

Neutrino Tridents at Near Detectors

Yuber F. Perez-Gonzalez

In collaboration with

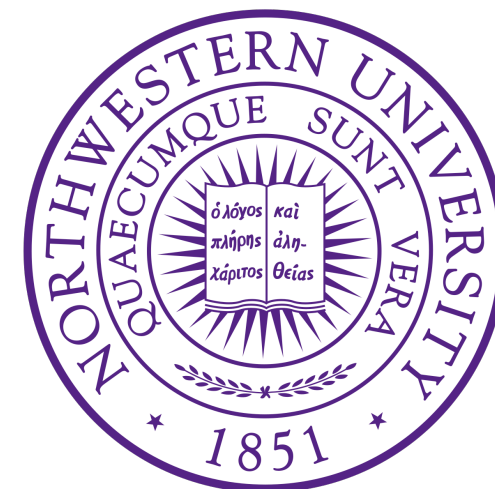
Peter Ballett - Matheus Hostert - Silvia Pascoli (Durham),
Zahra Tabrizi and Renata Zukanovich Funchal (U-São Paulo)

Physics Opportunities in the Near
DUNE Detector Hall
Fermilab, Nov 2018



Based on arXiv:1807.10973

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Outline

- Introduction
- Trident Cross Section
- Events in LAr Detectors (DUNE ND)
- Conclusions

Introduction

Near Detector Program

Neutrino sector
under scrutiny

Precision
measurements in
the neutrino sector

Reduce systematic
uncertainties

- ◆ Flux
- ◆ CC and NC cross sections
- ◆ Backgrounds

Near Detector Program

Neutrino sector
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- ♦ Flux
- ♦ CC and NC cross sections
- ♦ Backgrounds

High beam luminosity +
Large fiducial mass

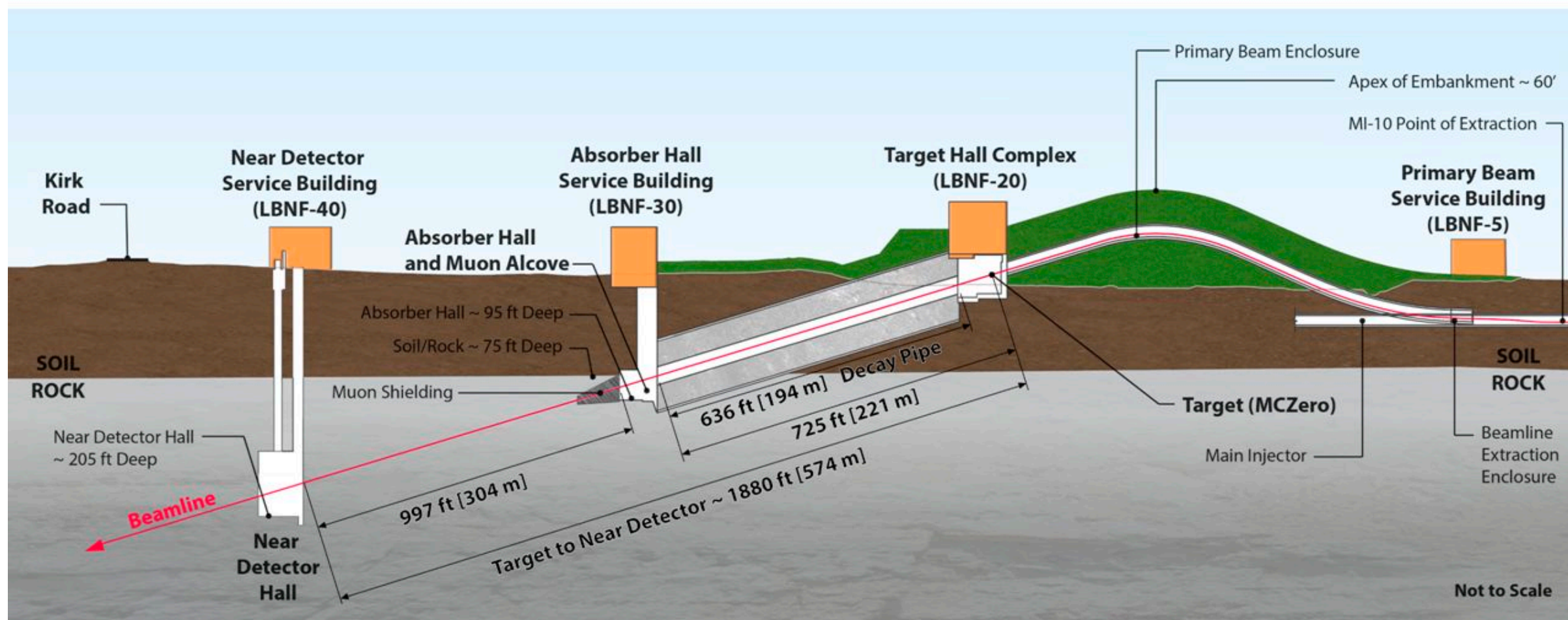
Ideal places to investigate rare
neutrino interactions

$$\sigma < 10^{-44} \text{ cm}^2$$

- ♦ Test SM predictions
- ♦ Search for BSM physics

DUNE ND

DUNE Beamline with Near Detector



Events per ton-year

$$\nu_{\mu} \text{ CC Total} \quad 1.64 \times 10^6$$

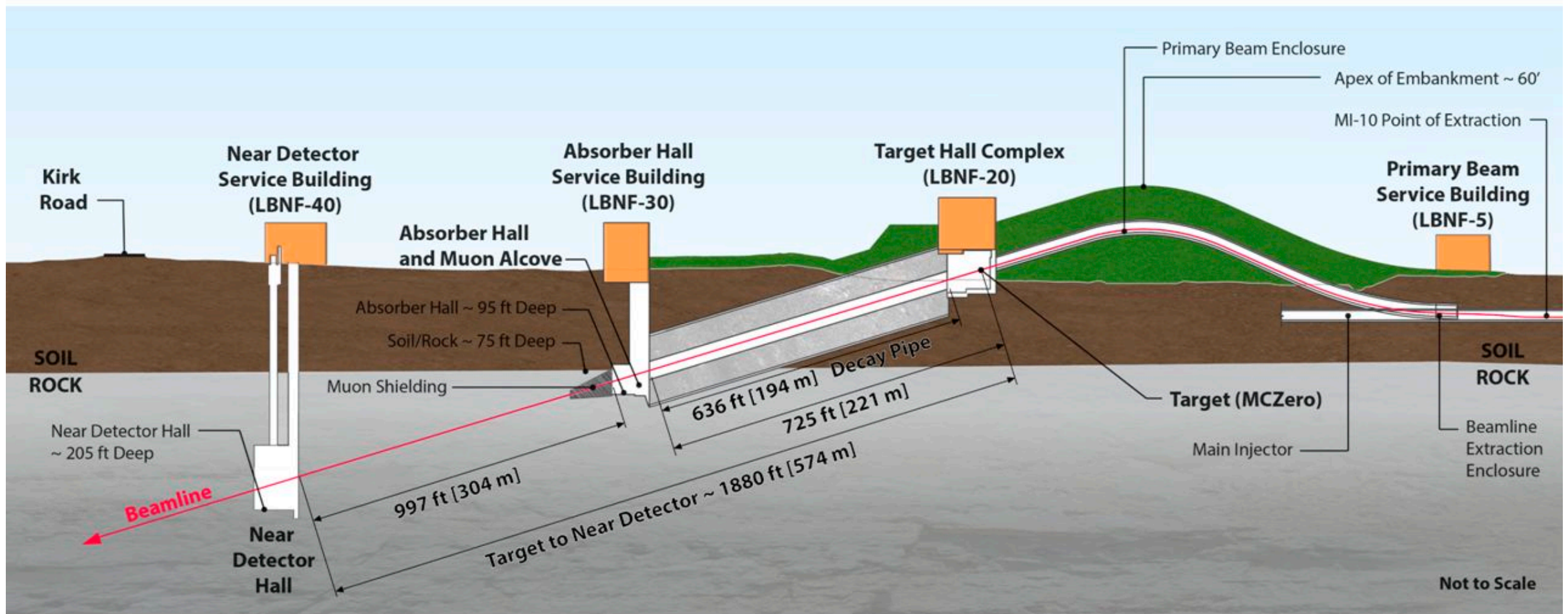
$$\nu_{\mu} \text{ NC Total} \quad 5.17 \times 10^5$$

$$\nu_{\mu} - e \quad 135$$

Alan's talk

DUNE ND

DUNE Beamline with Near Detector



Events per ton-year

ν_{μ} CC Total 1.64×10^6

ν_{μ} NC Total 5.17×10^5

$\nu_{\mu} - e$ 135

What about rare
neutrino
scatterings?

Alan's talk

Trident Inelastic Scattering

Production of a charged lepton pair from the inelastic neutrino scattering in the Coulomb field of the nucleus

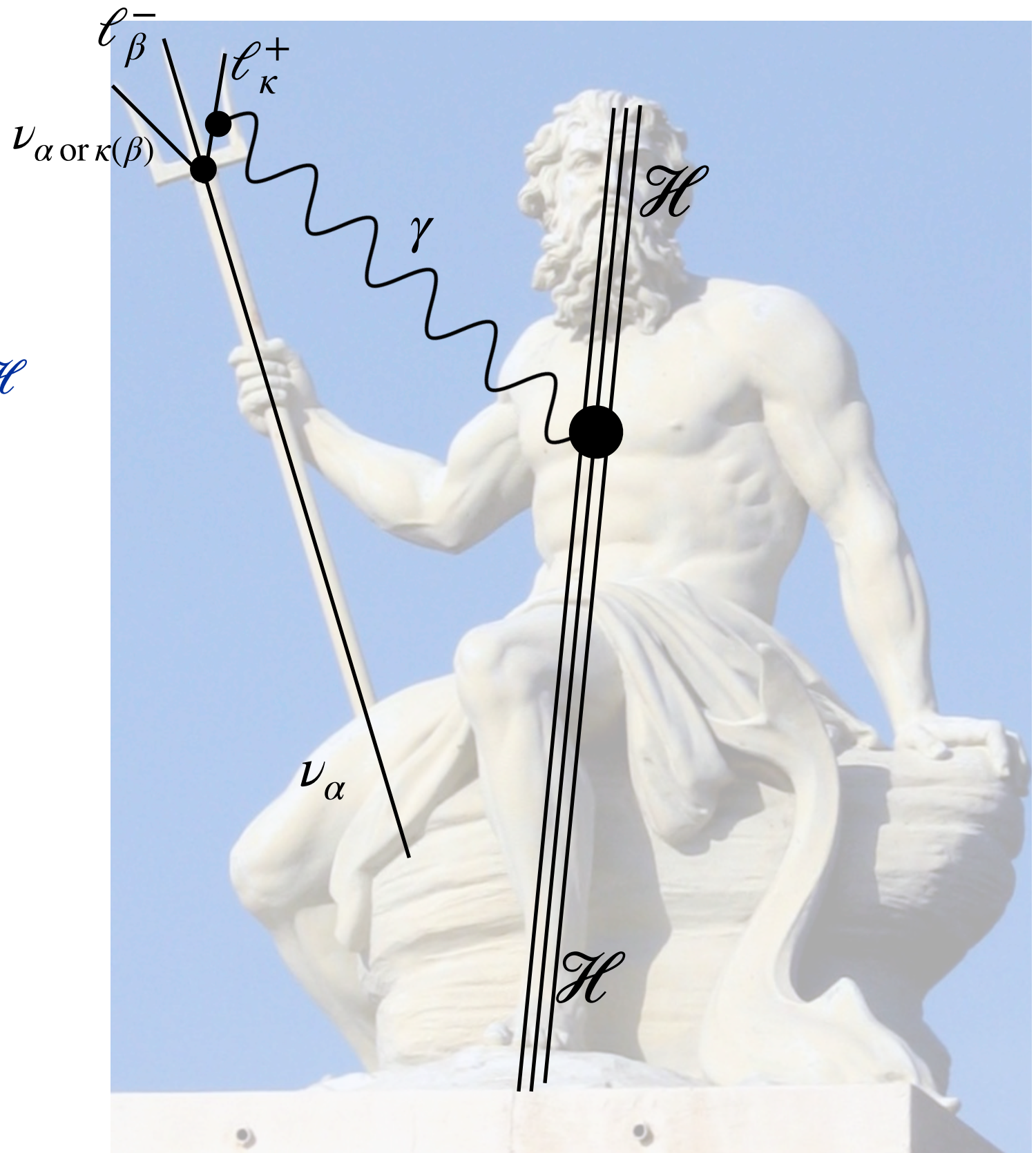
$$\nu_{\alpha} + \mathcal{H} \rightarrow \nu_{\alpha \text{ or } \kappa(\beta)} + \ell_{\beta}^{-} + \ell_{\kappa}^{+} + \mathcal{H}$$



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$$\nu_{\mu} \rightarrow \nu_{\mu} \mu^{+} \mu^{-}$$

CHARM II

PLB 245 (1990) 271

$$\frac{\sigma_{\text{CHARM II}}}{\sigma_{\text{SM}}} = 1.58 \pm 0.57$$

CCFR

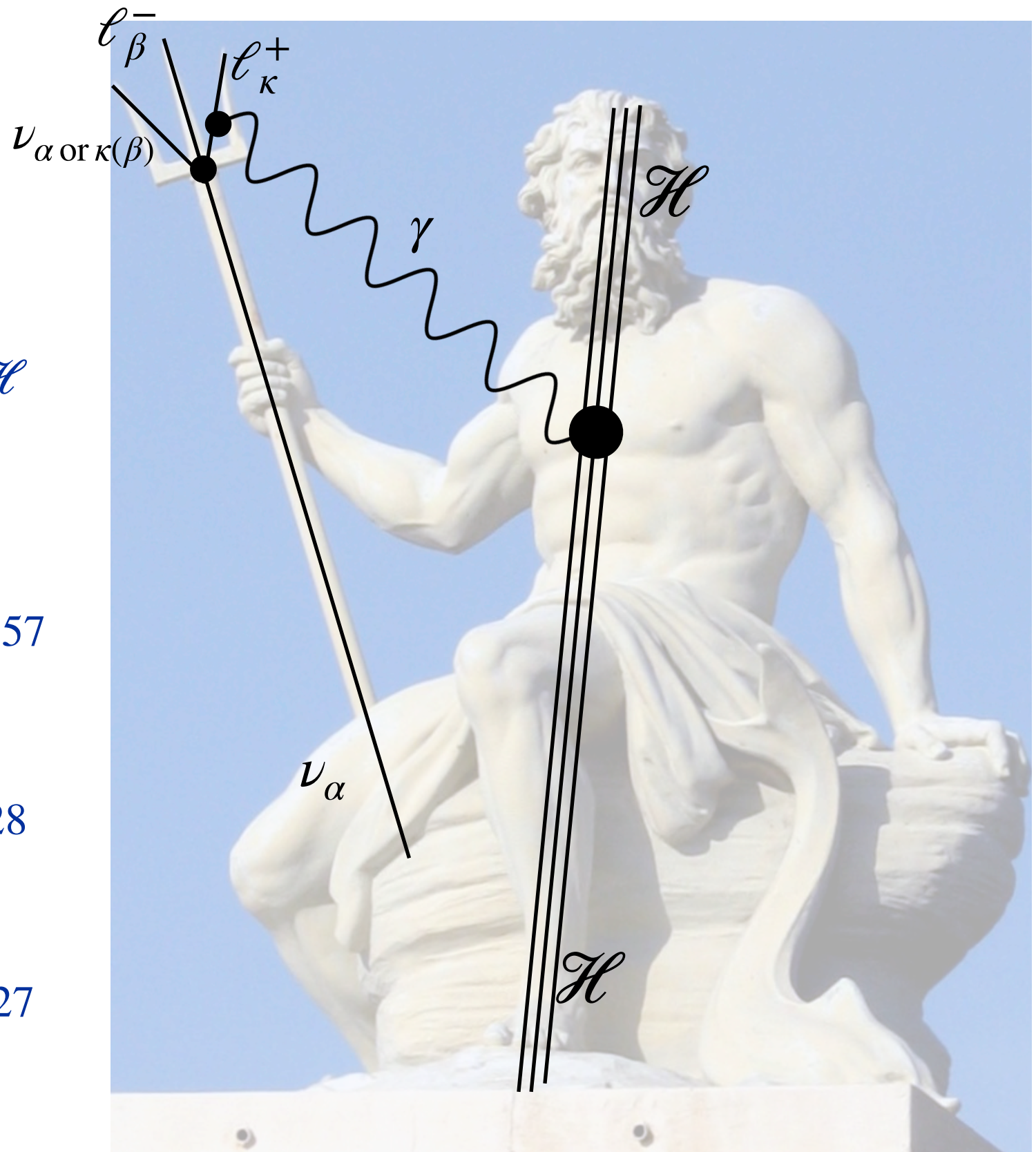
PRL 66 (1991) 3117

$$\frac{\sigma_{\text{CCFR}}}{\sigma_{\text{SM}}} = 0.82 \pm 0.28$$

NuTeV

Vancouver 1998, High energy physics, vol. 1

$$\frac{\sigma_{\text{NuTeV}}}{\sigma_{\text{SM}}} = 0.67 \pm 0.27$$



Trident Inelastic Scattering

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W. Altmannshoffer et. al., 2014

