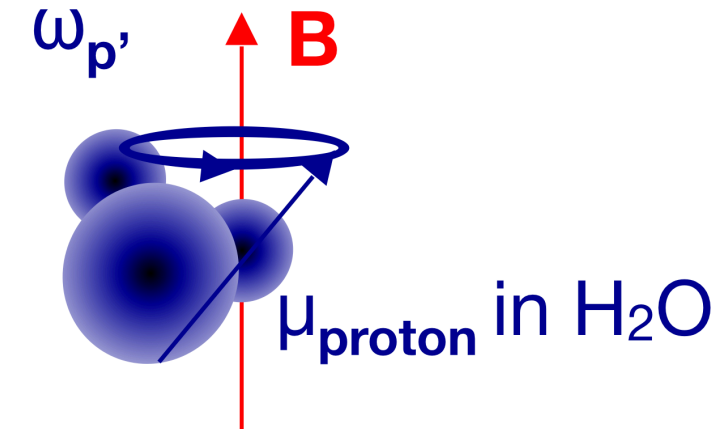
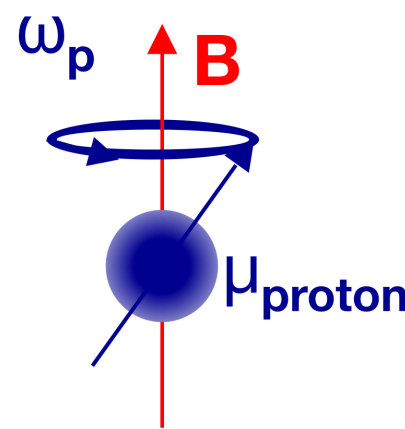


Run 1 Analysis Status — ω_p

$$a_{\mu} = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_{\mu}}{m_e} \frac{g_e}{2}$$

Run 1 Analysis Status: ω_p — Field Calibration

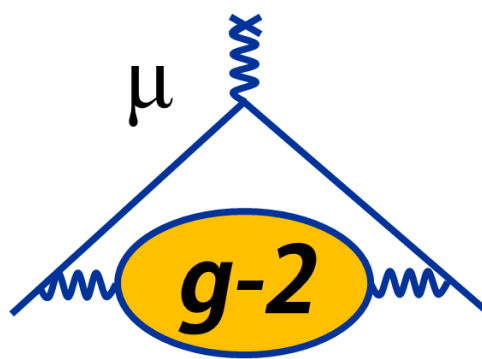
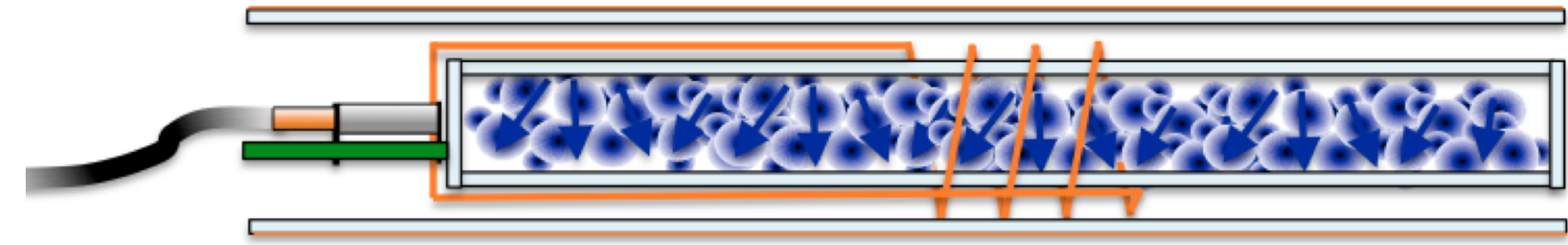
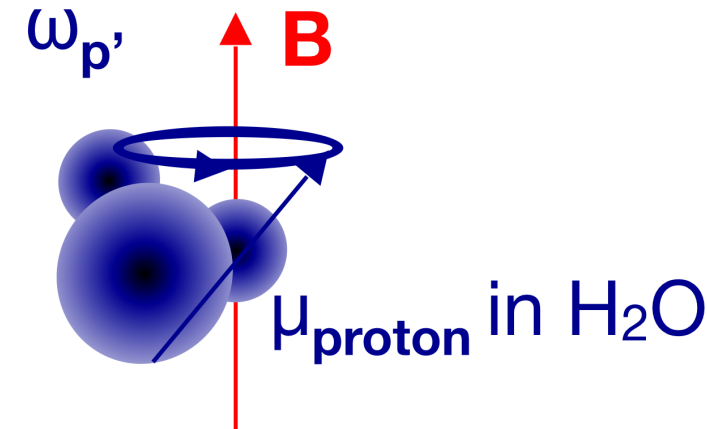
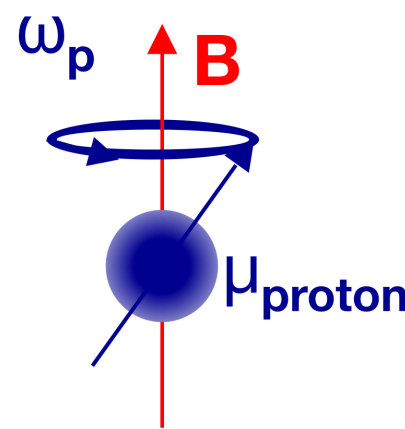
- In the experiment, need to extract ω_p ; however, don't have free protons
 - Need a calibration
- Field at the proton differs from the applied field



$$\omega_p^{\text{meas}} \approx \omega_p^{\text{free}}$$

Run 1 Analysis Status: ω_p — Field Calibration

- In the experiment, need to extract ω_p ; however, don't have free protons
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$$\omega_p^{\text{meas}} = \omega_p^{\text{free}} \left[1 - \sigma(\text{H}_2\text{O}, T) \right]$$

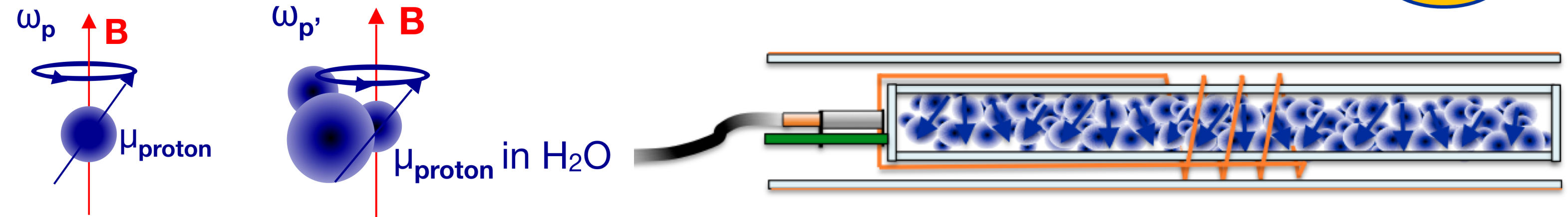
Protons in H_2O molecules, diamagnetism of electrons screens protons \Rightarrow local \mathbf{B} changes

- $\sigma = 25\,691(11) \times 10^{-9}$ at 25 deg C [P.J. Mohr et al, Rev. Mod. Phys. **84**, 1527 (2012)]

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$$\omega_p^{\text{meas}} = \omega_p^{\text{free}} \left[1 - \sigma(\text{H}_2\text{O}, T) - \left(\frac{\epsilon}{4\pi} - \frac{1}{3} \right) \chi(\text{H}_2\text{O}, T) \right]$$

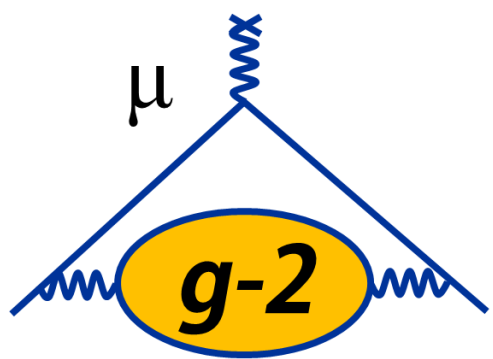
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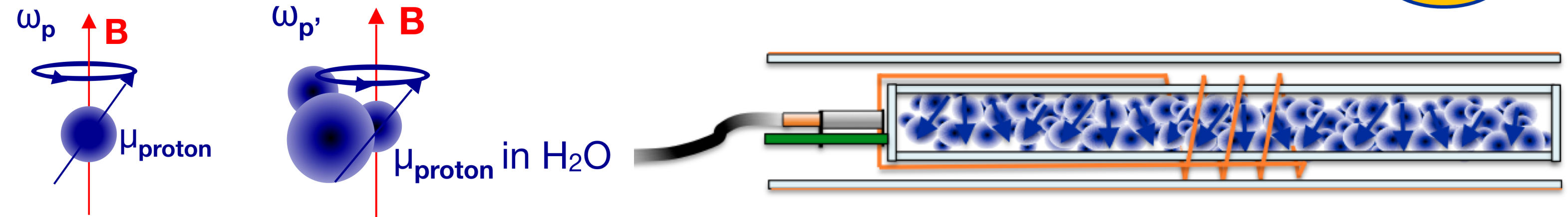
Magnetic susceptibility of water gives shape-dependent perturbation

- $\epsilon = 4\pi/3$ (perfect sphere)
- $\epsilon = 2\pi$ (infinite cylinder) when probe is perpendicular to B
- $\chi_{\text{H}_2\text{O}}(T = 20^\circ\text{C}) = -9049(9) \times 10^{-9}$ [world average]

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Protons in H_2O molecules, diamagnetism of electrons screens protons \Rightarrow local B changes

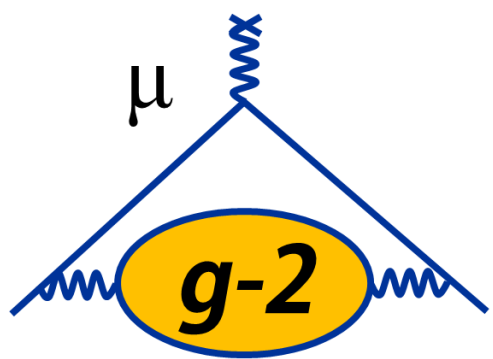
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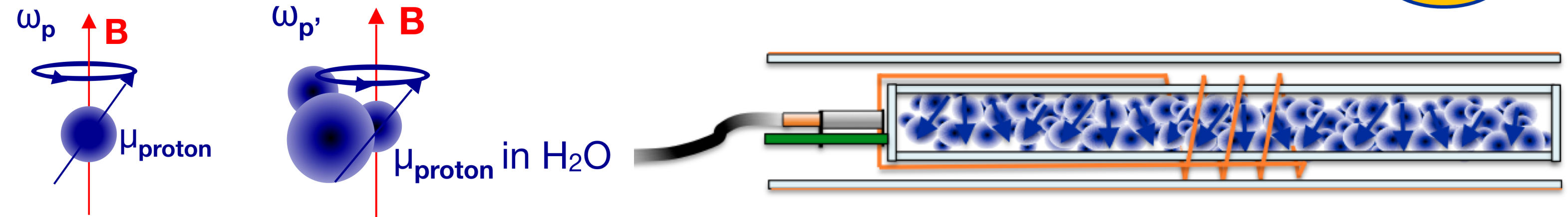
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paramagnetic impurities in water sample

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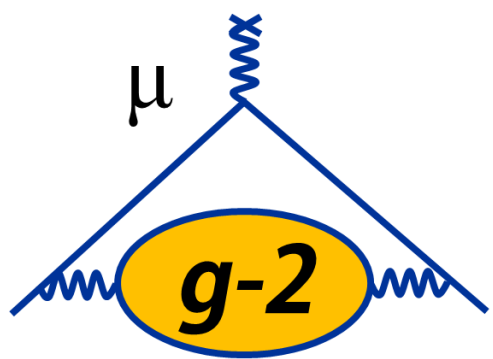
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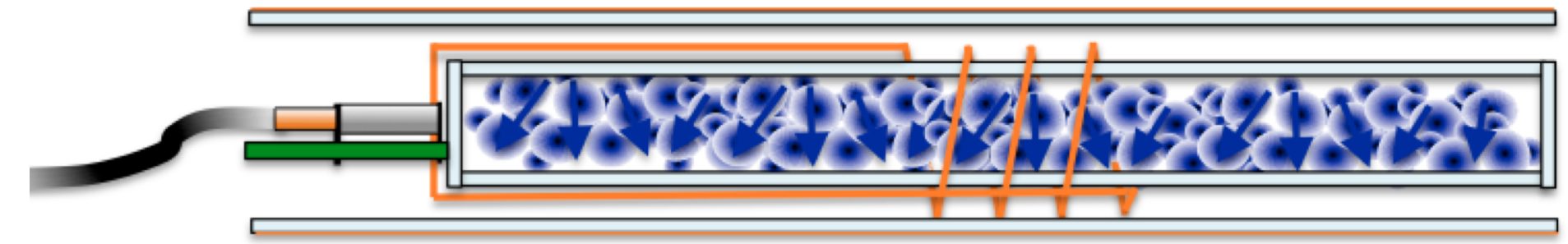
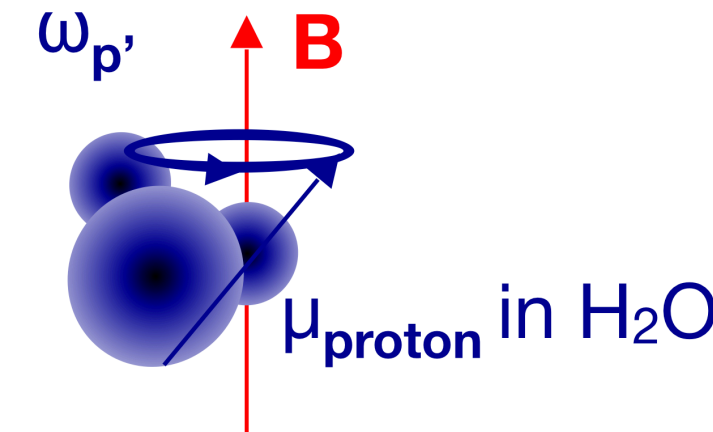
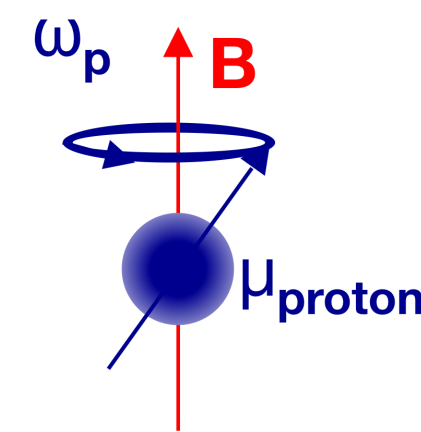
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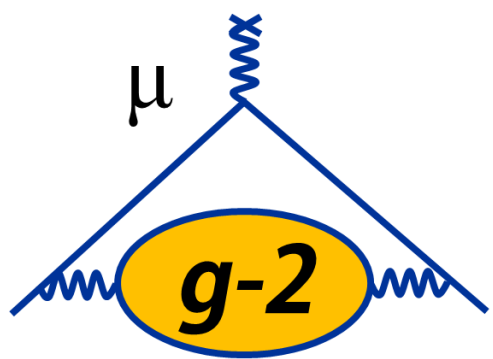
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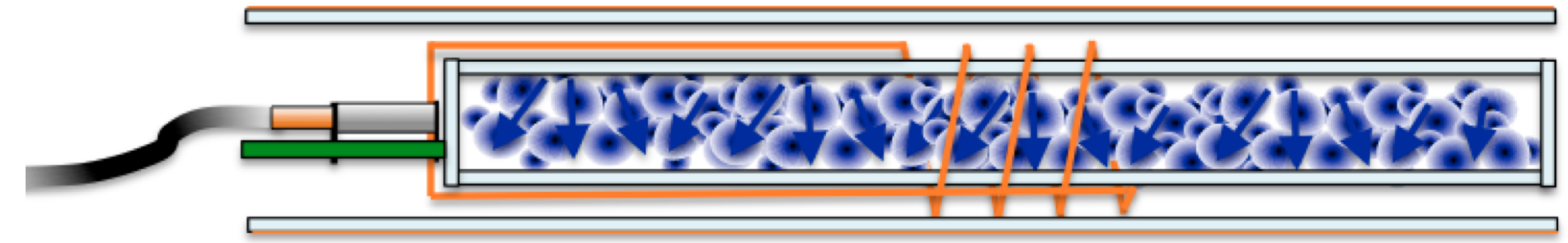
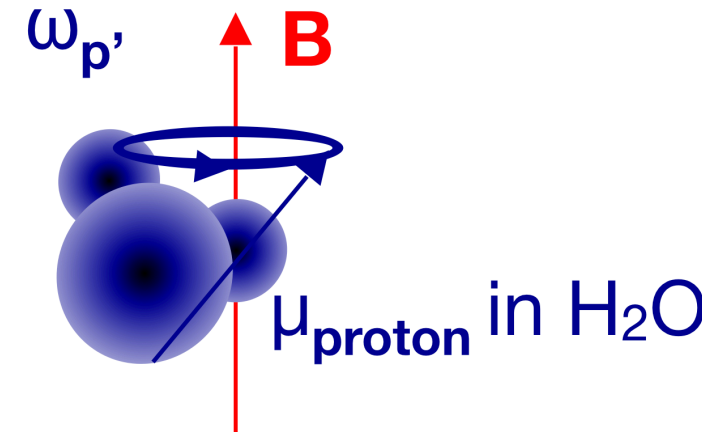
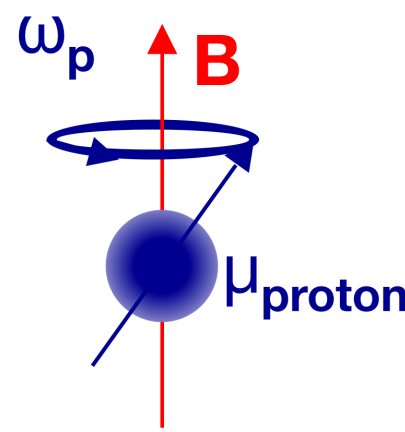
Magnetization of probe materials, geometry perturbs field experienced by protons

Dynamic effects: radiation damping, dipolar field from protons

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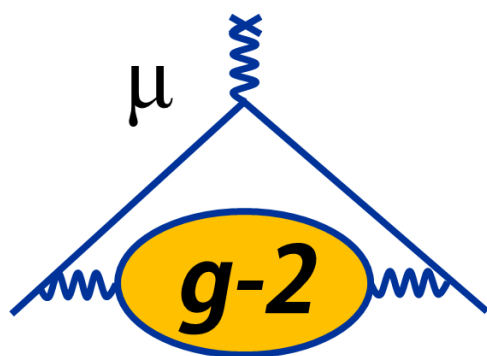
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Dynamic effects: radiation damping, dipolar field from protons



Goal: Determine total correction to ≤ 35 ppb accuracy

Run 1 Analysis Status: ω_p — Field Calibration

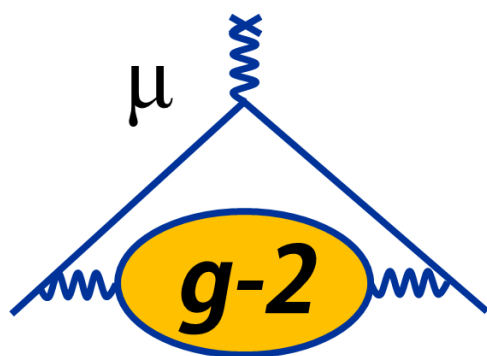


Plunging Probe

- Achieved **small perturbation of plunging probe** ($\delta_s + \delta_p + \delta_{RD} + \delta_d$): **(-0.2 ± 11.4) ppb**
- Quantified uncertainties on plunging probe material, dynamic effects — **under budget of 35 ppb by a factor of > 2**

Plunging Probe Perturbations		
Quantity	Symbol	Uncertainty (ppb)
Material Perturbation	δ_s	10.9
Paramagnetic Impurities	δ_p	1.1
Radiation Damping	δ_{RD}	2
Proton Dipolar Fields	δ_d	2.3
Bulk Magnetic Susceptibility	δ_b	2
Water Diamagnetic Shielding	σ	11
TOTAL		15.9

Run 1 Analysis Status: ω_p — Field Calibration



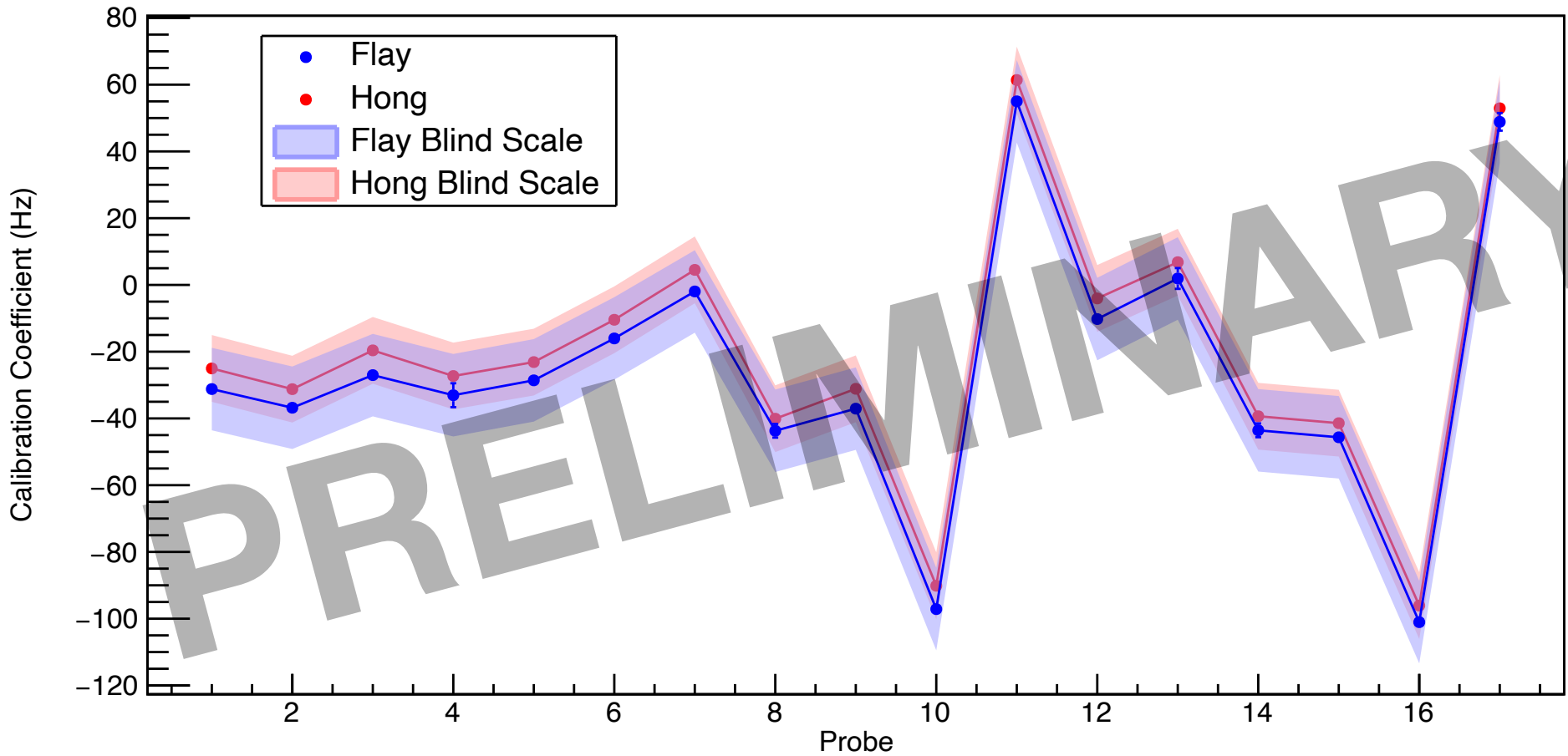
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TOTAL		15.9

Trolley Calibration

- Calibration of trolley probes under control**
- Factor of ≥ 2 improvement on uncertainties for nearly all probes compared to E821
- Uncertainty is $< \sim 40$ ppb on average per probe — **on budget**



Run 1 Analysis Status: ω_p — Field Interpolation



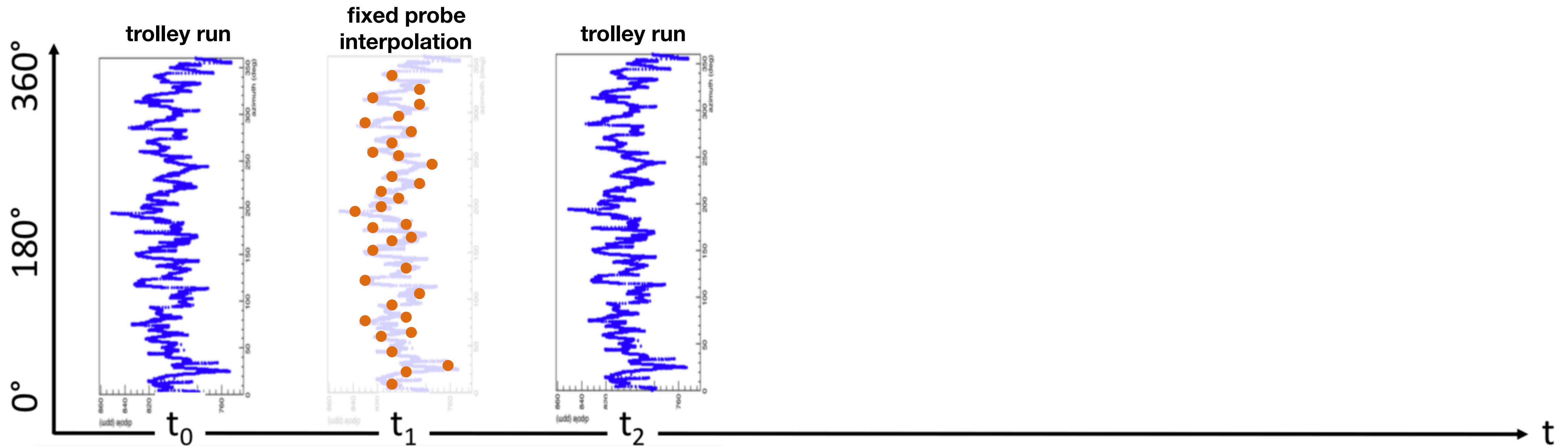
- Need to determine ω_p at all times while storing muons => interpolate between trolley maps using fixed probe data



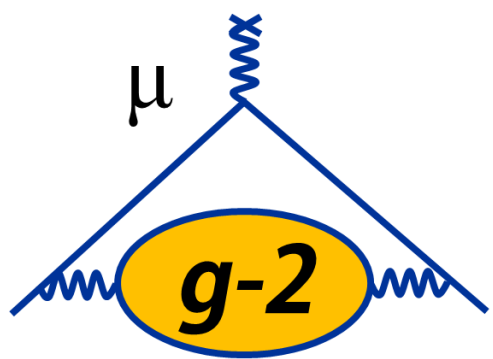
Run 1 Analysis Status: ω_p — Field Interpolation



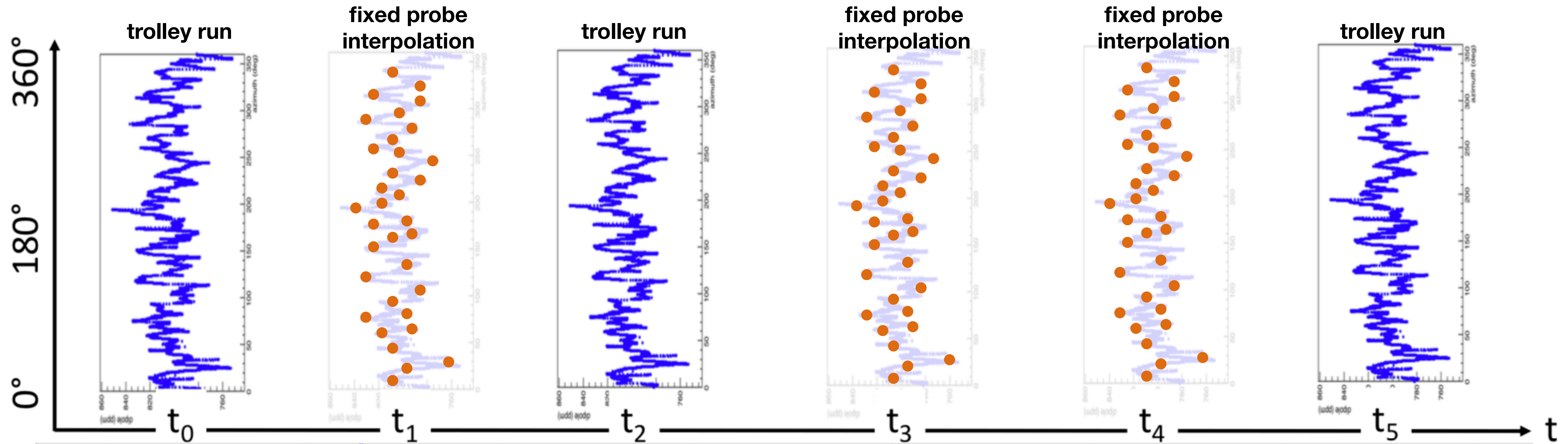
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Run 1 Analysis Status: ω_p — Field Interpolation



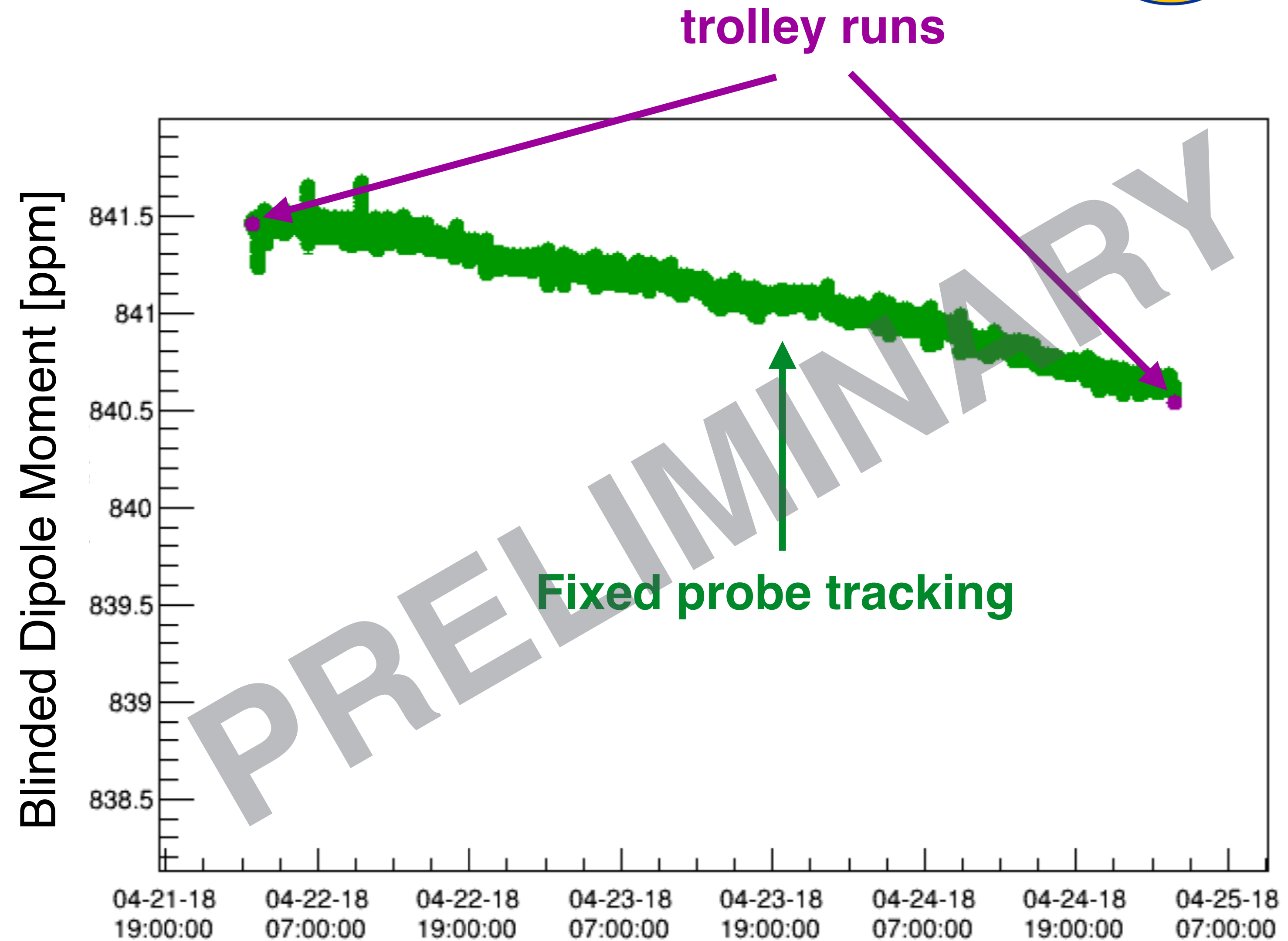
- Need to determine ω_p at all times while storing muons => interpolate between trolley maps using fixed probe data



Run 1 Analysis Status: ω_p — Field Interpolation



- Example from subset of data
- Tracking algorithms showing good agreement with trolley runs
- Also tracking higher-order multipole moments — important for extracting muon-weighted field $\tilde{\omega}_p$



Summary



- **The Muon g-2 Experiment** is a highly sensitive test of the SM
 - Discrepancy between theory and experiment for $a_\mu > \sim 3\sigma$
- ✓ Completed Run 1 in July 2018 (1.1x BNL statistics)
- ➡ Analyses are mature and progressing towards a result in **early 2020**
- ✓ Completed Run 2 in July 2019 (1.9x BNL statistics)
 - Starting to organize analysis efforts
- Run 3 starting this November: aiming to **triple** statistics to date

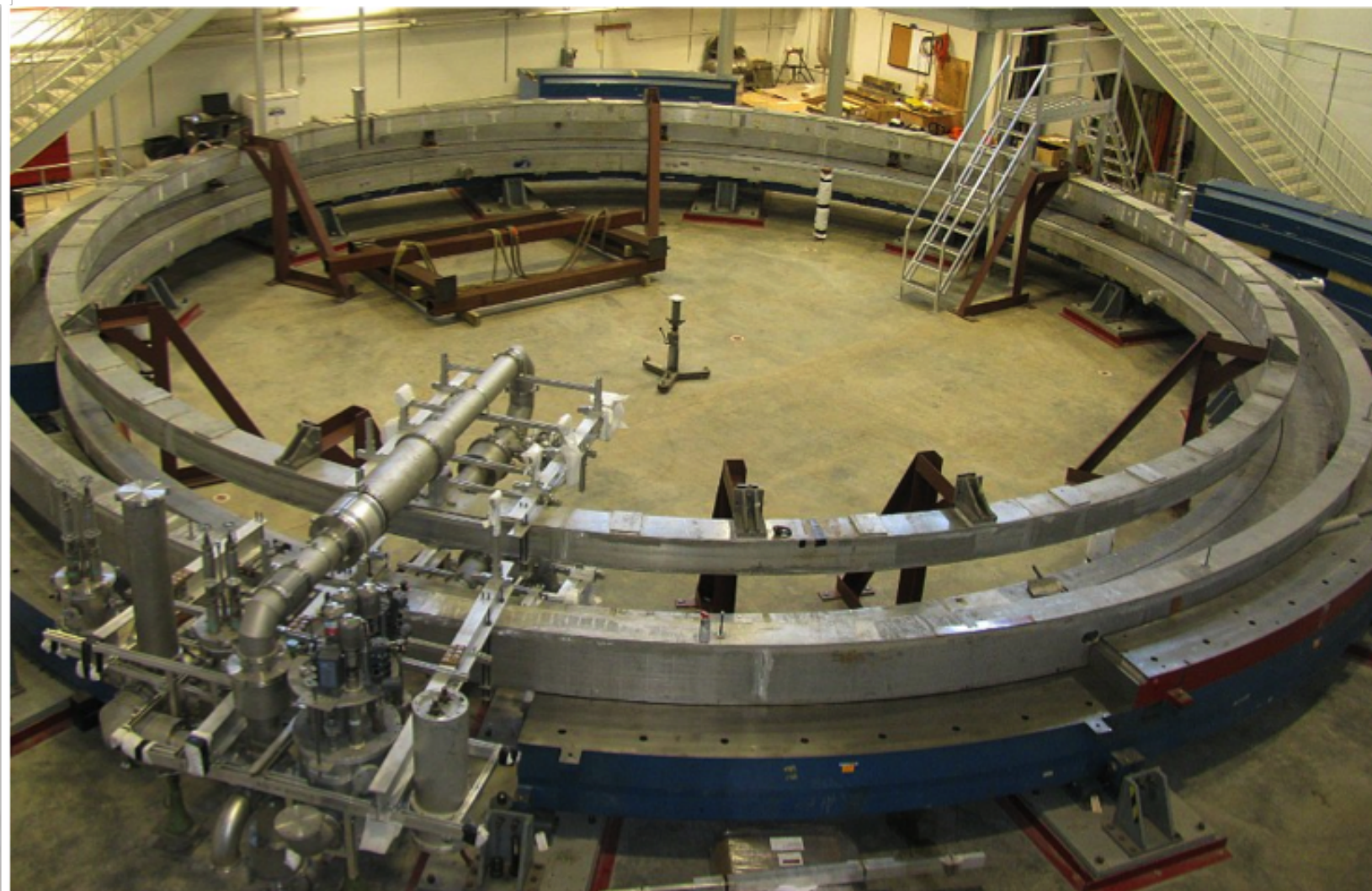
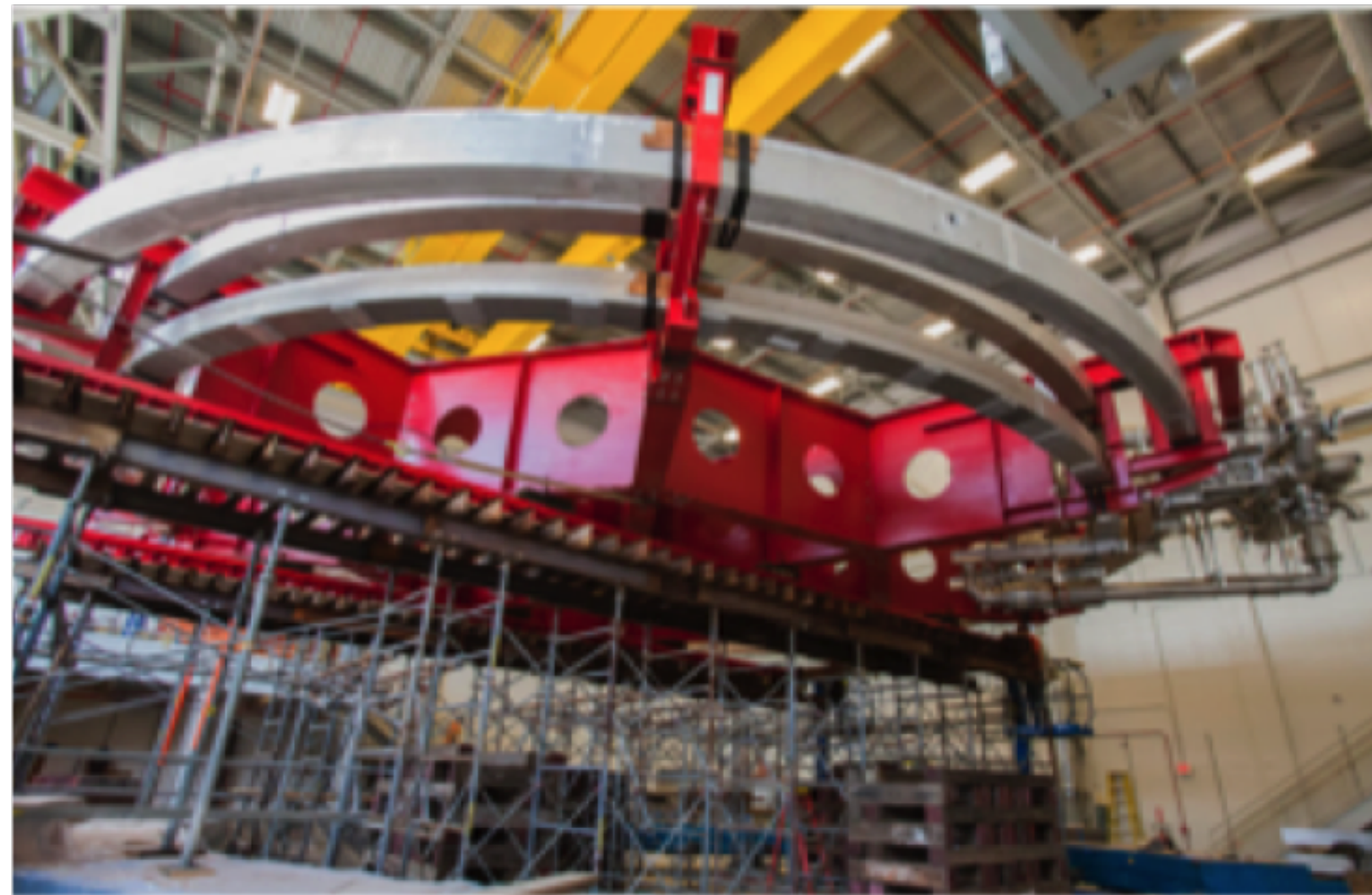
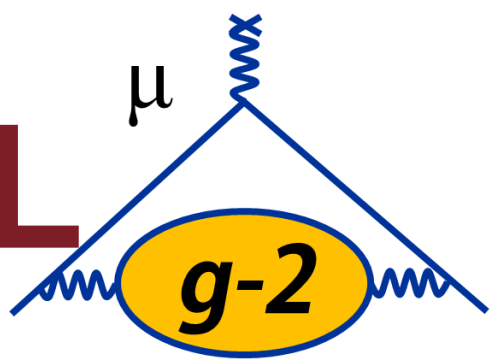
Thank You!