Z' constraints from trident

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Rates at DUNE ND

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Tridents produced by
atmospheric neutrinos

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Charged scalars

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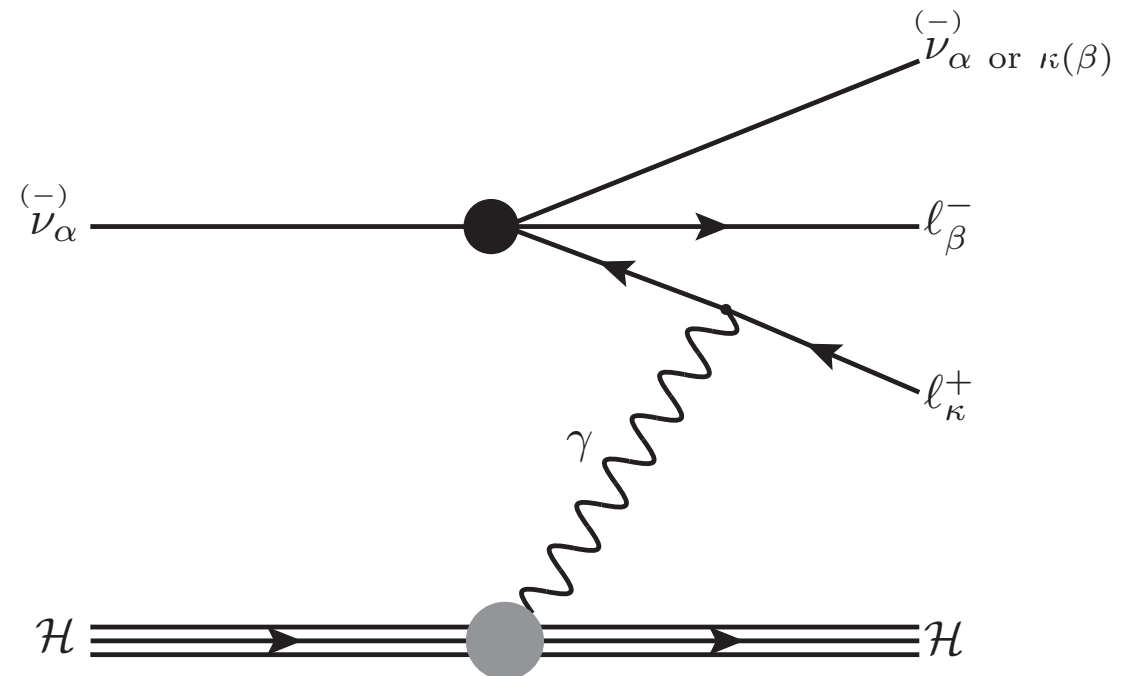
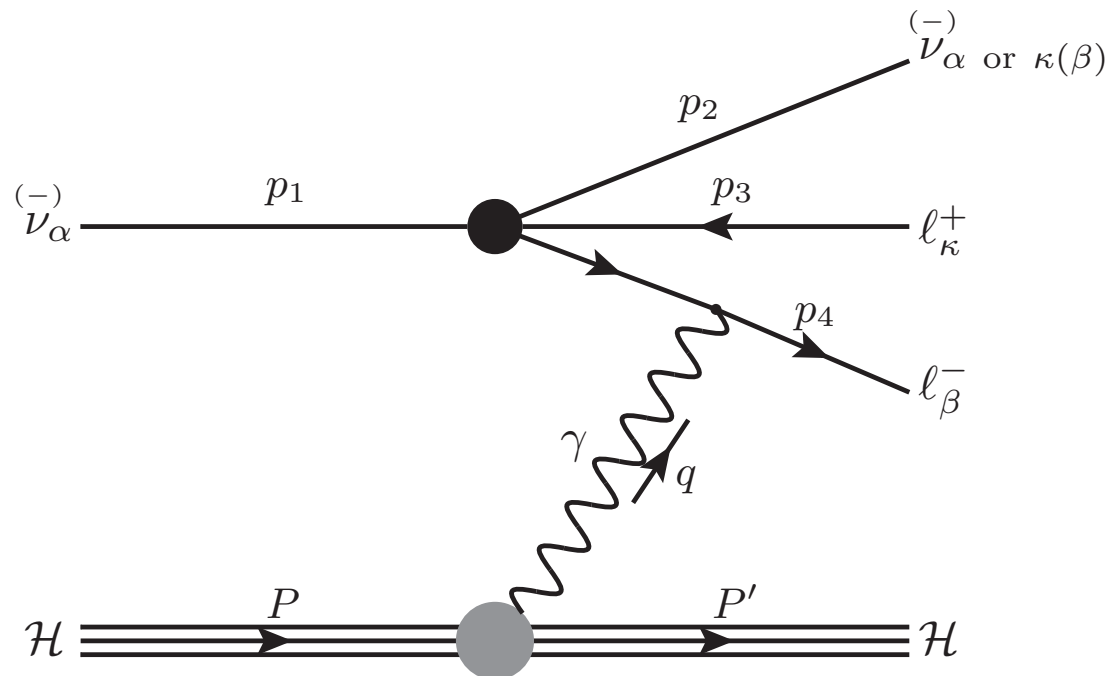
Equivalent Photon
Approximation
(EPA)



Overestimates the cross section in
some cases by more than 200%.

Trident Cross Sections

$$\nu_{\alpha}(p_1) + \mathcal{H}(P) \rightarrow \nu_{\alpha \text{ or } \kappa(\beta)}(p_2) + \ell_{\beta}^{-}(p_4) + \ell_{\kappa}^{+}(p_3) + \mathcal{H}(P')$$



(Anti)Neutrino	SM Contributions
$\bar{\nu}_{\mu} \mathcal{H} \rightarrow \bar{\nu}_{\mu} \mu^{-} \mu^{+} \mathcal{H}$	CC + NC
$\bar{\nu}_{\mu} \mathcal{H} \rightarrow \bar{\nu}_{e} e^{\pm} \mu^{\mp} \mathcal{H}$	CC
$\bar{\nu}_{\mu} \mathcal{H} \rightarrow \bar{\nu}_{\mu} e^{-} e^{+} \mathcal{H}$	NC
$\bar{\nu}_{e} \mathcal{H} \rightarrow \bar{\nu}_{e} e^{-} e^{+} \mathcal{H}$	CC + NC
$\bar{\nu}_{e} \mathcal{H} \rightarrow \bar{\nu}_{\mu} \mu^{\pm} e^{\mp} \mathcal{H}$	CC
$\bar{\nu}_{e} \mathcal{H} \rightarrow \bar{\nu}_{e} \mu^{-} \mu^{+} \mathcal{H}$	NC

4-point limit

→ Observed!

$$V_{\alpha\beta\kappa}(A_{\alpha\beta\kappa}) \equiv g_V^{\beta}(g_A^{\beta})\delta_{\beta\kappa} + \delta_{\alpha\beta}$$

Interference between
CC and NC

Trident Production


Differential
cross sections

$$\frac{d^2\sigma_{\nu X}}{dQ^2 d\hat{s}} = \frac{1}{32\pi^2(s - M_{\mathcal{H}}^2)^2} \frac{H_X^{\mu\nu} L_{\mu\nu}}{Q^4}$$

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Momentum Transfer 

$$s \equiv (p_1 + P)^2$$

Center-of-mass energy
of the total system

$$\hat{s} \equiv (p_1 + q)^2$$

Center-of-mass energy
of the neutrino-photon
system

Trident Production

Differential
cross sections

It depends on scattering regime

Hadronic Tensor

$$\frac{d^2\sigma_{\nu X}}{dQ^2 d\hat{s}} = \frac{1}{32\pi^2(s - M_{\mathcal{H}}^2)^2} \frac{H_X^{\mu\nu} L_{\mu\nu}}{Q^4}$$

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Center-of-mass energy
of the neutrino-photon
system

We can separate
photon contributions

$$\frac{d^2\sigma_{\nu X}}{dQ^2 d\hat{s}} = \frac{1}{32\pi^2} \frac{1}{\hat{s} Q^2} \left[h_X^T(Q^2, \hat{s}) \sigma_{\nu\gamma}^T(Q^2, \hat{s}) + h_X^L(Q^2, \hat{s}) \sigma_{\nu\gamma}^L(Q^2, \hat{s}) \right]$$

Transversal

Longitudinal

Scattering Regimes

Coherent

Scattering Regimes

Coherent

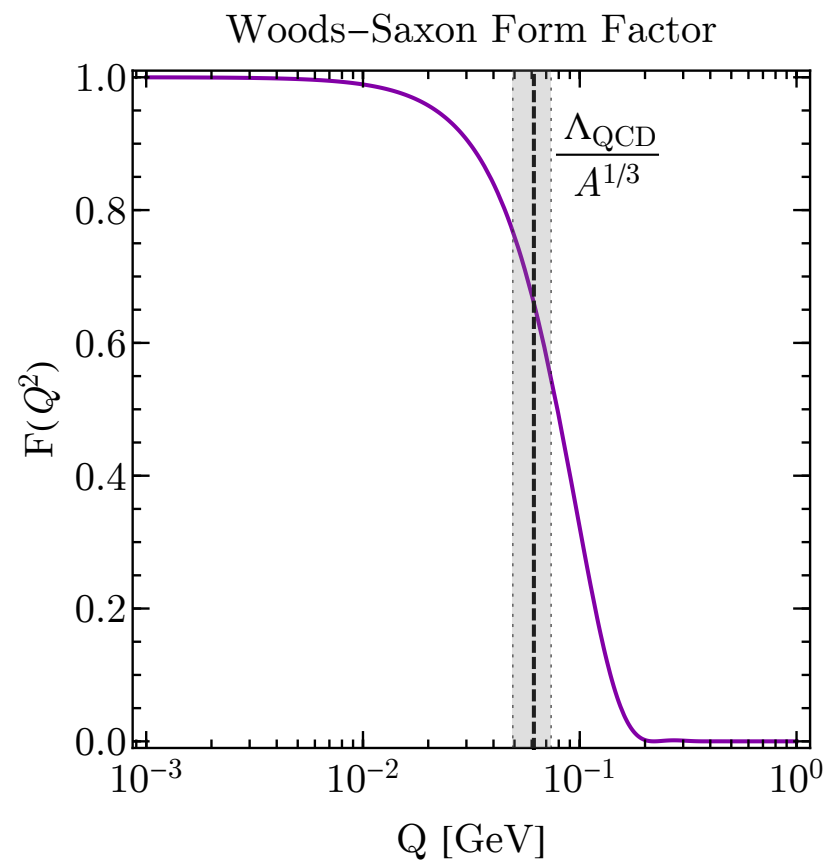
$$H_c^{\mu\nu} = 4Z^2 e^2 \left| F(Q^2) \right|^2 \left(P^\mu - \frac{q^\mu}{2} \right) \left(P^\nu - \frac{q^\nu}{2} \right)$$

Scattering Regimes

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Cut on Q

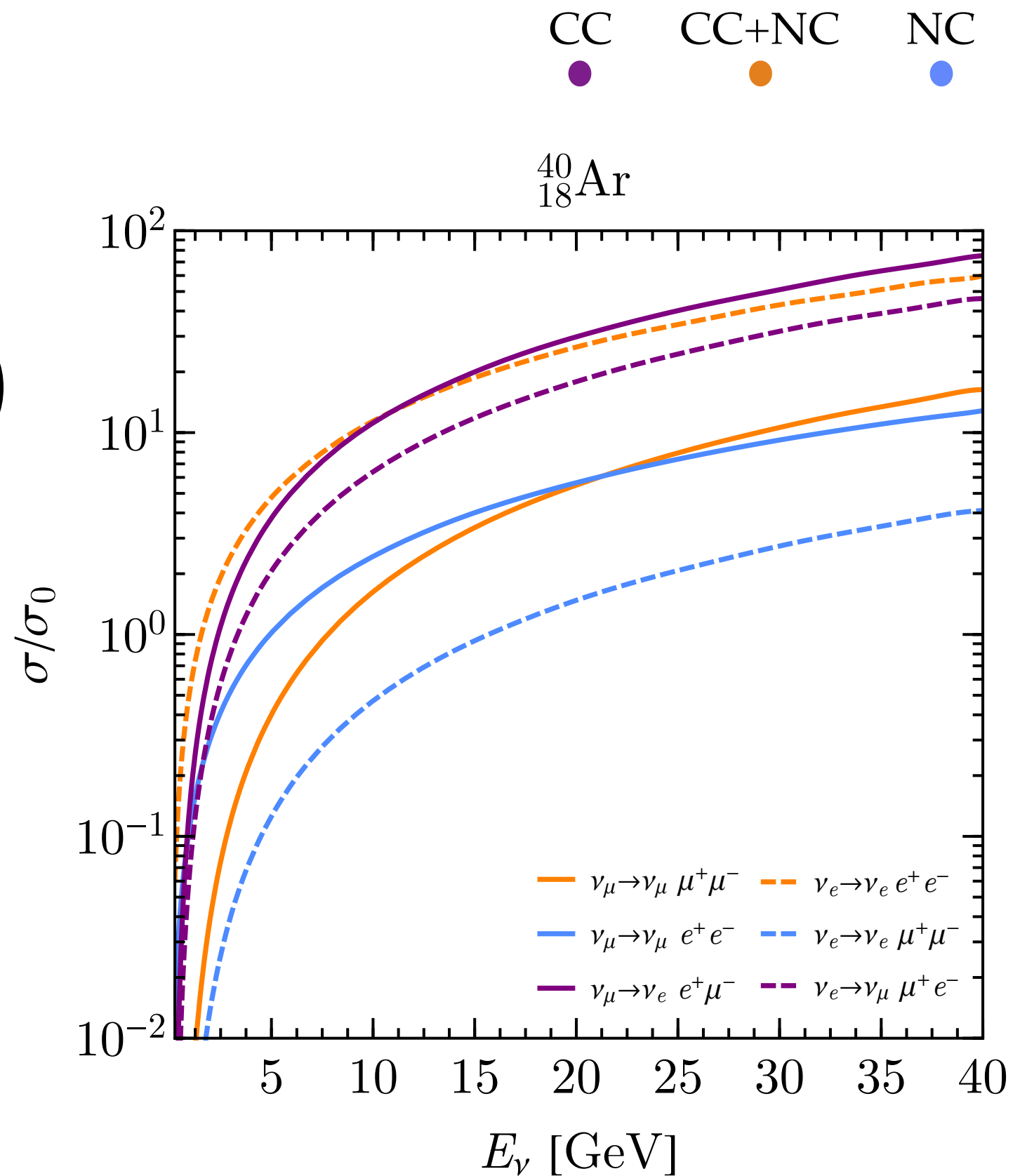
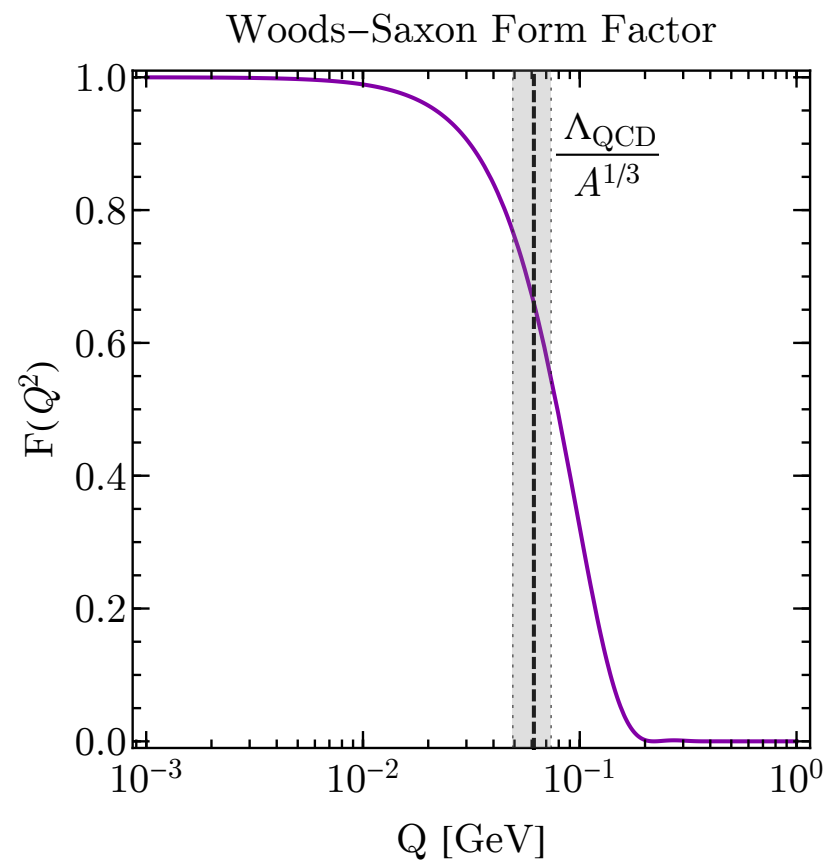


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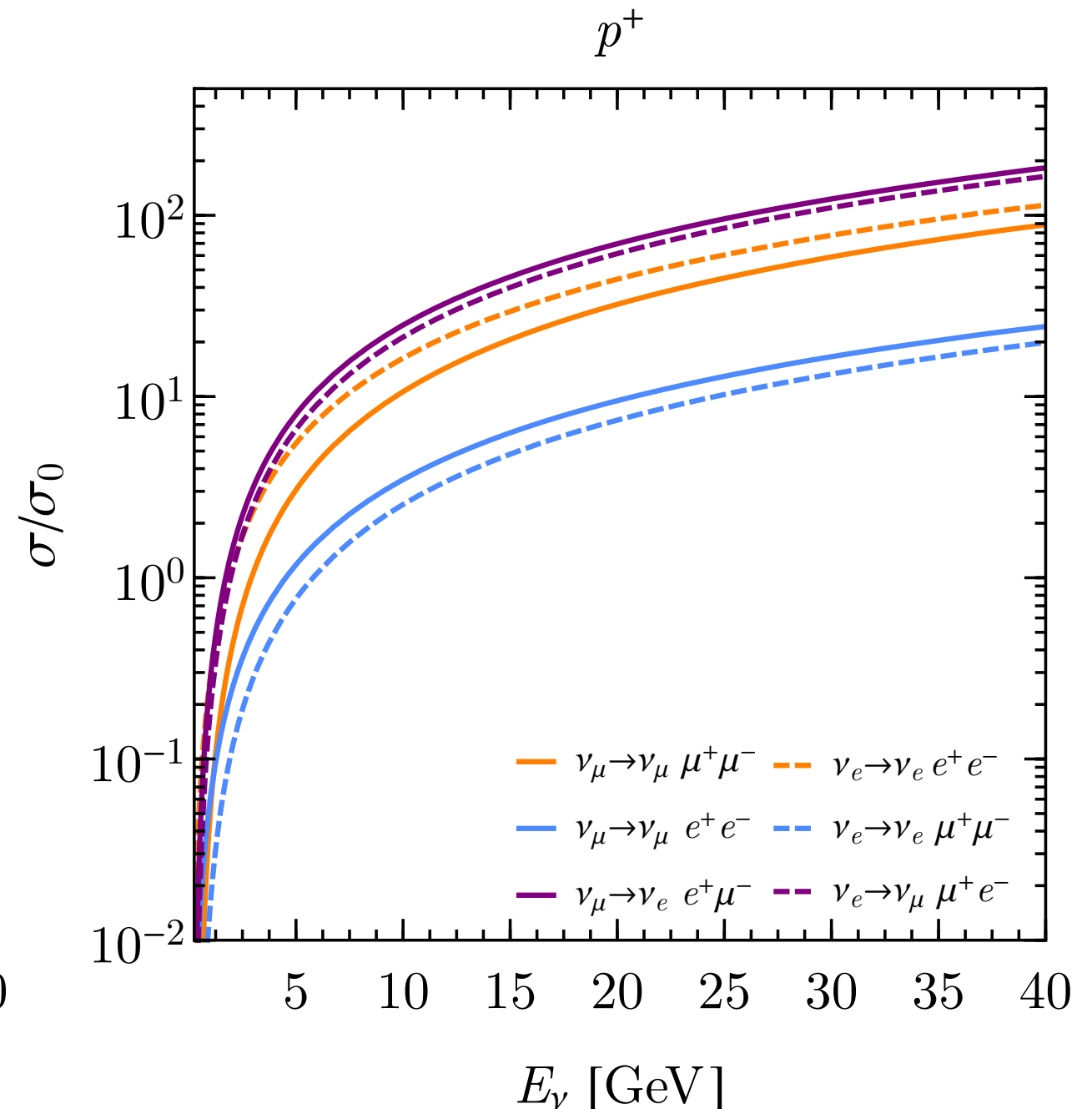
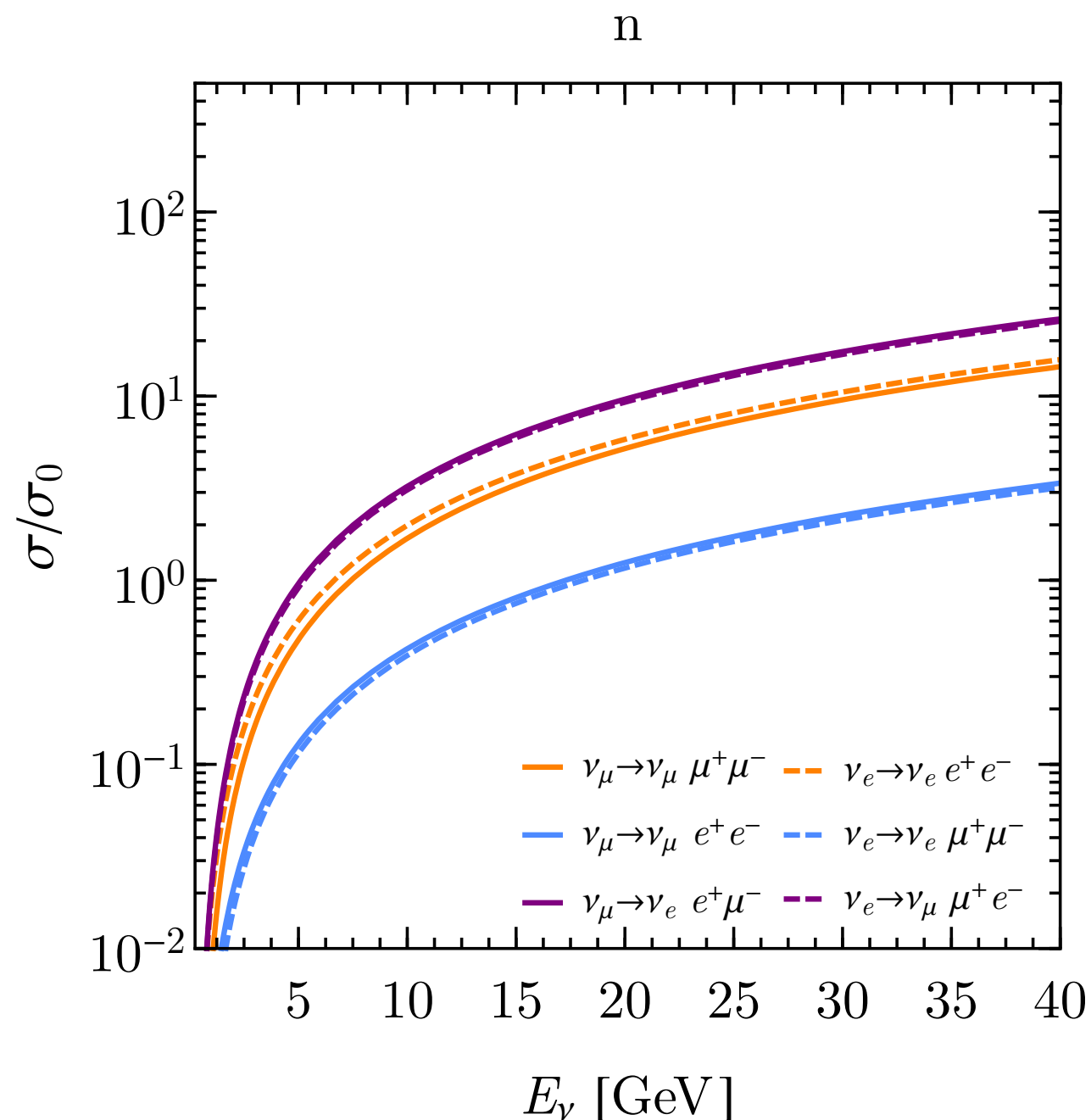
$$\sigma_0 = Z^2 \cdot 10^{-44} \text{ cm}^2$$

Scattering Regimes

Diffraction

$$\frac{d^2\sigma_{\nu d}}{dQ^2 d\hat{s}} \rightarrow f(|\vec{q}|) \frac{d^2\sigma_{\nu d}}{dQ^2 d\hat{s}}$$

$$H_d^{\mu\nu}(P, P') = Z H_p^{\mu\nu}(P, P') + (A - Z) H_n^{\mu\nu}(P, P')$$

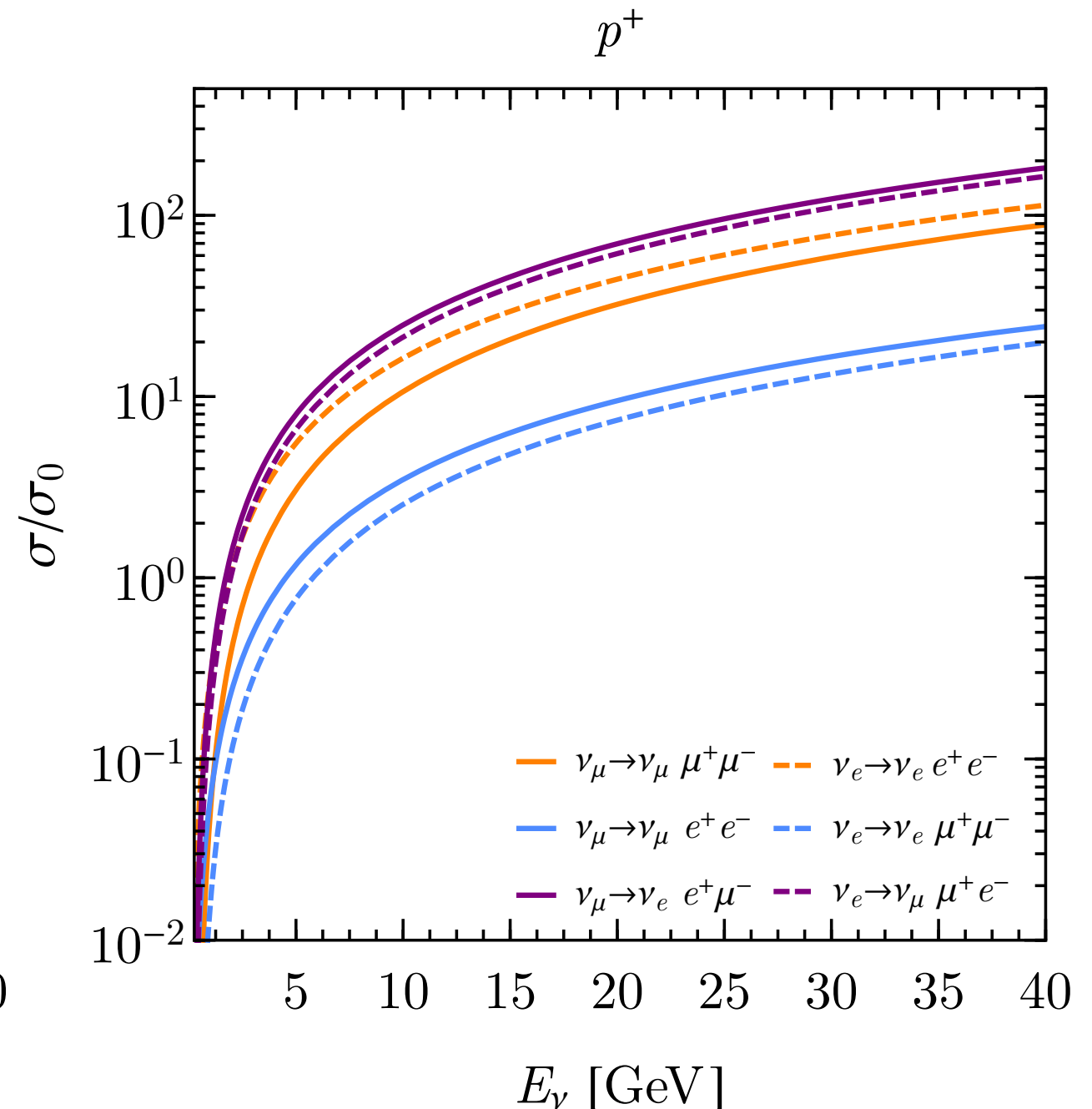
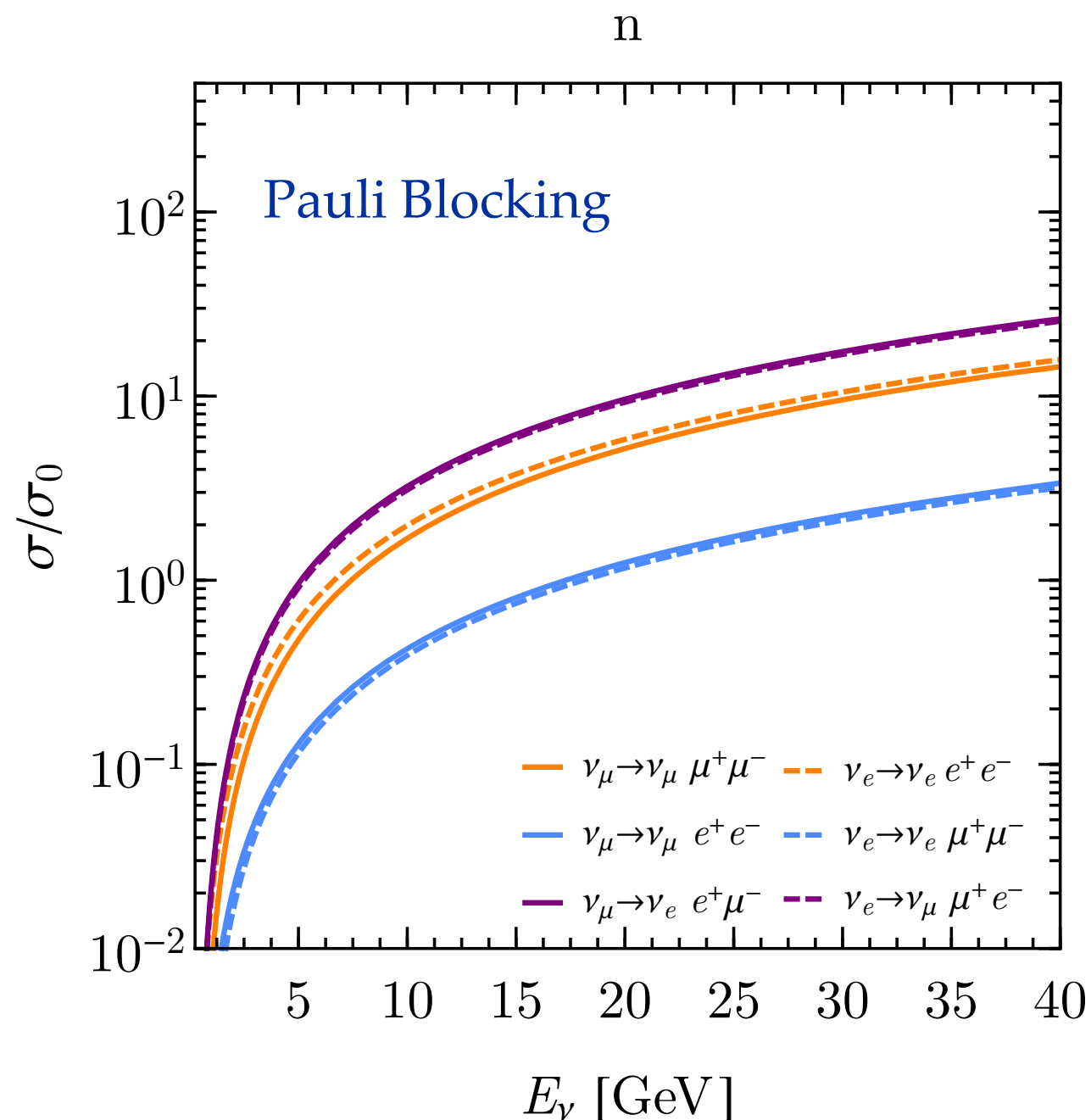


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Equivalent Photon Approximation (EPA)

Can we obtain an approximate cross section?