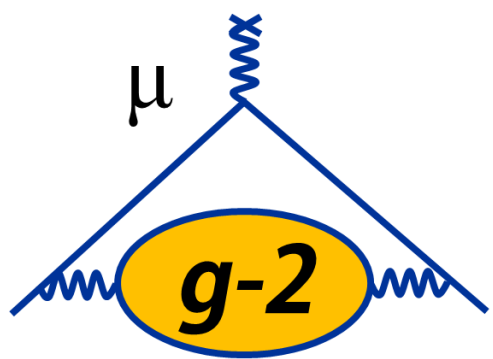
Backup

The Big Move: Transporting the Ring from BNL to FNAL

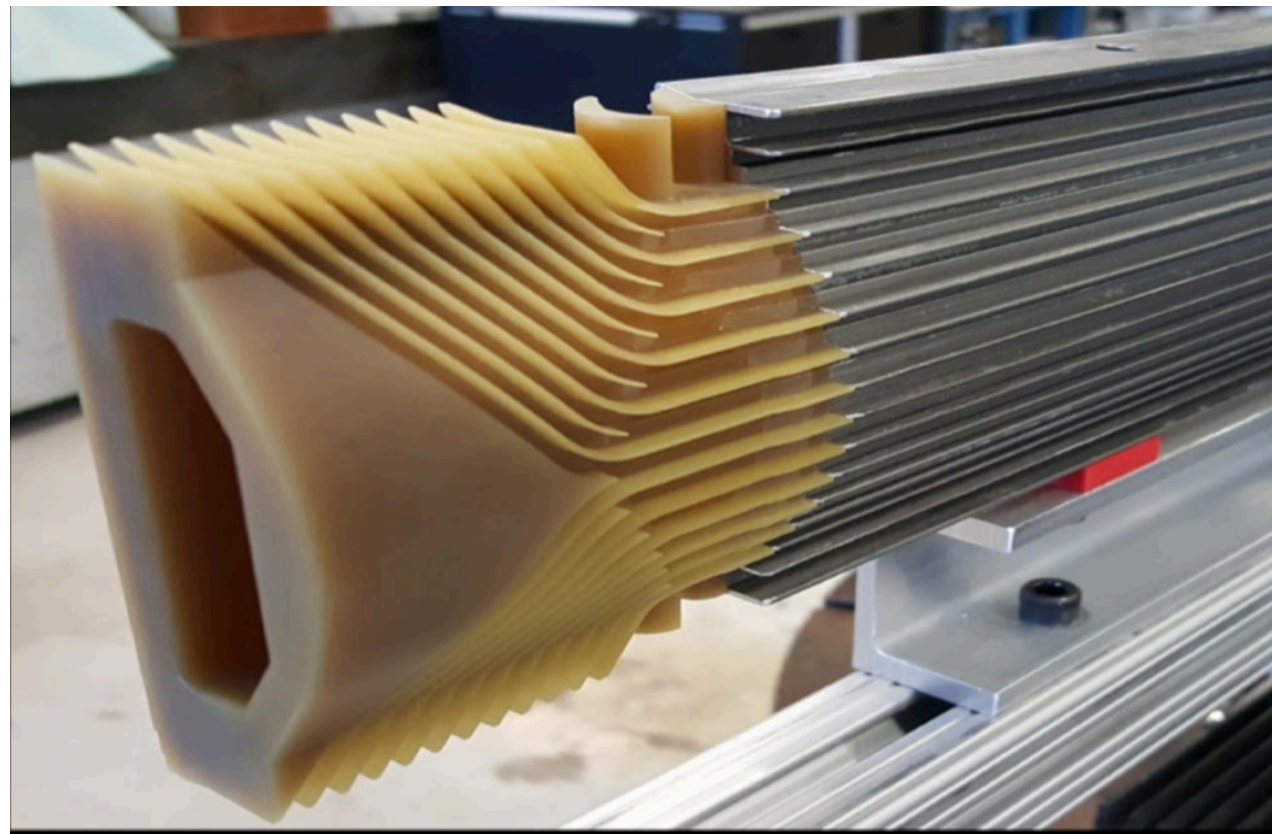
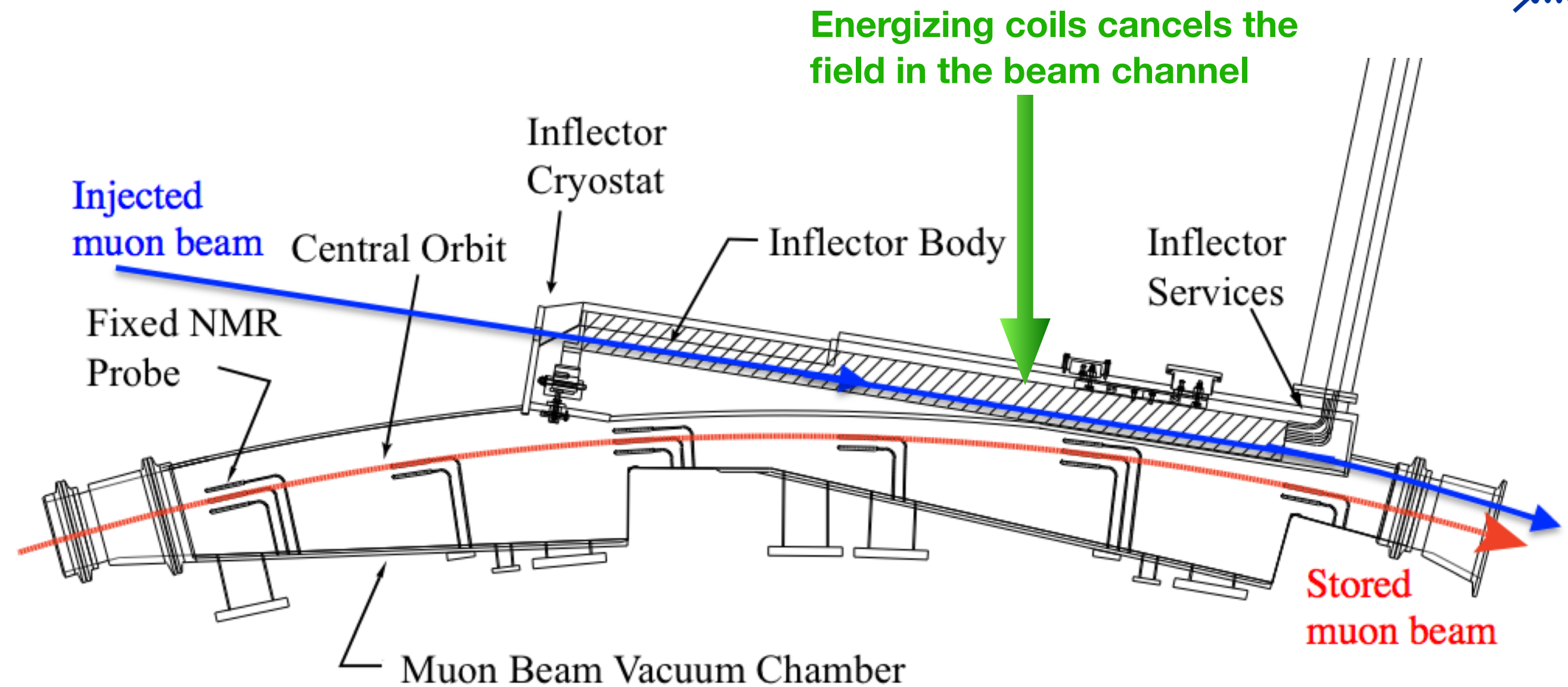


- June 2013—June 2015
- Ring deconstructed at BNL, transported by barge/flatbed trailer
- Reassembled at FNAL
- Ring successfully cooled and powered to 1.45 T in September 2015 — remarkable achievement!

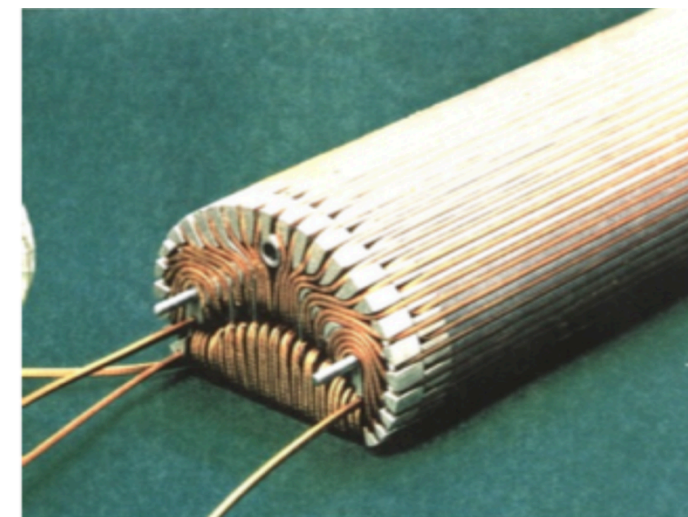
Getting Muons Into the Ring: Inflector Magnet



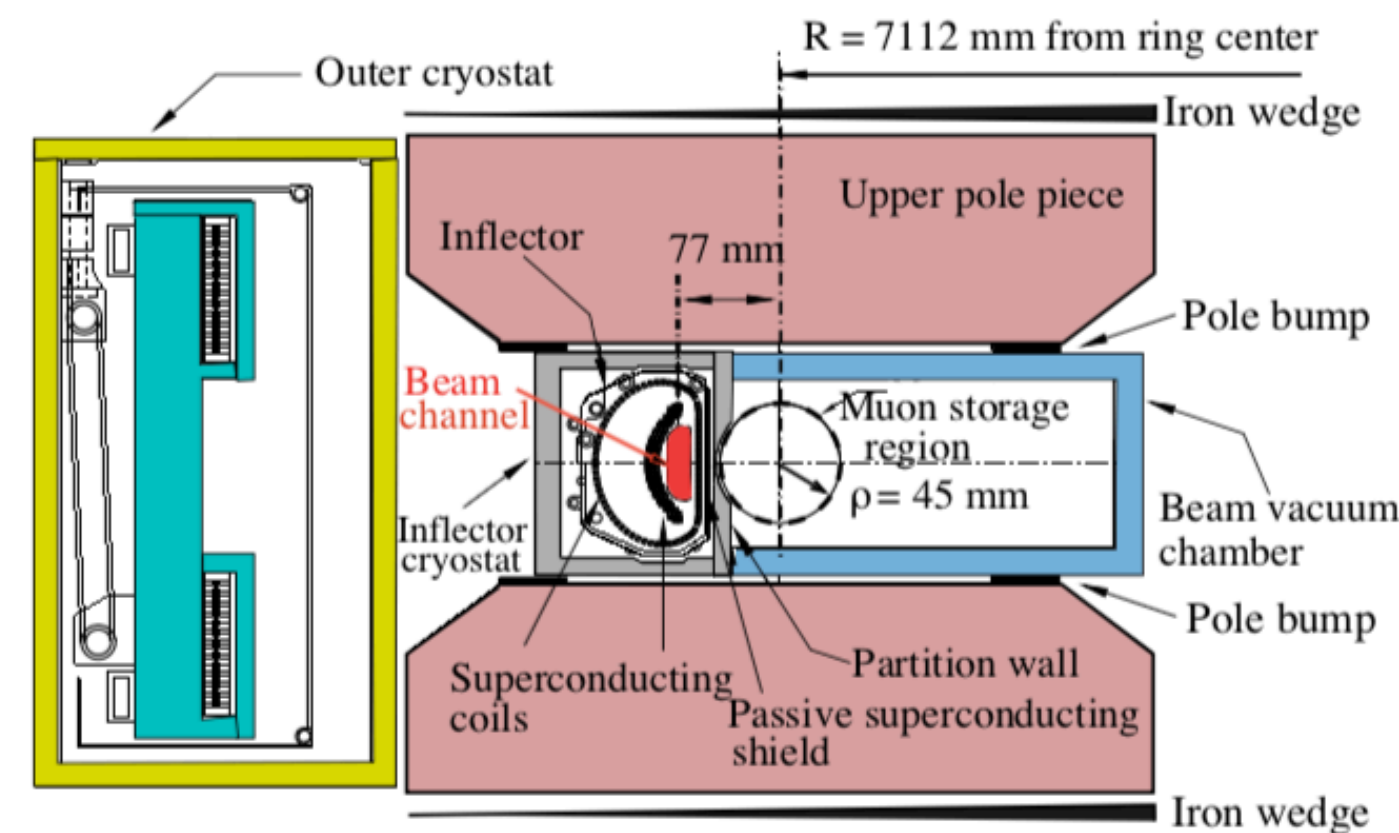
- Outside ring: $B = 0$ T, inside: $B = 1.45$ T
- Need to cancel field in order to get muons in (strong deflection otherwise)
- No perturbation to field outside shield
- New inflector design with higher transmission under development
- **Improve injection by 40%**



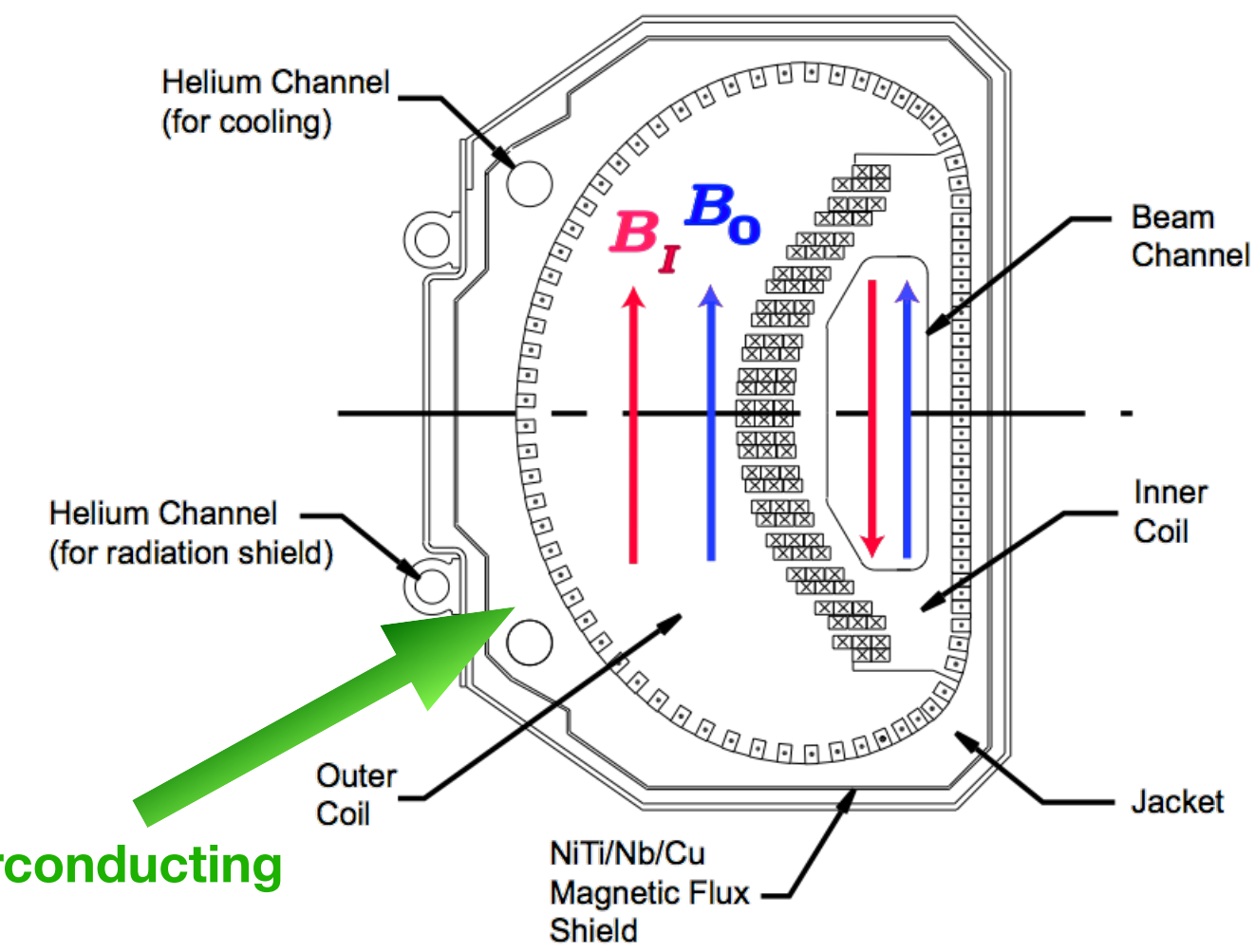
New inflector coil winding mount



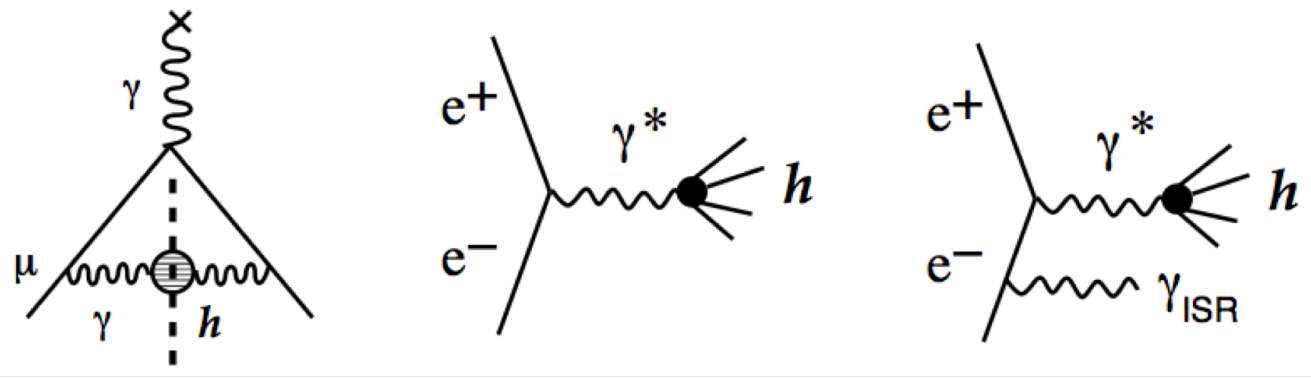
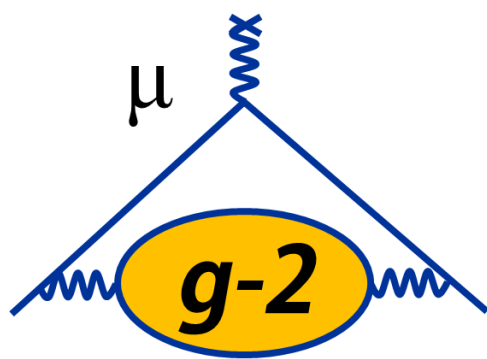
Present inflector



Super currents in passive superconducting shield prevents flux leakage

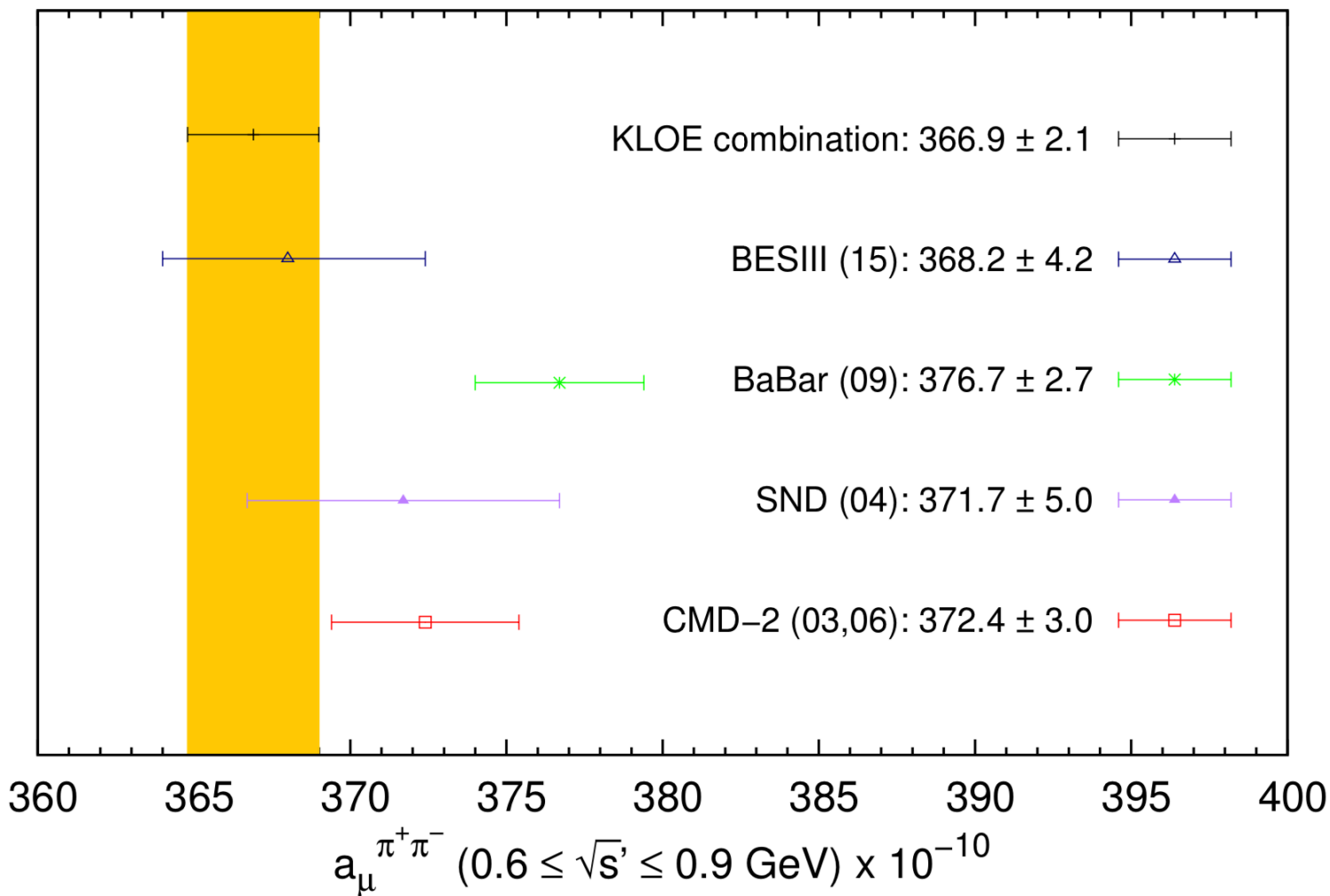


Theory Status of Hadronic Contribution to a_μ

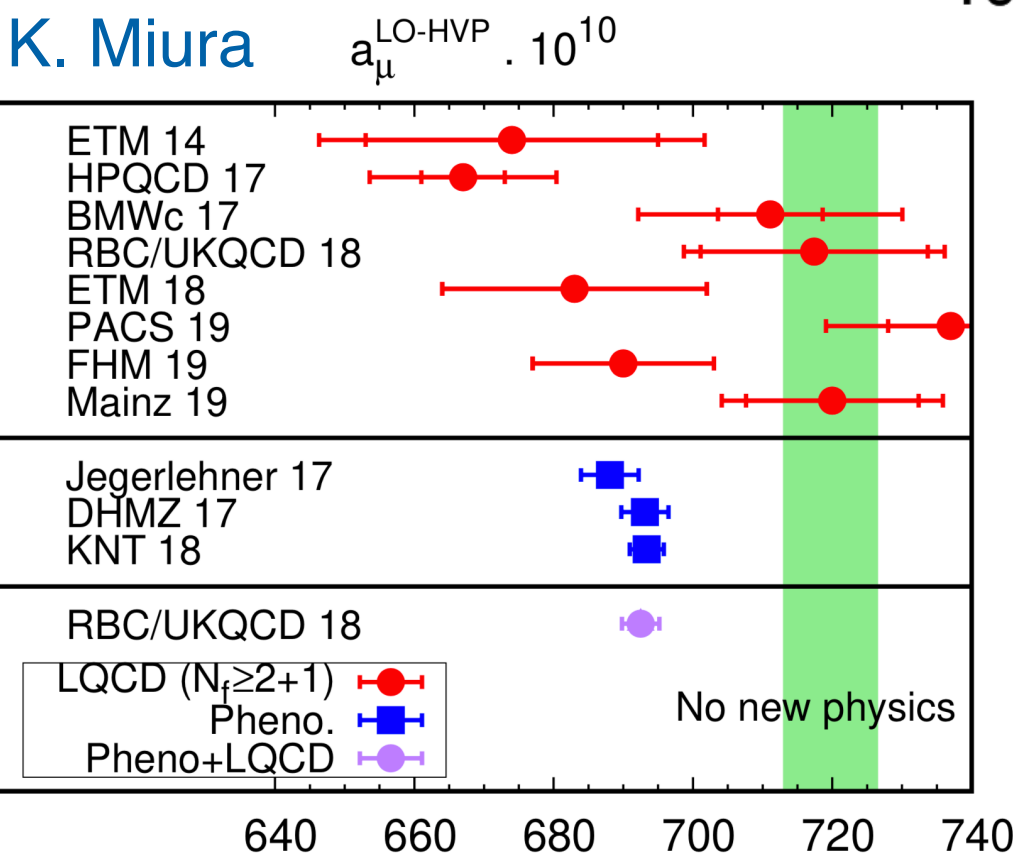


$$a_\mu^{\text{had;LO}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2 \int_{m_\pi^2}^\infty \frac{ds}{s^2} K(s) R(s)$$
$$R \equiv \frac{\sigma_{\text{tot}}(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

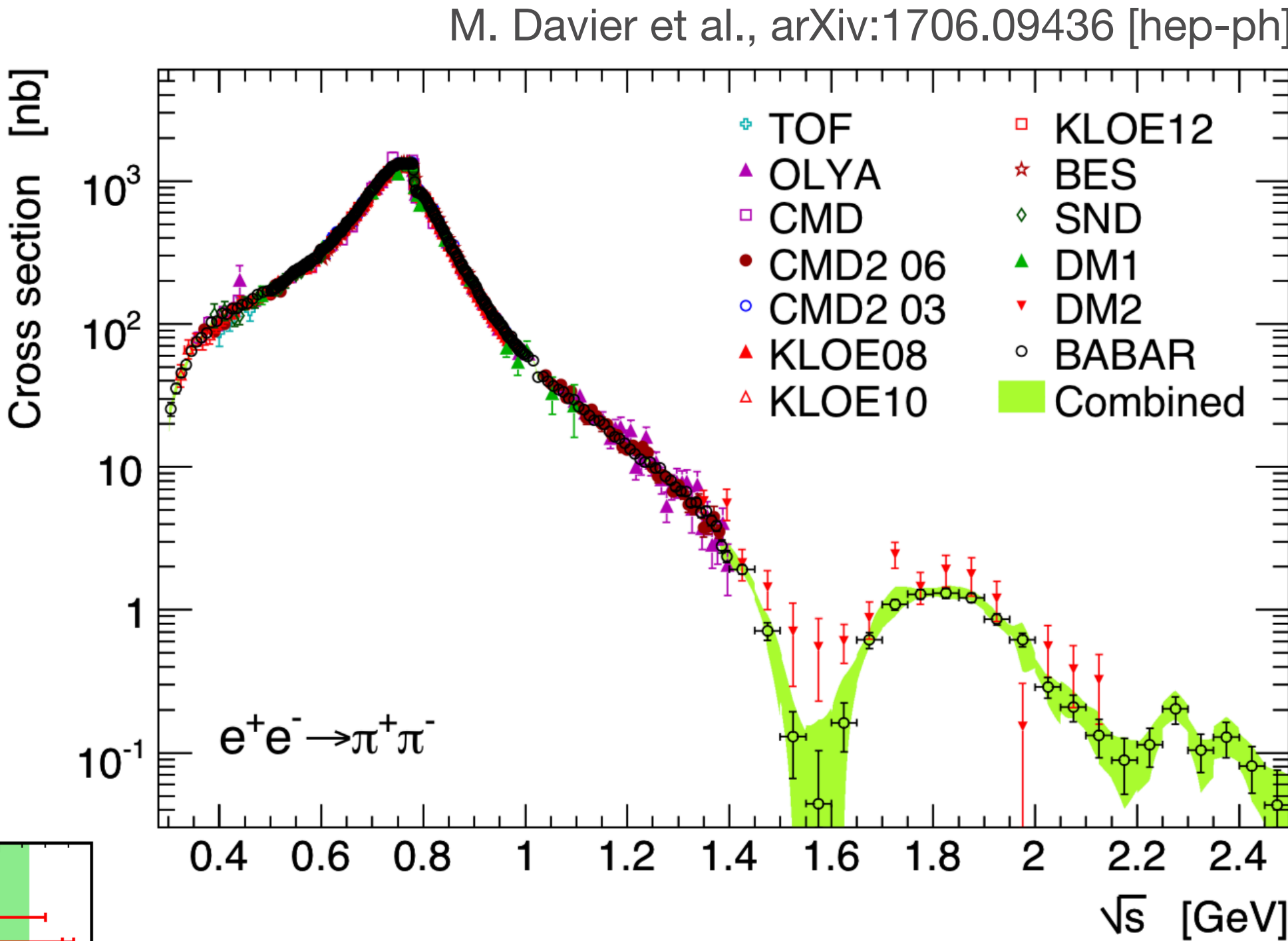
- **Critical input to HVP** from e+e- colliders (SND, CMD3, BaBar, KLOE, Belle, BESIII)
- **BESIII**: 3x more data available, luminosity measurement improvements
- **VEPP-2000**: Aiming for 0.3% (fractional) uncertainty; radiative return + energy scan
- **CMD3**: Will measure up to 2 GeV (energy scan, ISR — good cross check)



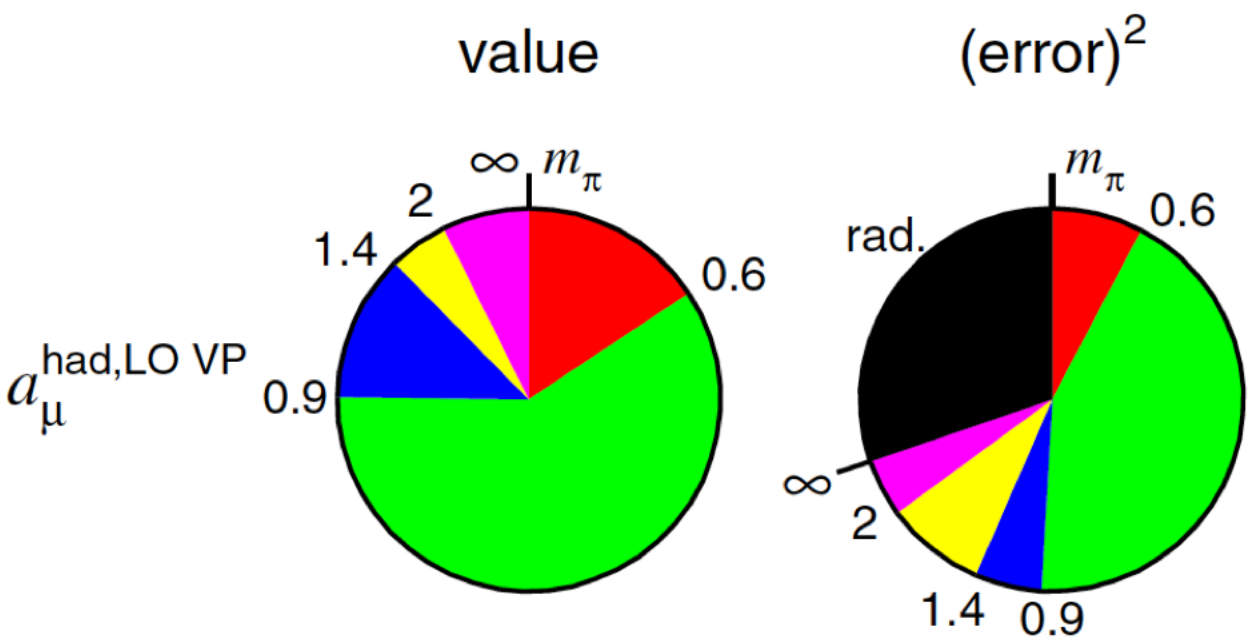
A. Anastasi et al., arXiv:1711.03085 [hep-ex]



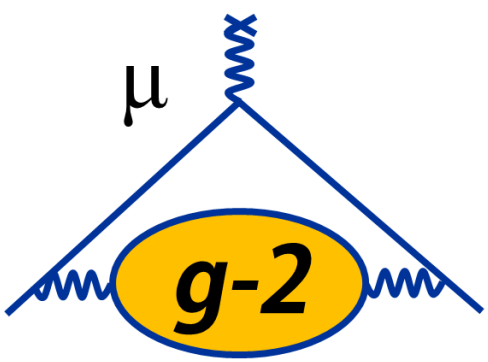
- **Lattice calculations** of a_μ^{HVP} to 1% soon, 30% for HLbL in 3–5 years



M. Davier et al., arXiv:1706.09436 [hep-ph]

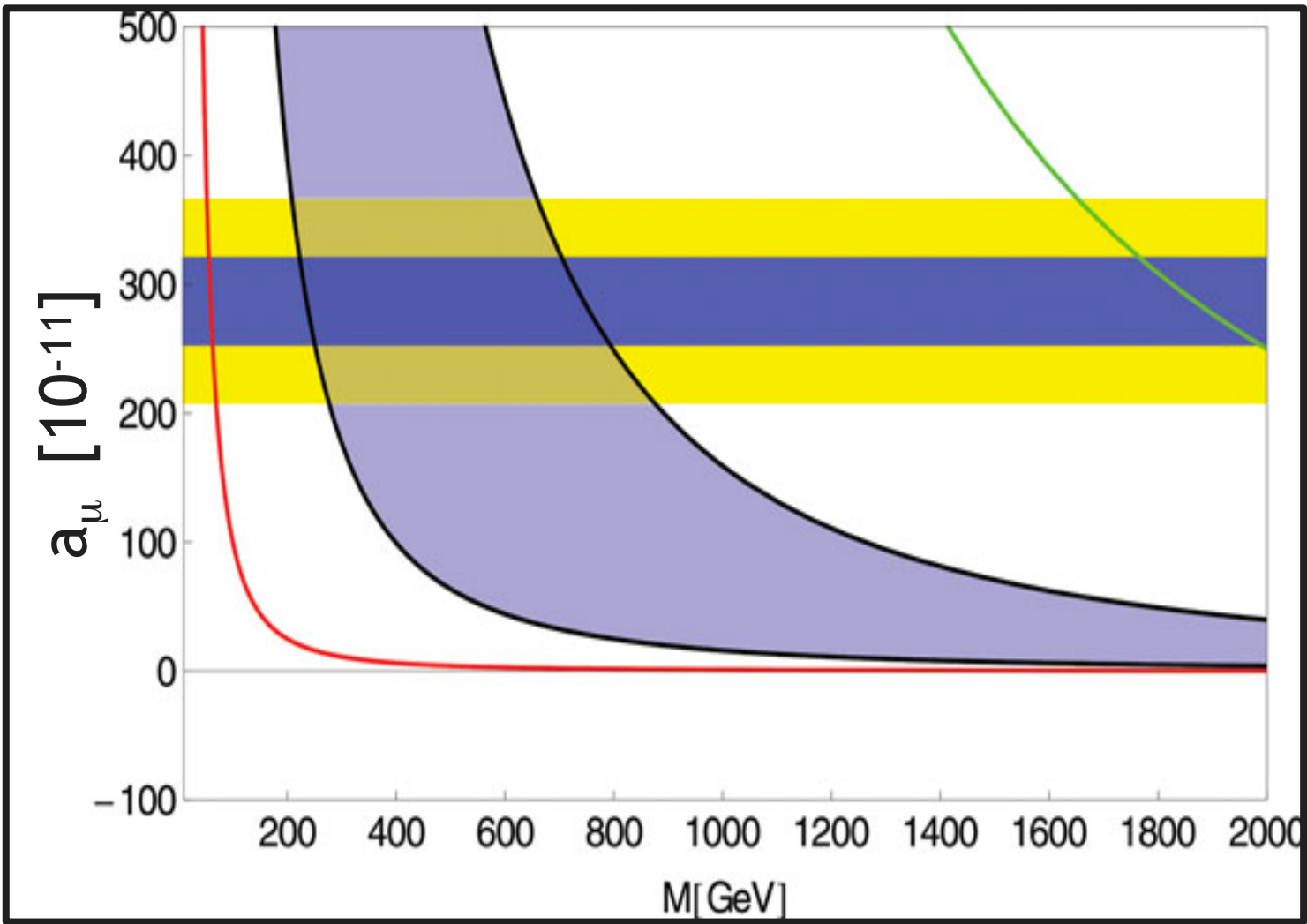
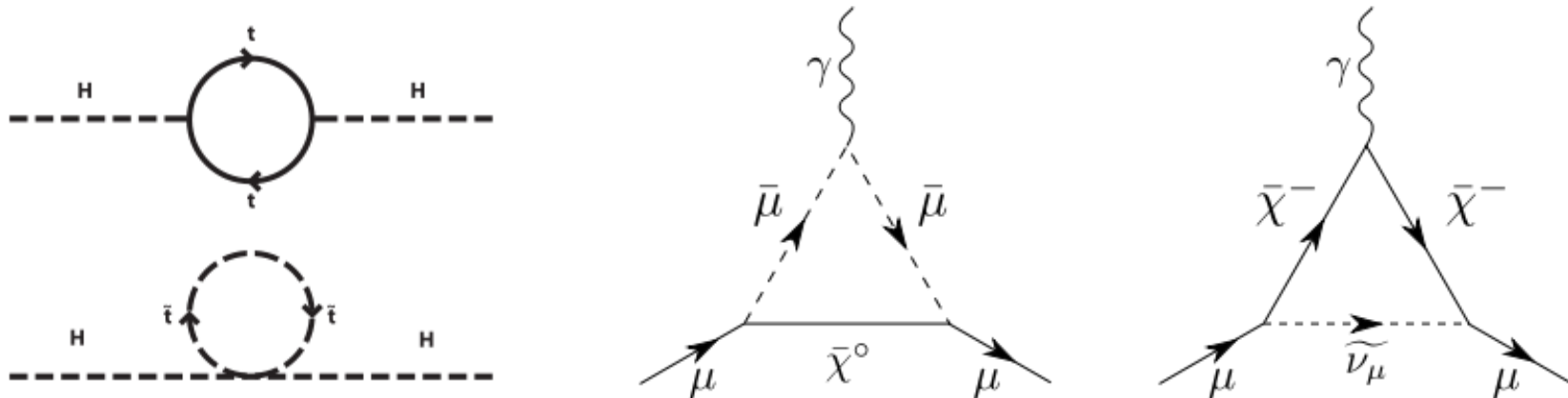


Physics Beyond the Standard Model?