Fermi, 1924
Weizsacker, Williams, 1934

EPA

$$\sigma_t(P_i + C_s \rightarrow P_f + C_s) \approx \int dP(Q^2, \hat{s}) \sigma_\gamma(P_i + \gamma \rightarrow P_f; \hat{s}, Q^2 = 0)$$

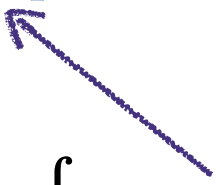
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Energy spectrum of
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Incoming particle
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Assumptions

- * Longitudinal contribution can be neglected
- * Transverse contribution can be taken as on-shell

$$\sigma_{\nu\gamma}^L(Q^2, \hat{s}) \approx 0$$

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Trident case

Photon spectrum

$$\sigma_{\text{EPA}} = \int_{m_L^2}^{\hat{s}_{\text{max}}} \int_{(\hat{s}/2E_\nu)^2}^{Q_{\text{max}}^2} \sigma_{\nu\gamma}^T(0, \hat{s}) dP(Q^2, \hat{s})$$

$$dP(Q^2, \hat{s}) = \frac{Z^2 e^2}{4\pi^2} |F(Q^2)|^2 \frac{d\hat{s}}{\hat{s}} \frac{dQ^2}{Q^2}$$

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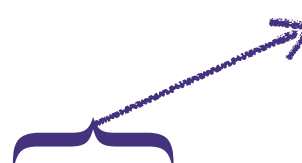
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Kinematical Limits



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QED

Fermi Limit of the SM

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QED

$$\sigma_\gamma^{\text{QED}}(P_i + \gamma \rightarrow P_f; \hat{s}, 0) \propto \frac{1}{\hat{s}}$$

Decreases with
increasing transferred
four-momentum

On-shell \gg off-shell

Fermi Limit of the SM

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Fermi Limit of the SM

$$\sigma_\gamma^{\text{FL}}(P_i + \gamma \rightarrow P_f; \hat{s}, 0) \propto G_F^2 \hat{s}$$

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On-shell \ll off-shell

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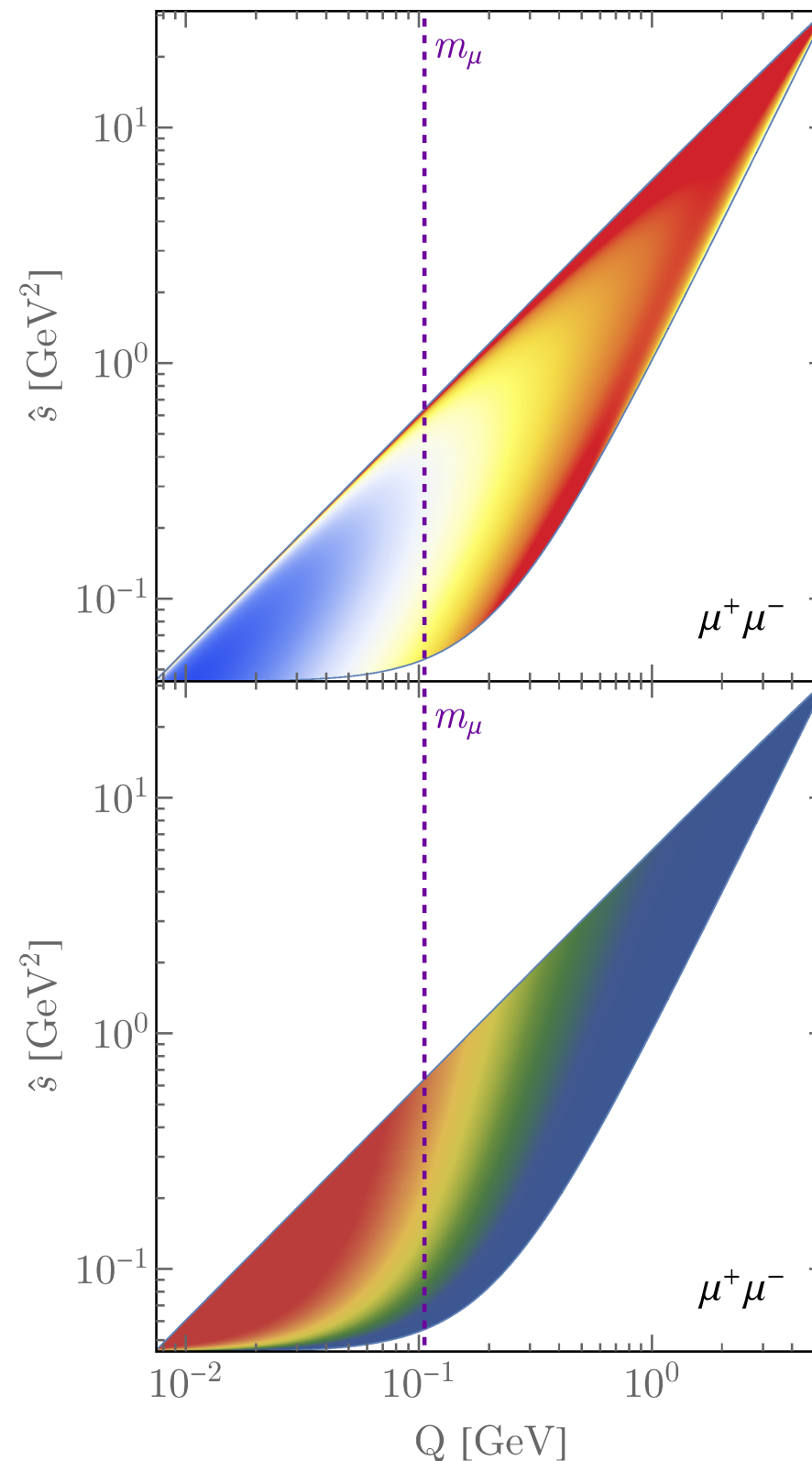
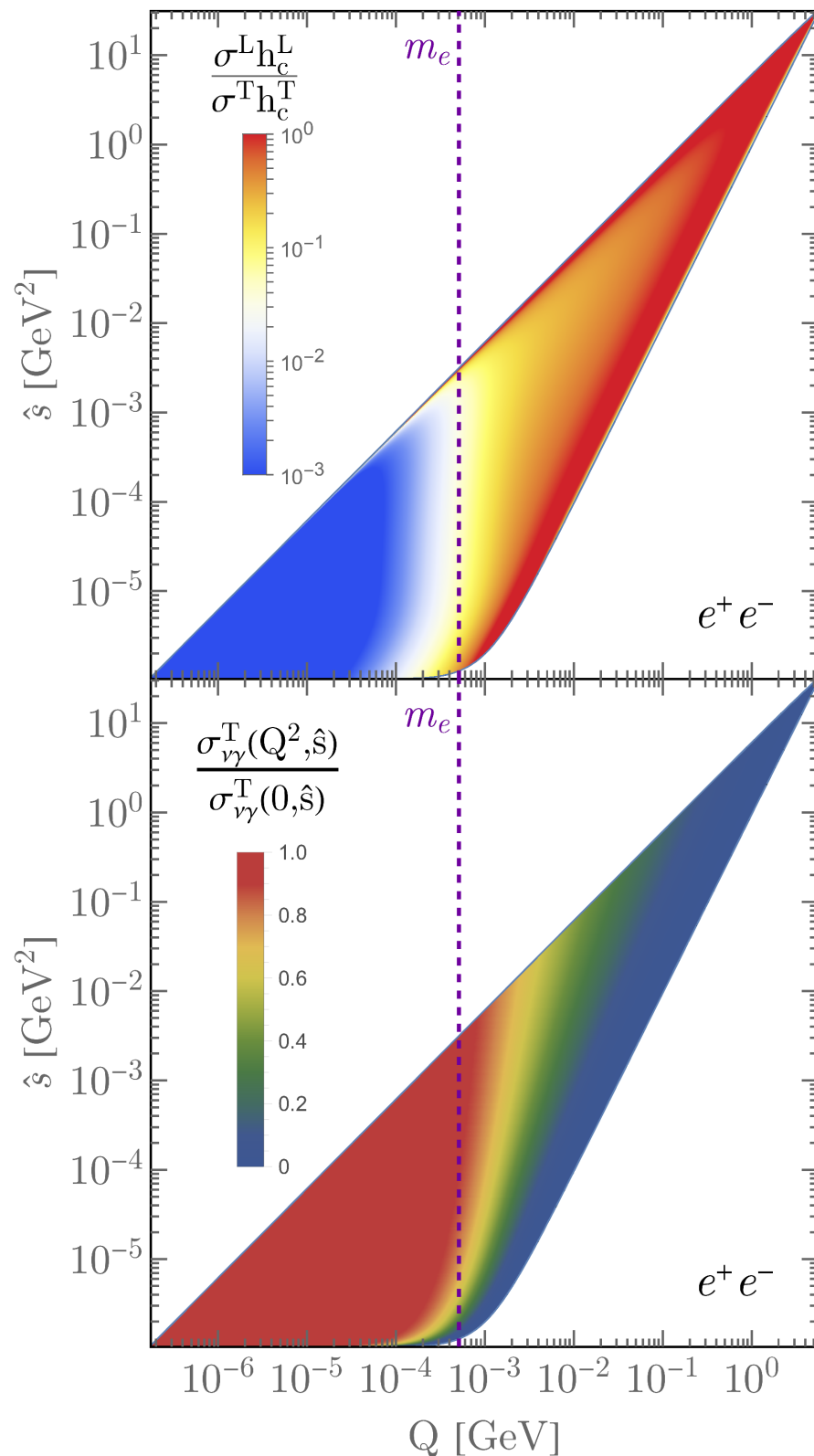
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EPA not valid for trident, right?

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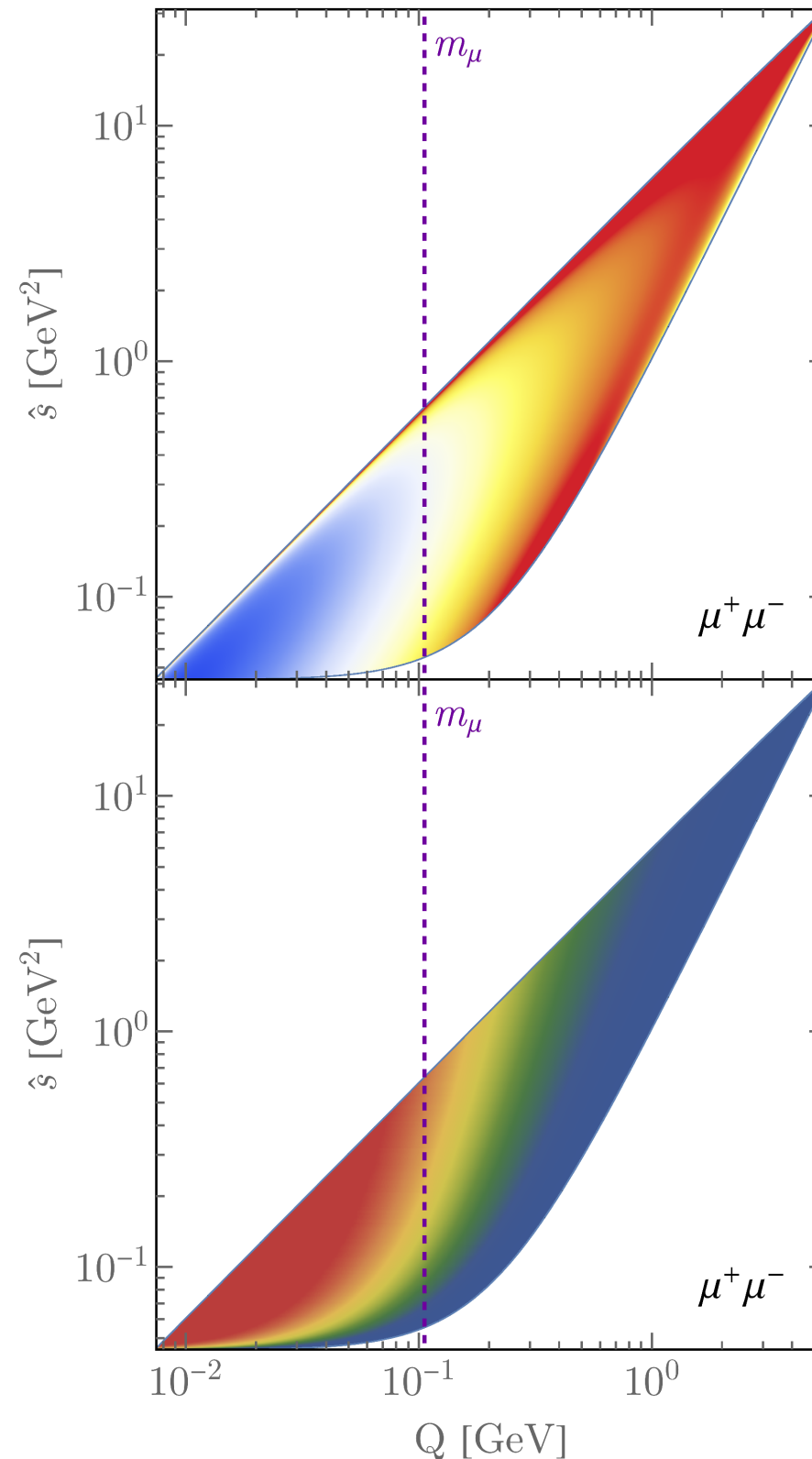
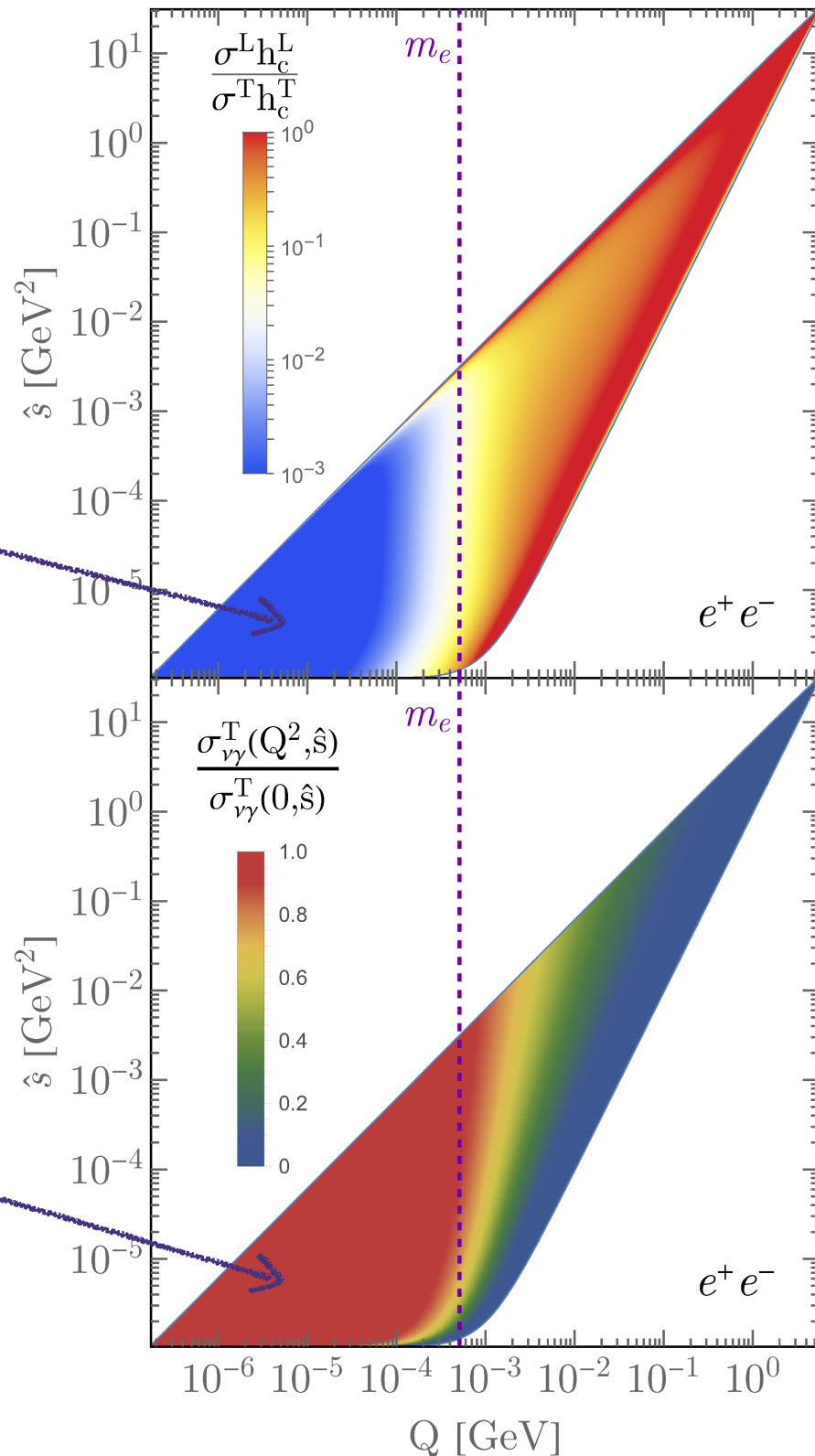
$$\frac{\sigma^L(Q^2, \hat{s}) h_c^L(Q^2, \hat{s})}{\sigma^T(Q^2, \hat{s}) h_c^T(Q^2, \hat{s})}$$

Kinematically
allowed region

$$E_\nu = 3 \text{ GeV}$$

EPA

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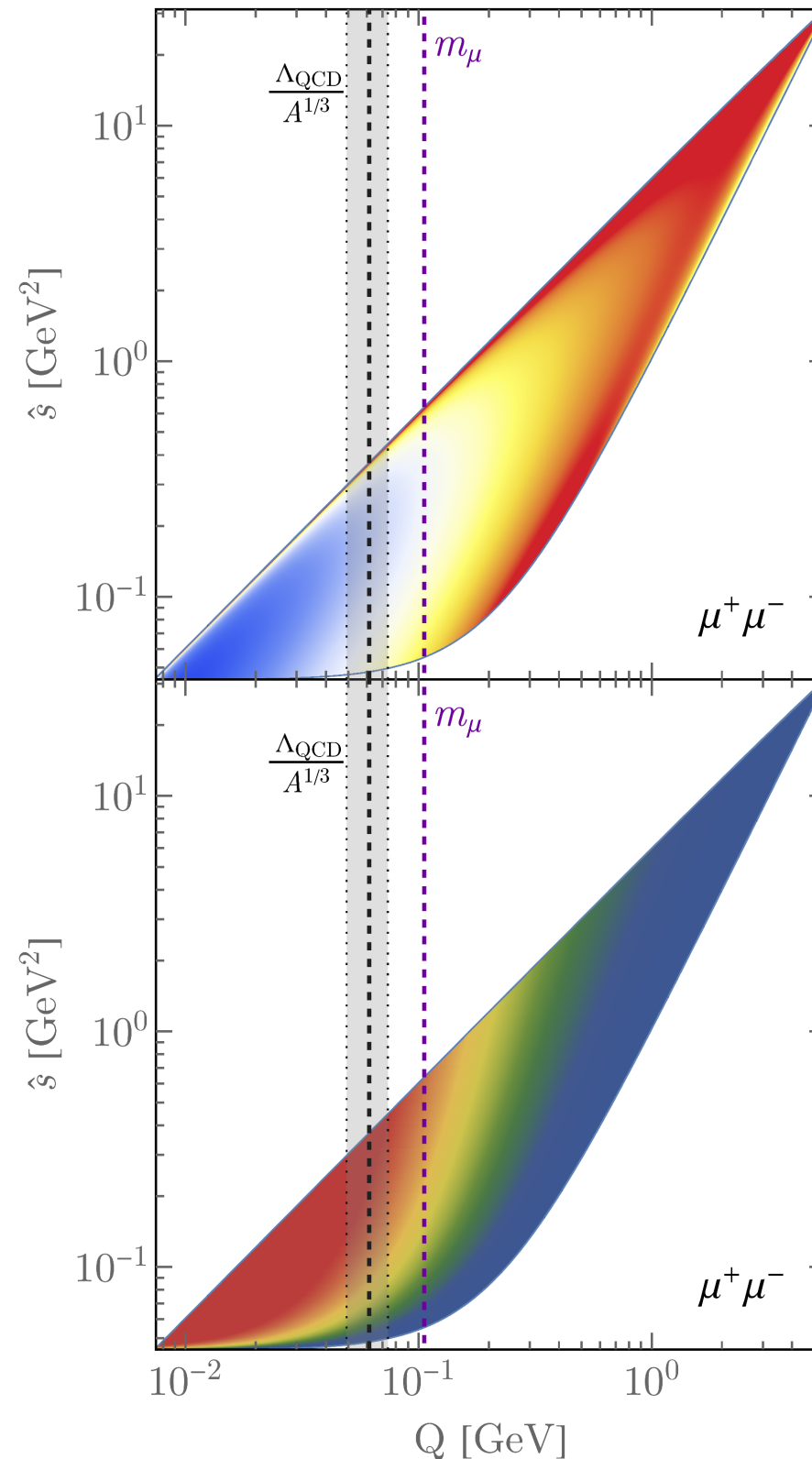
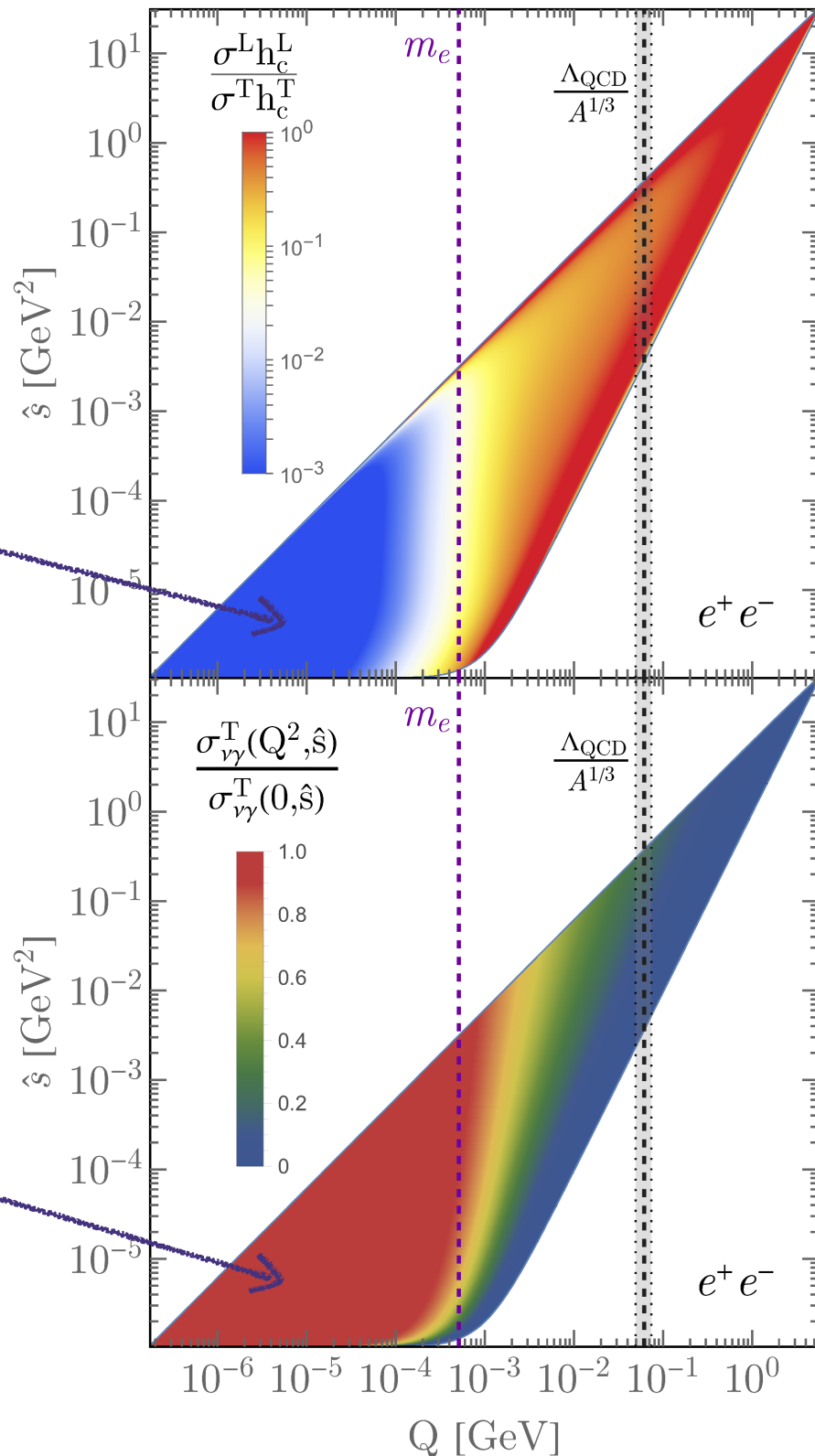
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Physical cutoff on the
momentum integration

$$\frac{\sigma_{\nu\gamma}^T(Q^2, \hat{s})}{\sigma_{\nu\gamma}^T(0, \hat{s})}$$

$$\frac{\sigma^L(Q^2, \hat{s}) h_c^L(Q^2, \hat{s})}{\sigma^T(Q^2, \hat{s}) h_c^T(Q^2, \hat{s})}$$

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allowed region



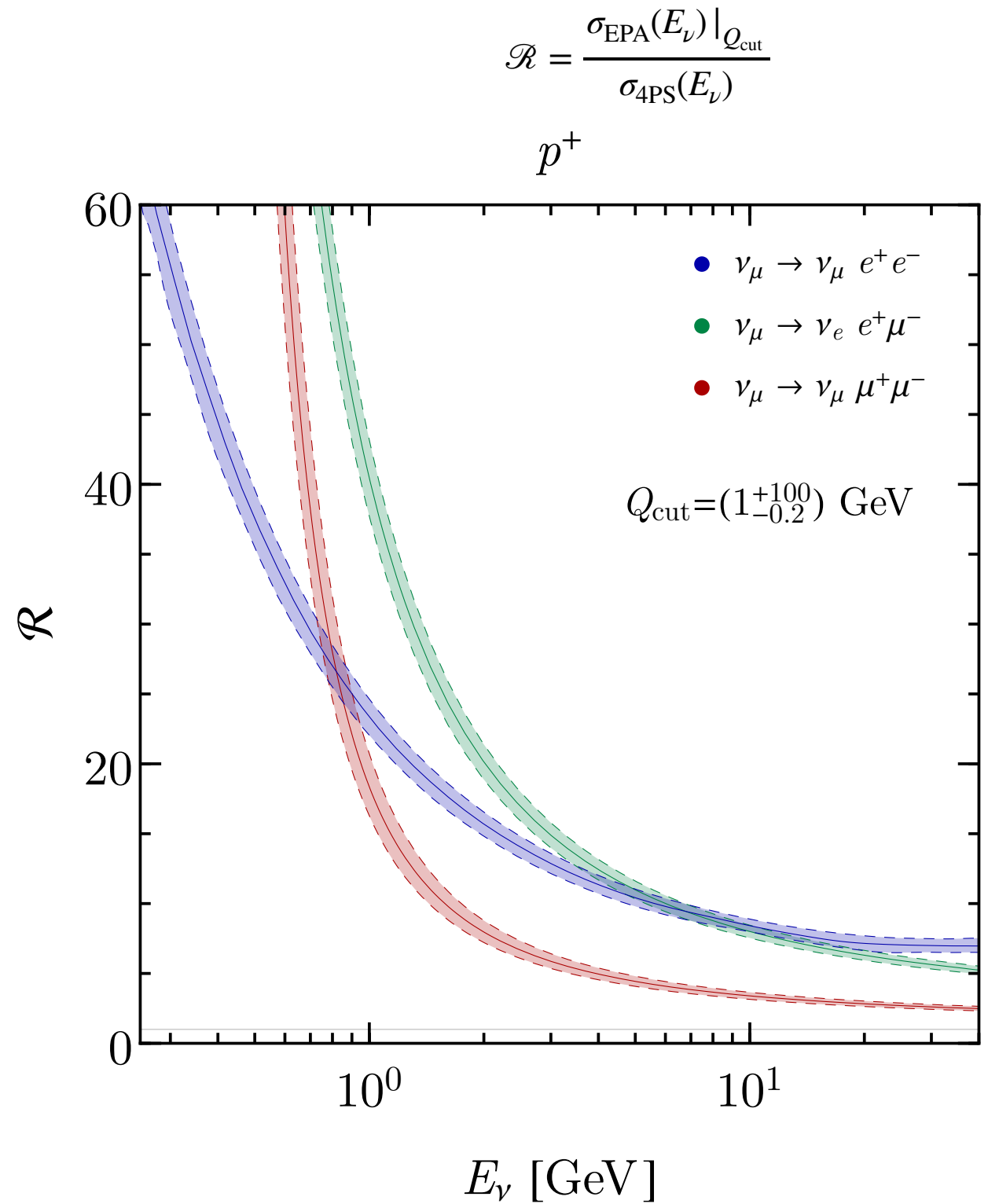
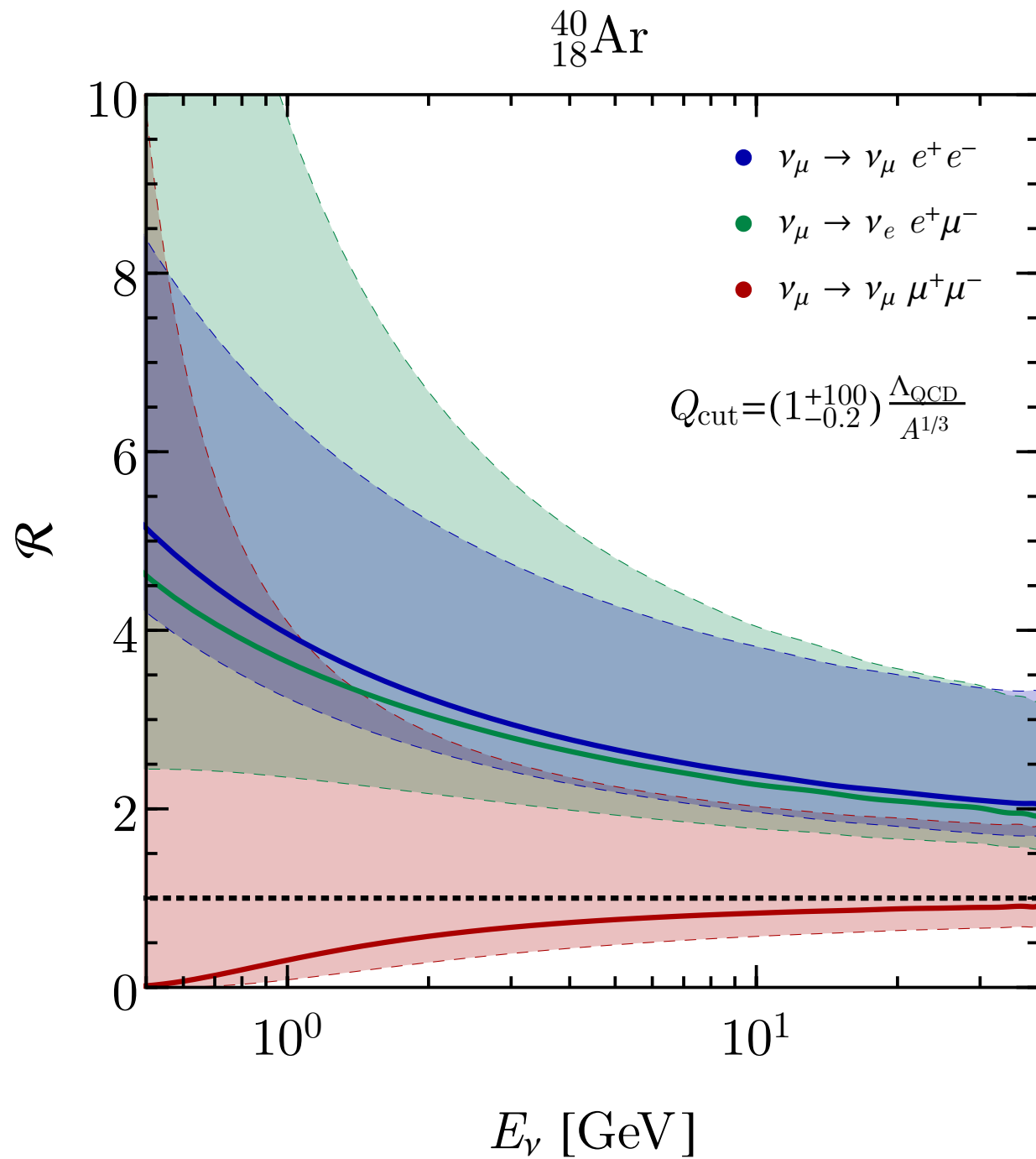
Form factor cuts
the Q integration
in a region which
EPA works
reasonably

$$E_\nu = 3 \text{ GeV}$$

EPA

EPA not valid for trident, right?