SUSY, TeV-Scale Models

- Higgs measured at the LHC to be 125 GeV
- Theory: Higgs should acquire much heavier mass from loops with heavy SM particles (e.g., top quark)
 - Supersymmetry: new class of particles** that enters such loops and **cancels this contribution**

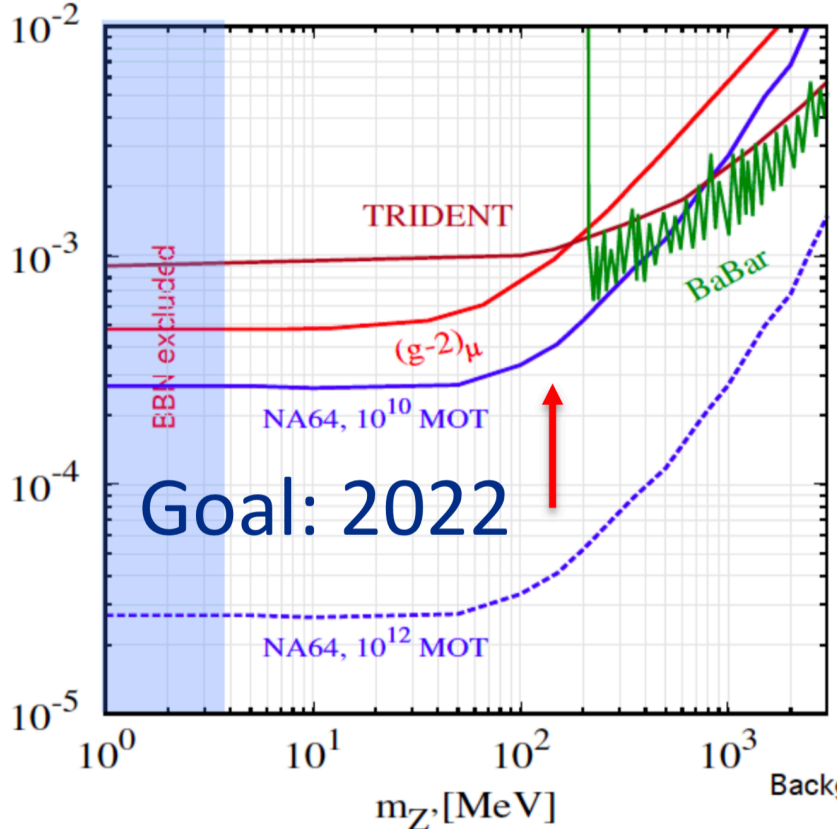
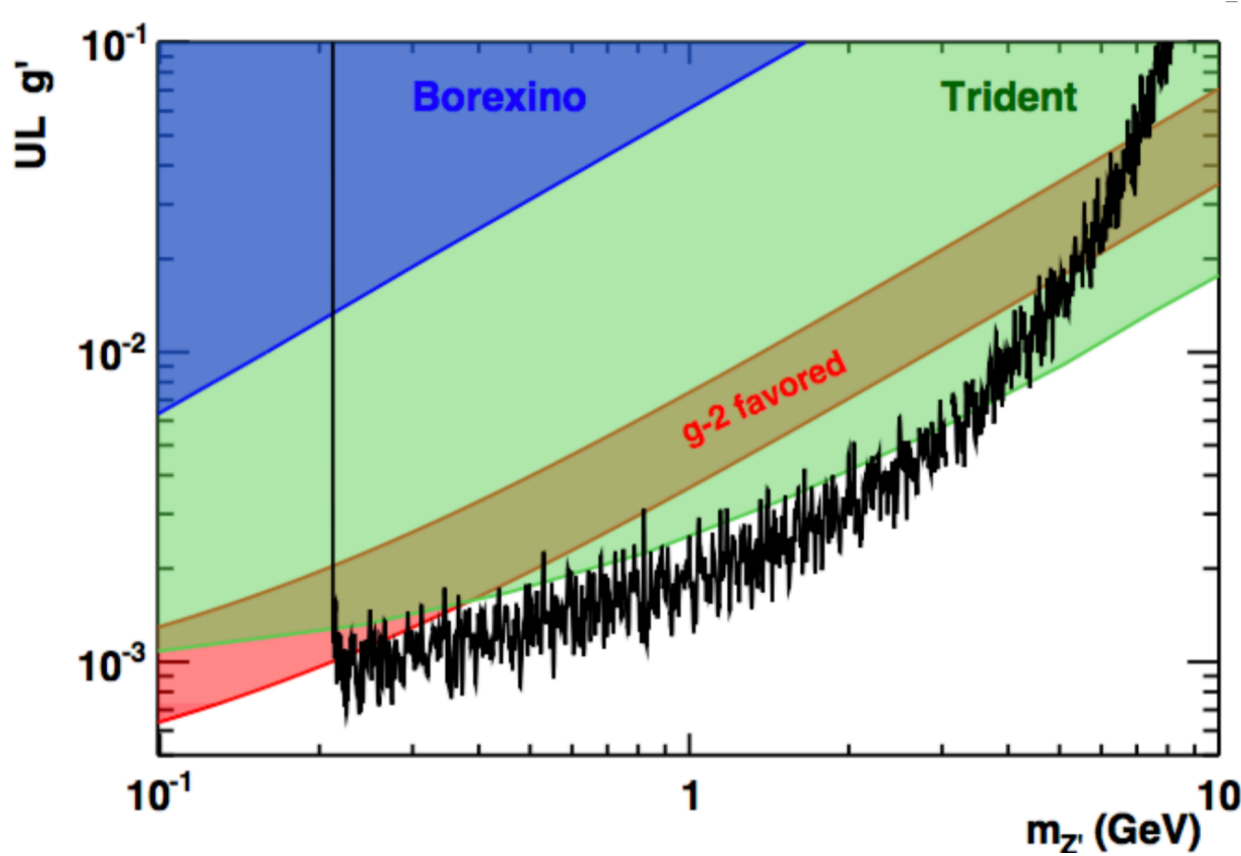


D. Hertzog, Ann. Phys. (Berlin), 2015, courtesy D. Stockinger

- Complementary to direct searches at the LHC**
 - Sensitivity to $\text{sgn}(\mu)$, $\tan(\beta)$
 - Contributions to a_μ arise from charginos, sleptons
 - LHC searches sensitive to squarks, gluinos

- Z', W', UED, Littlest Higgs**
 - Assumes typical weak coupling
- Radiative muon mass generation**
- Unparticles, Extra Dimension Models, SUSY ($\tan \beta = 5$ to 50)**

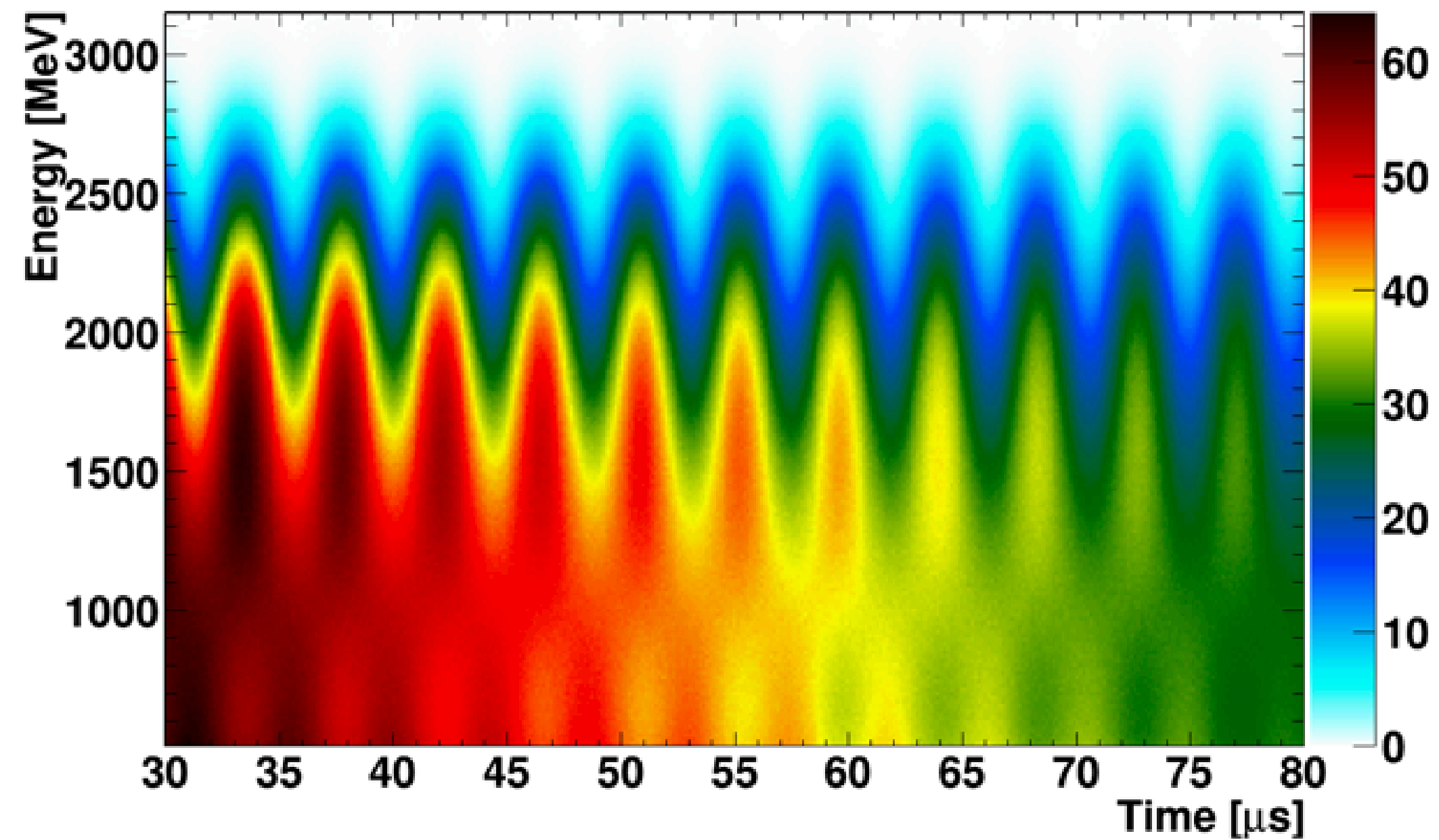
Z' Possibilities



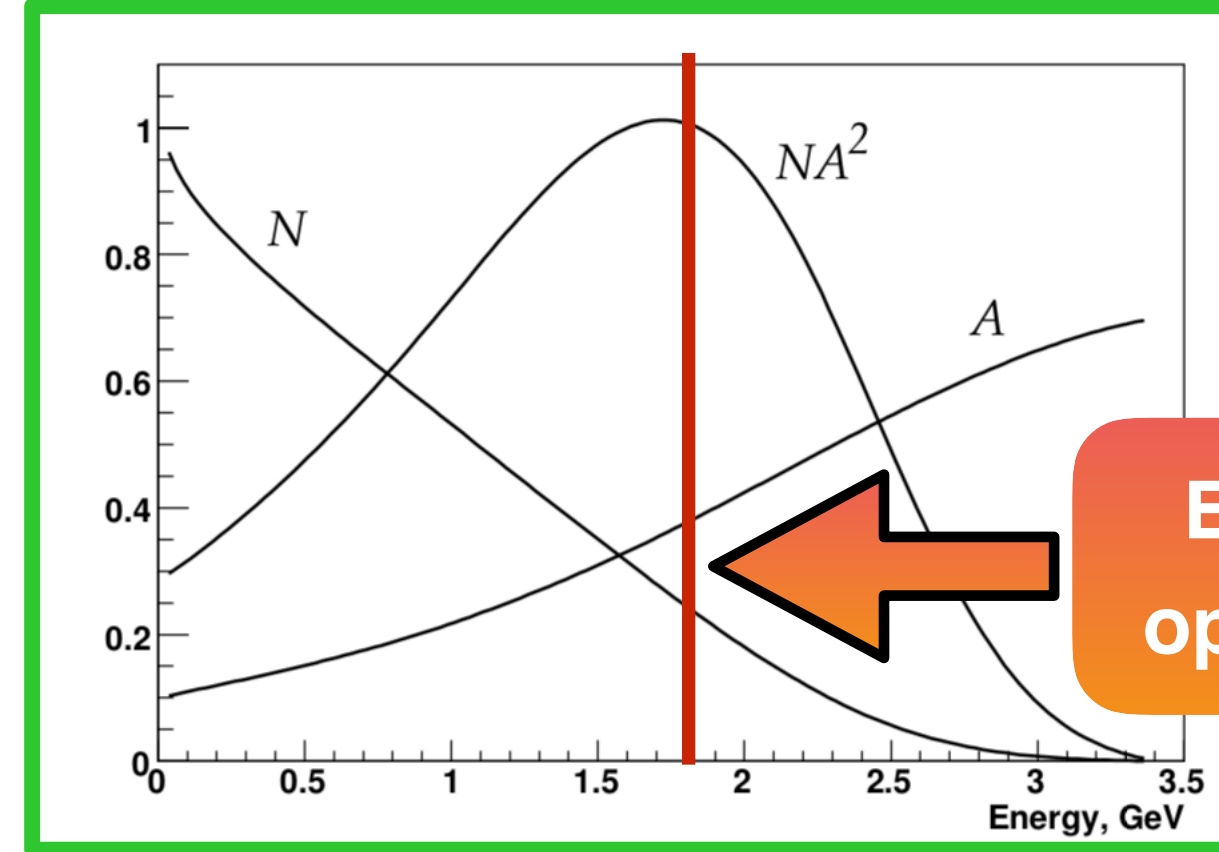
Others...

- Axion-like particles
- Dark photons (invisible)
- Extended Higgs/leptoquarks

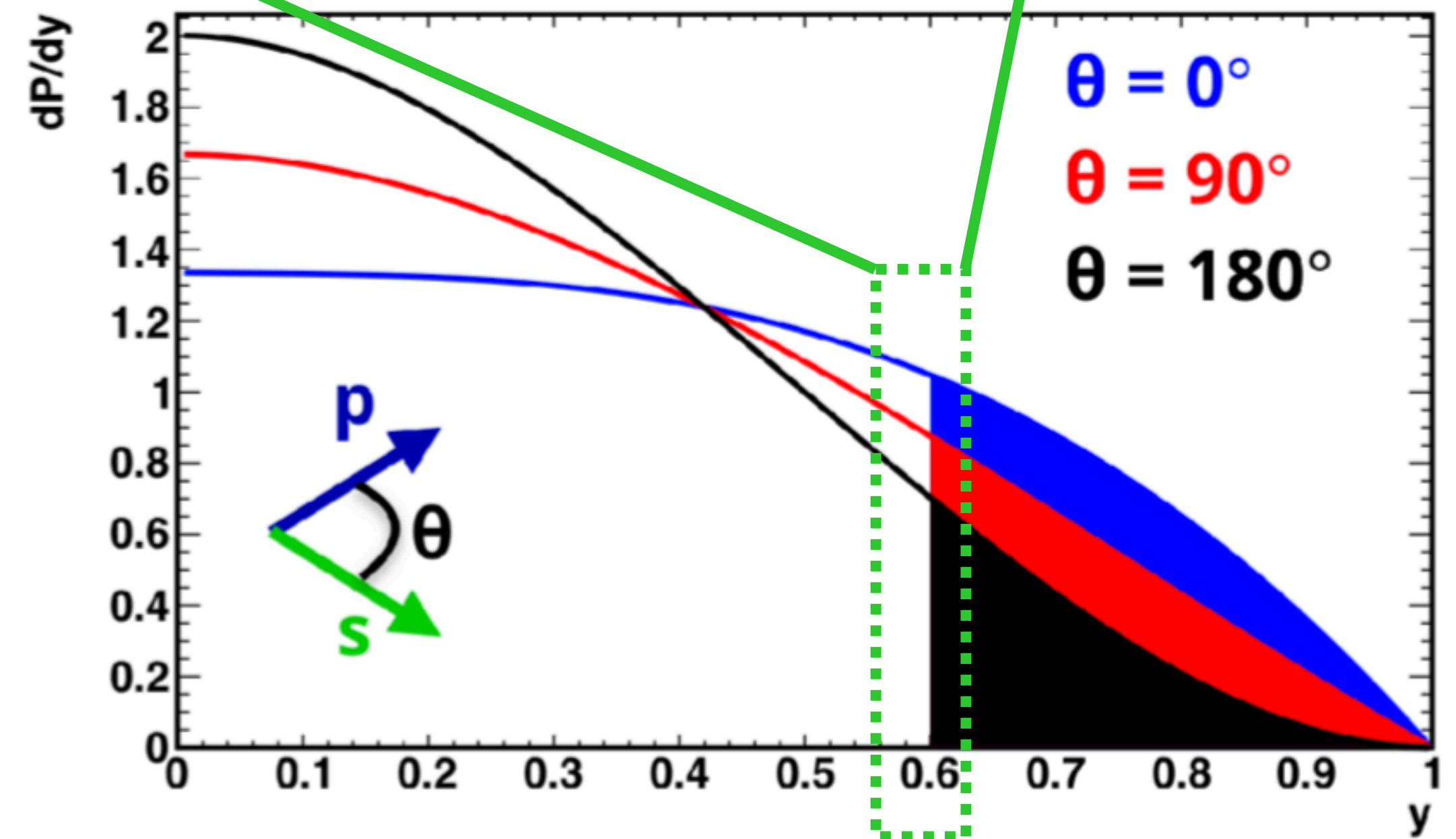
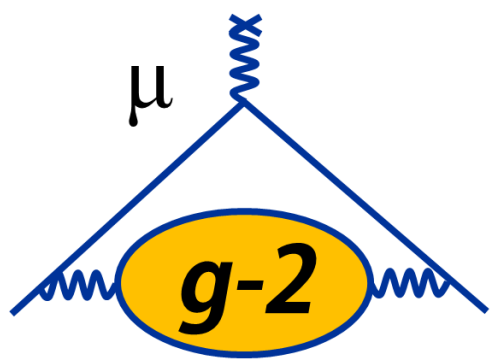
Analysis Details: ω_a



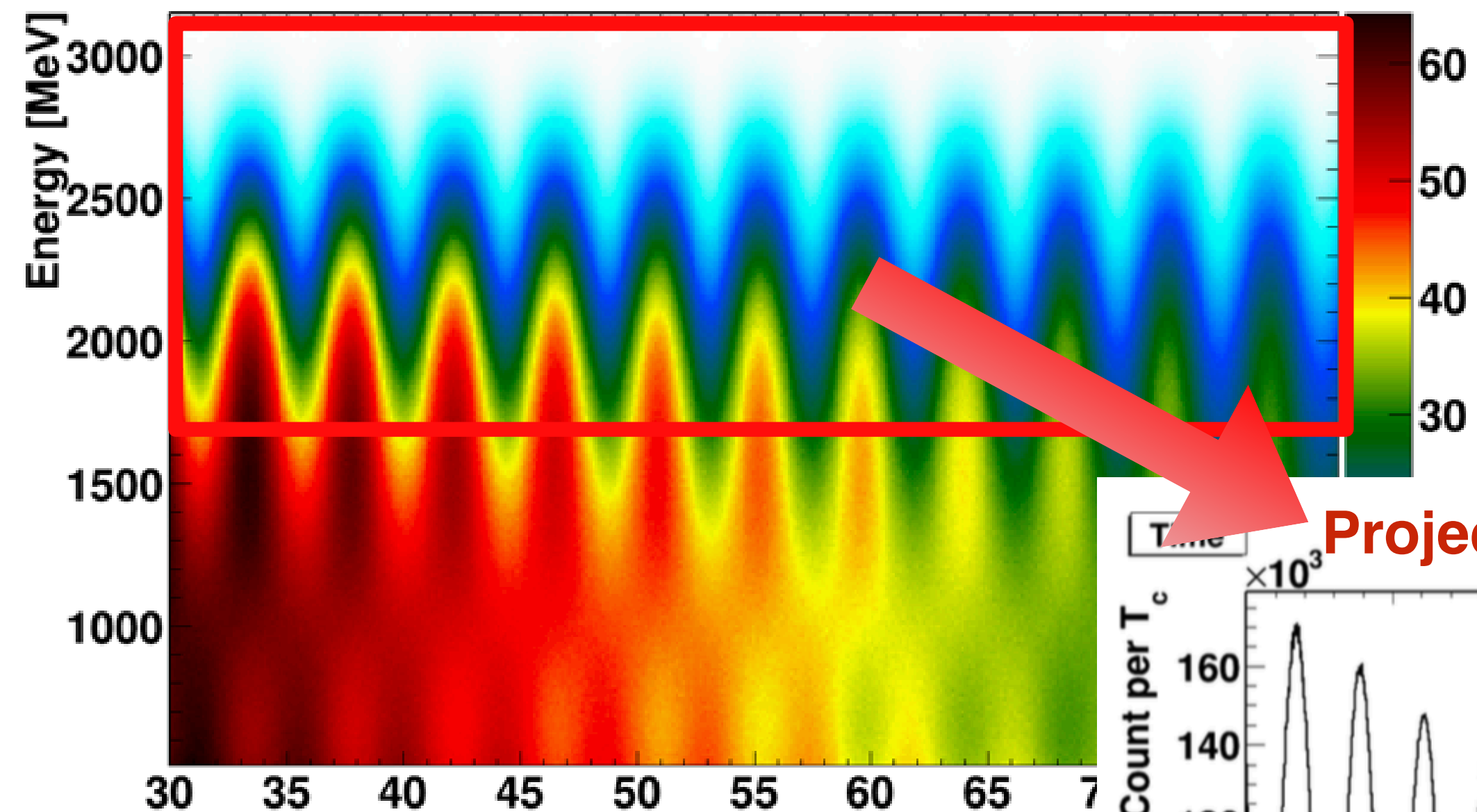
$$\frac{\delta\omega_a}{\omega_a} = \frac{\sqrt{2}}{2\pi f_a \tau_\mu N^{\frac{1}{2}} A}$$



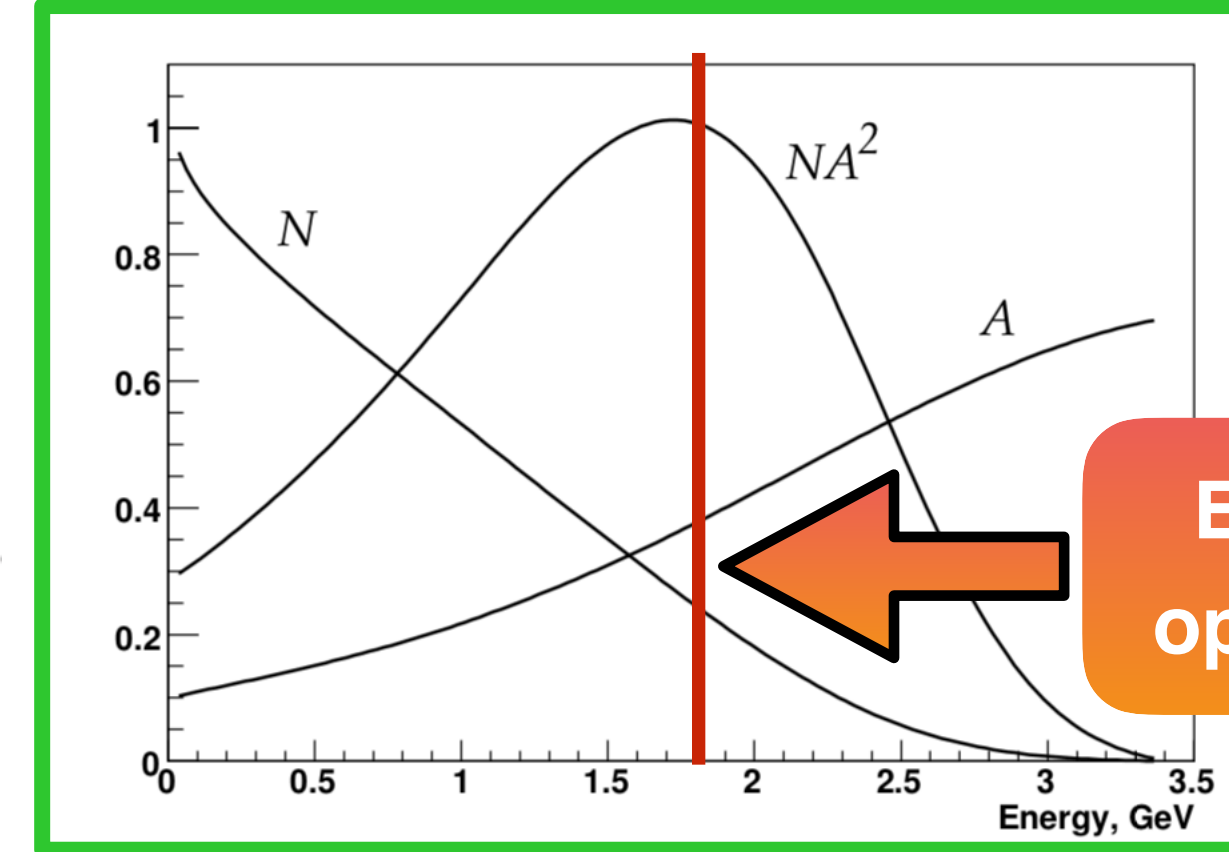
Energy cut chosen to optimize figure-of-merit



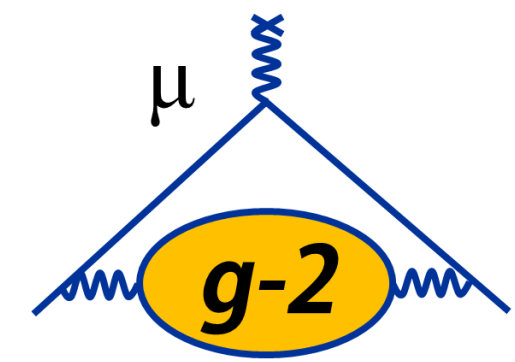
Analysis Details: ω_a



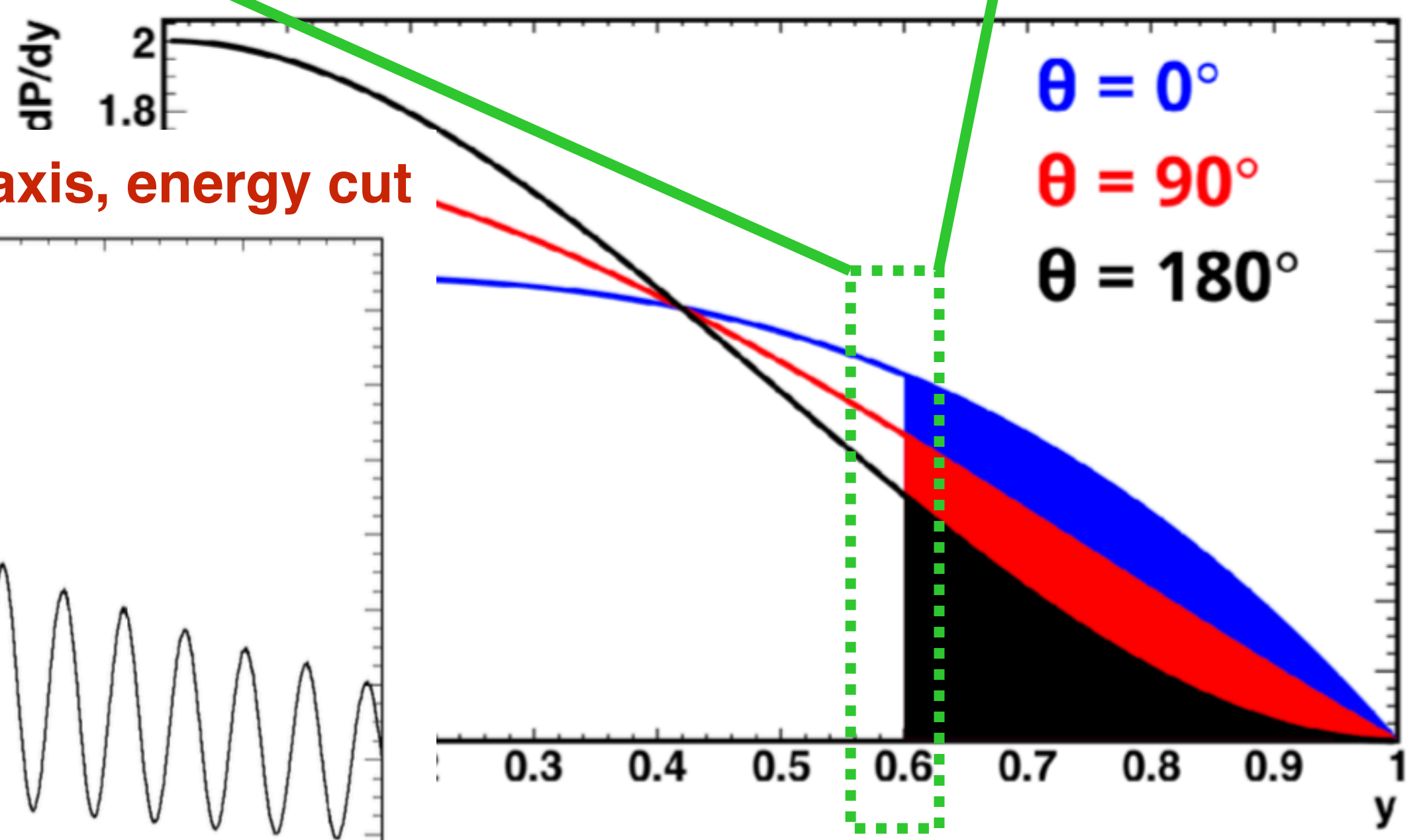
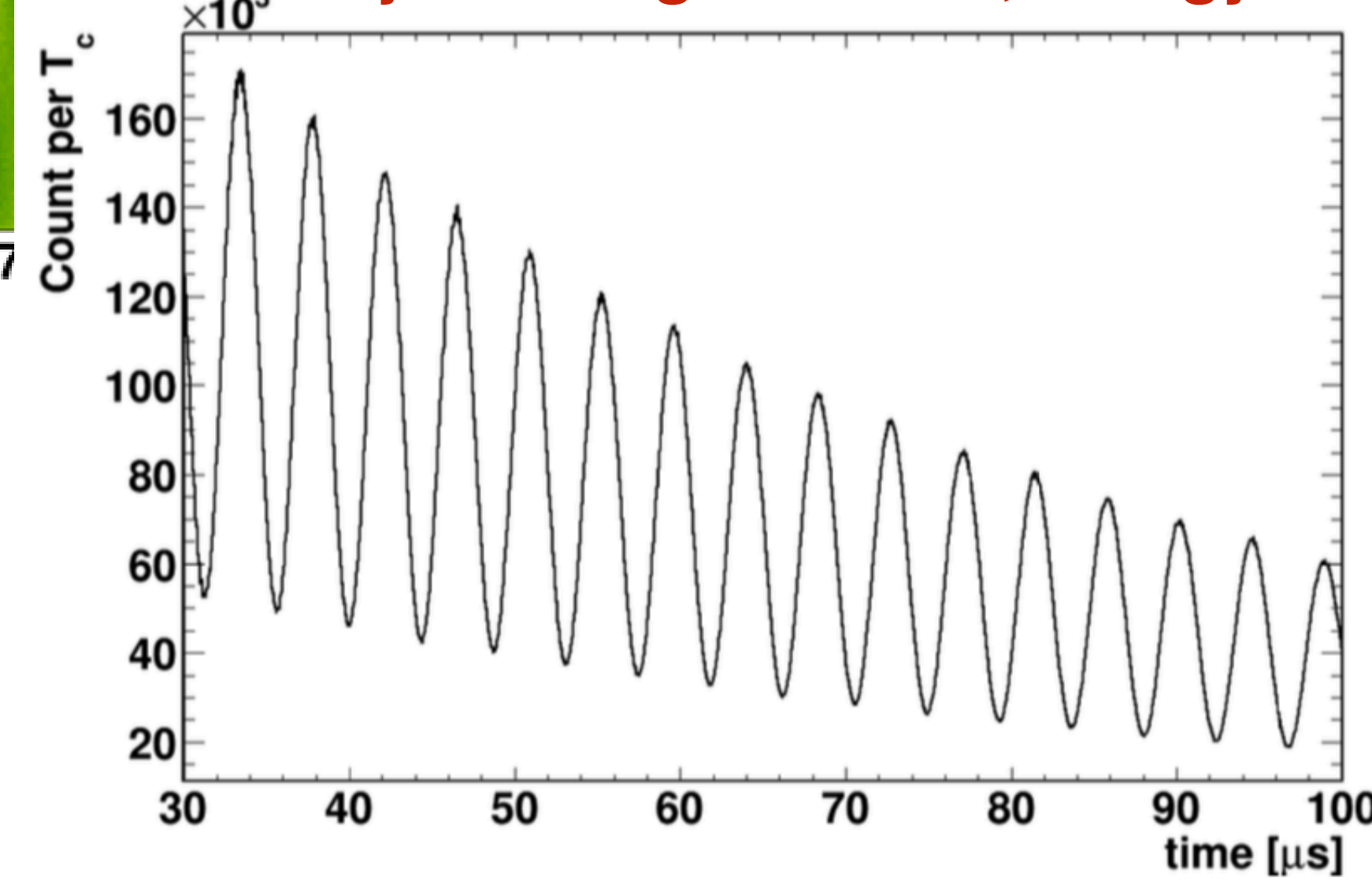
$$\frac{\delta\omega_a}{\omega_a} = \frac{\sqrt{2}}{2\pi f_a \tau_\mu N^{\frac{1}{2}} A}$$



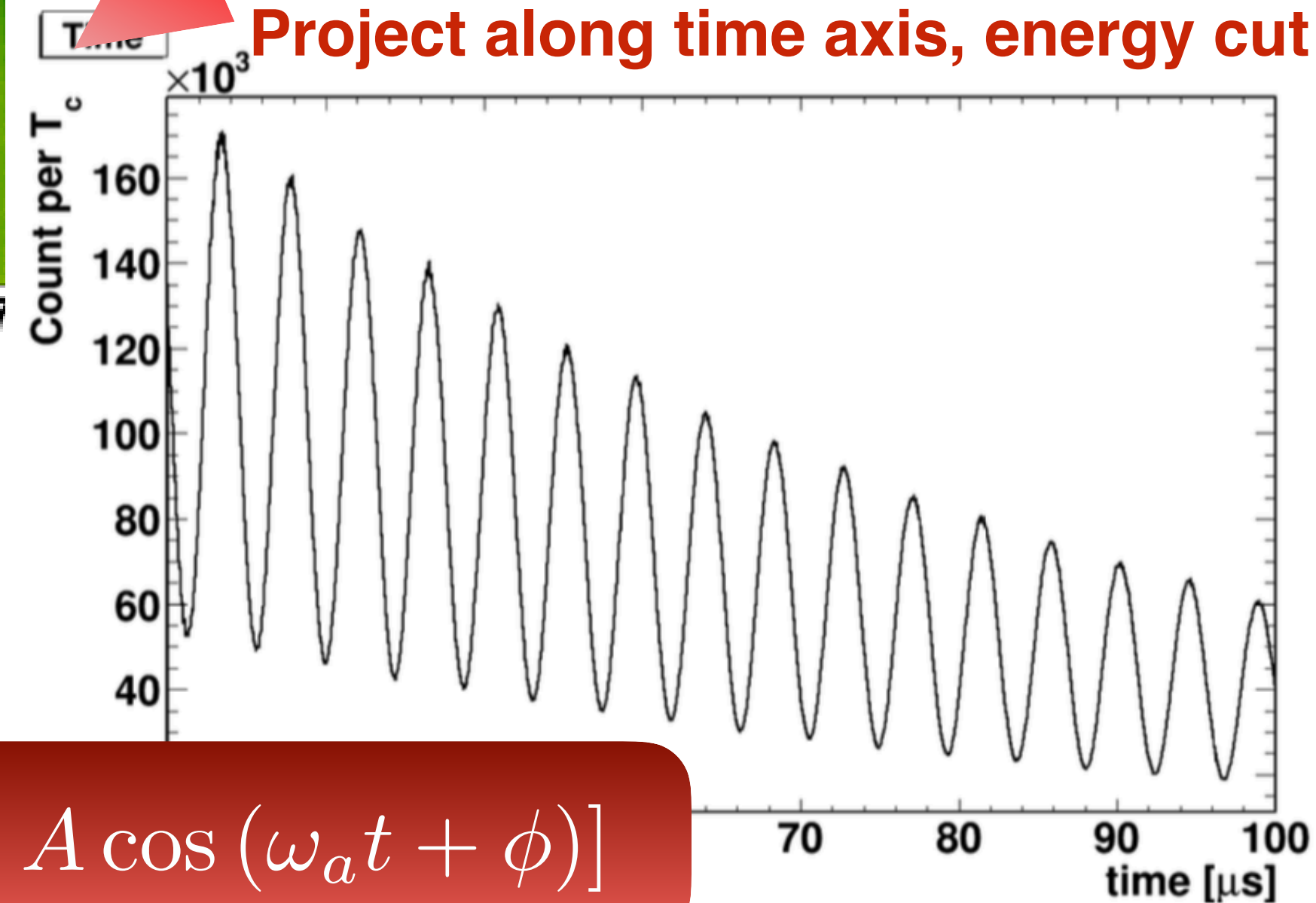
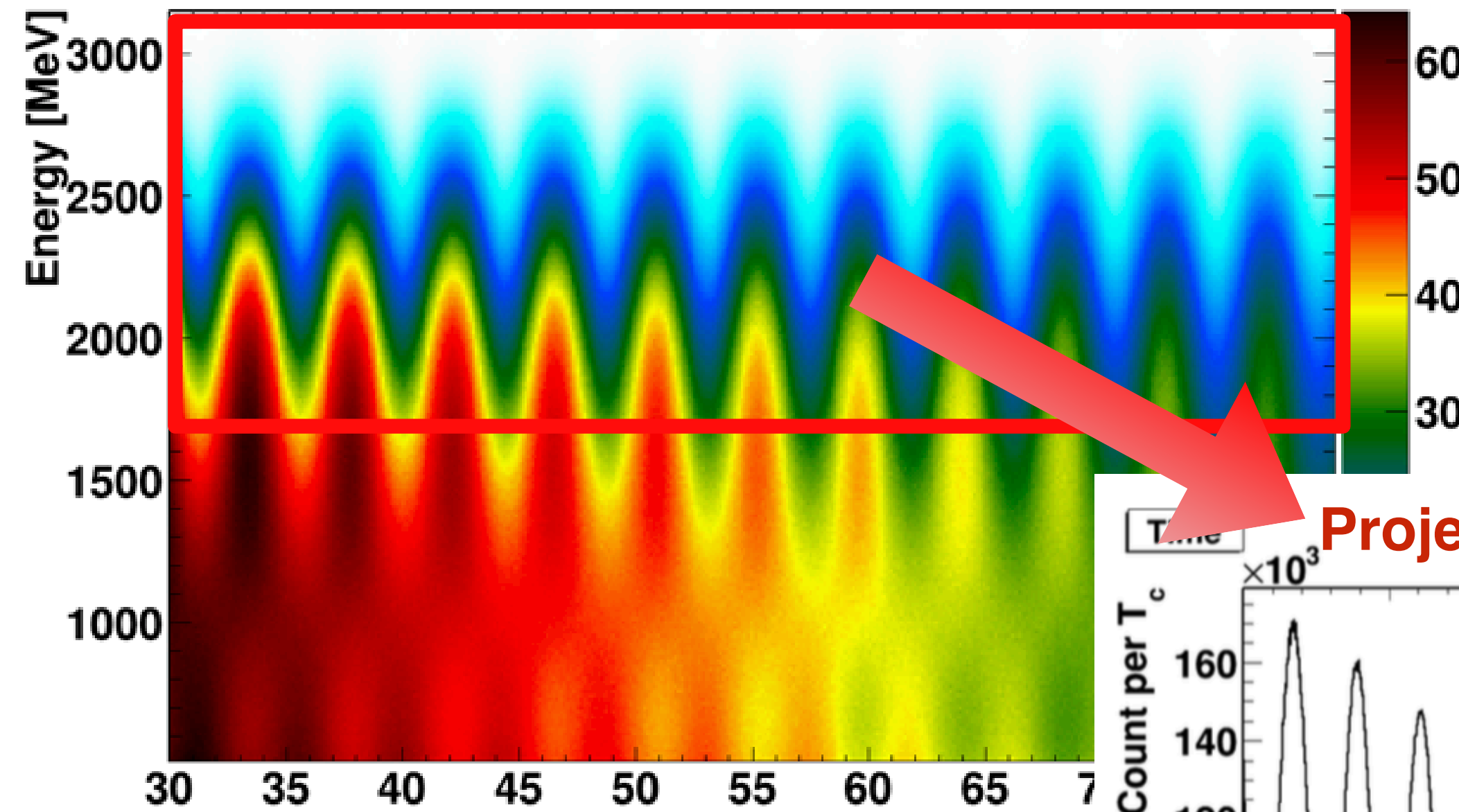
Energy cut chosen to optimize figure-of-merit



Project along time axis, energy cut

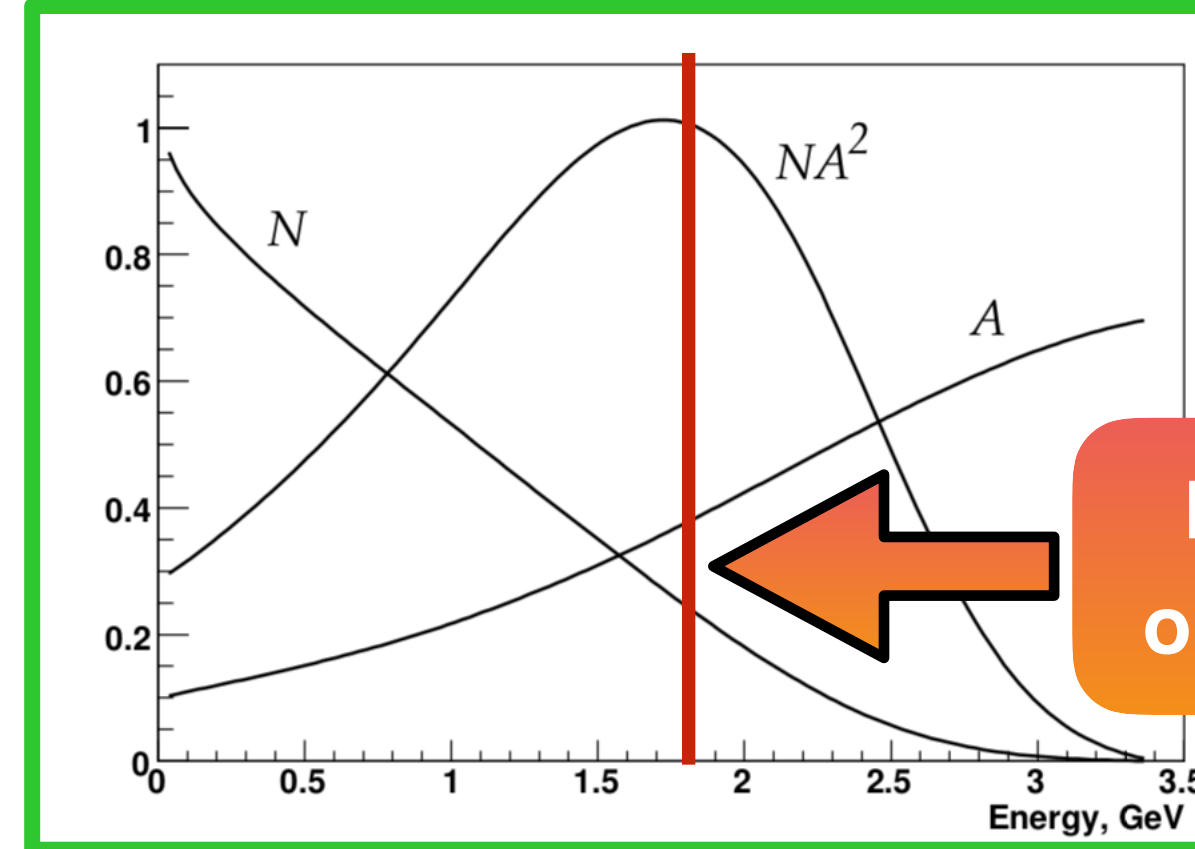


Analysis Details: ω_a

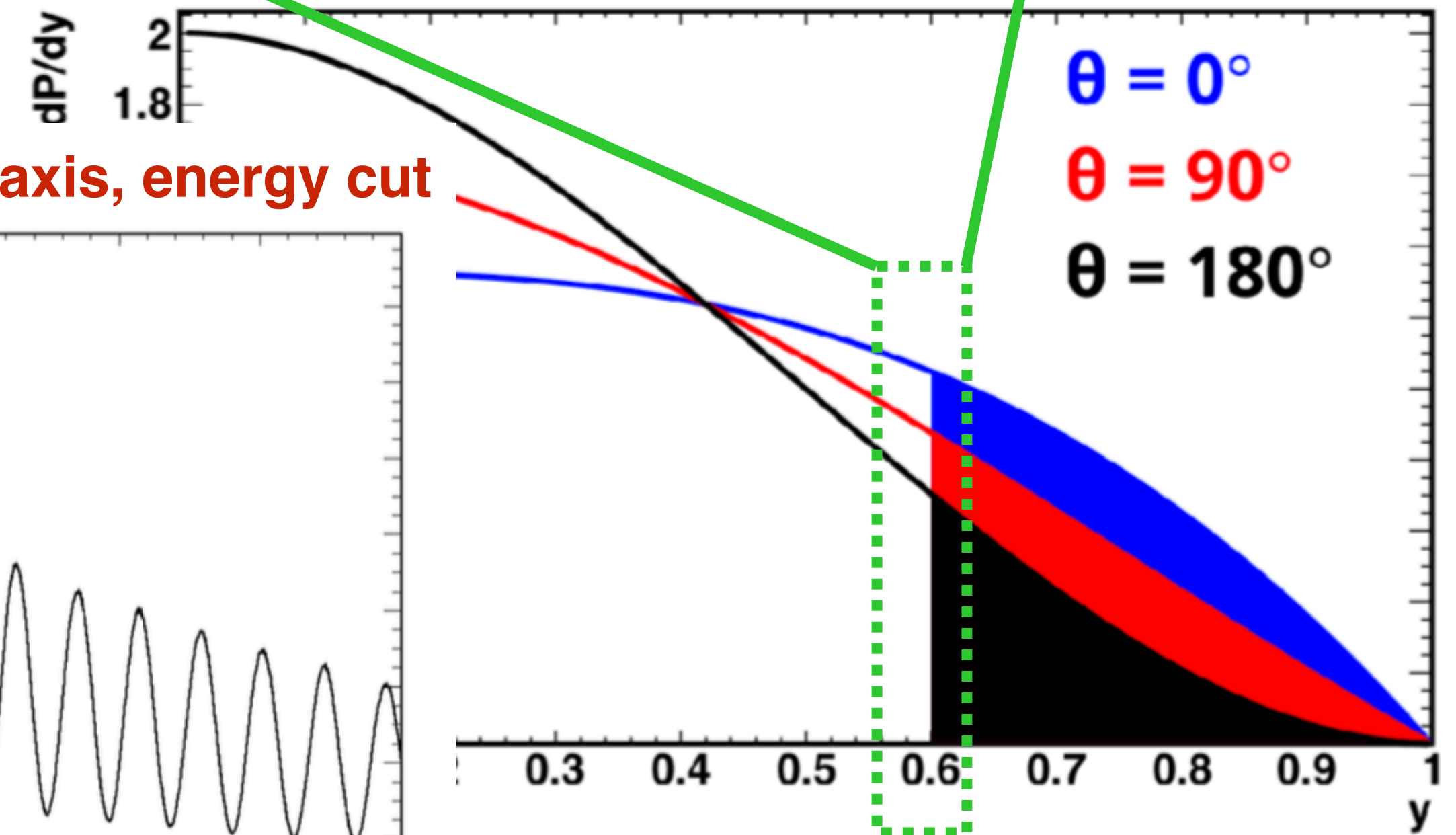
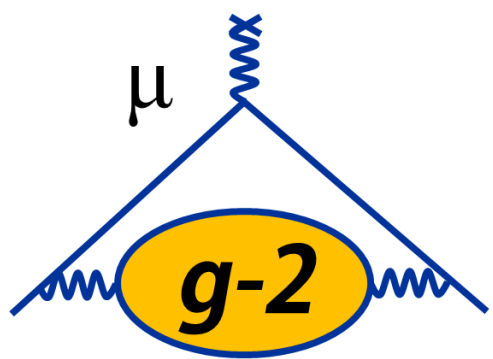


Fit to: $N(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi)]$

$$\frac{\delta\omega_a}{\omega_a} = \frac{\sqrt{2}}{2\pi f_a \tau_\mu N^{\frac{1}{2}} A}$$



Energy cut chosen to optimize figure-of-merit



Beam Dynamics Corrections



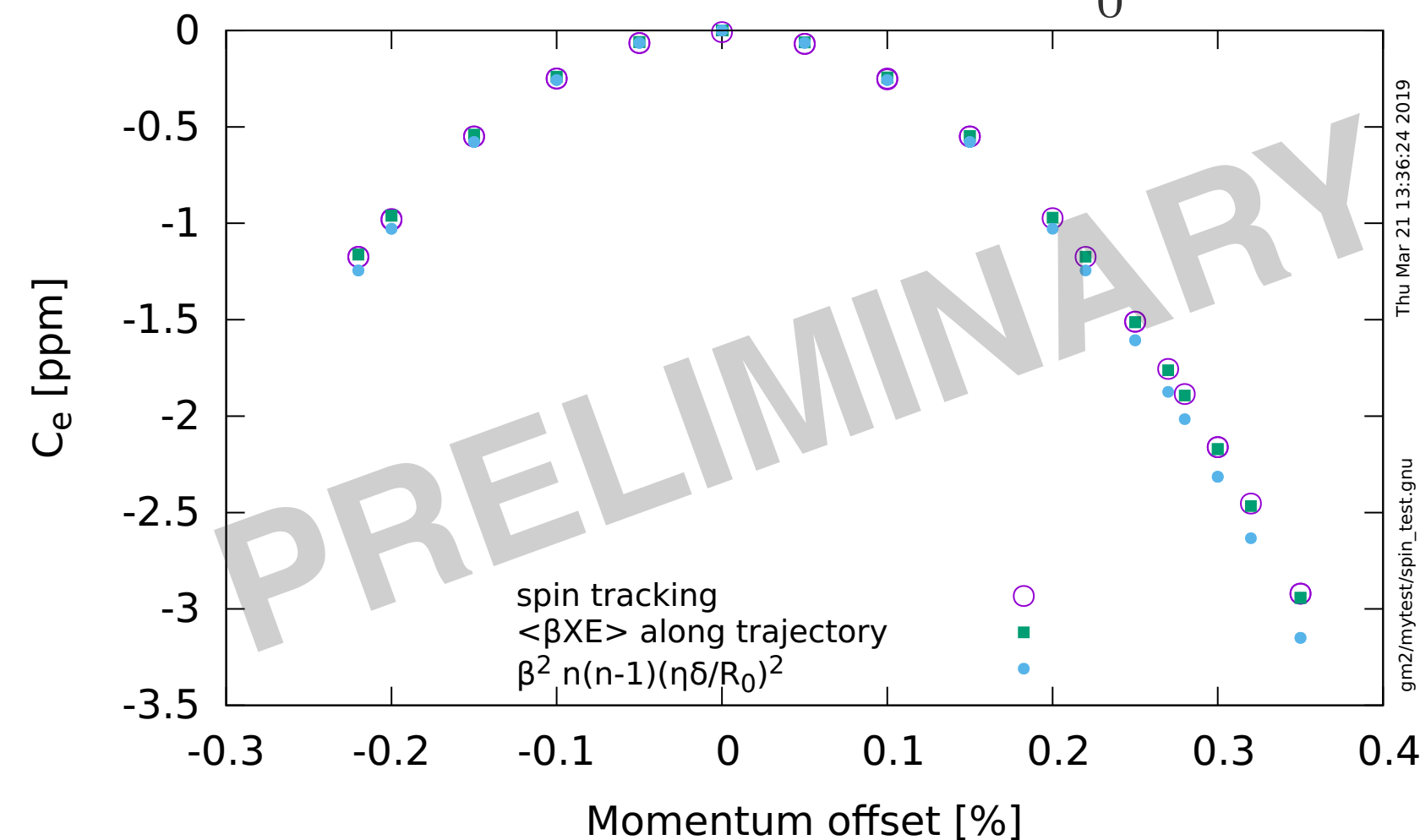
- Full expression for $\vec{\omega}_a$:

$$\vec{\omega}_a = \vec{\omega}_S - \vec{\omega}_C = -\frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

- Choose $\gamma = 29.3$ ($p_\mu = 3.094$ GeV/c)

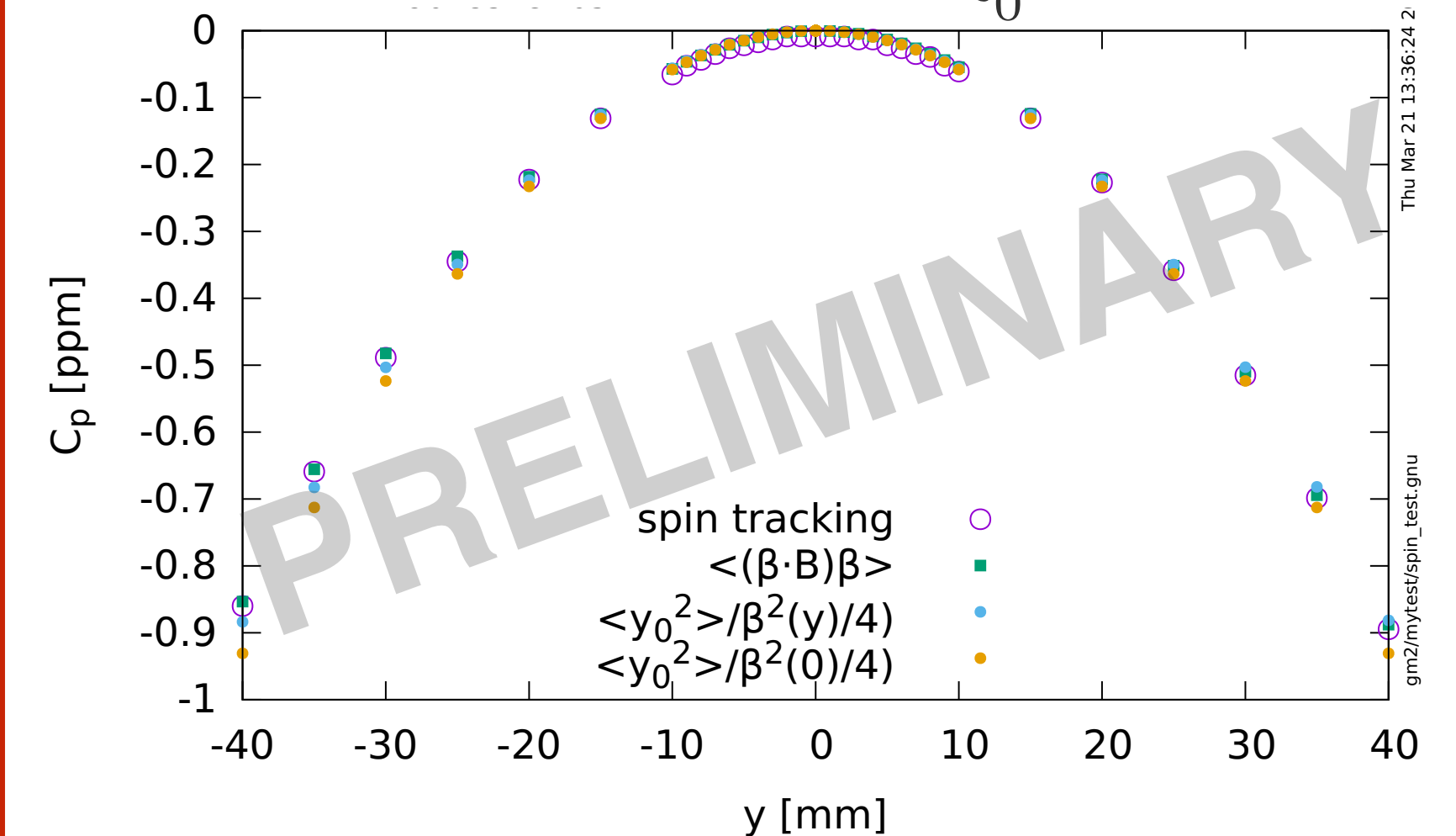
Not all μ^+ at this $\gamma \rightarrow$ **E-field correction**

$$C_E = -2n(1-n)\beta^2 \frac{\langle x_e^2 \rangle}{R_0^2}$$



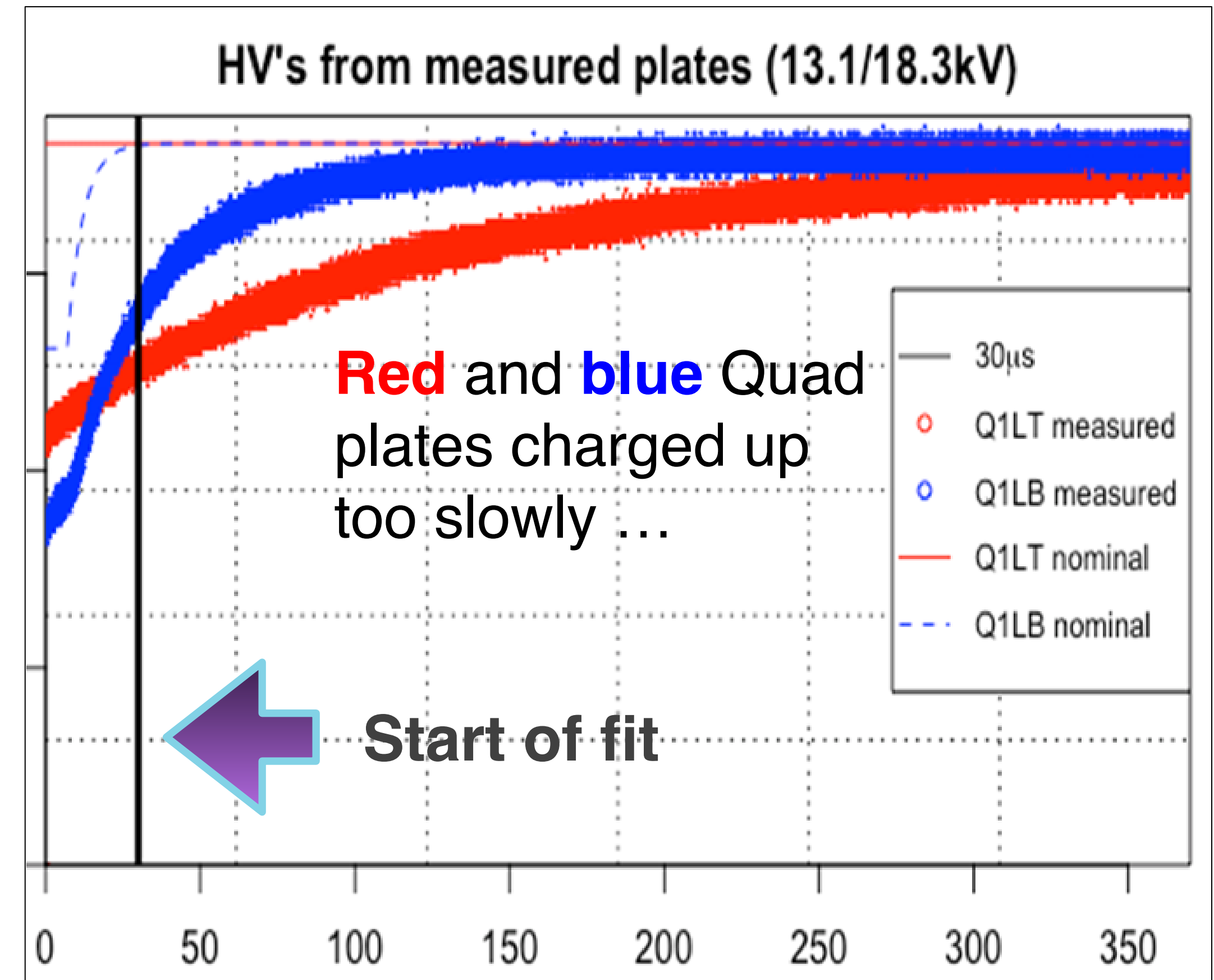
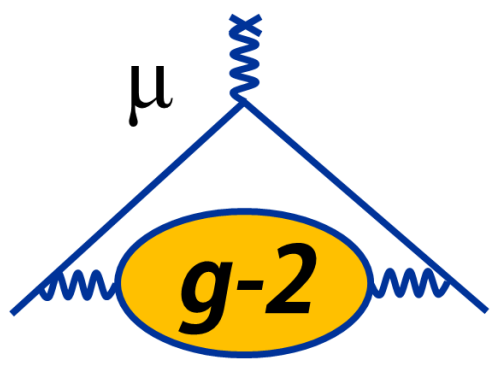
Vertical beam oscillations \rightarrow **pitch correction**

$$C_p = -\frac{n}{4} \frac{\langle y^2 \rangle}{R_0^2}$$



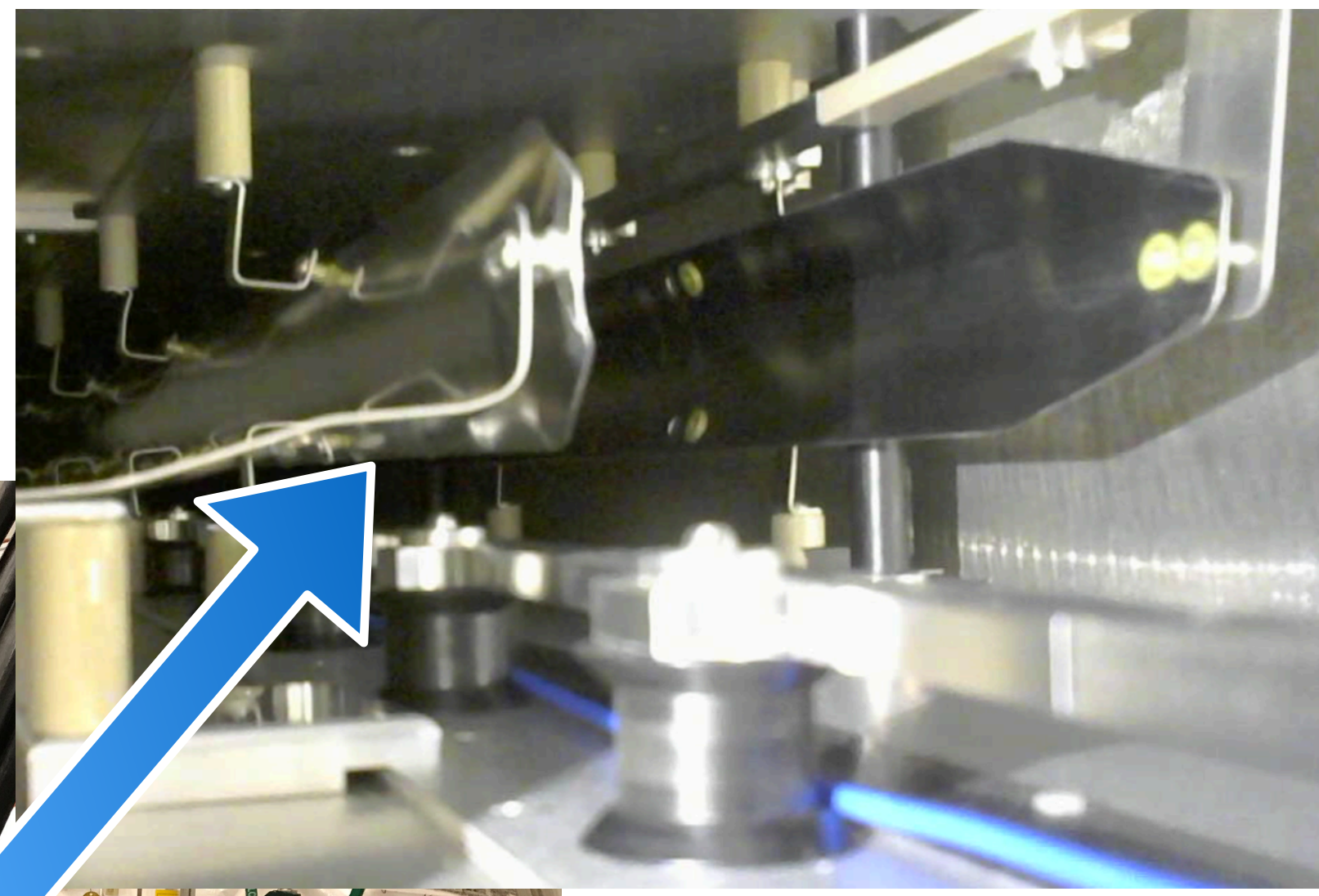
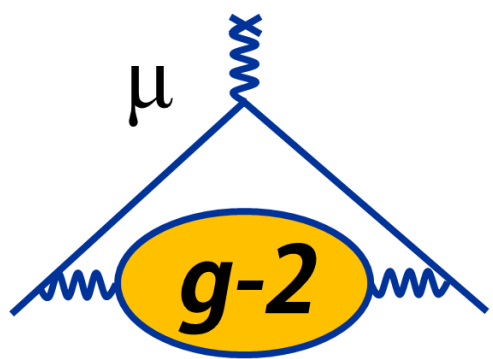
Quad Challenges (Run 1)

- 2 of 32 HV resistors on quad plates were flawed => did not stabilize in time for “fit start time”
- Mean of vertical muon distribution moves down by 0.6 mm
 - Investigating impact on ω_a (calculations, systematic measurements)
- Problem was fixed for Run 2

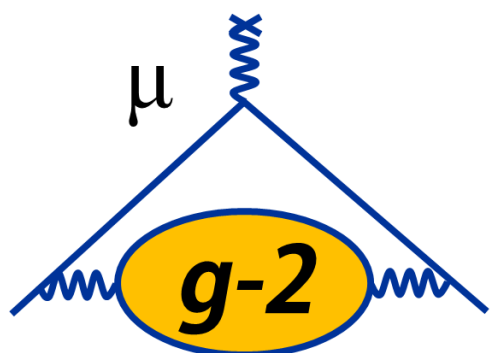


How is a Kick Made?

- A **charging power supply** charges up a **capacitor bank** to 700 V
- Capacitors are **discharged** through a transformer into a **Blumlein** (a HV capacitor up to 55 kV)
- Current in **Blumlein** is discharged into a **resistive load** ($Z = 12.5 \Omega$)
- Current delivered to plates, producing a **~ 200 G magnetic field**, rotating muon's momentum vector

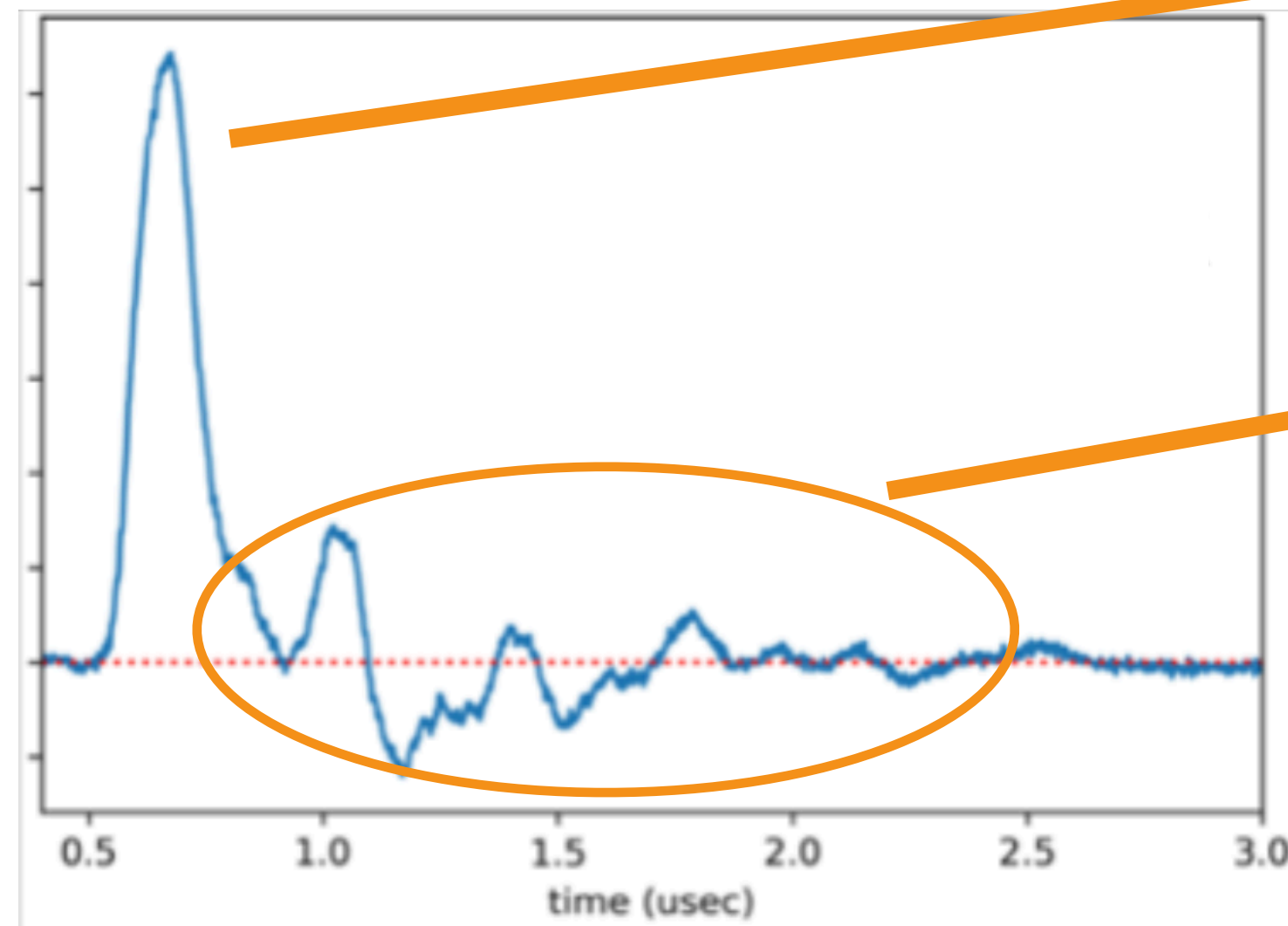


What Affects the Beam Shape?