EPA not valid for trident, right?

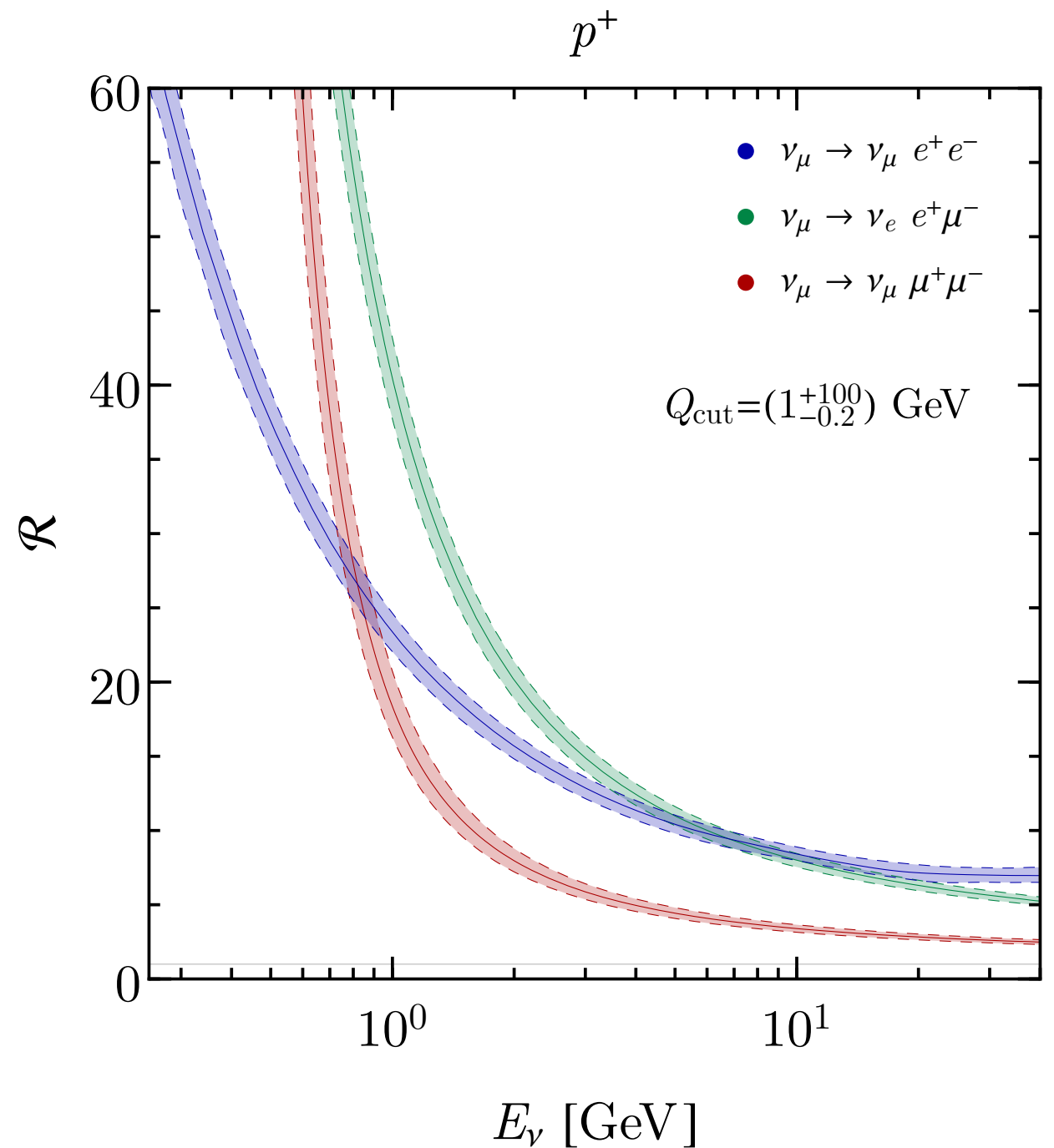
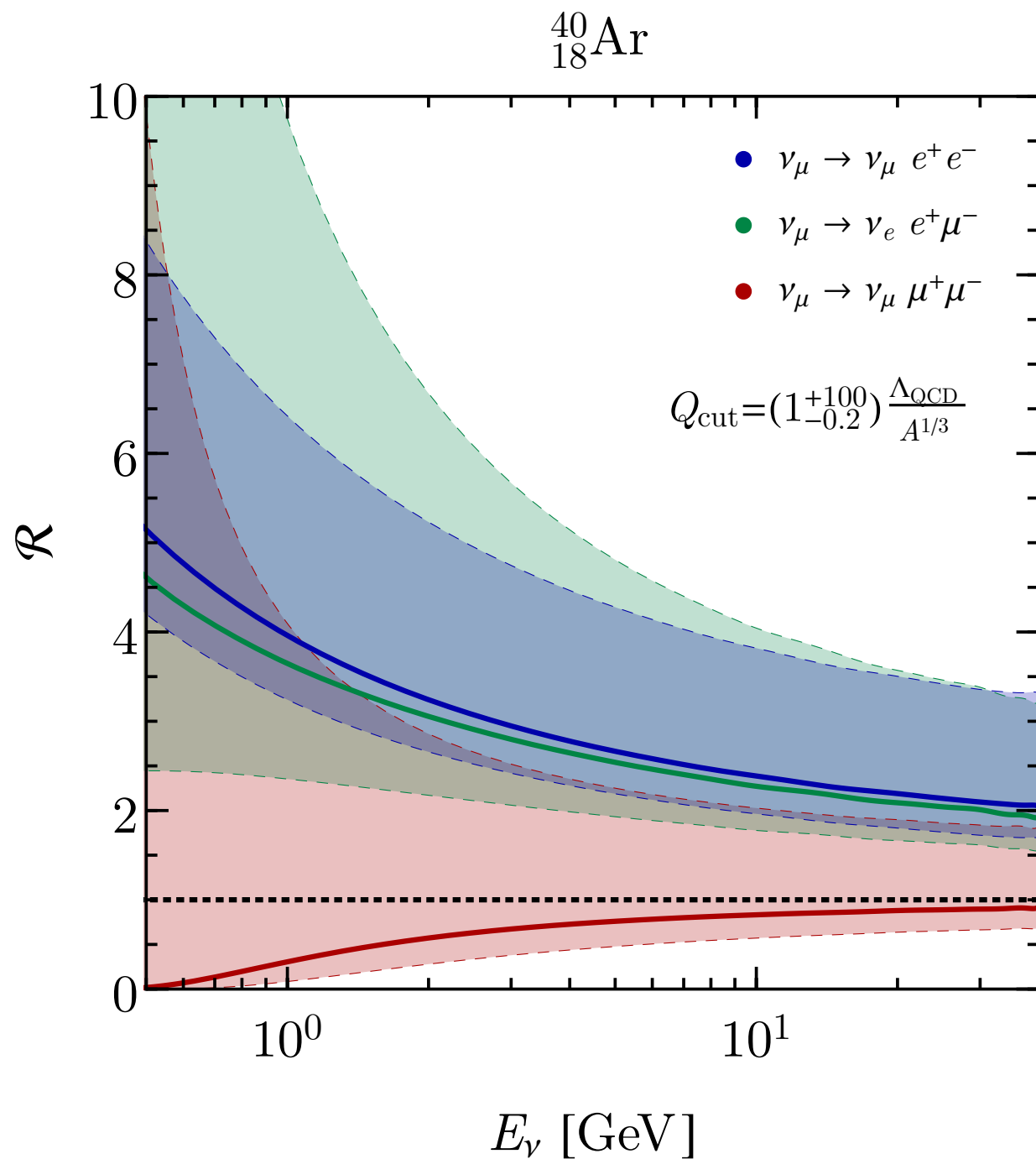


EPA

EPA not valid for trident, right?

Yes.

$$\mathcal{R} = \frac{\sigma_{\text{EPA}}(E_\nu)|_{Q_{\text{cut}}}}{\sigma_{4\text{PS}}(E_\nu)}$$



Compute the full 4PS cross sections

Events in LAr Detectors

Rates

$$N_X^\psi = \text{Norm} \times \int dE_\nu \sigma_{\nu X}(E_\nu) \frac{d\phi_\nu(E_\nu)}{dE_\nu} \epsilon(E_\nu)$$

Rates

$$N_X^\psi = \text{Norm} \times \int dE_\nu \sigma_{\nu X}(E_\nu) \frac{d\phi_\nu(E_\nu)}{dE_\nu} \epsilon(E_\nu)$$


$$\text{Exposure [POT]} \times \frac{\text{Fiducial Detector Mass} \times N_A}{m_T} [\text{target particles}]$$

Assuming 100%
efficiency

Rates

$$N_X^\psi = \text{Norm} \times \int dE_\nu \sigma_{\nu X}(E_\nu) \frac{d\phi_\nu(E_\nu)}{dE_\nu} \epsilon(E_\nu)$$

$\text{Exposure [POT]} \times \frac{\text{Fiducial Detector Mass} \times N_A}{m_T} \text{ [target particles]}$

Assuming 100% efficiency

Experiment	Baseline (m)	Total Exposure (POT)	Fiducial Mass (t)	E _ν (GeV)
SBND	110	6.6 × 10 ²⁰	112	0 – 3
μBooNE	470	1.32 × 10 ²¹	89	0 – 3
ICARUS	600	6.6 × 10 ²⁰	476	0 – 3
DUNE	574	12.81 (12.81) × 10 ²¹	50	0 – 40
νSTORM	50	10 ²¹	100	0 – 6

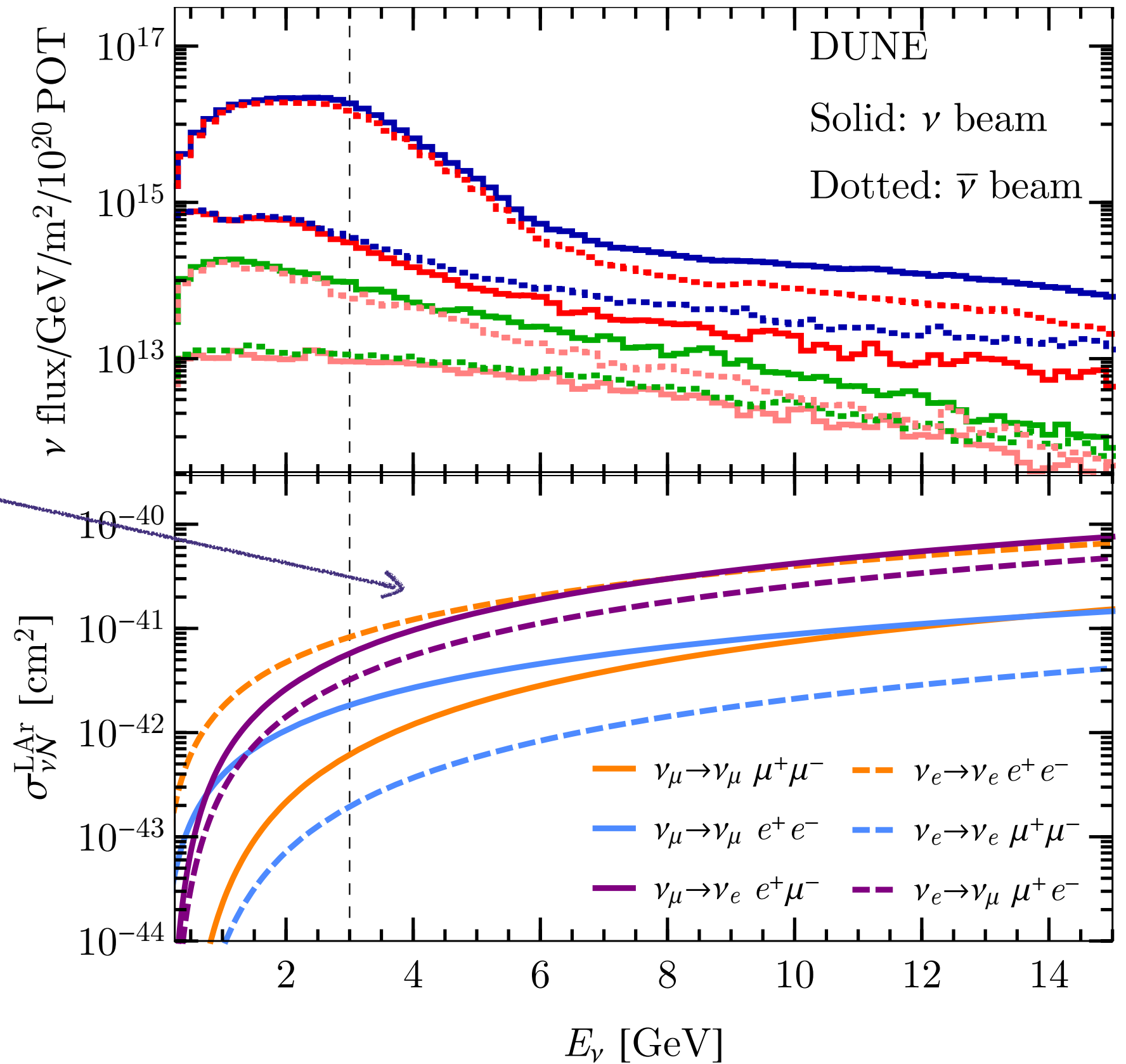
ν mode

ν̄ mode

$\text{Exposure} = 1.83 \times (3 + 2 \times 2) \times 10^{21} \text{ POT}$

Rates

Higher cross sections at DUNE ND



Rates

Exposure = $1.83 \times (3 + 2 \times 2) \times 10^{21}$ POT

Not seen yet

Not seen yet

Channel	SBND	μ BooNE	ICARUS	DUNE ND	ν STORM ND
Total $e^\pm \mu^\mp$	10	0.7	1	2993 (2307)	191
	1	0.1	0.1	391 (299)	23
Total $e^+ e^-$	6	0.4	0.7	1007 (800)	114
	0.2	0.0	0.02	64 (49)	6
Total $\mu^+ \mu^-$	0.4	0.0	0.0	286 (210)	11
	0.3	0.0	0.0	143 (108)	6