- **Kicker pulse** strength, shape affects structure of beam
- **Beam width** affected by dynamics

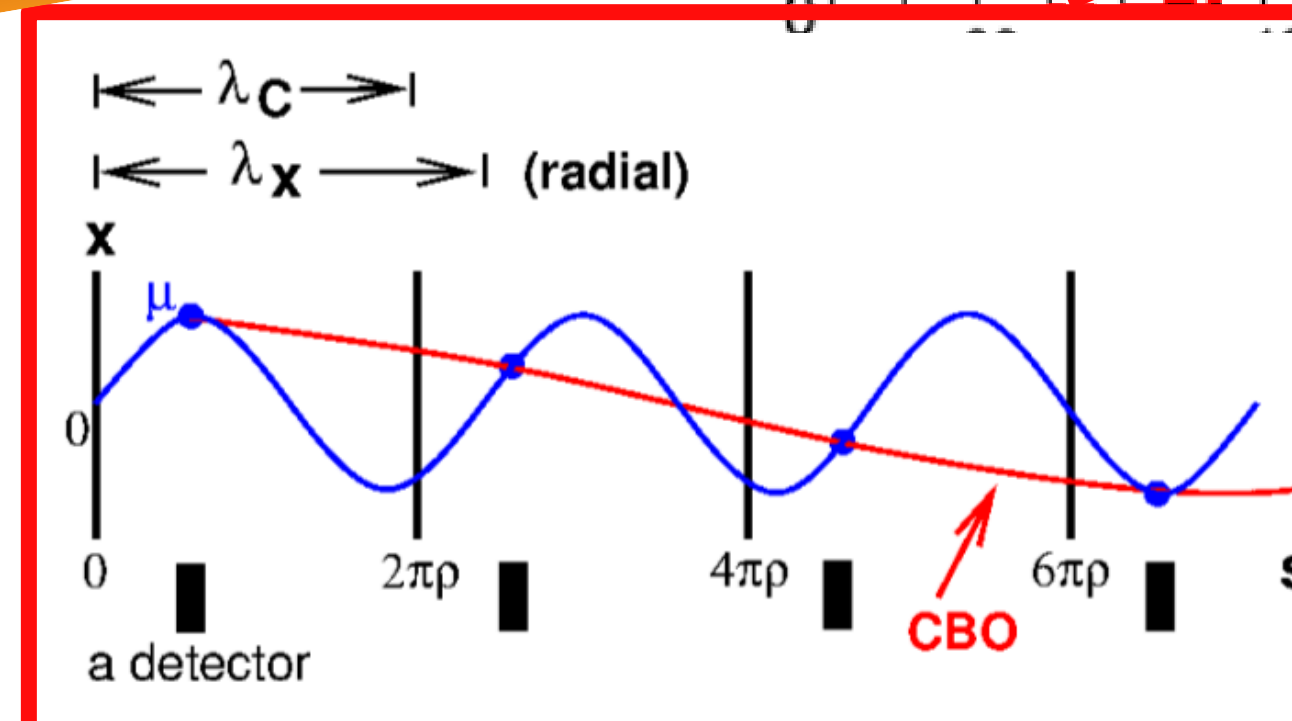
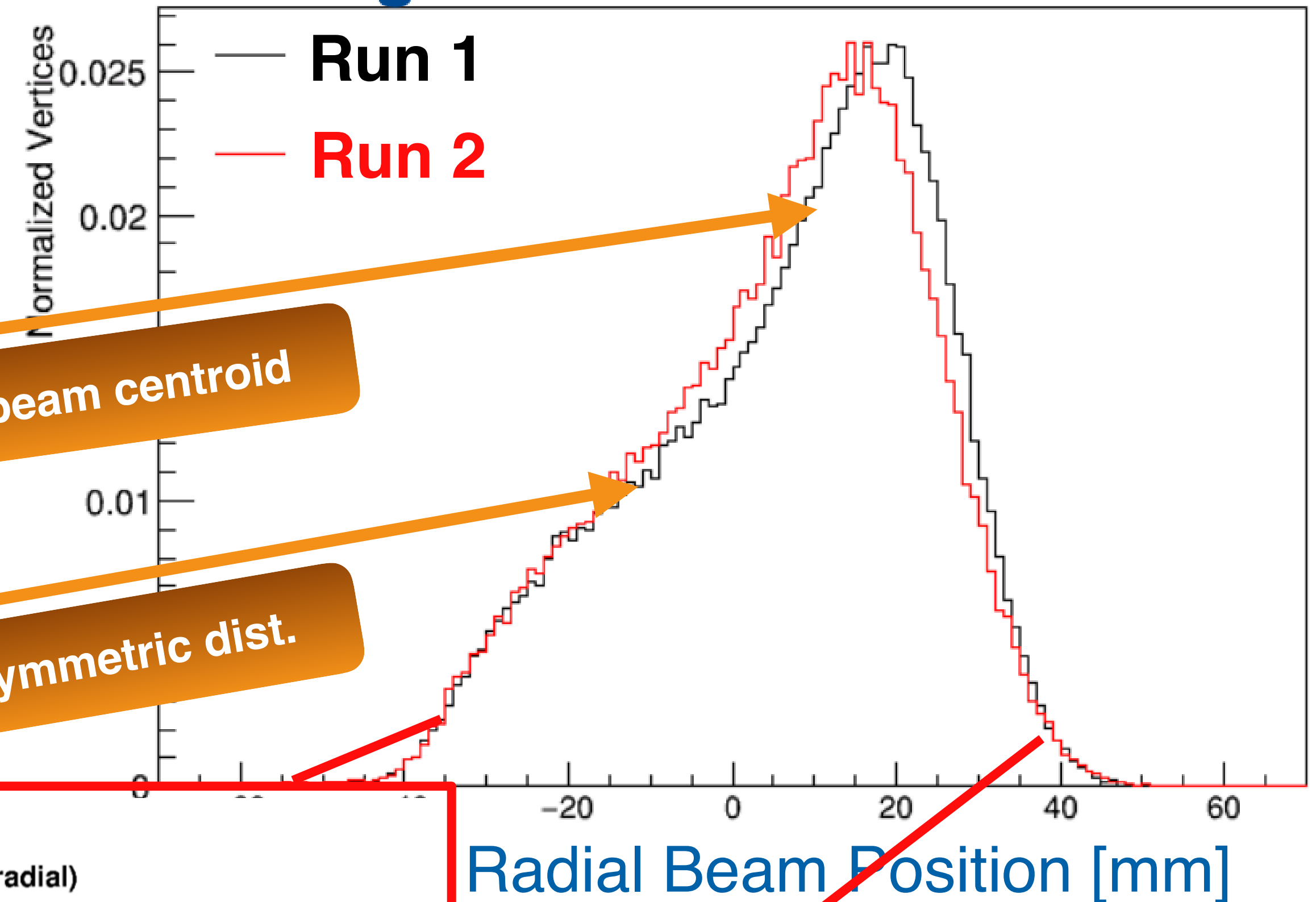
Kicker Pulse



kick strength \rightarrow beam centroid

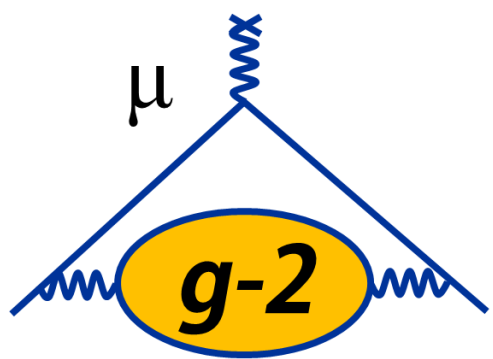
ringing kick \rightarrow asymmetric dist.

Higher kick \rightarrow lower radius

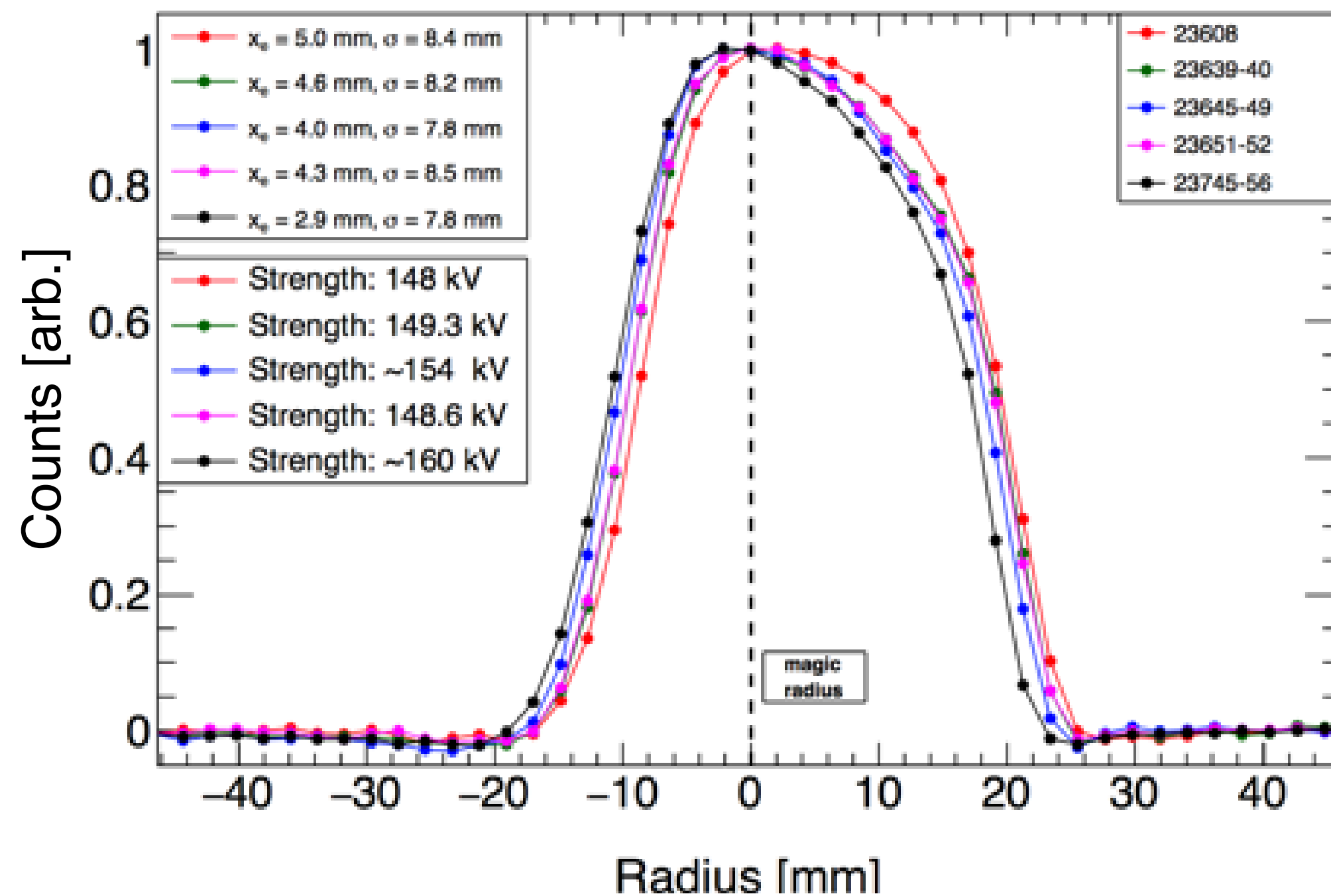


CBO \rightarrow radial width

Kicker Challenges

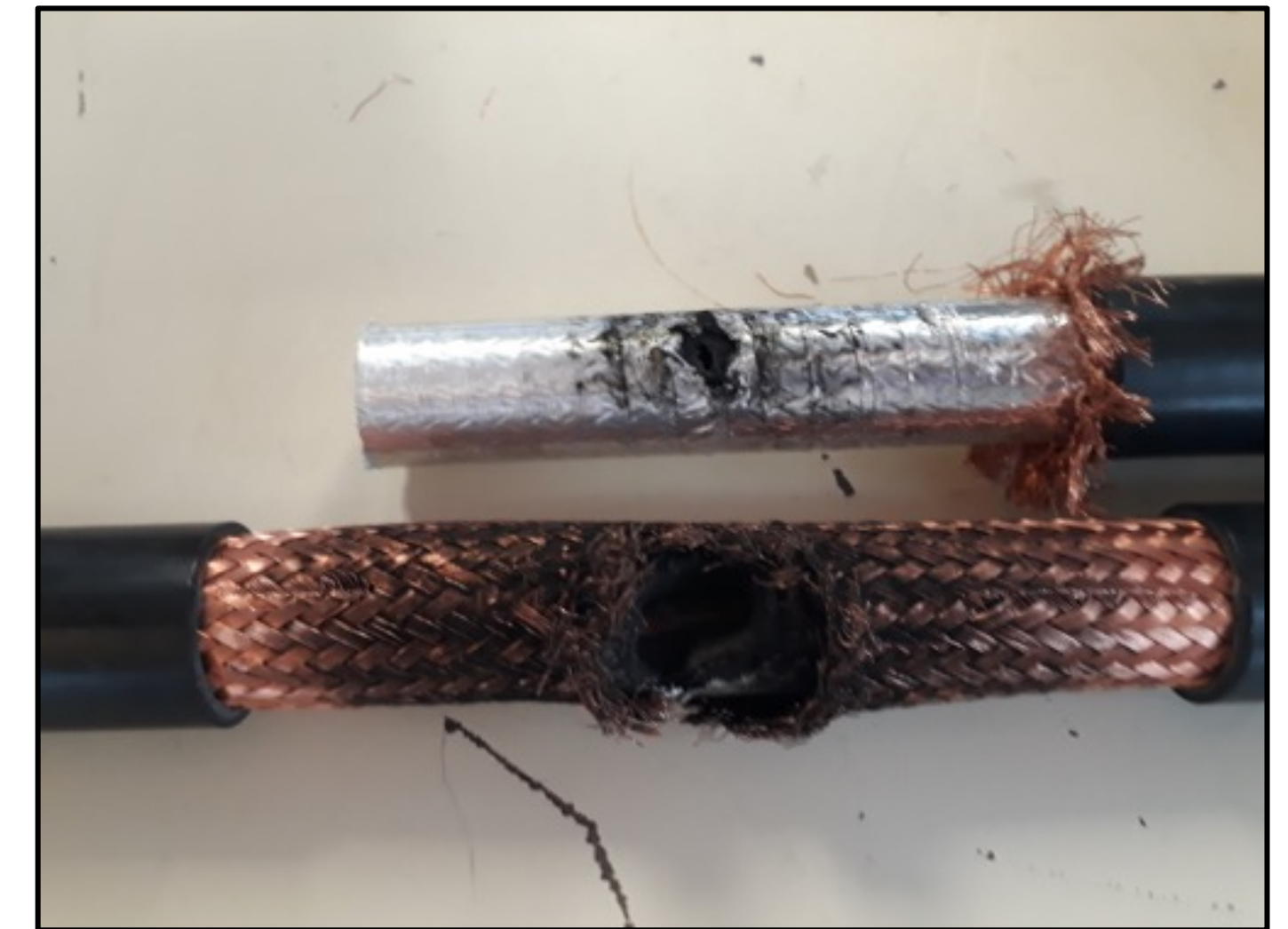


- Need to deflect beam by ~ 11 mrad on the first turn, then turn off (< 149 ns)
- Engineering challenge never fully realized at BNL — and not yet for us



Radial beam distribution at various kick strengths; well contained, but not centered

Take away: we can live with this — we just don't like it



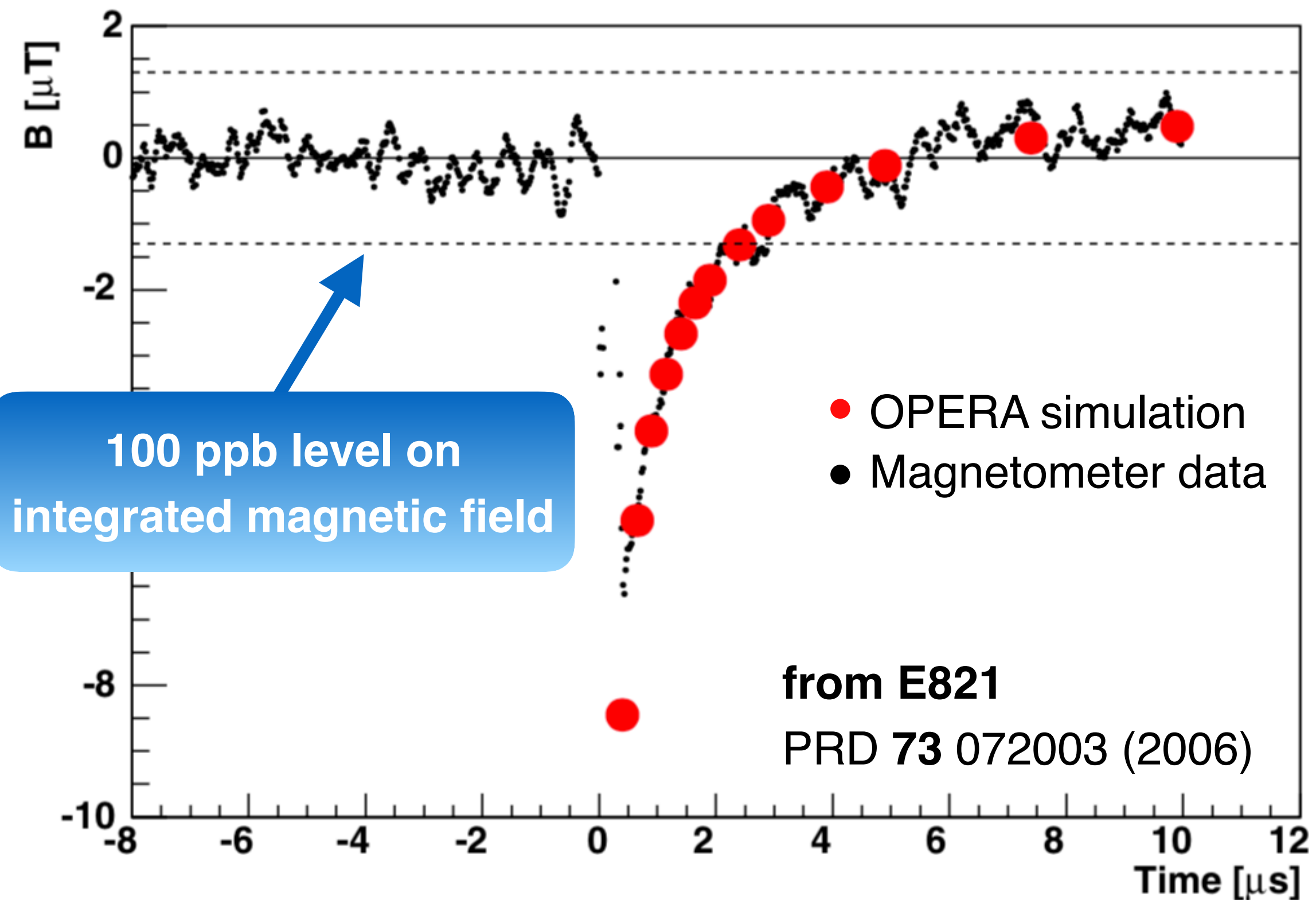
Running at high voltages (~ 50 kV) burns out cables — had to back off

What Drives the ω_a Fit Start Time?

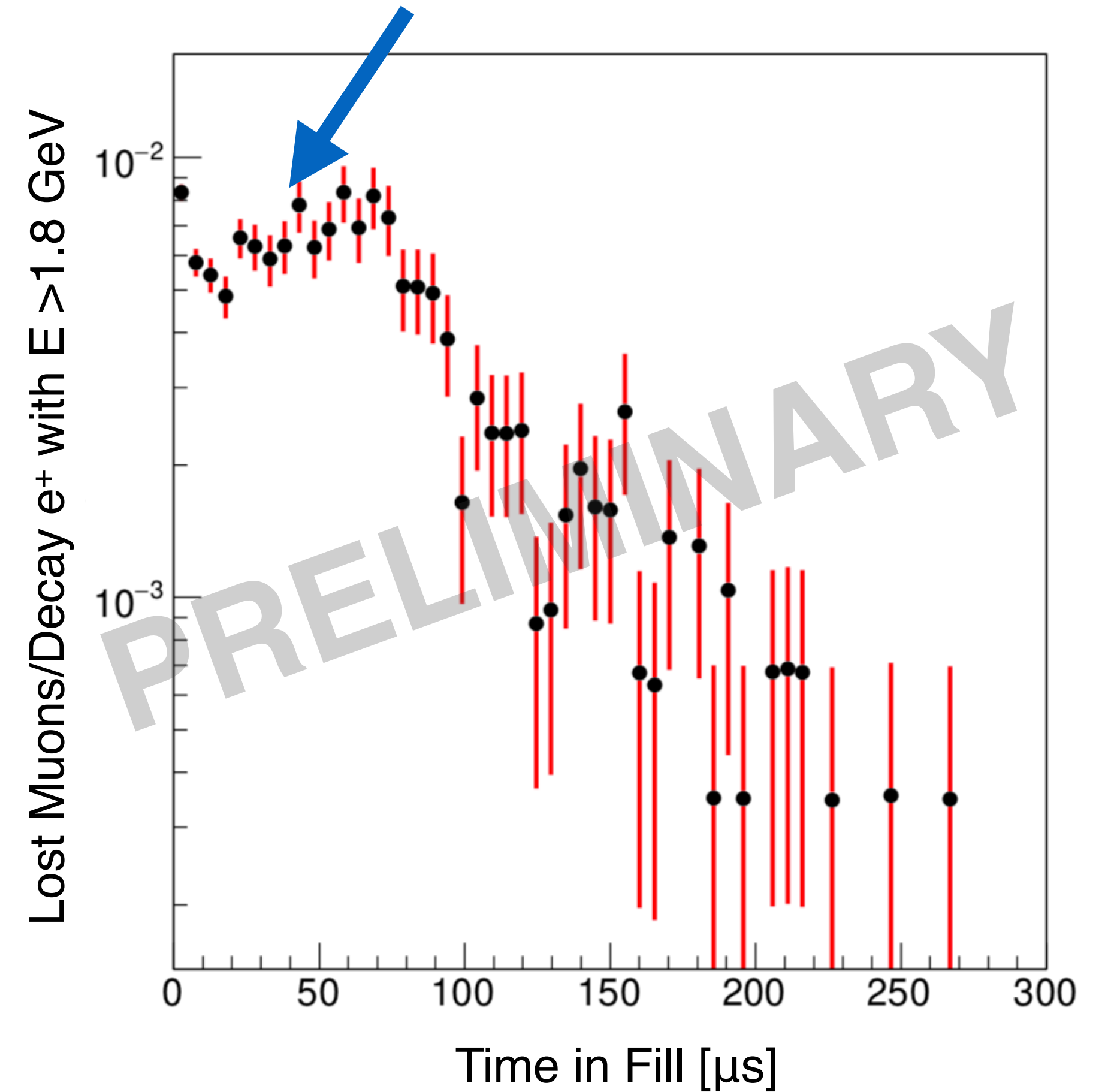


- Start fit window to extract ω_a at $\sim 30 \mu\text{s}$ to avoid:

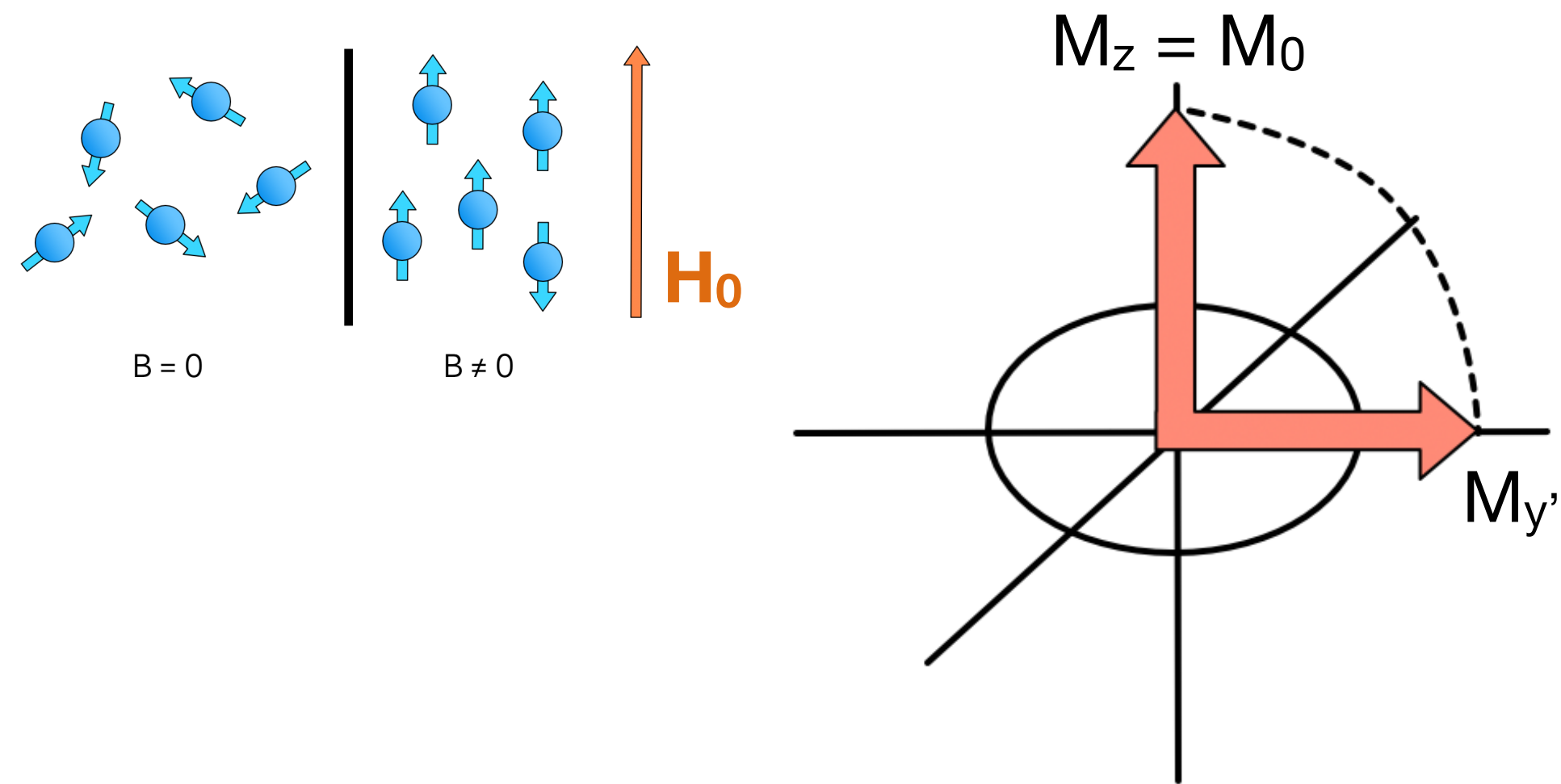
Kicker eddy currents affect the magnetic field



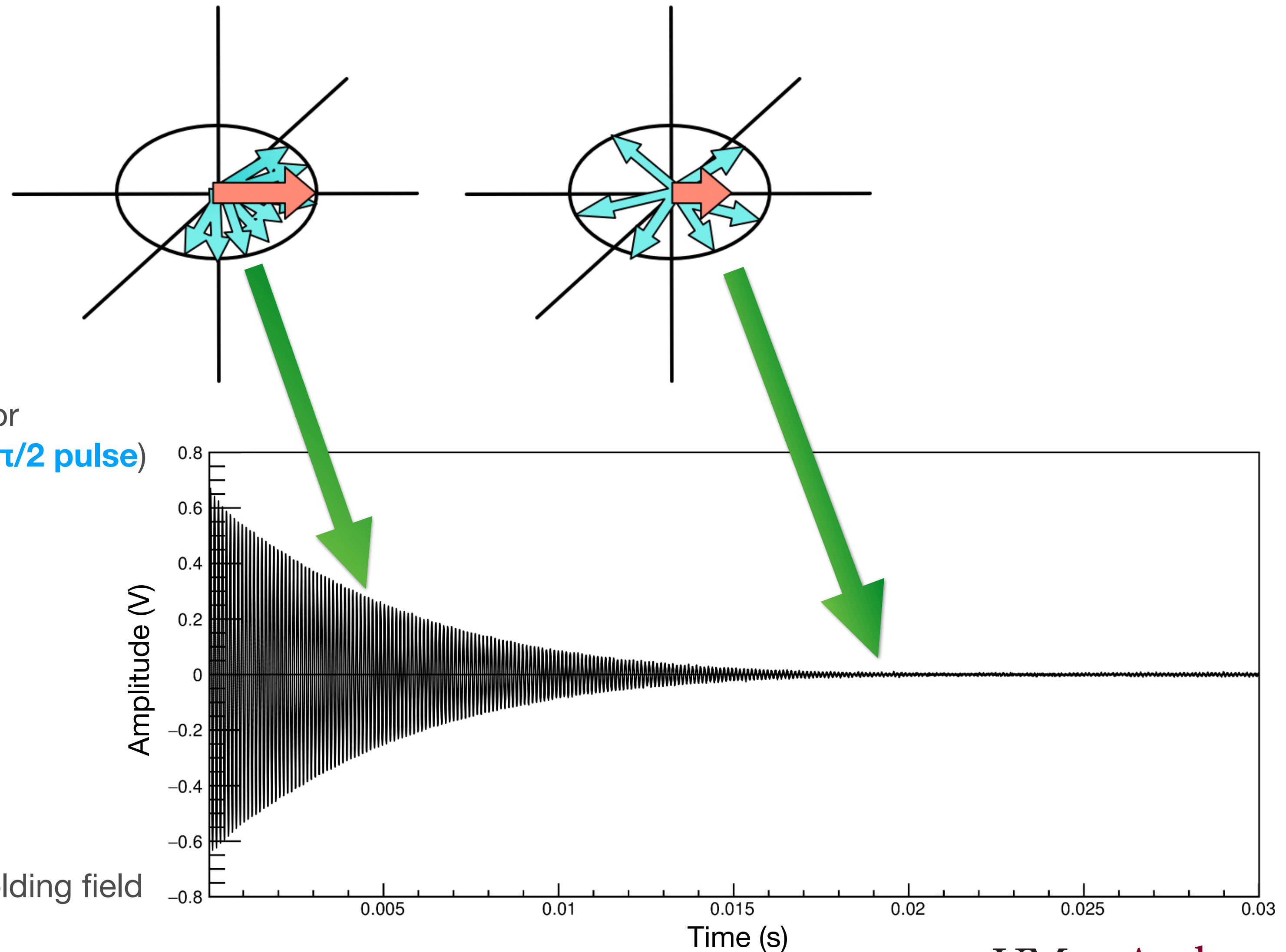
Quad scraping at early times to reduce losses



Pulsed Nuclear Magnetic Resonance



- Apply an RF pulse for a short time to the sample at Larmor frequency — tips spins perpendicular to external B field (**$\pi/2$ pulse**)
- Spin precession induces an EMF in the pickup coil
 - So-called **Free-Induction Decay (FID)**
- Decay of signal driven by:
 - Spin-spin interactions (dephasing) (pure T_2)
 - Field inhomogeneities (T_2^*)
 - Simultaneously, spins relax back to alignment with holding field (spin-lattice relaxation, T_1)



Magnetic Circuits

$$\mathcal{E} = \oint \vec{f}_s \cdot d\vec{\ell} = V = IR$$

Can write a similar equation for magnets

$$\mathcal{F} = \oint \vec{H} \cdot d\vec{\ell} = NI$$

Magnetomotive Force (mmf)

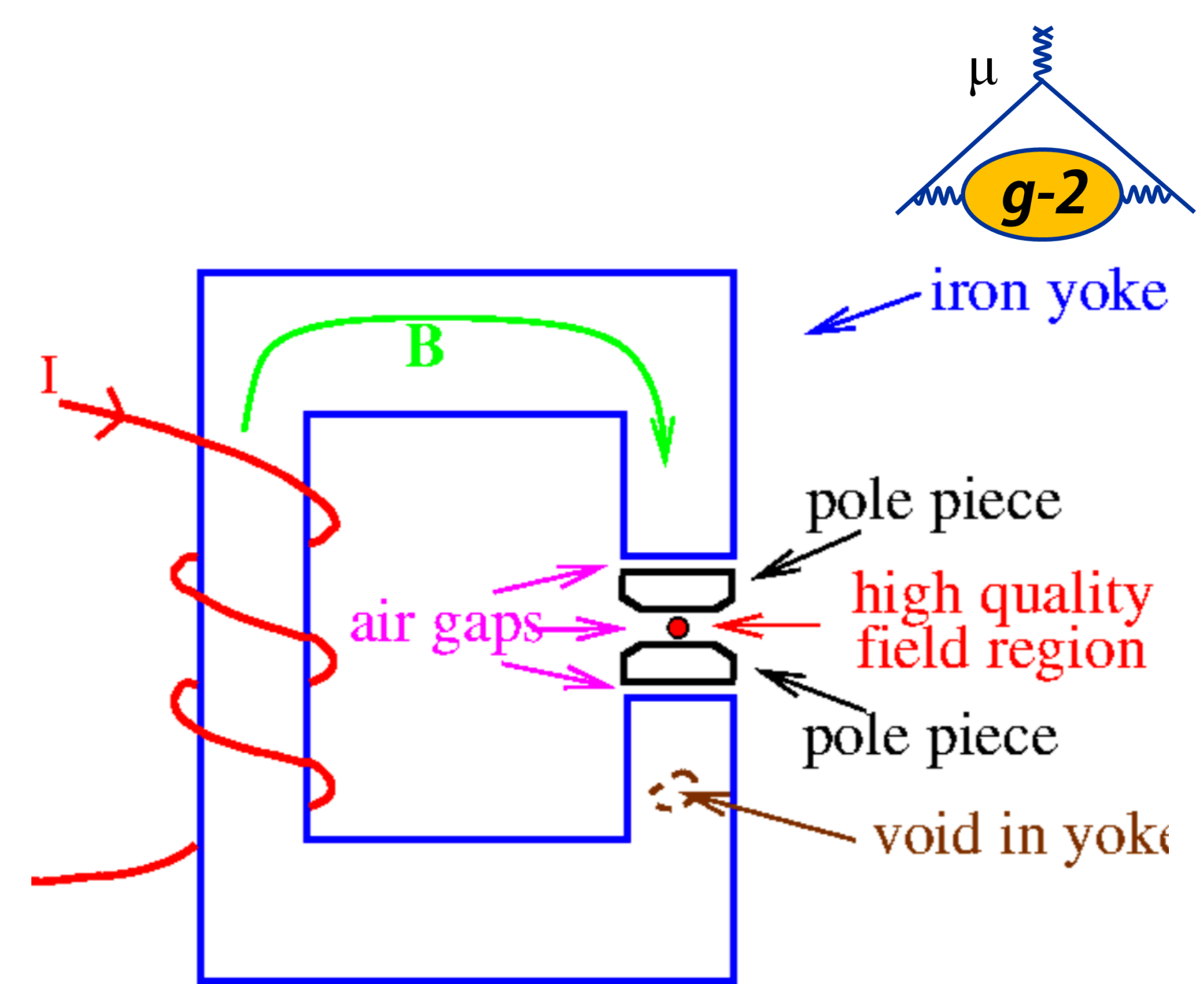
$$\vec{B} = \mu_0 (1 + \chi_m) \vec{H} = \mu \vec{H}$$

Rewrite H in terms of B

$$\Phi = \vec{B} \cdot \vec{A} = \mu \vec{H} \cdot \vec{A}$$

Consider magnetic flux

$$\Phi \oint \frac{d\ell}{\mu A} = \mathcal{F} \Rightarrow \boxed{\mathcal{R} = \oint \frac{d\ell}{\mu A} = \frac{\mathcal{F}}{\Phi}}$$



Magnetic Reluctance

- Analogous to resistance in an electrical circuit

$$\boxed{V = IR \Leftrightarrow \mathcal{F} = \Phi \mathcal{R}}$$

- Current flows along a path of least resistance while field lines will take a path of least reluctance
- While the emf drives electric charges (Ohm's Law), the mmf "drives" magnetic field lines (Hopkinson's Law)

Magnet Anatomy

- For E821, Gordon Danby had a brilliant magnet design

$B = 1.45 \text{ T}$ (~5200 A)

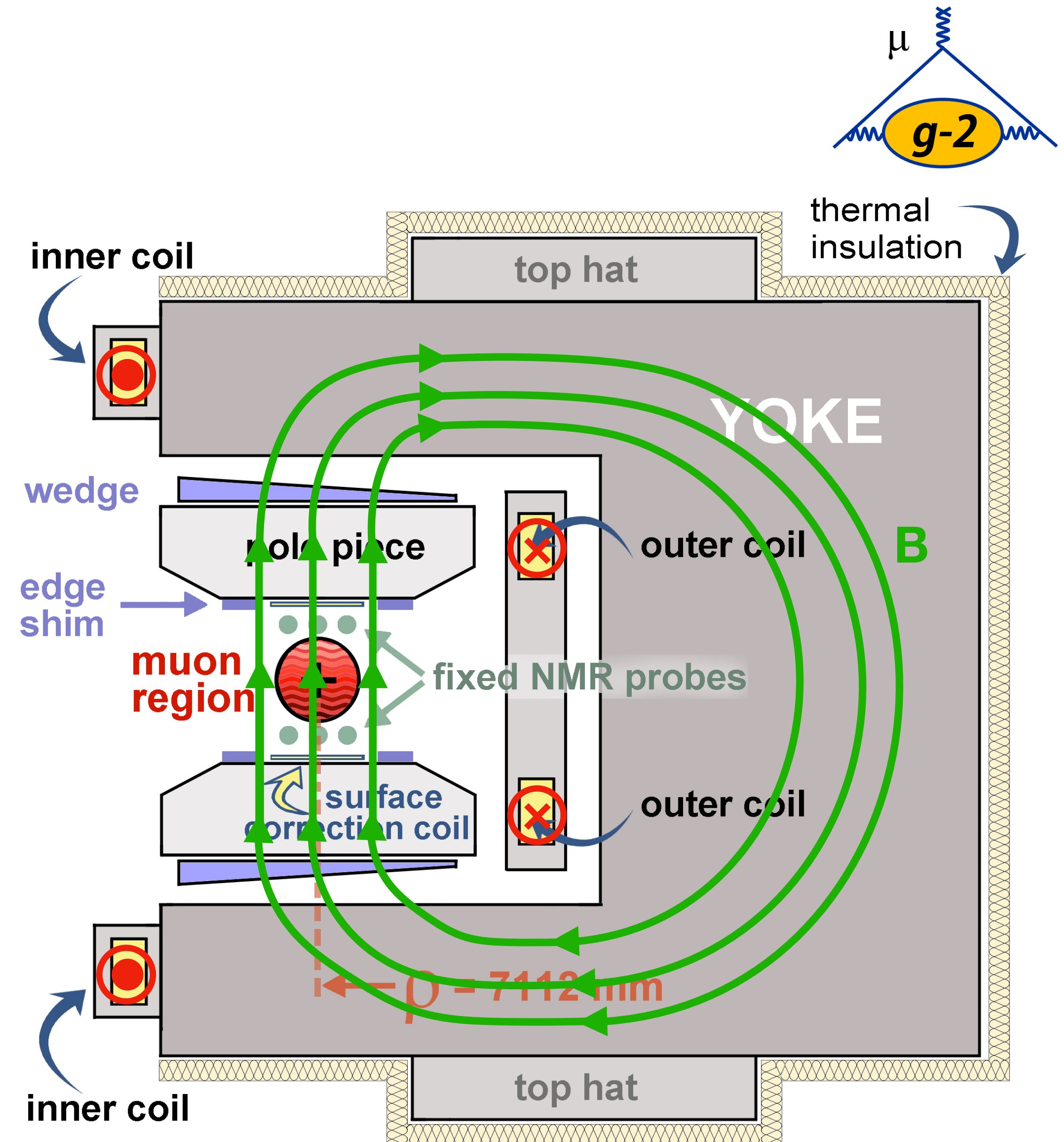
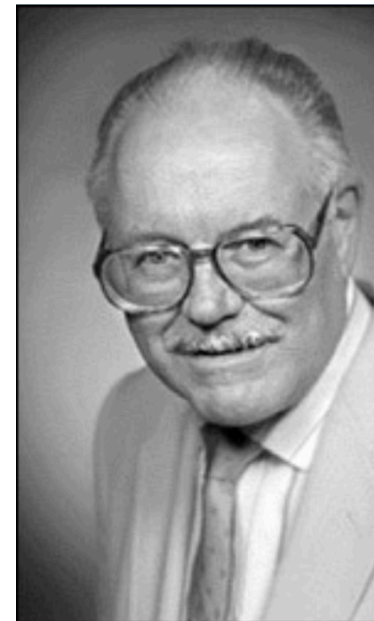
- Non-persistent current: fine-tuning of field in real time

12 C-shaped yokes

- 3 upper and 3 lower poles per yoke
- 72 total poles

Shimming knobs

- Pole separation determines field: pole tilts, non-flatness affect uniformity
- Top hats (30 deg effect, dipole)
- Wedges (10 deg effect, dipole, quadrupole)
- Edge shims (10 deg effect, dipole, quadrupole, sextupole)
- Laminations (1 deg effect, dipole, quadrupole, sextupole)
- Surface coils (360 deg effect, quadrupole, sextupole,...)

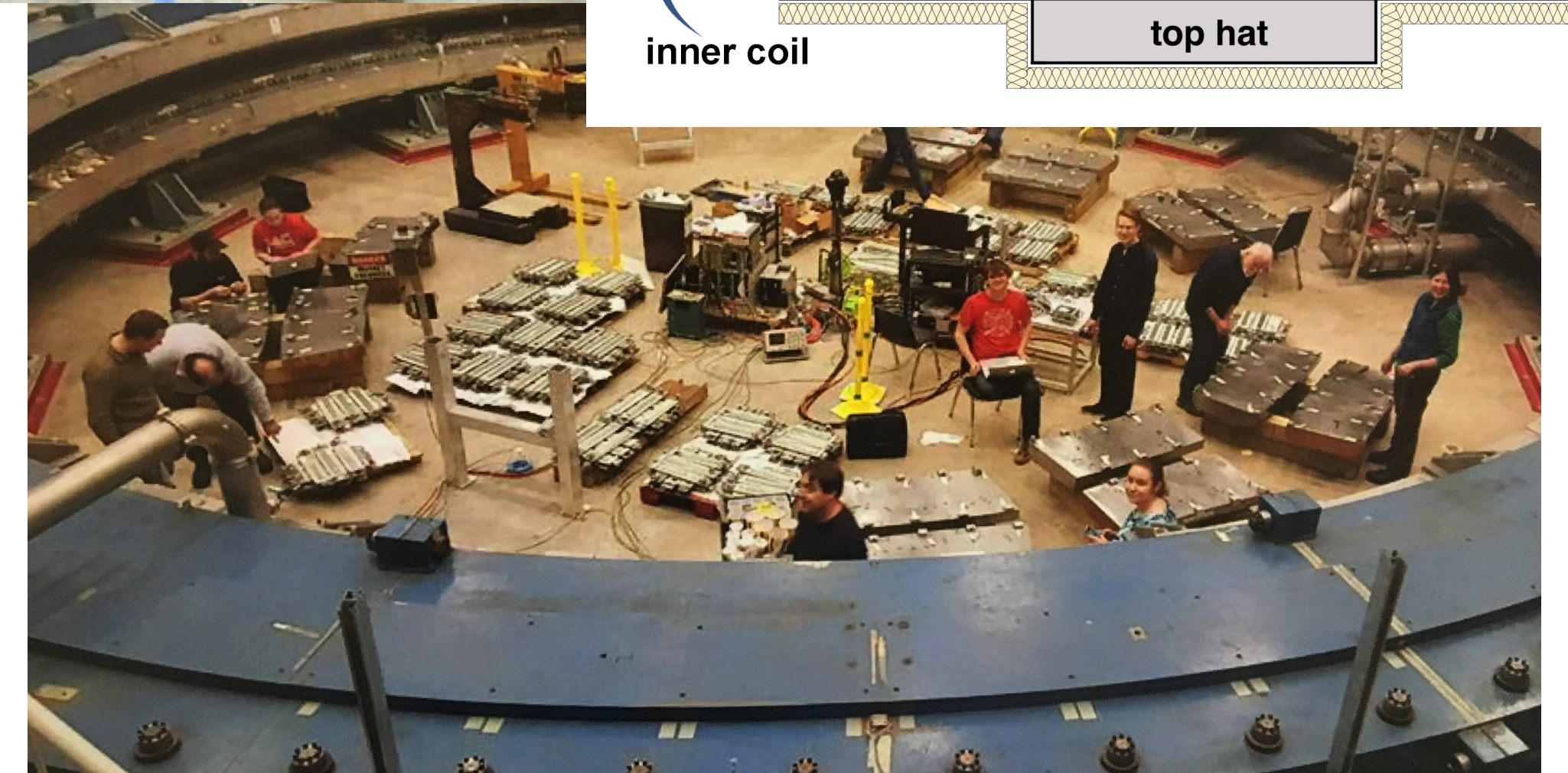
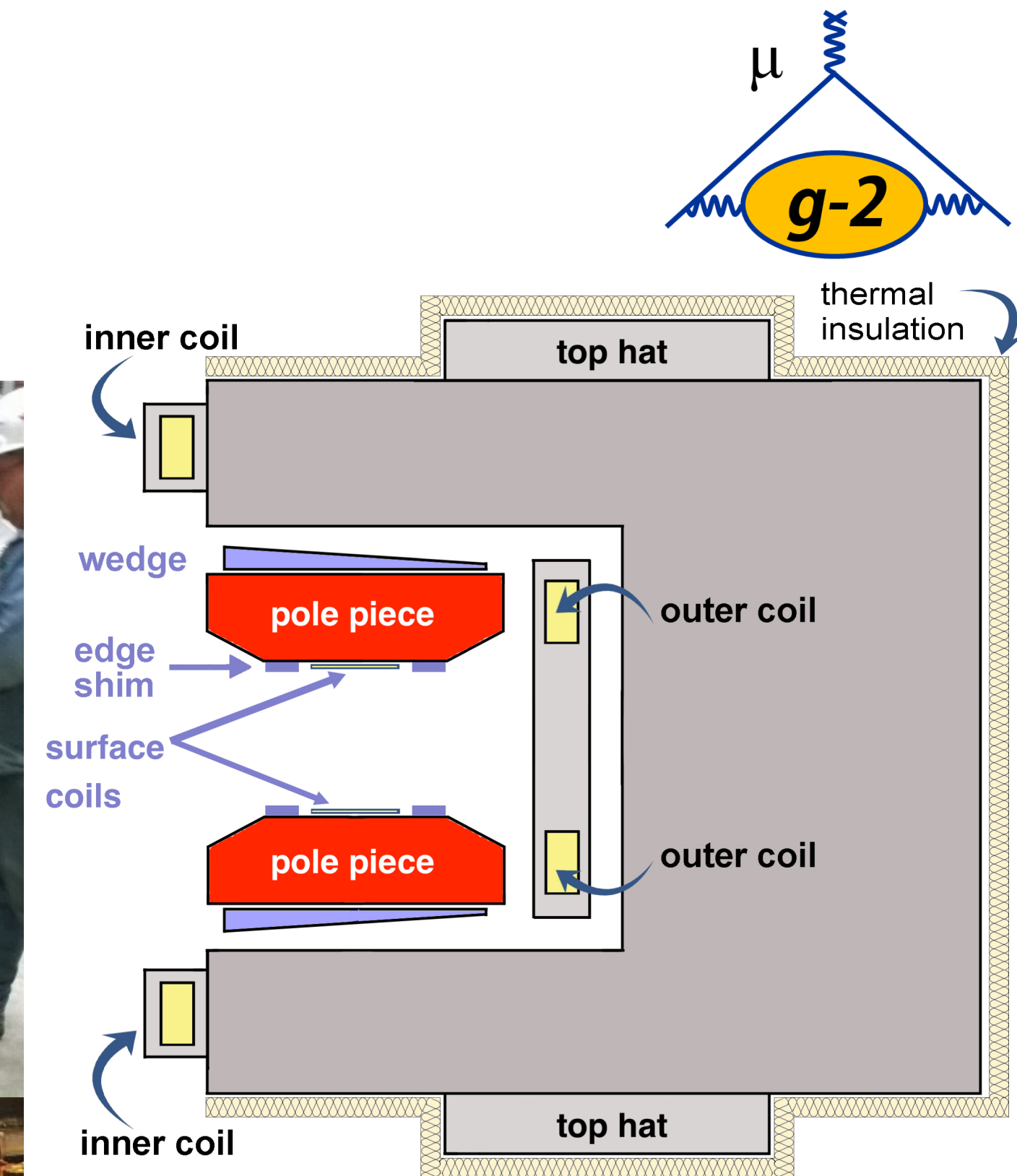


Current direction indicated by red markers

UMassAmherst

Optimizing the Dipole Moment

- Want to optimize the vertical component of the field
- Step and tilt discontinuities in pole surfaces yield large variations in the field
- To reduce/remove such effects, make adjustments to pole feet, which changes the magnet gaps and tilts
 - Use 0.001—0.010” thick shims
 - Requires removal of poles from the ring
- Informed by a computer model that optimizes the pole configurations
 - Requires global continuity between pole surfaces
 - Allows only three adjacent poles to be moved at a time (preserves alignment)



Minimizing the Quad, Sext, Octu Moments

Calibrated shimming knobs

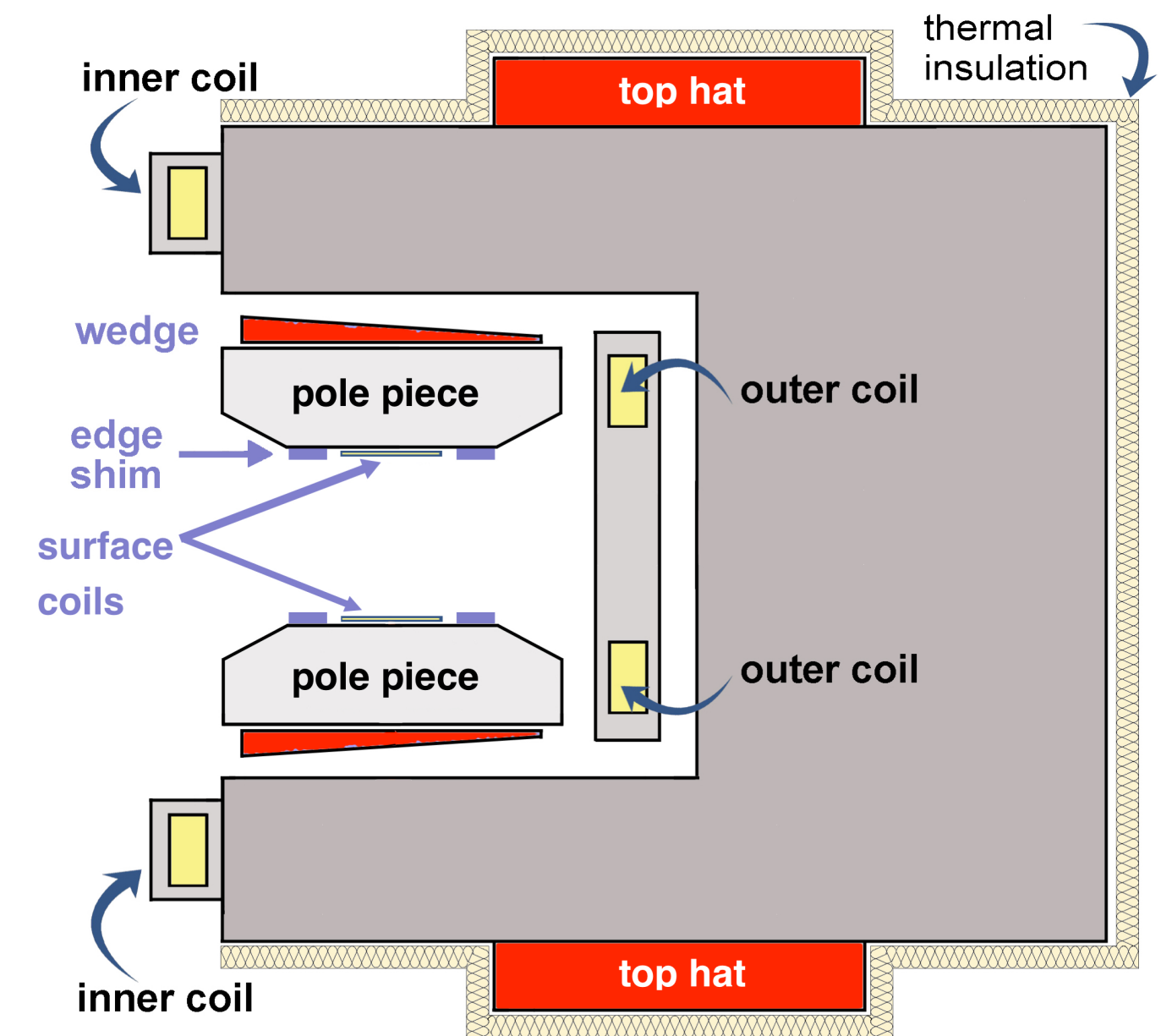
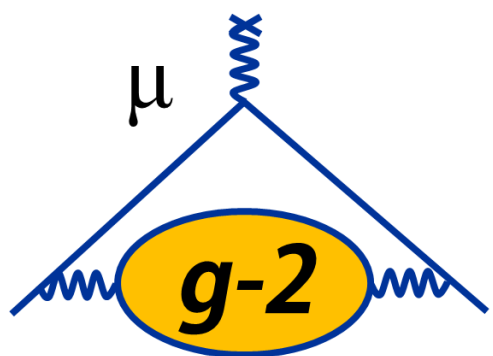
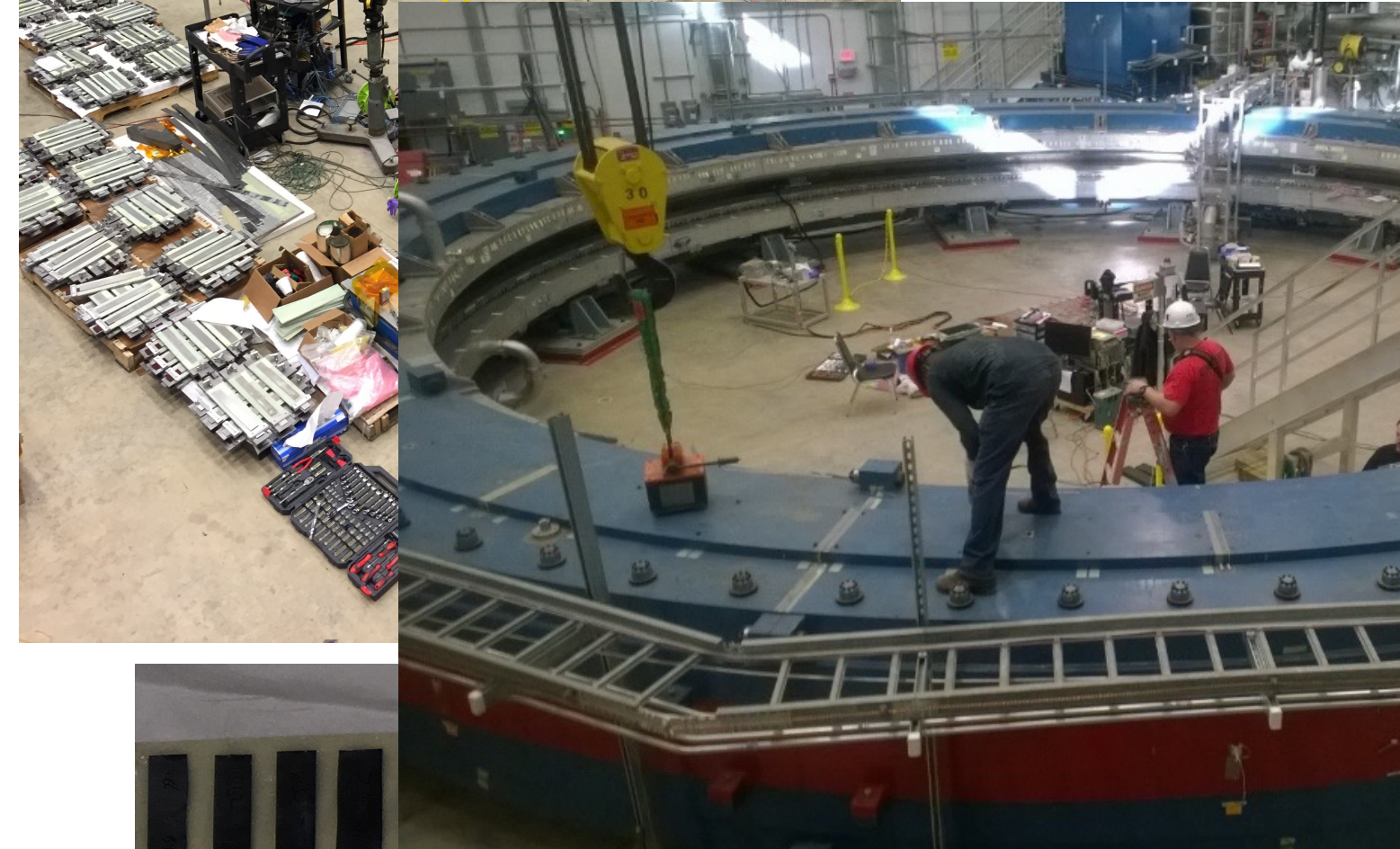
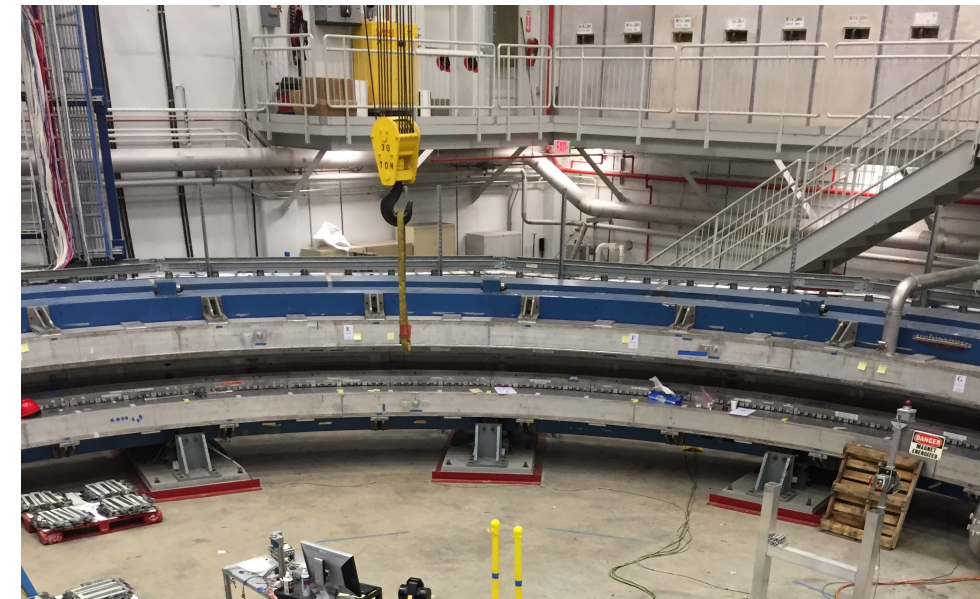
- 48 top hats
- 864 wedges
- ~8400 iron foils (on pole surfaces)

Coarse tuning: top hat & wedge adjustments (**dipole, quadrupole**)

- Least-squares fit to field maps predicts top hat and wedge positions

Fine tuning: iron foils (**quadrupole, sextupole,...**)

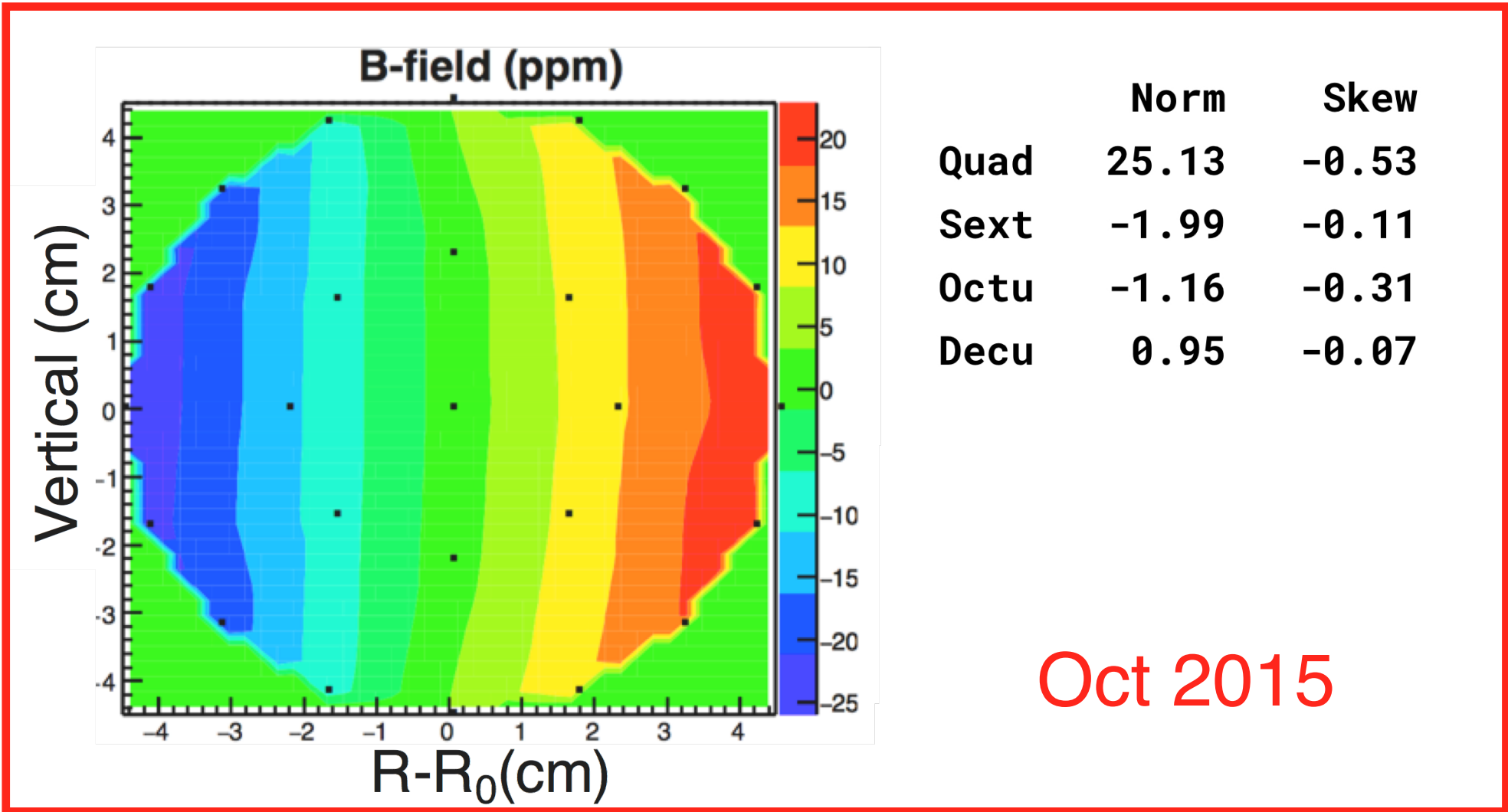
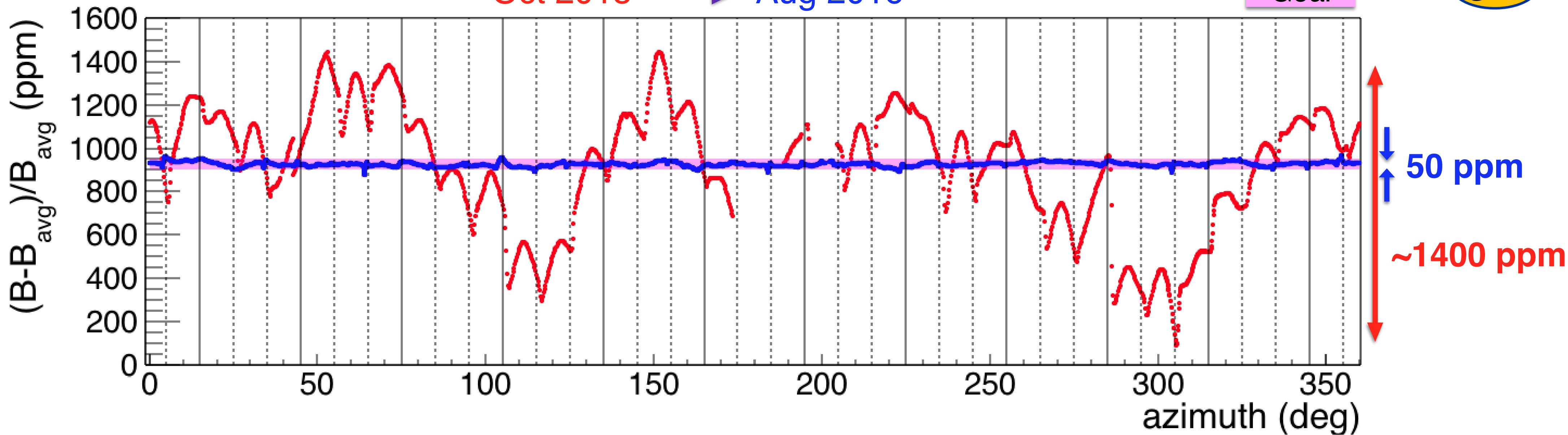
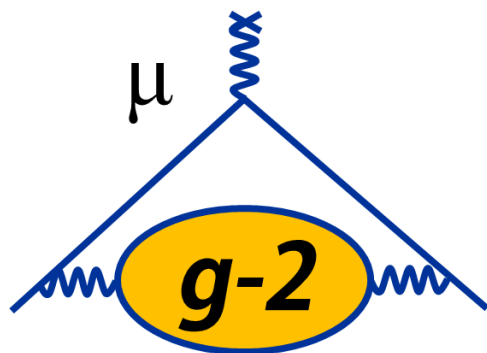
- Modeled as saturated dipoles in 1.45 T field
- Computer code predicts foil width (mass) distribution to fill in the valleys of the field map



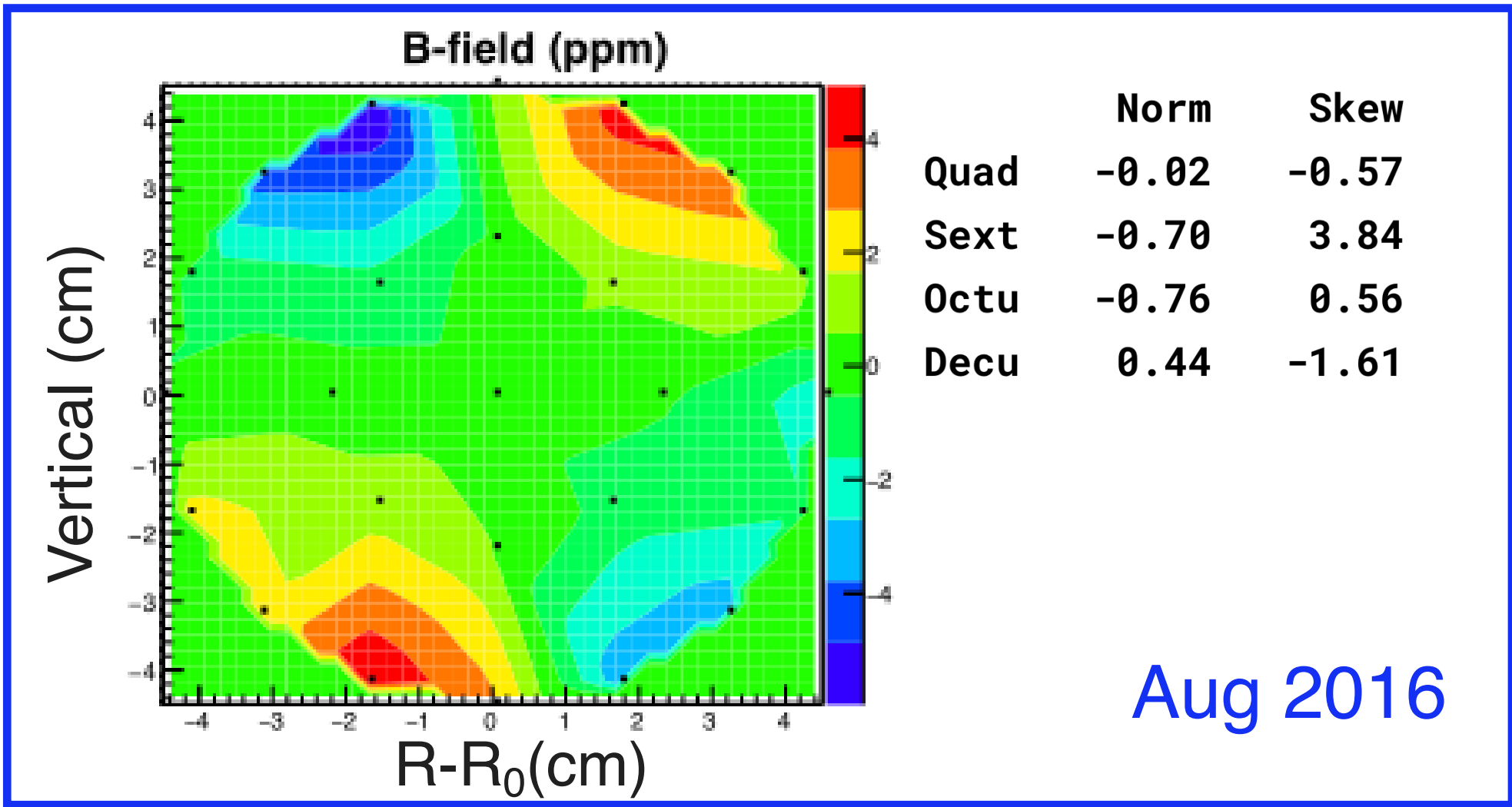
Rough Shimming Results

Oct 2015 → Aug 2016

Goal



Oct 2015

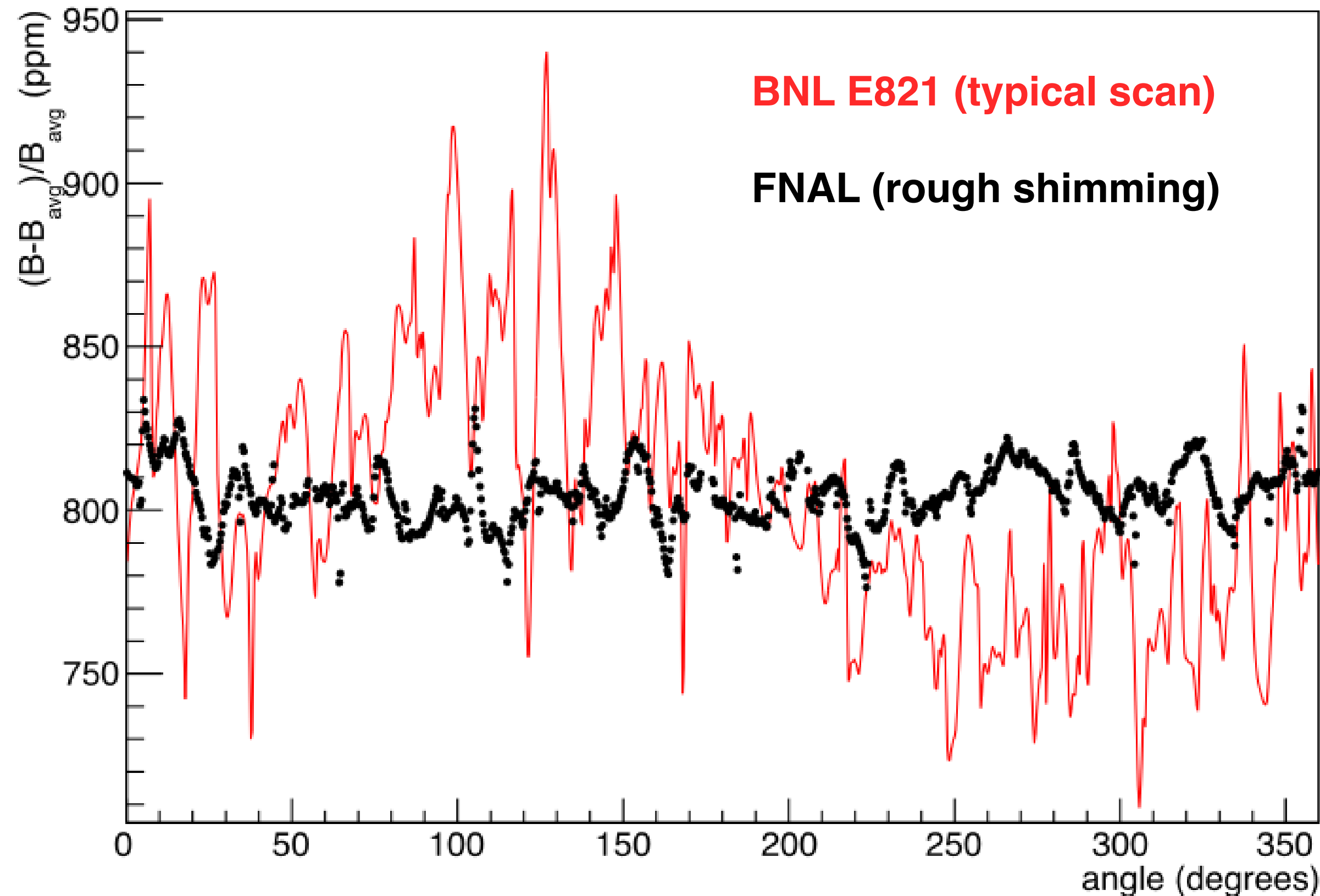


Aug 2016

Magnetic Field Comparison: BNL 821 and FNAL E989



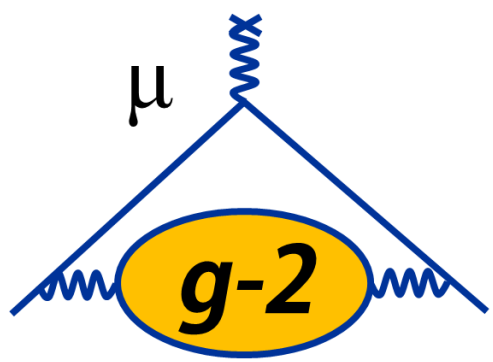
Dipole Vs Azimuth



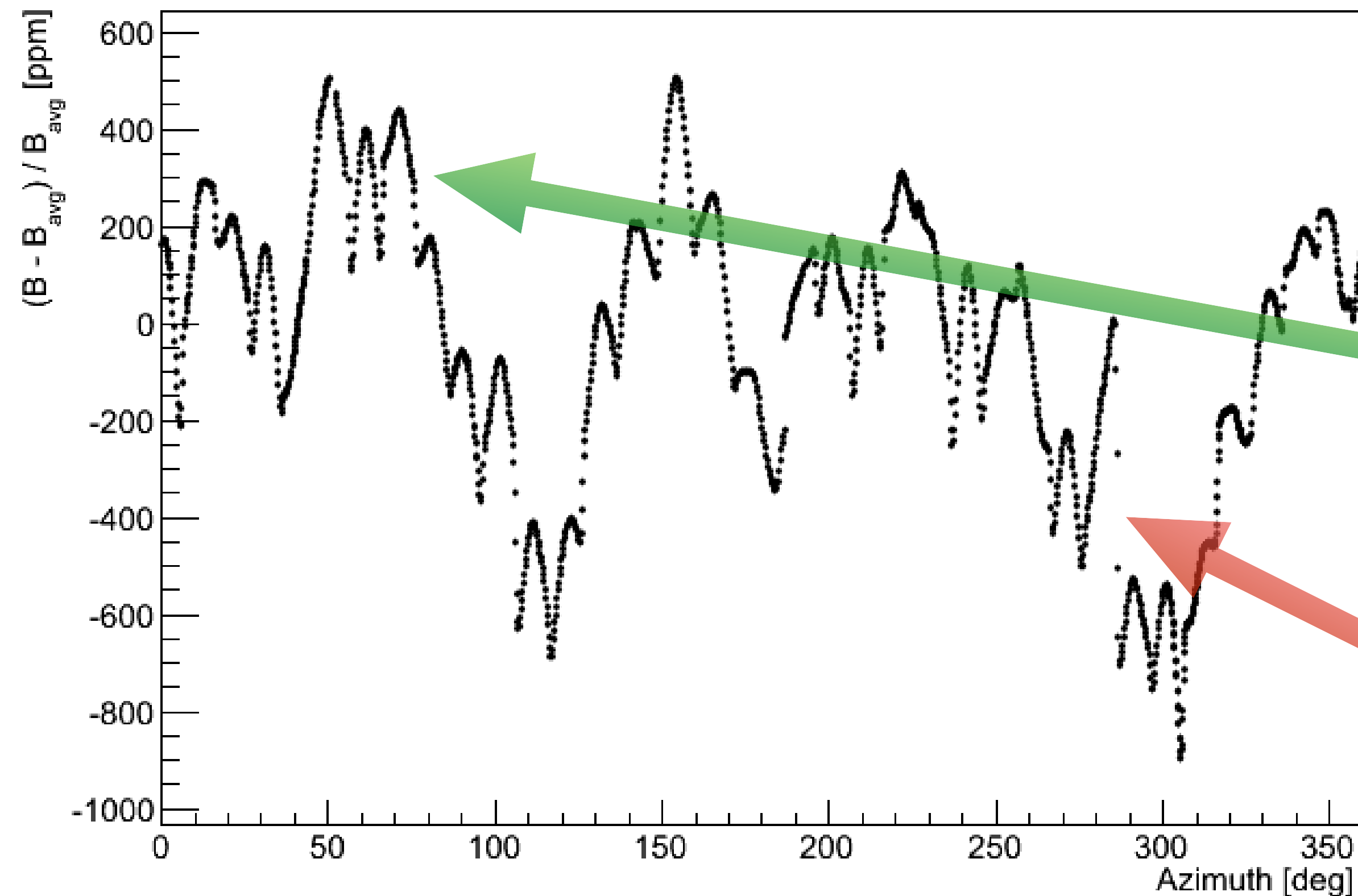
- Laminations very successful in reducing field variations


- **BNL E821: 39 ppm RMS (dipole), 230 ppm peak-to-peak**
- FNAL rough shimming: 10 ppm RMS (dipole), 75 ppm peak-to-peak