Coherent

Diffractive

Large contributions of diffractive events

Compare order of magnitudes

ν mode $\bar{\nu}$ mode

Kinematical Distributions

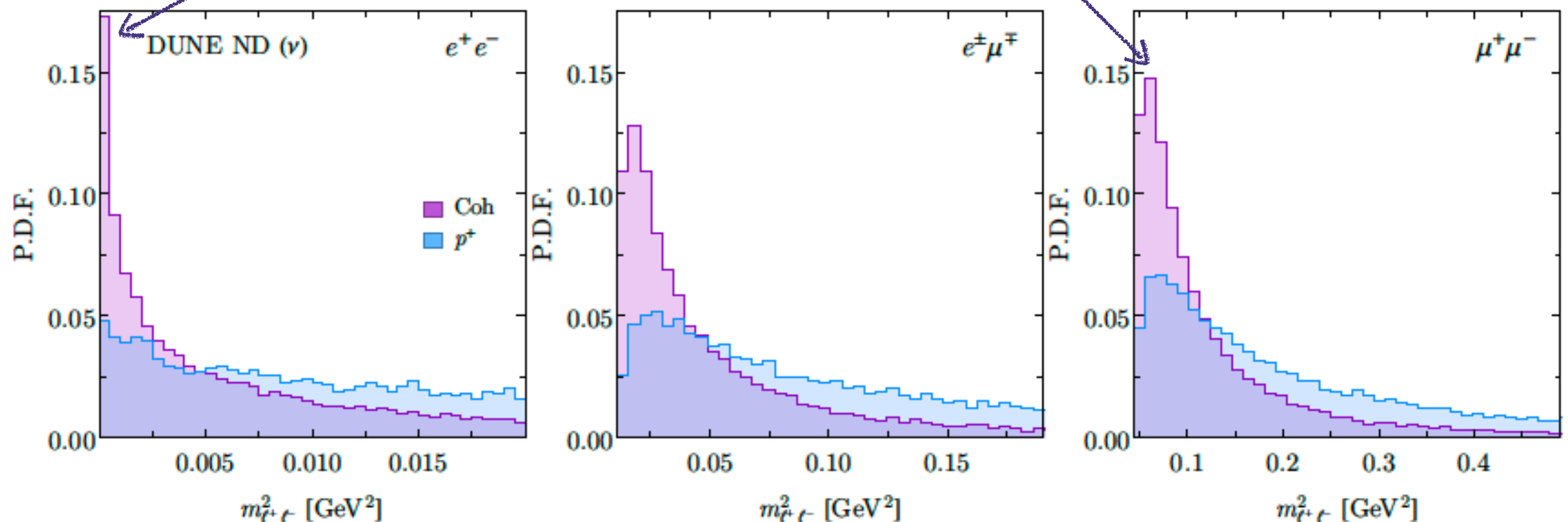
EPA gives also
misleading
distributions

Invariant charged lepton masses $m_{\ell^+\ell^-}^2$

Tool for
background
suppression

Small values

Peaked distributions

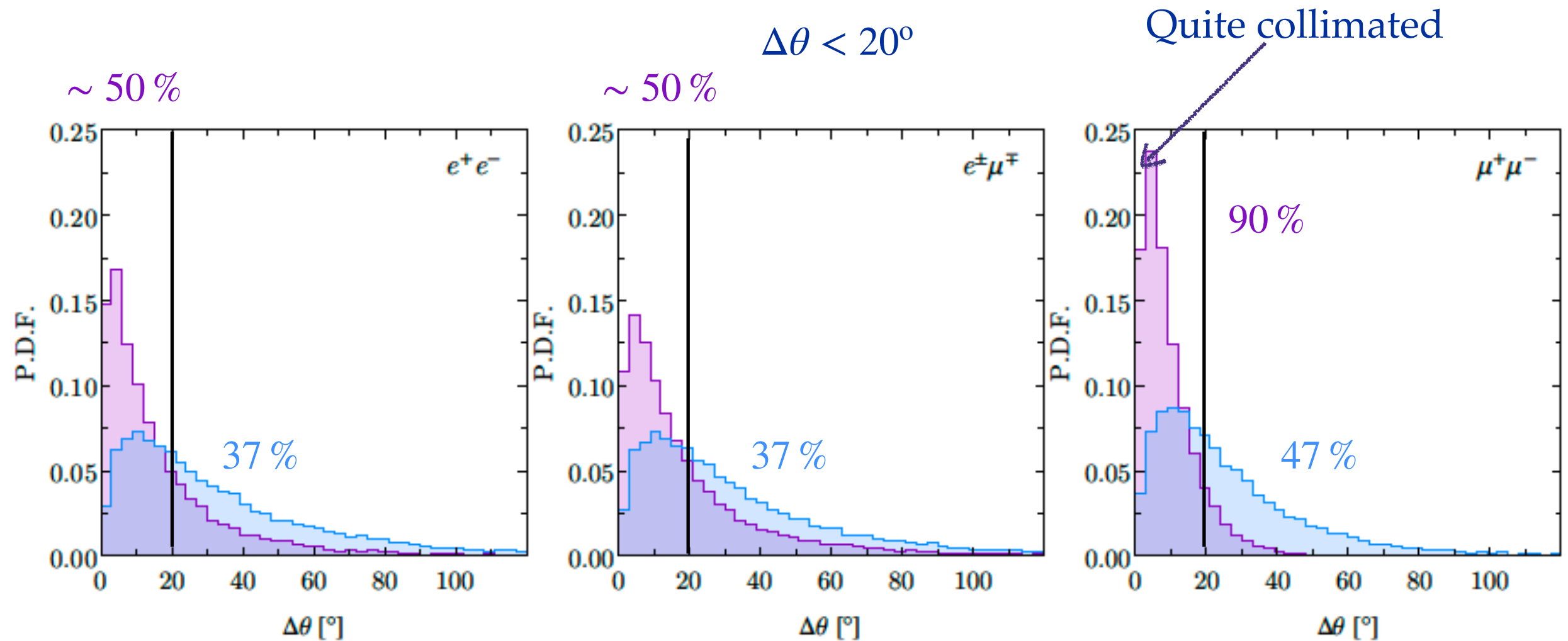


Flux integrated
distributions

All channels with same
final states

Kinematical Distributions

Separation angle $\Delta\theta$



Flux integrated
distributions

All channels with same
final states

Backgrounds?

Goal: Reach suppressions of order $\mathcal{O}(10^{-6} - 10^{-5})$

♦ misID

$$\mu^+\mu^- \longrightarrow \nu_\mu \text{CC} 1\pi^\pm$$

$$e^+e^- \longrightarrow \text{NC}\pi^0$$

$$e^\pm\mu^\mp \longrightarrow \text{CC}\pi^0$$

$$\text{CC}\gamma \quad \nu_e \text{CC}\pi^\pm$$

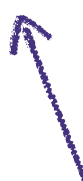
misID	Rate
γ as e^\pm	0.05
γ as e^+e^-	0.1 (w/ vertex) 1 (no vertex + overlapping)
π^\pm as μ^\pm	0.1

♦ No hadronic activity

♦ $m_{\mu^+\mu^-}^2 < 0.2 \text{ GeV}^2, \Delta\theta < 20^\circ, \theta_\pm < 15^\circ$

A more
careful
analysis is
needed

Channel	$N_B^{\text{misID}}/N_{\text{CC}}$	$N_B^{\text{had}}/N_{\text{CC}}$	$N_B^{\text{kin}}/N_{\text{CC}}$	$\epsilon_{\text{sig}}^{\text{coh}}$	$\epsilon_{\text{sig}}^{\text{dif}}$ ¹¹
$e^\pm\mu^\mp$	$1.67 (1.62) \times 10^{-4}$	$2.68 (4.31) \times 10^{-5}$	$4.40 (3.17) \times 10^{-7}$	0.61 (0.61)	0.39 (0.39)
e^+e^-	$2.83 (4.19) \times 10^{-4}$	$1.30 (2.41) \times 10^{-4}$	$6.54 (14.1) \times 10^{-6}$	0.48 (0.47)	0.21 (0.21)
$\mu^+\mu^-$	$2.66 (2.73) \times 10^{-3}$	$10.4 (9.75) \times 10^{-4}$	$3.36 (3.10) \times 10^{-8}$	0.66 (0.67)	0.17 (0.16)



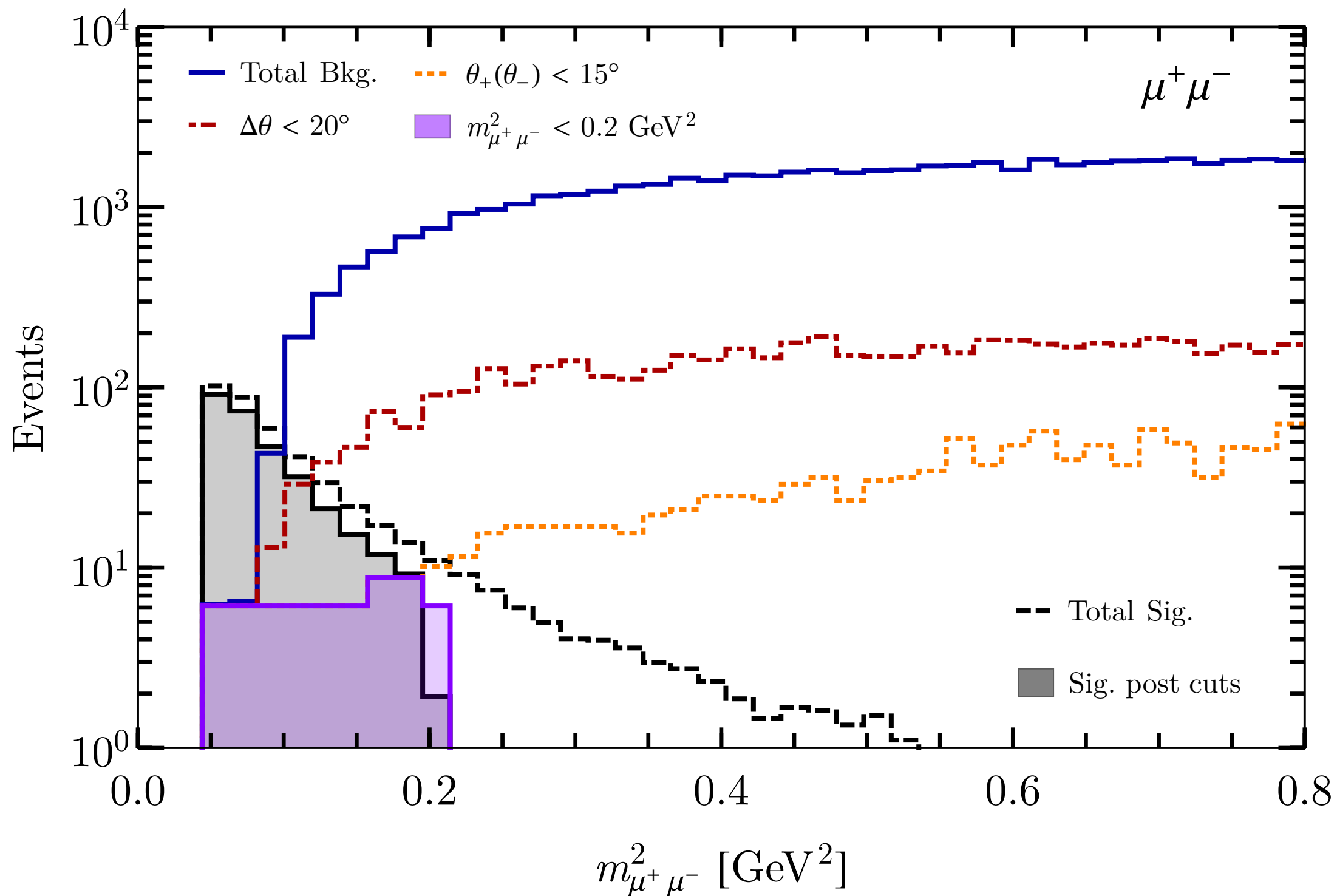
Efficiencies after cuts

misID

Hadronic veto

Kinematic cuts

Backgrounds?



GENIE simulation

Conclusions

- A full 4PS computation is required to obtain correct estimates for the number of events.
- Let us stress that the EPA gives a reasonable result for the dimuon channel due to a serendipitous behavior of the Form Factor.
- EPA can artificially suppress the coherent scattering contribution and increase the diffractive one giving rise to an incorrect rate and distributions of observable quantities.
- We have estimated the background for each trident channel via a Monte Carlo simulation using GENIE, and identified the dominant contributions arising primarily from particle misidentification.
- Reduction of ~ 6 orders of magnitude \times CC in the background is necessary to observe trident events at DUNE ND. A more careful analysis is needed.

Thank you!



Backup

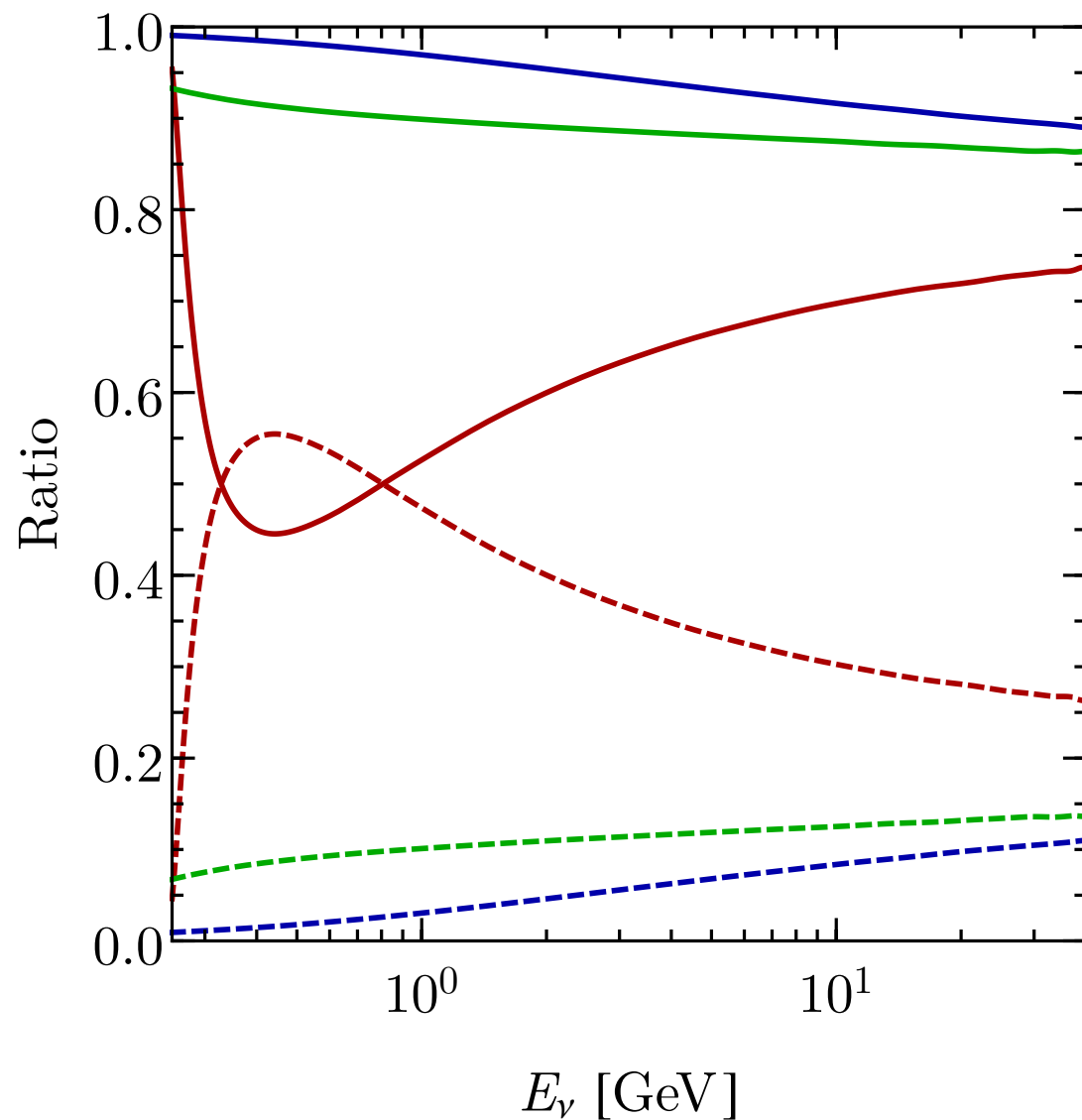
Coherent vs Diffractive

$$\sigma_{\nu\mathcal{N}} = \sigma_{\nu c} + \sigma_{\nu d}$$

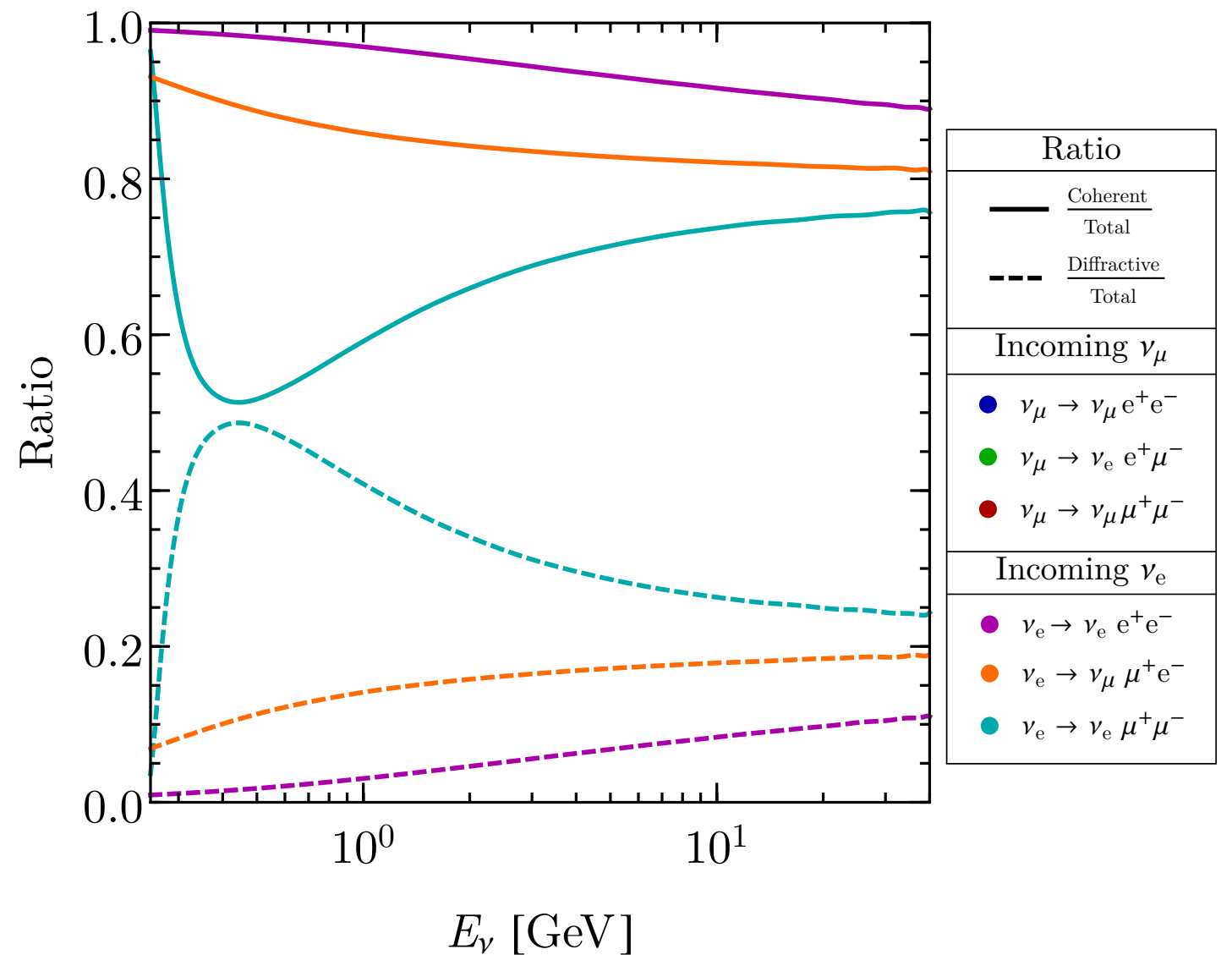
$$\frac{\sigma_{\nu c}}{\sigma_{\nu\mathcal{N}}}$$