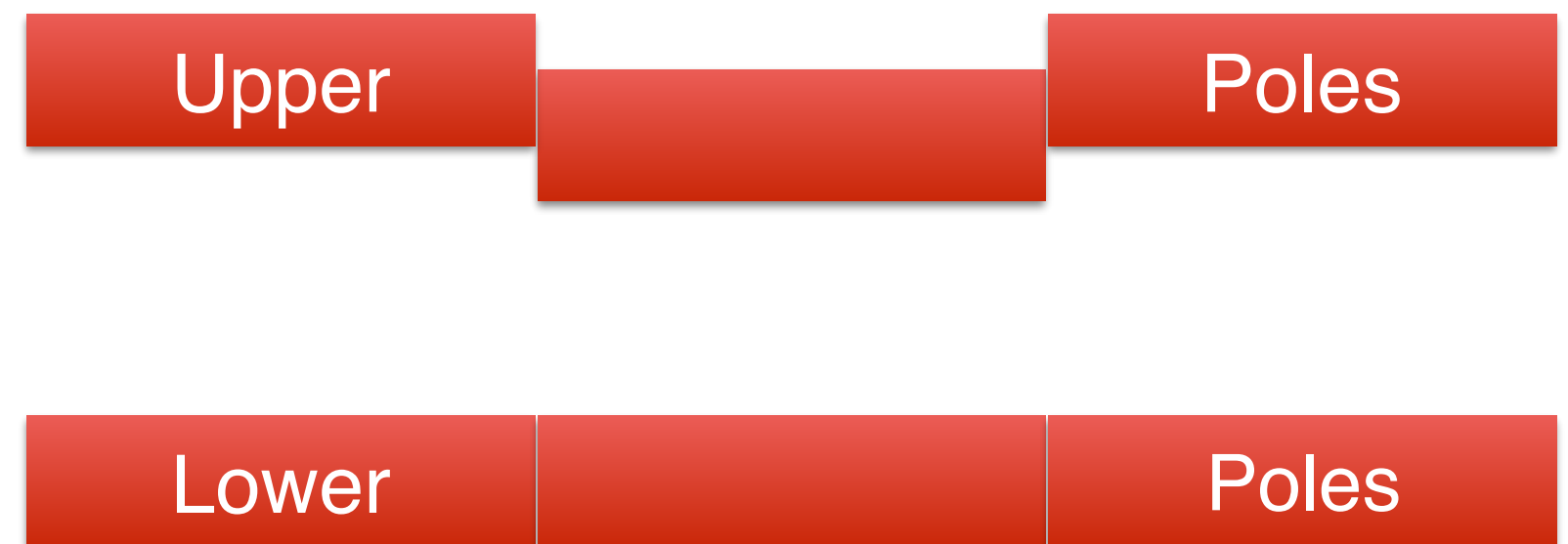
Magnetic Field Variations



First Magnetic Field Map, Oct 14 2015



- Gradual drift from materials, pole gap changes
- 36 pairs of poles \rightarrow 10-degree structure
- Pole shape: 
- Pole-to-pole discontinuities

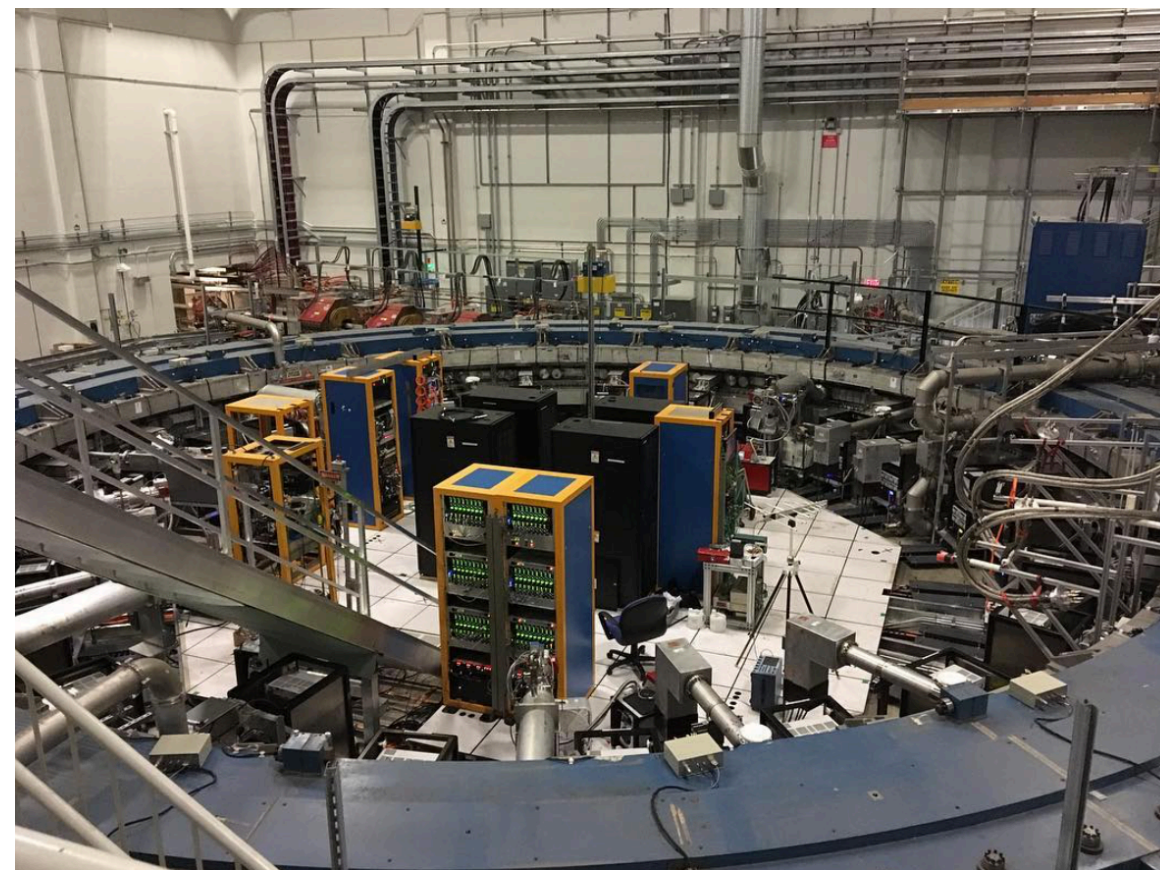
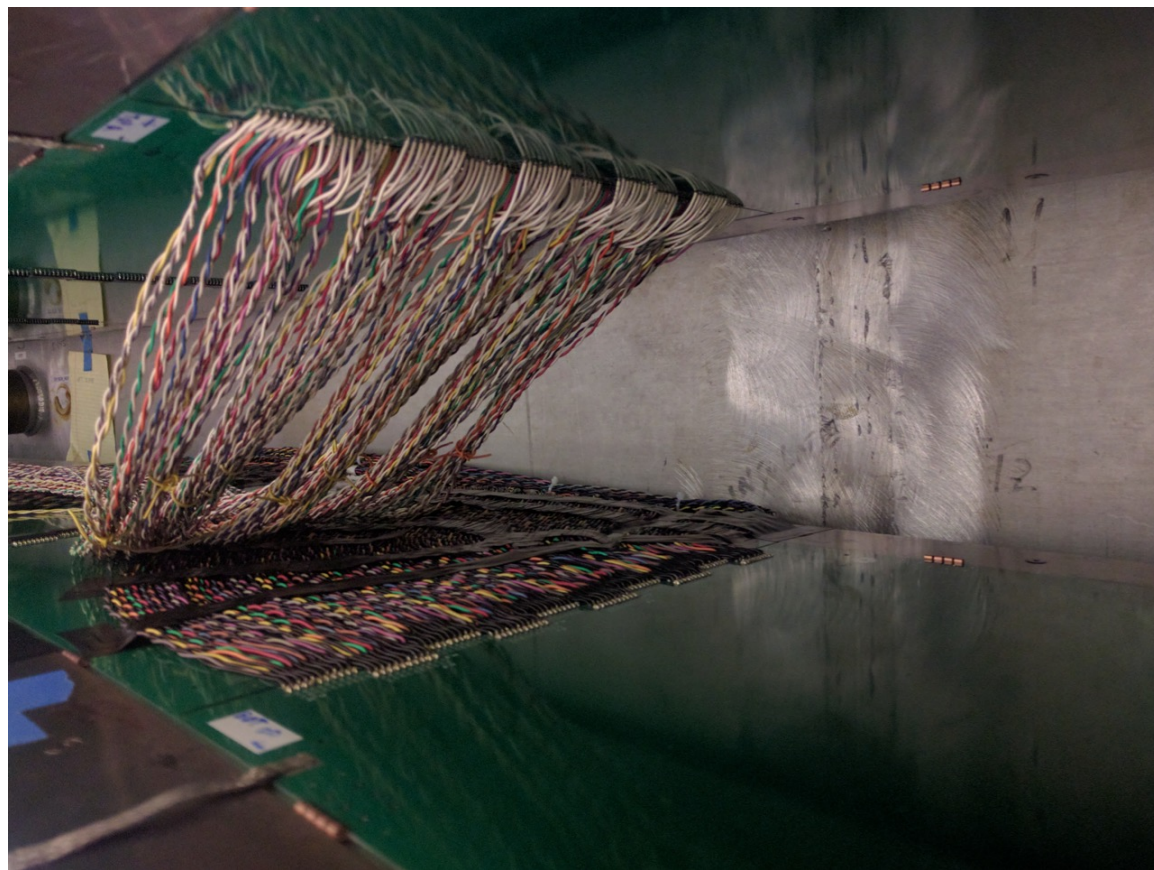


Auxiliary Field Systems



Surface Correction Coils

- Continuous PCB traces going around the ring on pole surfaces
- 100 concentric traces on upper poles, 100 on lower poles
- Programmable range: ± 20 ppm on the field
- Used to cancel higher-order multipole moments in the magnetic field (on average)



Power Supply Feedback

- Programmable current source with a range of ± 5 ppm on the field
- Uses data from **fixed probe** system to stabilize the field at a specified set point

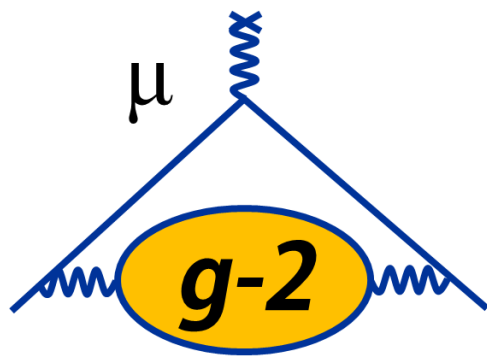


Fluxgates

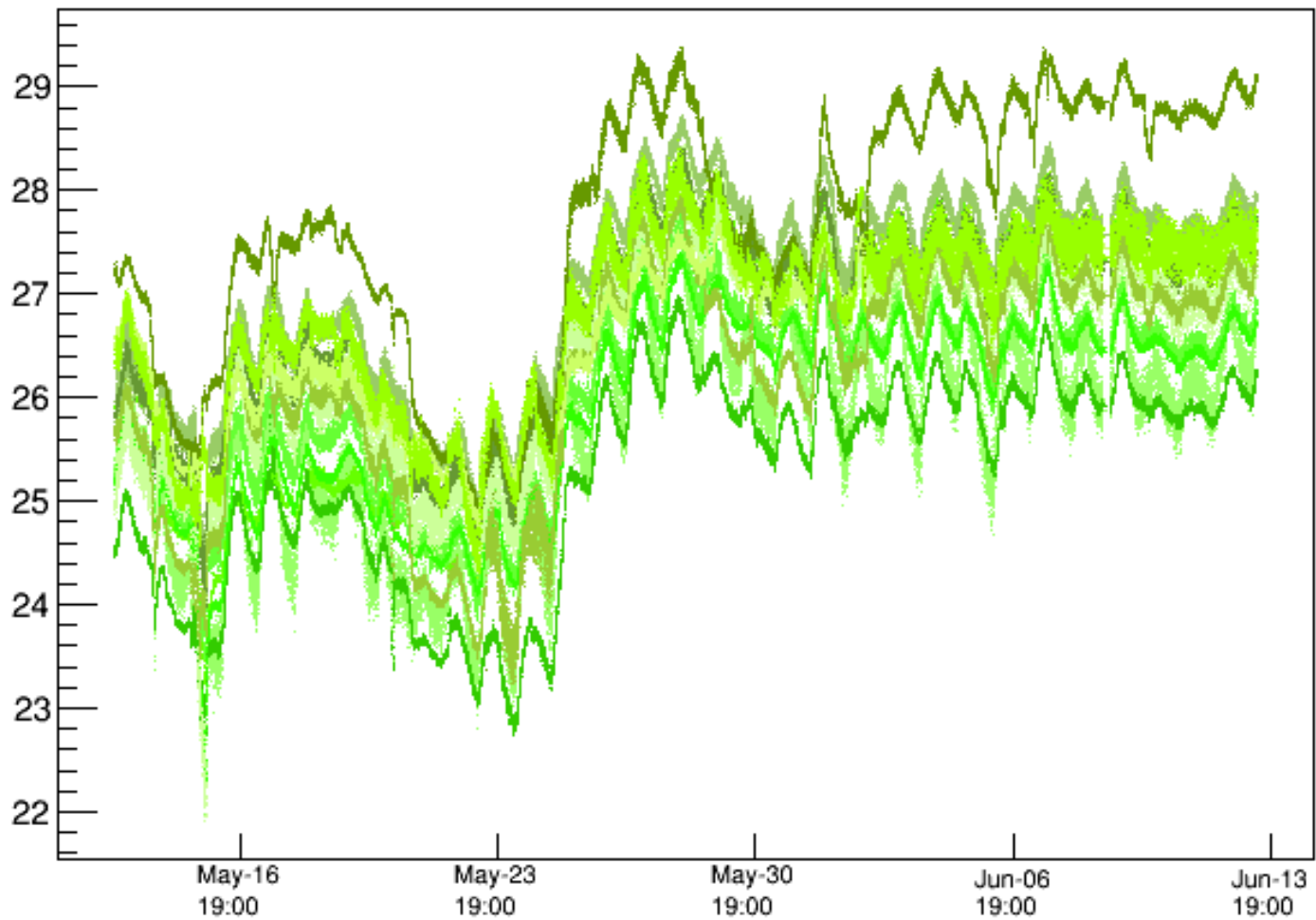
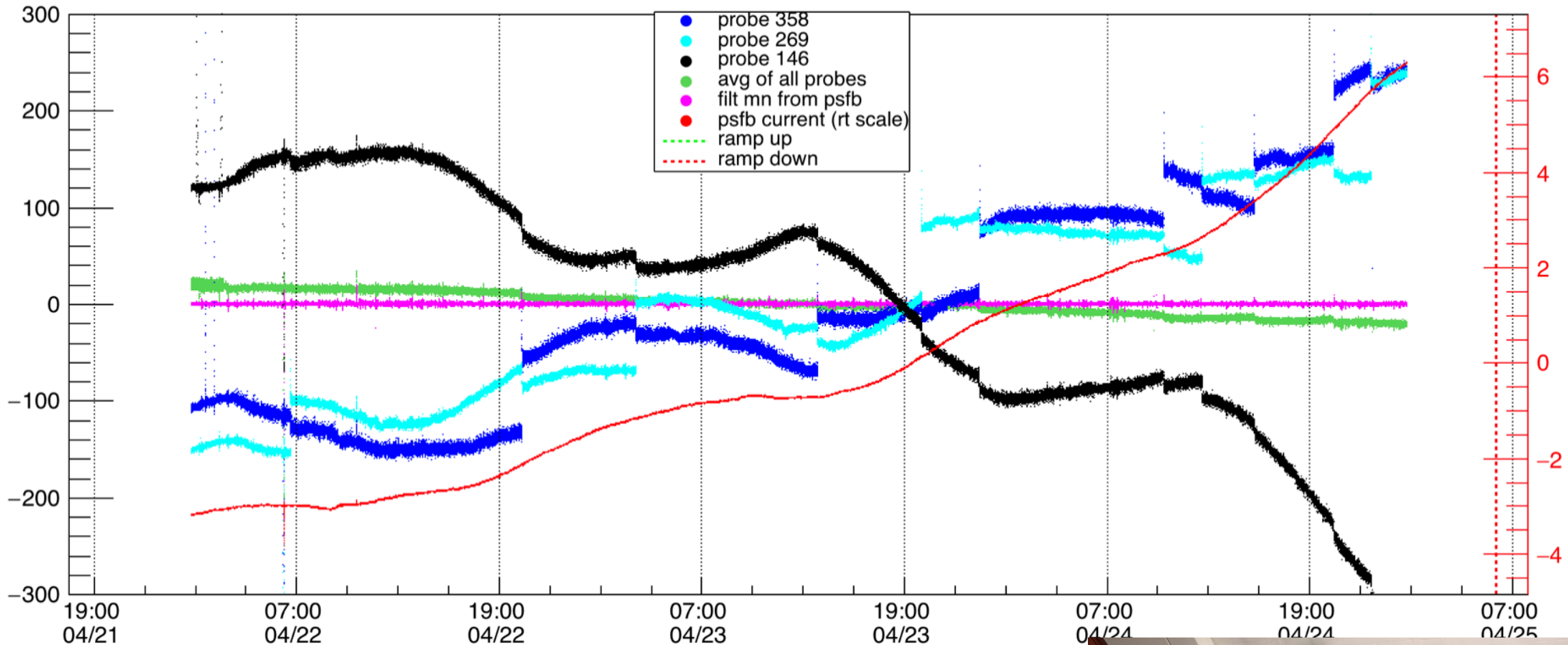
- Measure (x,y,z) components of transient fields in the hall
- Sensitive down to 10^{-9} T (DC or AC) fields
- Bandwidth up to 1 kHz



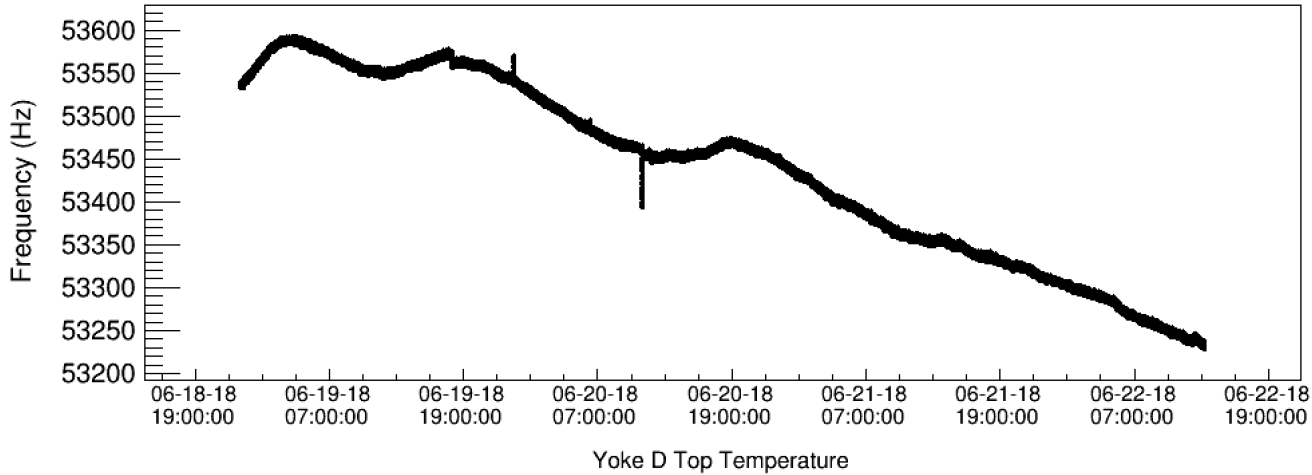
Magnet Insulation



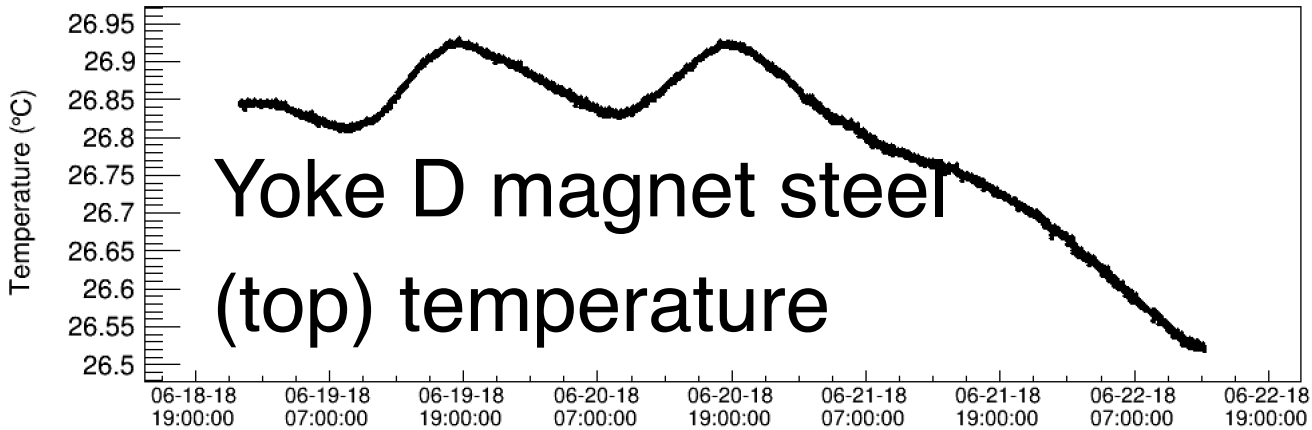
- Temperature variations in the hall affect the quality of the magnetic field
 - Observed ~ 20 ppm/deg C effects on the dipole moment during the run
 - Also affects ability to track higher-order multipoles
- Two main issues
 - Large changes in average temperature over time (2–3°C)
 - Differential changes across the magnet (~3°C)
- Two-pronged solution:
 - Improved cooling system in the hall
 - Install fiberglass insulation blanket on magnet steel



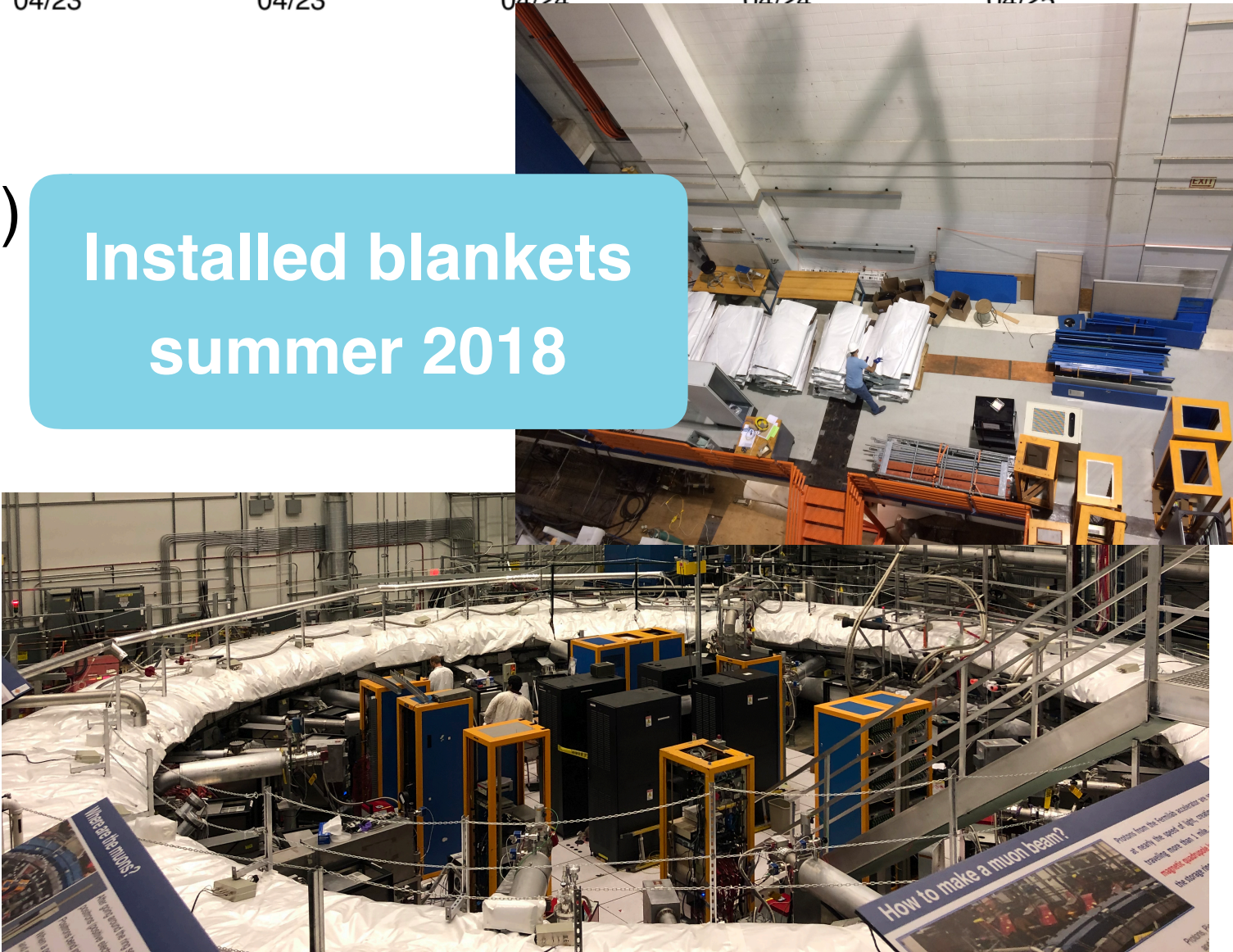
Fixed probe on yoke D vacuum chamber (top)



Yoke D magnet steel (top) temperature



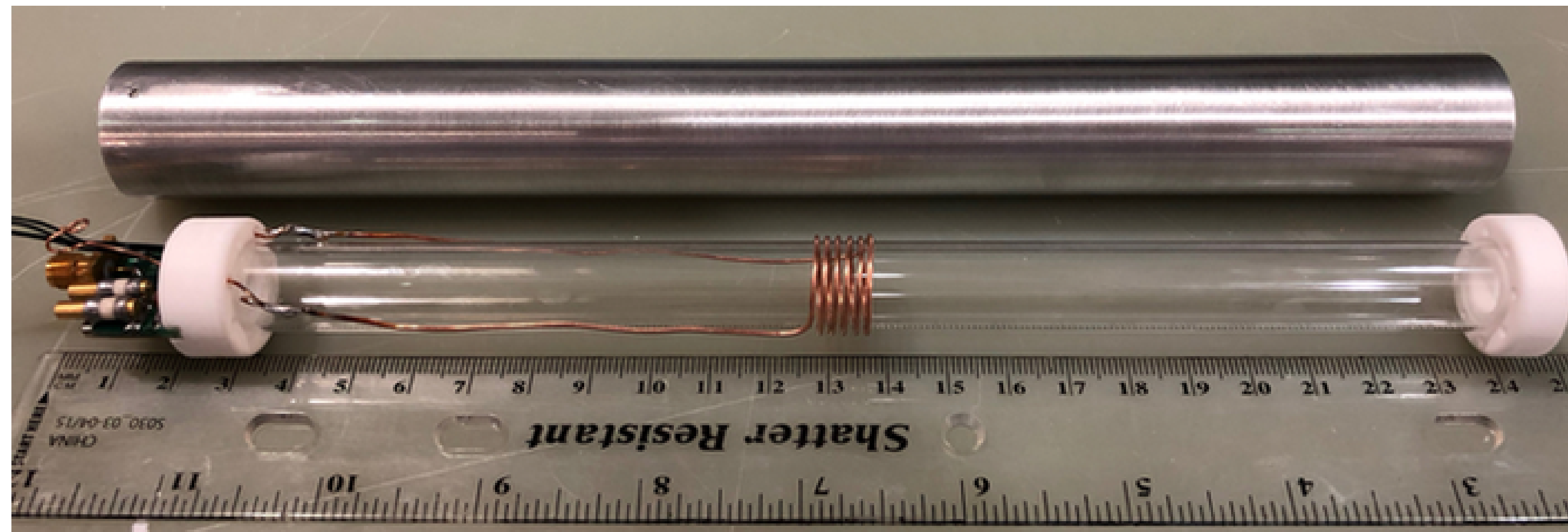
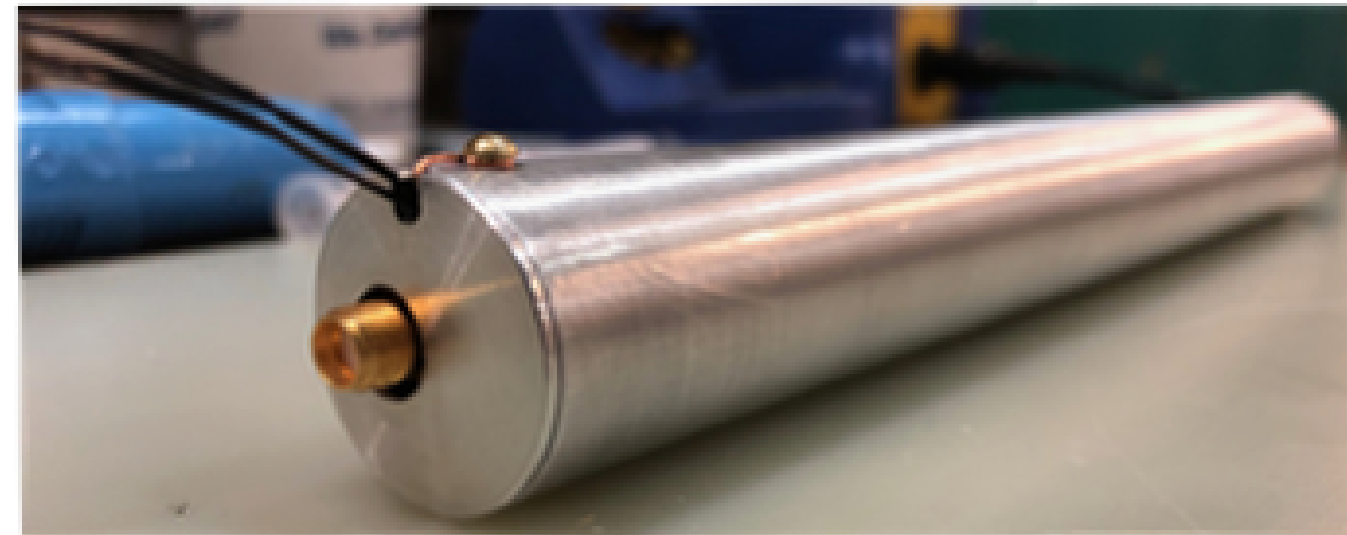
Installed blankets
summer 2018



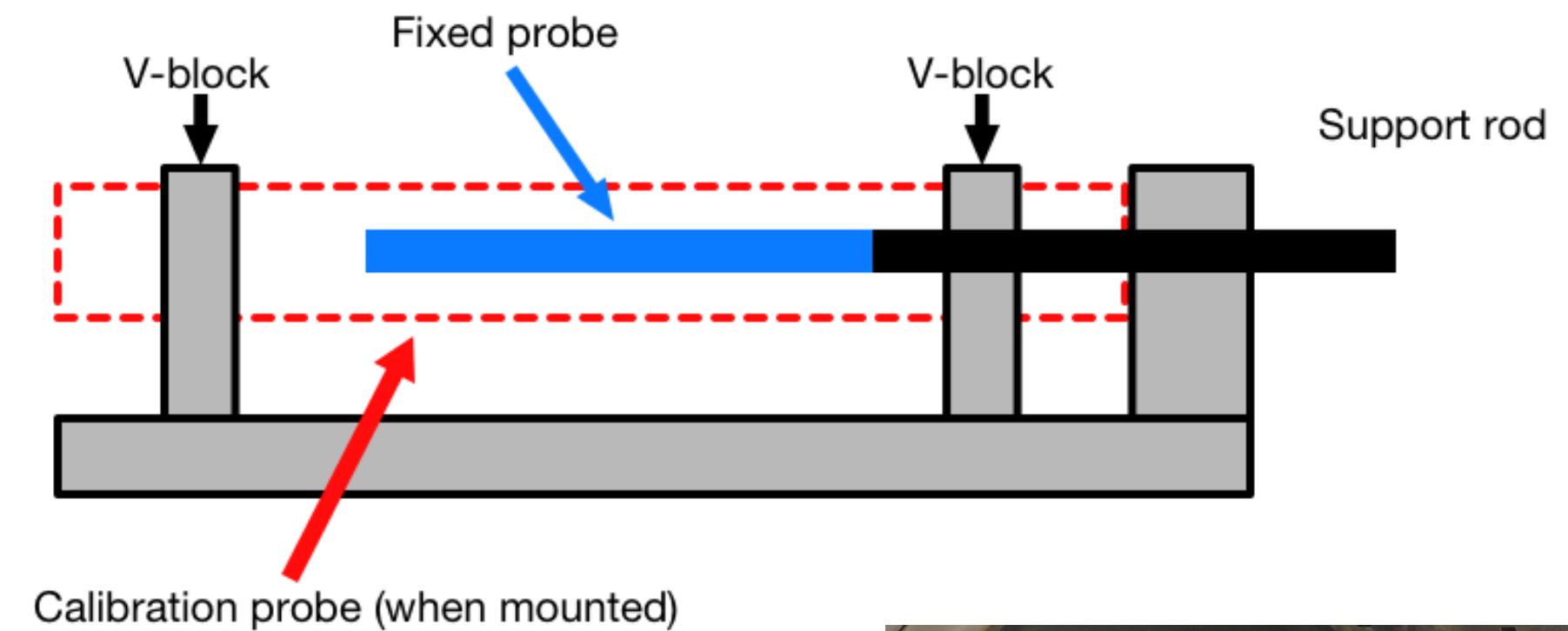
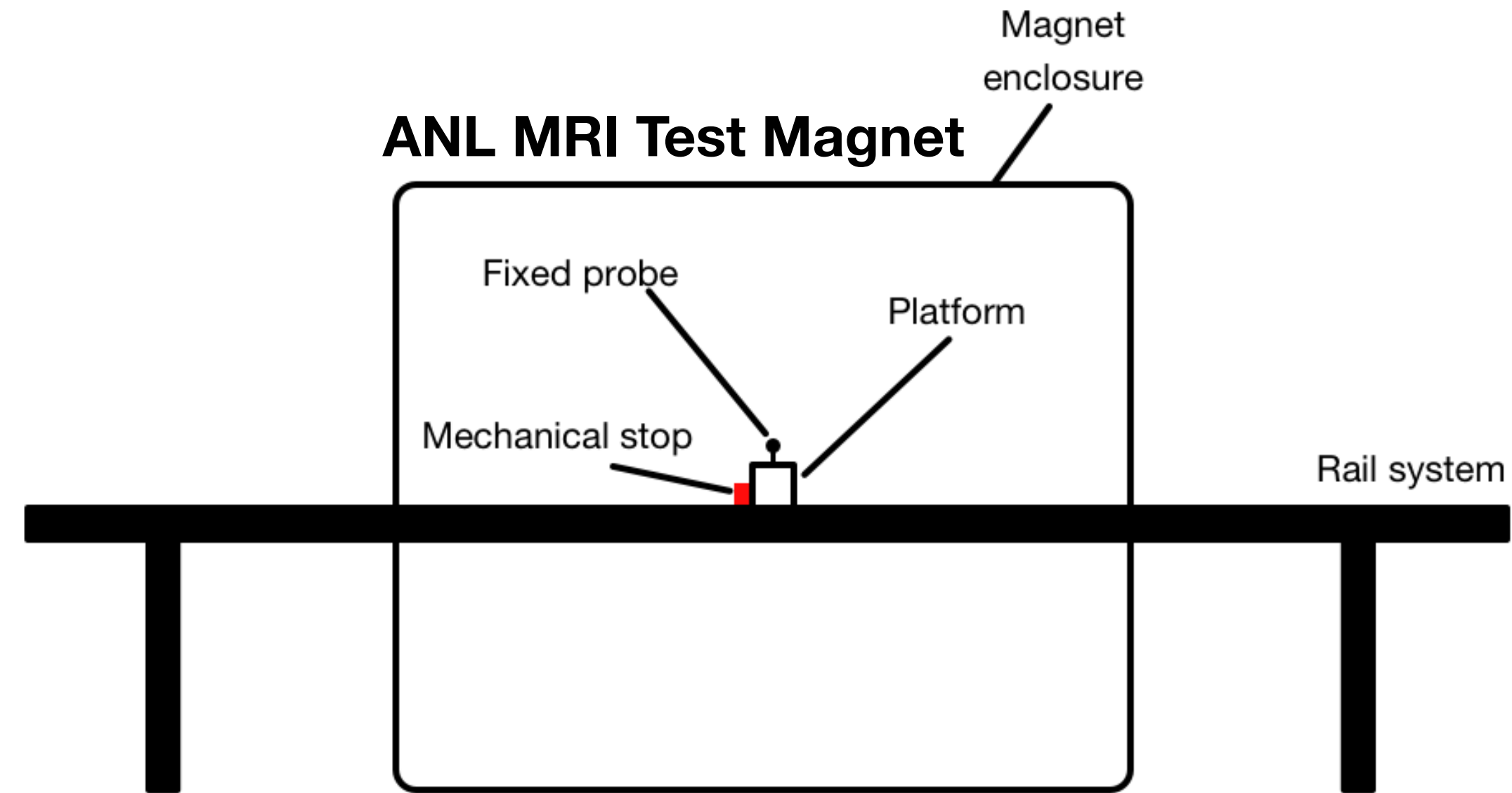
Plunging Probe Design



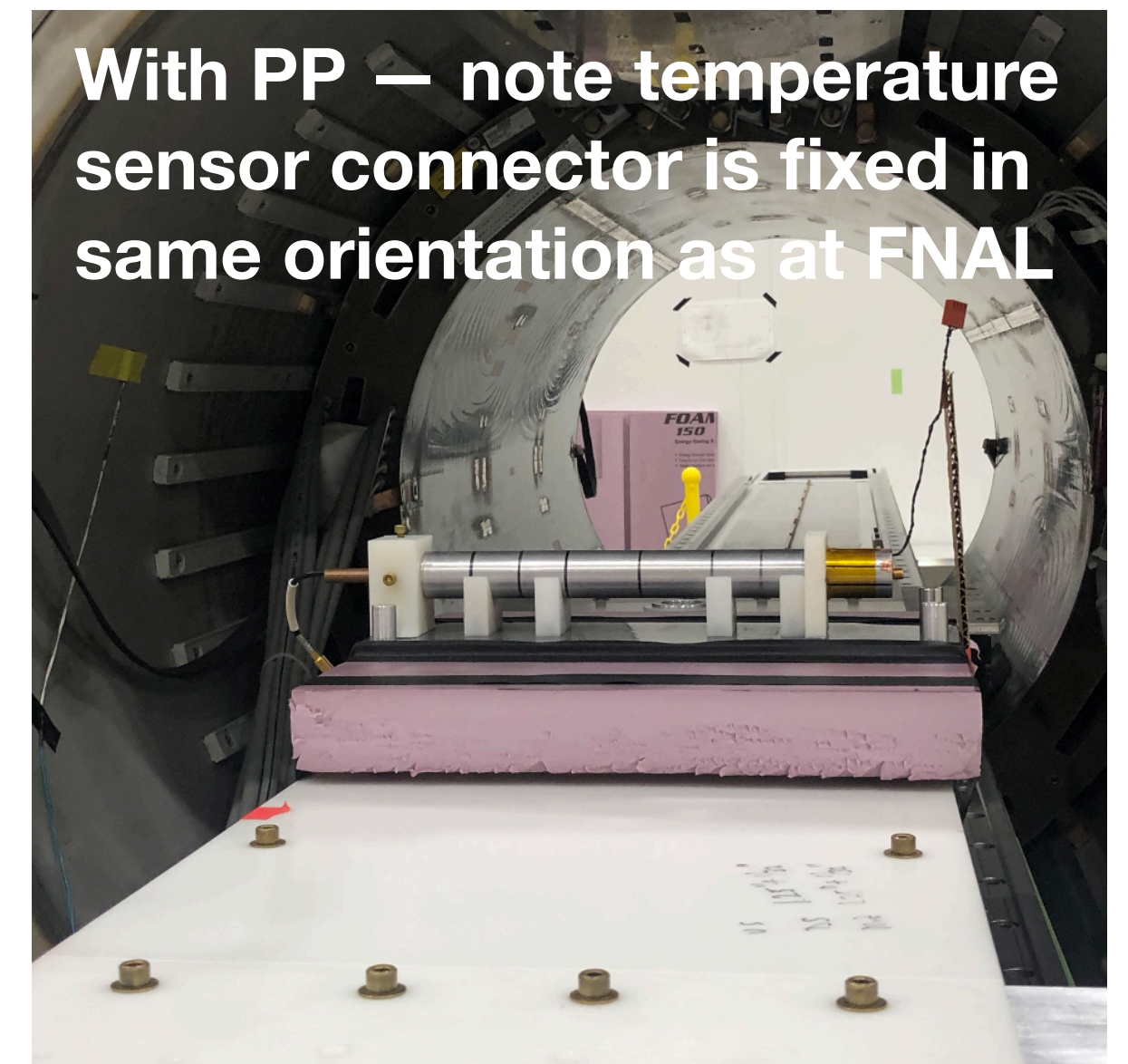
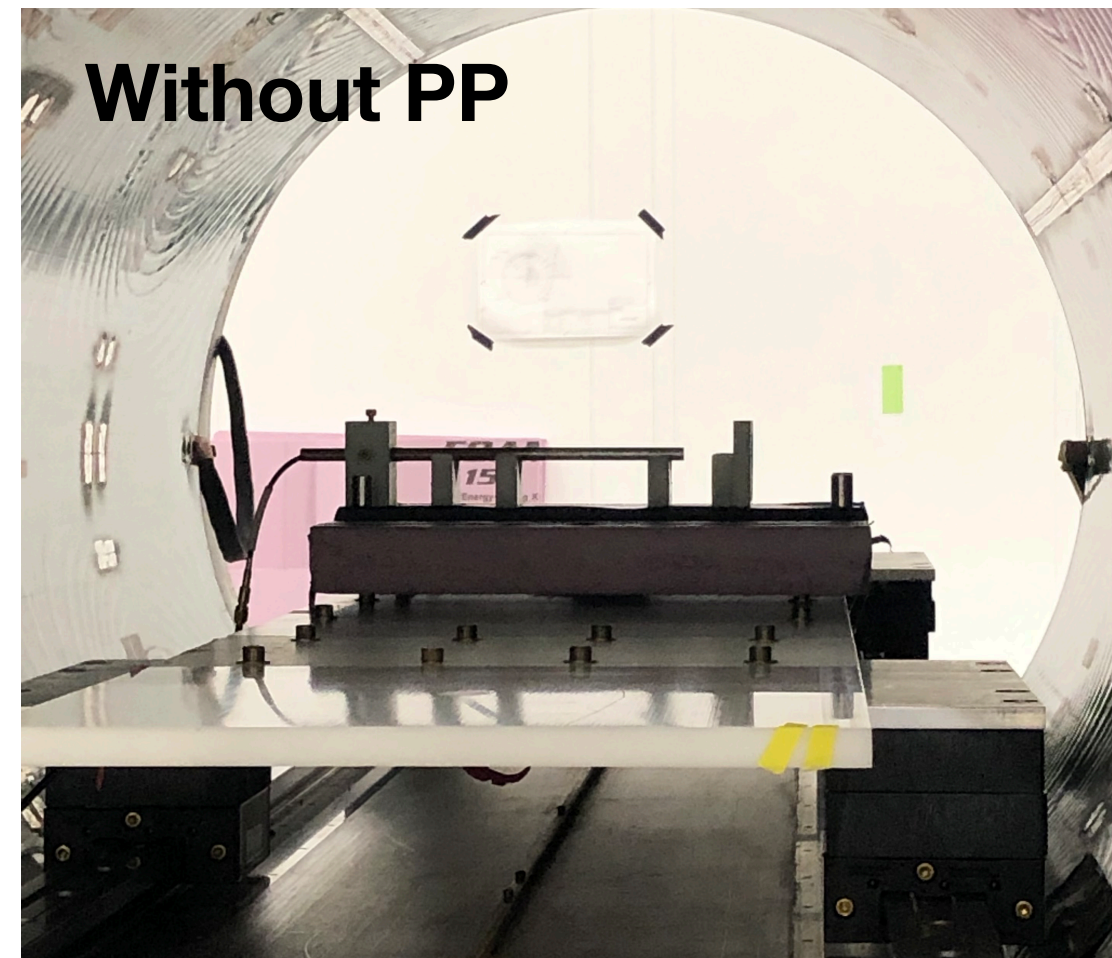
- Used to calibrate the **trolley** NMR probes
- **Symmetry** is very important => minimizes field perturbations => reduced systematic uncertainties
- **RF coil support**: 15-mm OD high-precision glass cylinder
 - Macor supports ensure alignment of **RF coil** (zero- χ 0.97-mm OD wire)
- **Ground shield**: 1" OD, 1-mm wall 2024-T3 Al
 - Stabilizes probe tune, reduces noise pickup
- **Vacuum compatible**



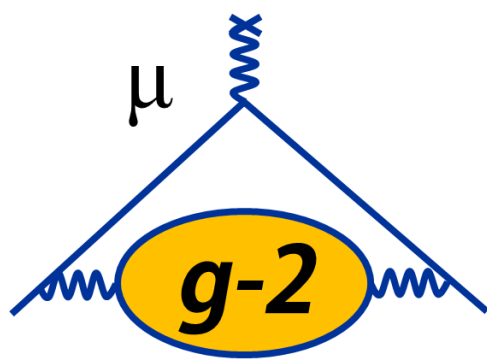
Plunging Probe: Measuring Perturbations



- Take measurements of the field using the **fixed probe**
- **Compare** measurements without and with the PP mounted on stand
- **Difference** with and without gives the effect



Plunging Probe: Material Perturbations (δ_s)



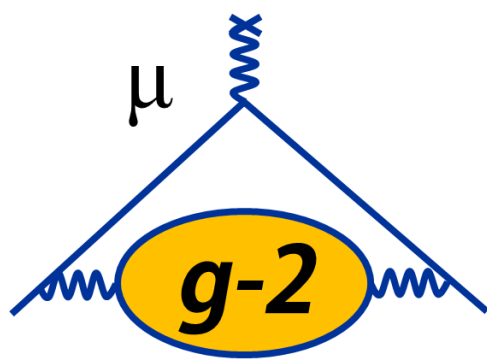
How much the field changes due to probe presence

Quantity	Symbol	Value (ppb)
General Material Perturbation	$\delta_{\text{mat}} + \delta_{\text{mag}}$	4.2 ± 8.0
SMA Cable Perturbation	δ_{cable}	$-1.4 + 3.0$
Probe Temperature	δ_T	0 ± 5
Roll Effect	δ_{roll}	0 ± 1
Pitch Effect	δ_{pitch}	-4.4 ± 4.4
Water Sample Camber	δ_c	0 ± 1
TOTAL	δ_s	-1.6 ± 10.9

Dependence on orientation about its **own** axis

Dependence on orientation relative to **field** axis

Plunging Probe: Water Purity (δ_p)



- Impurities in the water sample will perturb the field — paramagnetic contamination (e.g., dissolved oxygen) will increase the field
- Conduct two tests:
 1. Measure field when we **degas** the water — that is, heat up water just enough so that oxygen escapes. Compare to nominal water sample (at same T)
 2. Compare field measurements using water from two different vendors

• Define $\delta_p = \delta_{O_2} + \delta_w$

dissolved oxygen

different vendors

Quantity	Symbol	Value (ppb)
Oxygen Contamination	δ_{O_2}	1.4 ± 0.5
Different Vendors	δ_w	0 ± 1
TOTAL	δ_p	1.4 ± 1.1

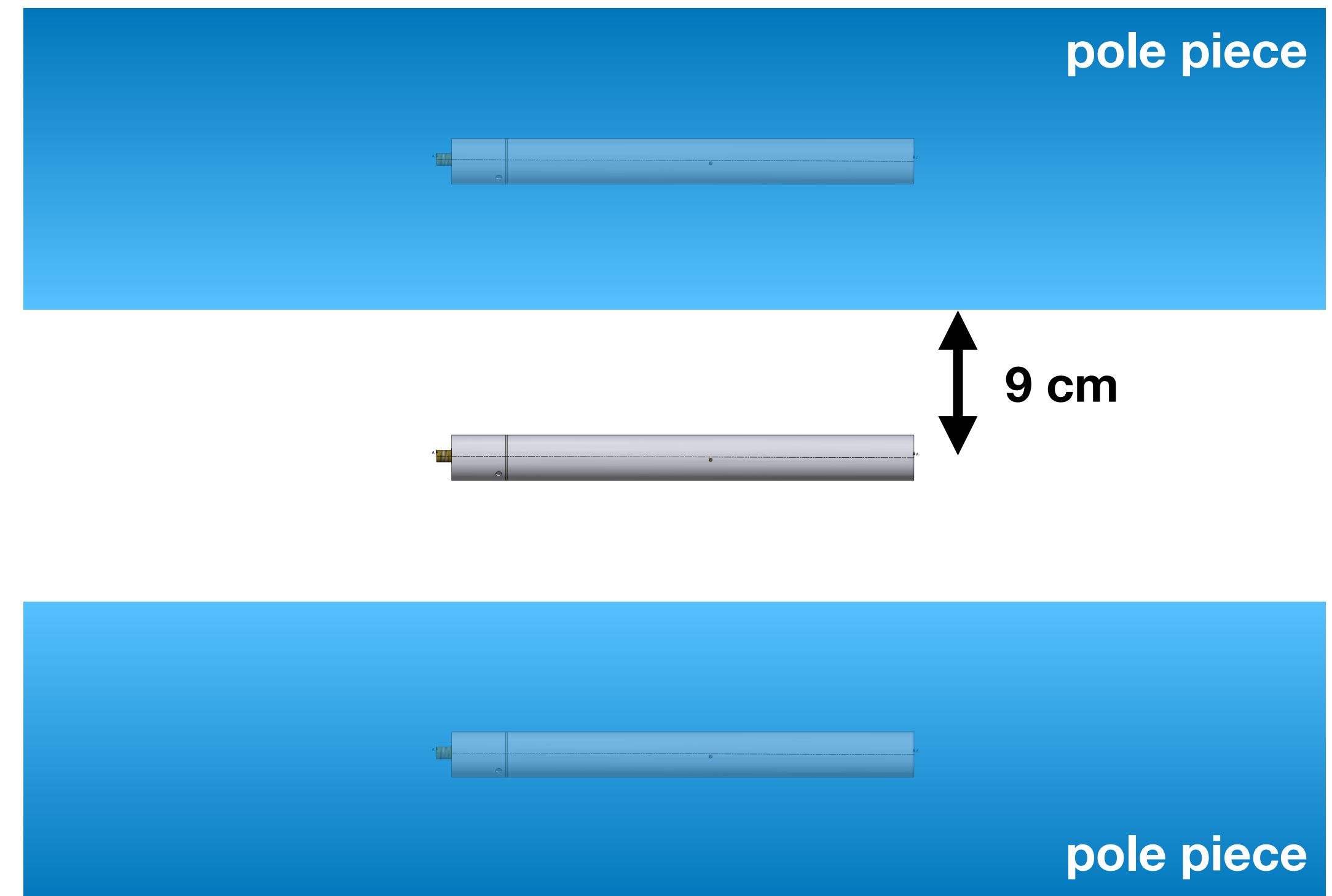
Plunging Probe: Magnetic Images



- Need to account for the effect due to magnetic images of the PP in the pole pieces
- For an infinite plane with magnetic permeability μ_r , the field due to the image of a material with perturbation ΔB is:

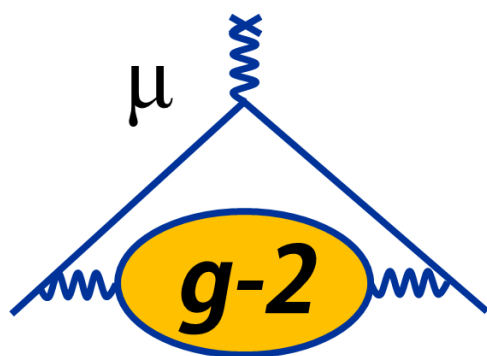
$$\Delta B' \approx \left(\frac{\mu_r - 1}{\mu_r + 1} \right) \Delta B (x, y, z') \approx \Delta B (x, y, z')$$

For $\mu_r \gg 1$ (~ 1450 for the magnet*). Evaluate the perturbation at image distance z' in the pole piece (upper **and** lower)

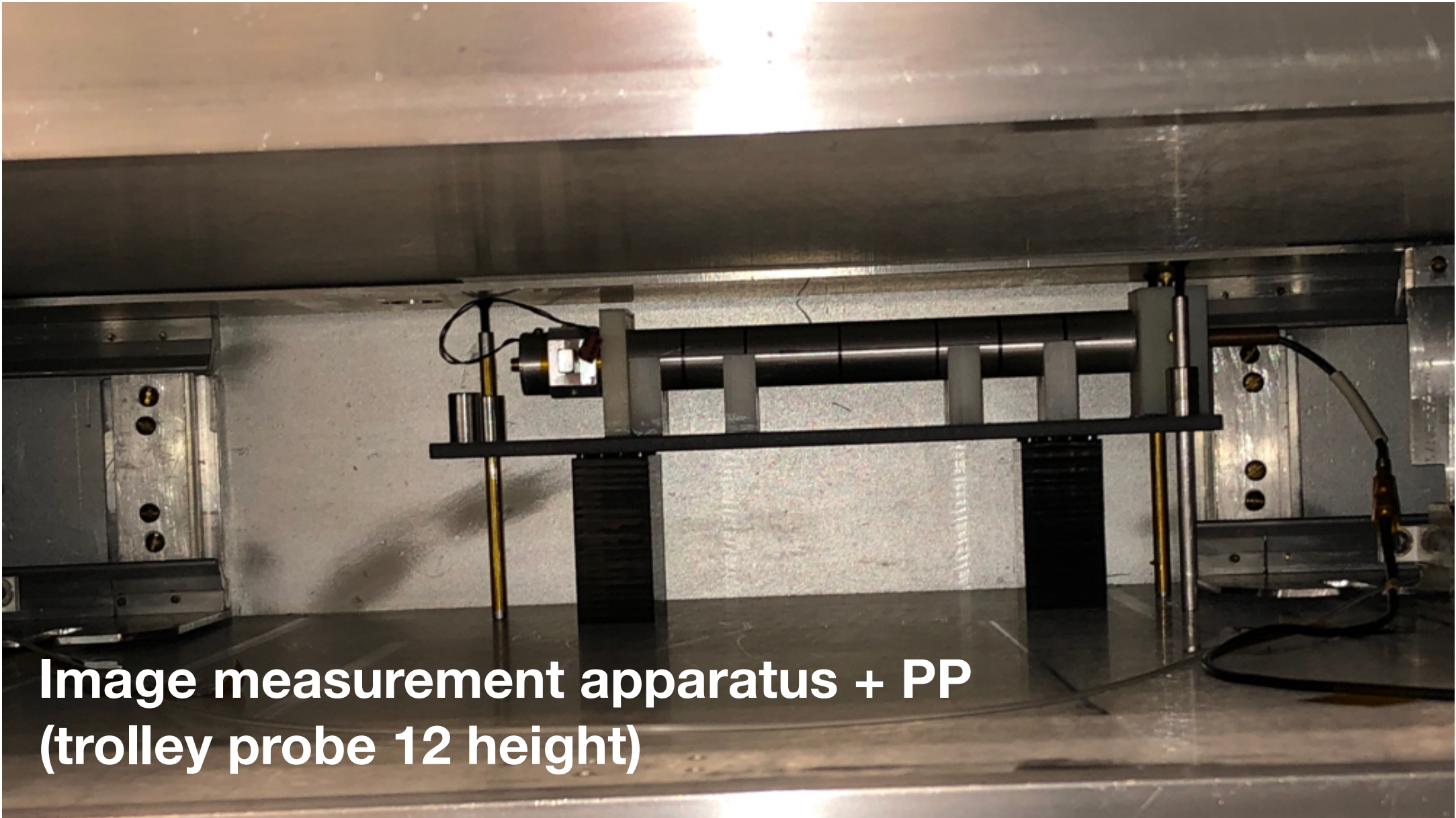


* G. T. Danby *et al*, Nucl. Instrum. Meth. A **457**, 151 (2001)

Plunging Probe: Measuring the Images at FNAL



- Use a stage to mount a fixed probe along the axis of the PP, which can slide over fixed probe
- Compare field measurements with and without PP installed on the stage
- Repeat measurements at height of center trolley probe, and highest trolley probe location
- Also conduct measurements with PP mounting rod attached/detached



Rod composition may not be pure aluminum (typically up to ~20% variation in $\chi \Rightarrow$ imperfect predictions)

Type	Height Above Midplane (mm)	Image + Perturbation (ppb)	Calculated Prediction (ppb)
PP + Holder	35.5	6.1 ± 4.4	11.5
PP + Holder	-0.2	7.3 ± 5.4	4.3
Rod	-0.2	-3.1 ± 5.6	-12.9
TOTAL	-0.2	4.2 ± 8.0	-8.6

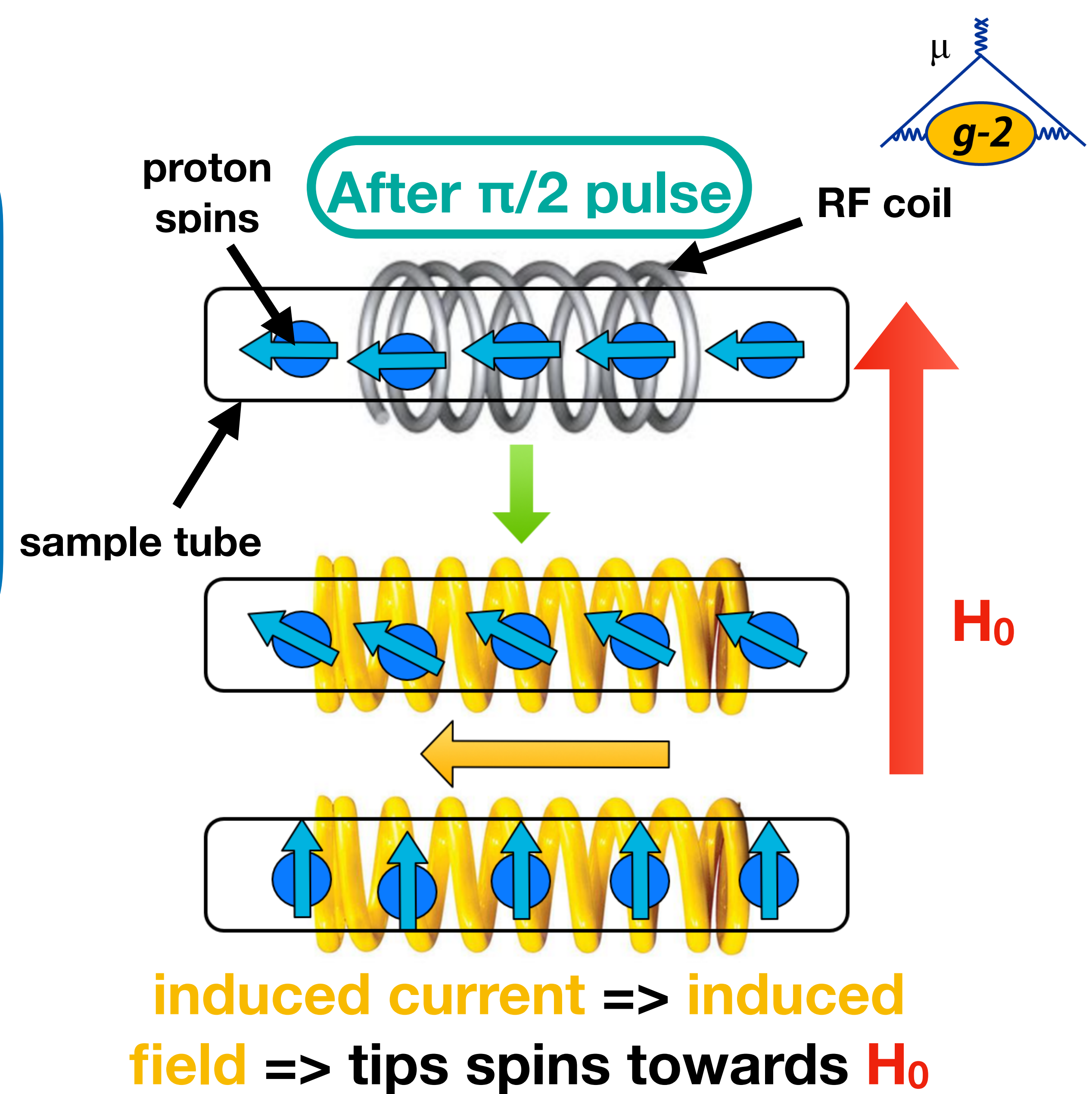
Radiation Damping

What is it?

- Precessing spins induce emf in pickup coil; this in turn generates an alternating magnetic field that **acts to rotate spins back towards the main field**
- **Size of effect:** $\delta_{RD} \sim [(f_0 - f_L)/f_0] \eta Q M_z(t) / \tau_{RD}$
 - f_0 = resonant frequency of circuit; f_L = Larmor frequency
 - η = filling factor; Q = quality factor of circuit
 - $M_z(t)$ = magnetization of sample, τ_{RD} = time scale of effect

How to quantify?

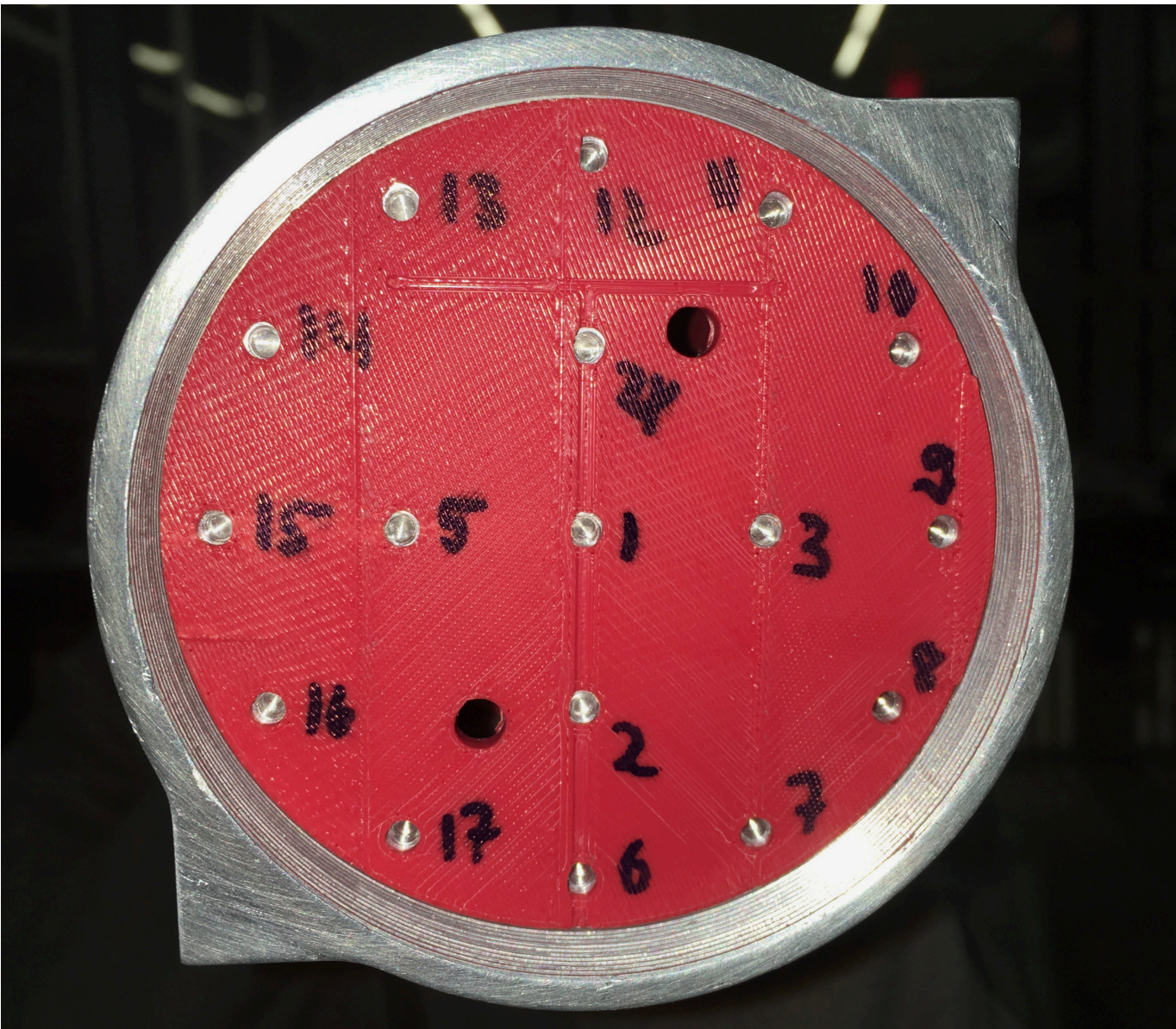
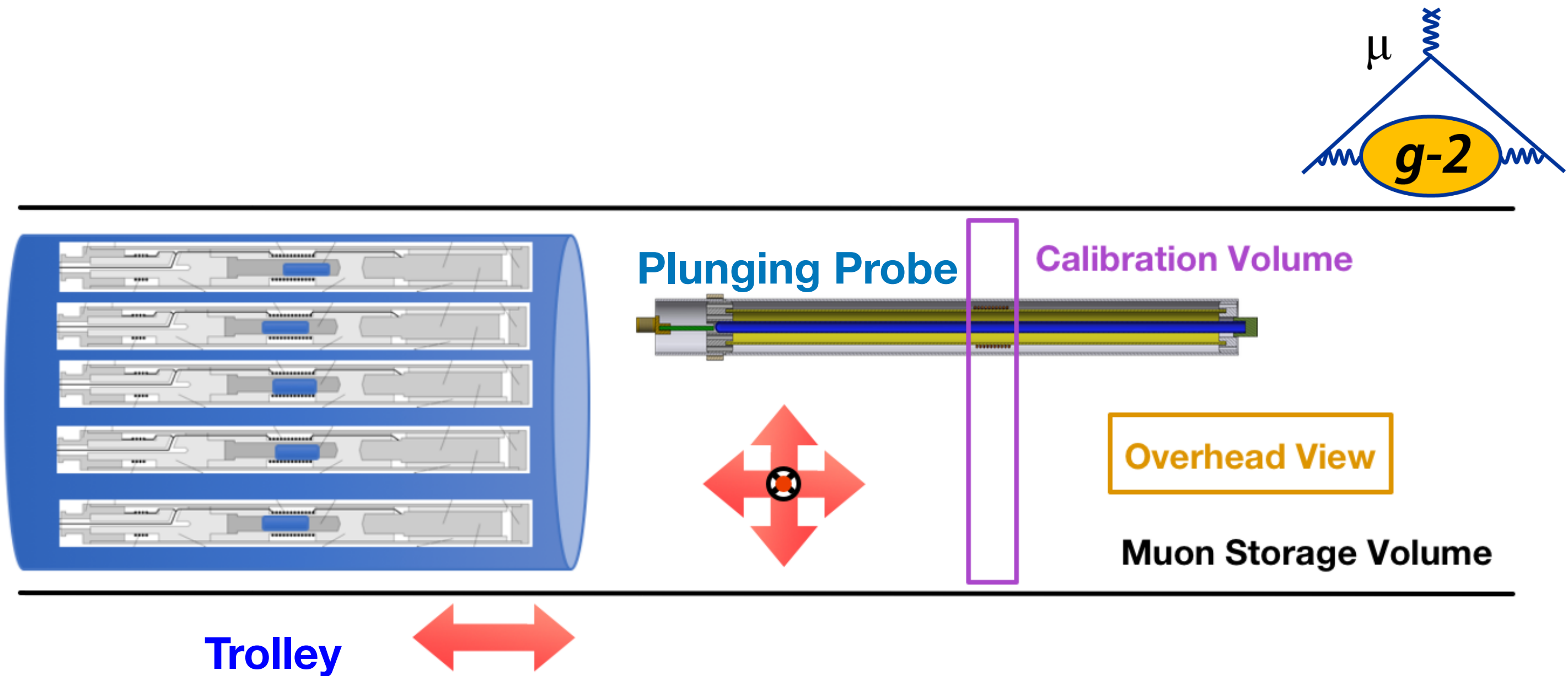
- Use coils to produce a longitudinal field
 - Precise control over main field to mimic damping effect
- Vary $\pi/2$ pulse \Rightarrow vary $M_z(t)$ \Rightarrow changes δ_{RD}



Calibrating the Trolley

Procedure

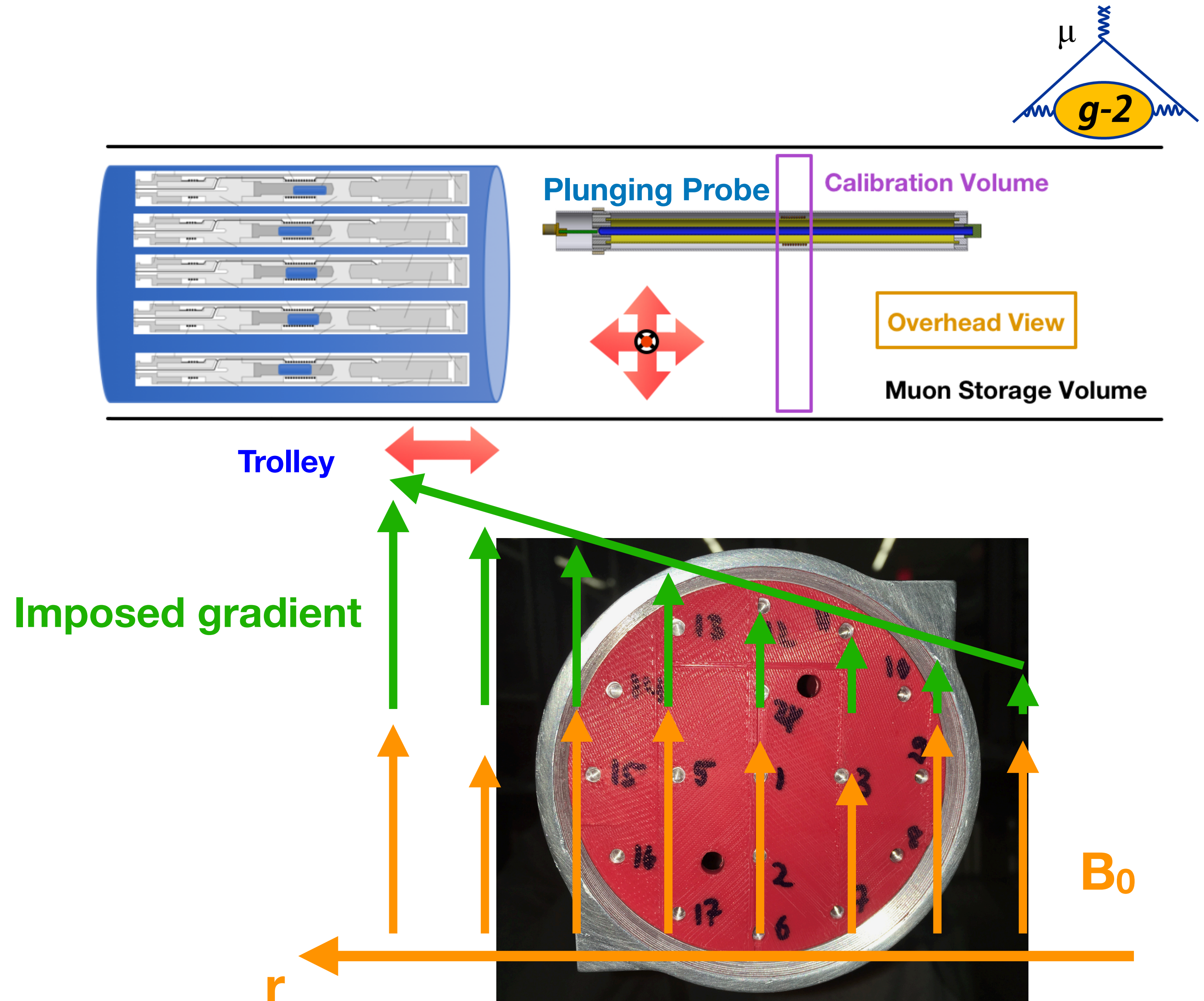
- Select **trolley** probe to calibrate



Calibrating the Trolley

Procedure

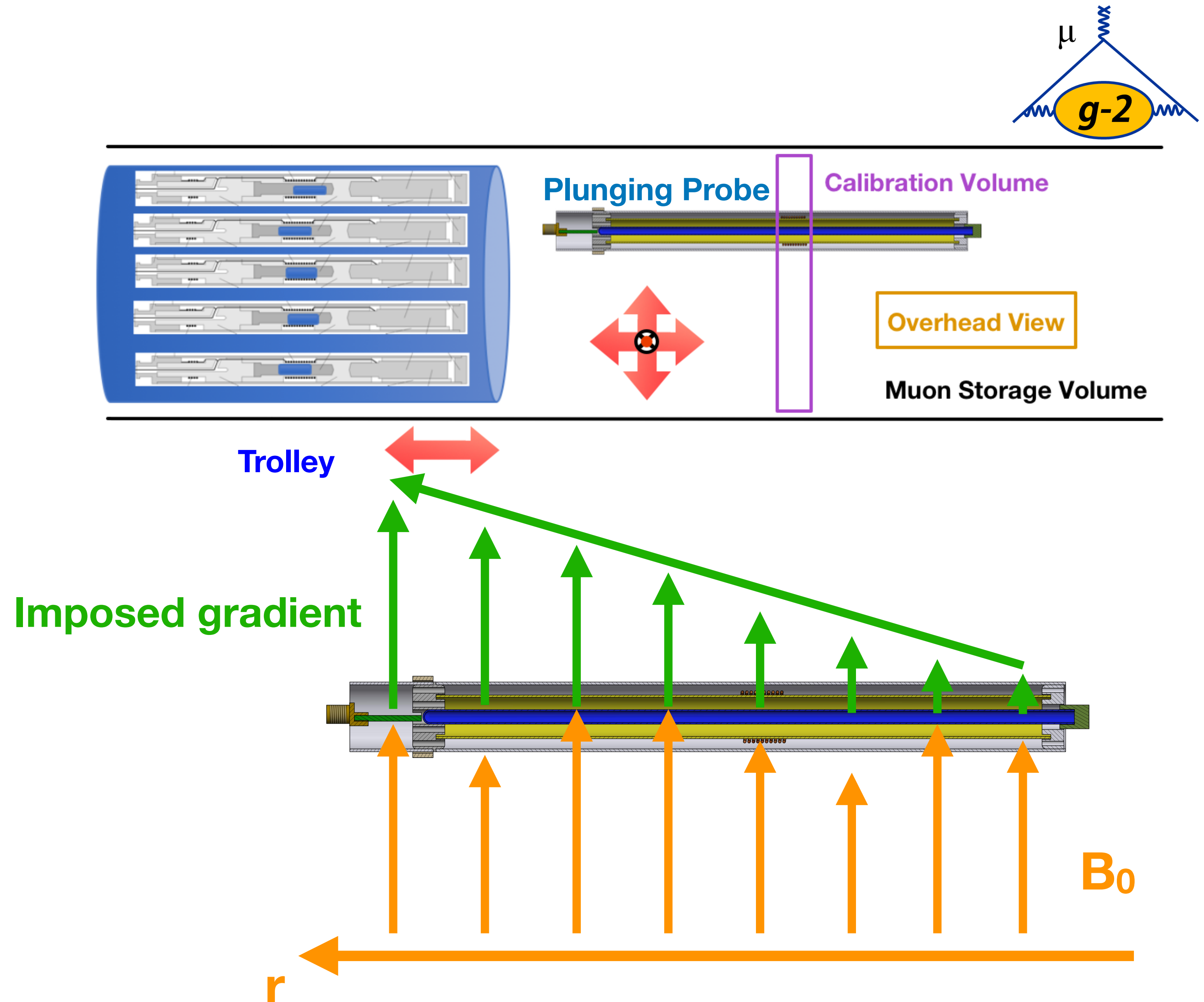
- Select **trolley** probe to calibrate
- Impose a **known gradient** across the trolley; compare to **bare field** B_0 . Define $\Delta B = B(I \neq 0) - B(I = 0)$
- Unique ΔB for each **trolley** probe gives position



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- Move **plunging probe** into volume; measure ΔB and determine distance to move **plunging probe**
- Iterate until **plunging probe** ΔB matches **trolley** probe ΔB
- Perform for radial, vertical, azimuthal coordinates



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- Perform for radial, vertical, azimuthal coordinates
- Shim the field to be highly uniform, and measure using the **PP** and the **trolley** (rapid swapping)