$$\frac{\sigma_{\nu d}}{\sigma_{\nu\mathcal{N}}}$$

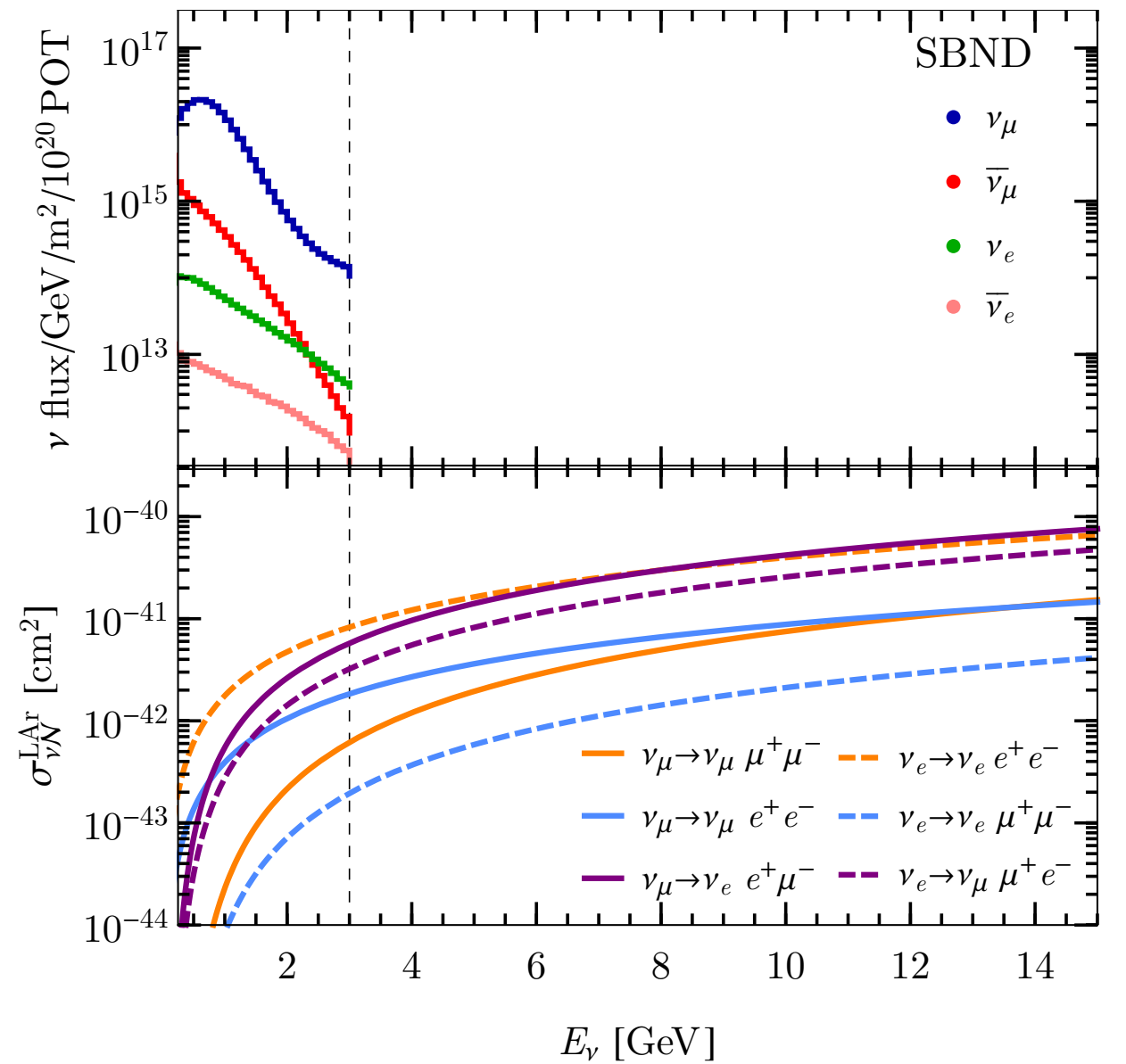
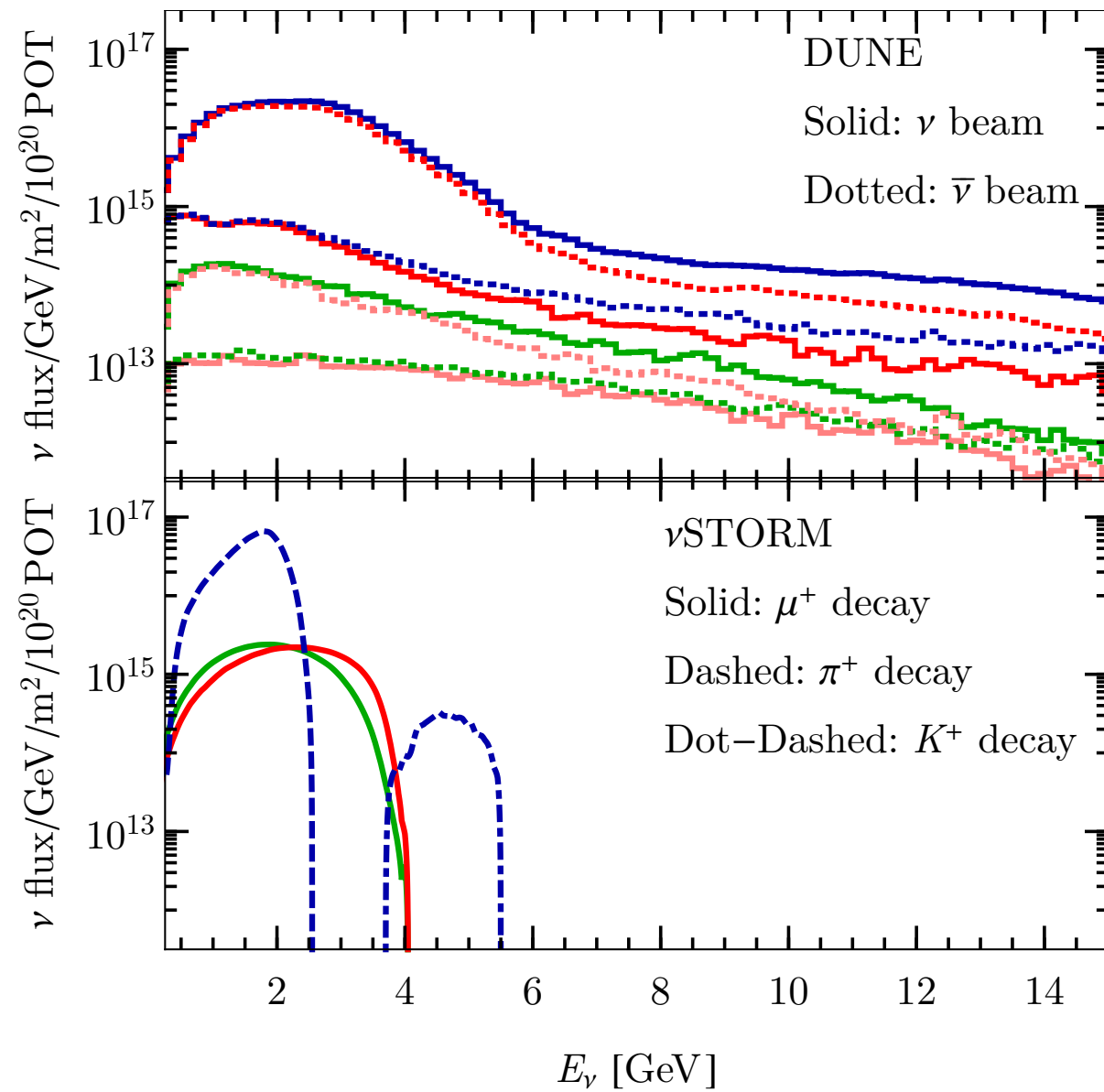
Incoming ν_μ flux, $^{40}_{18}\text{Ar}$



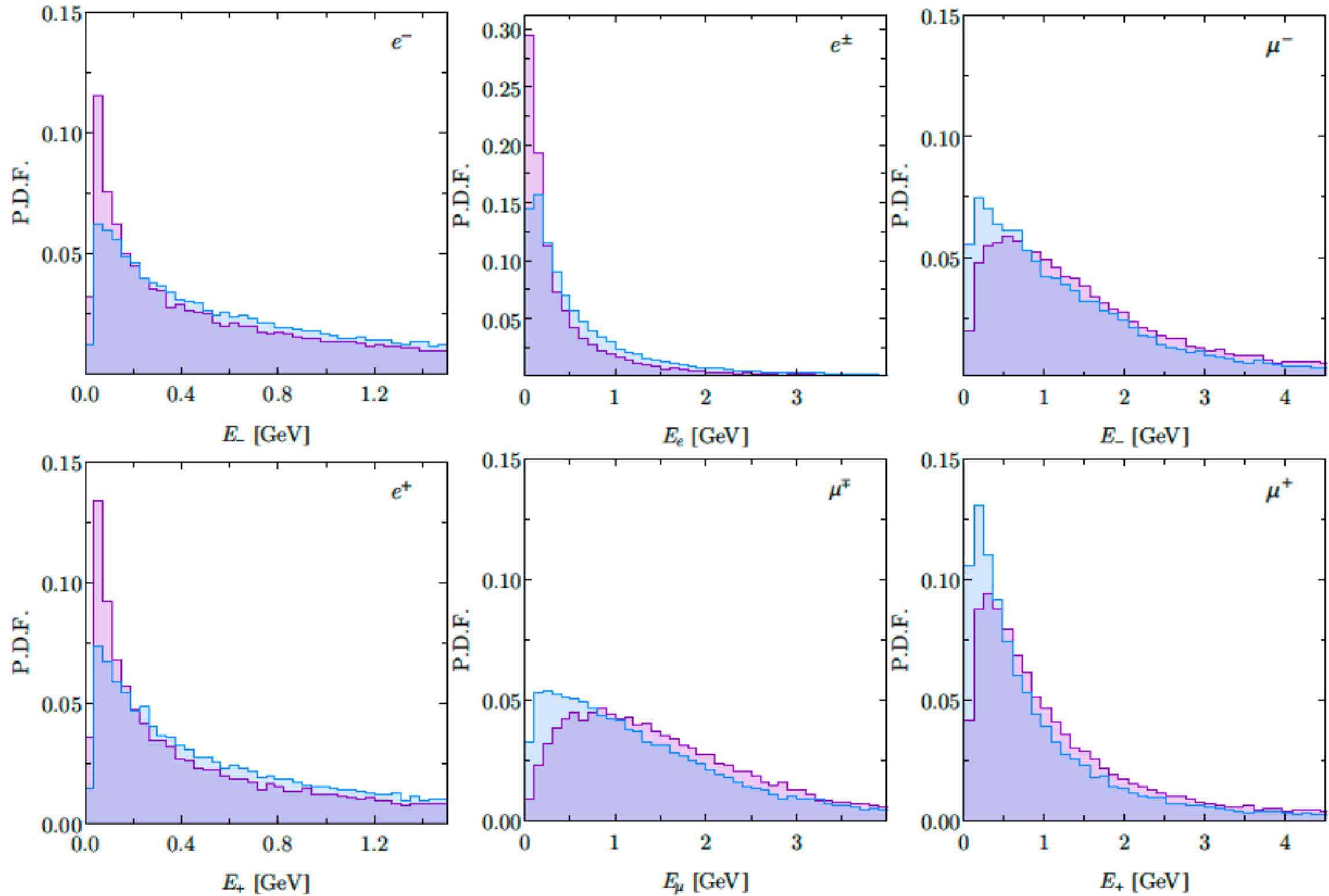
Incoming ν_e flux, $^{40}_{18}\text{Ar}$



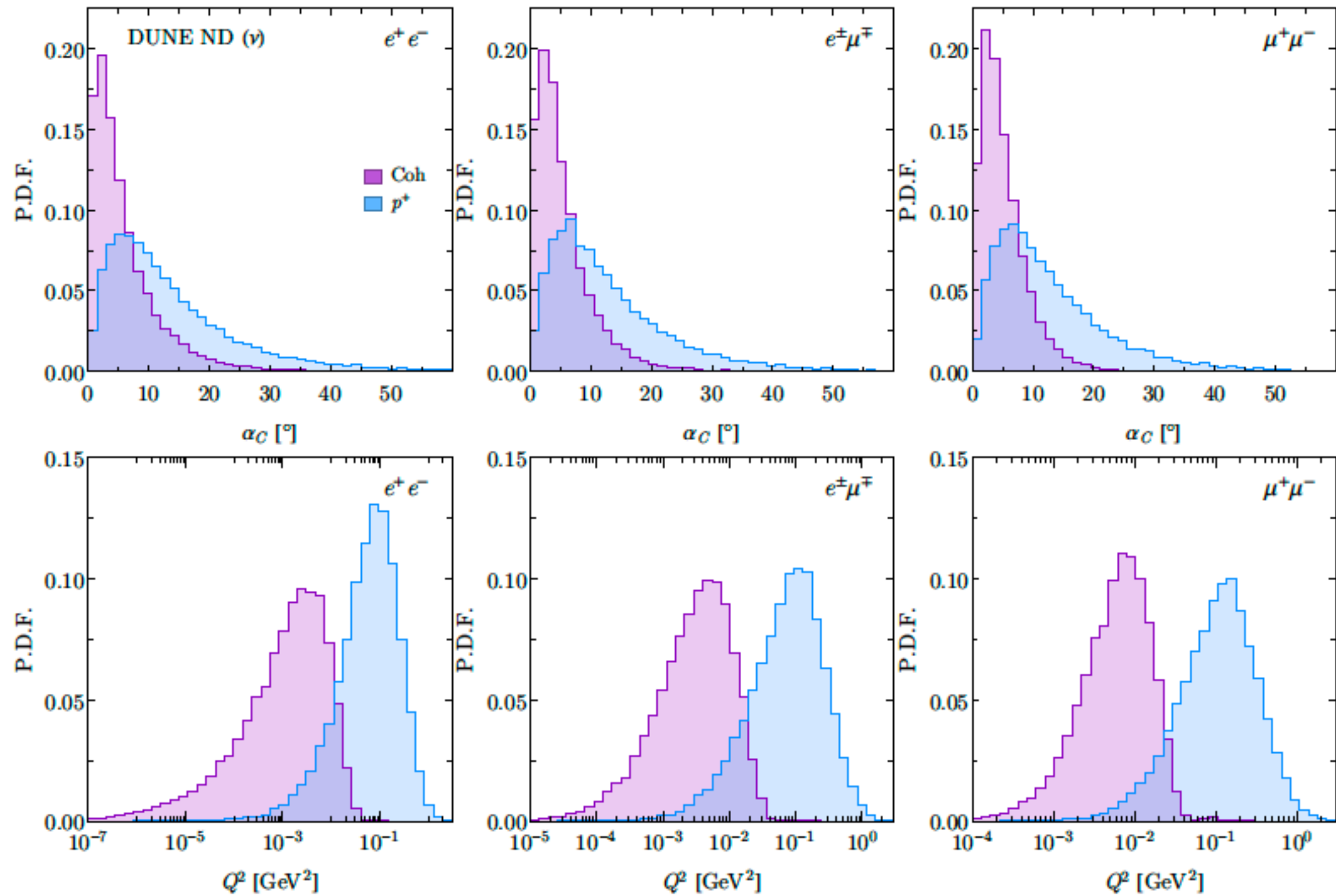
Rates



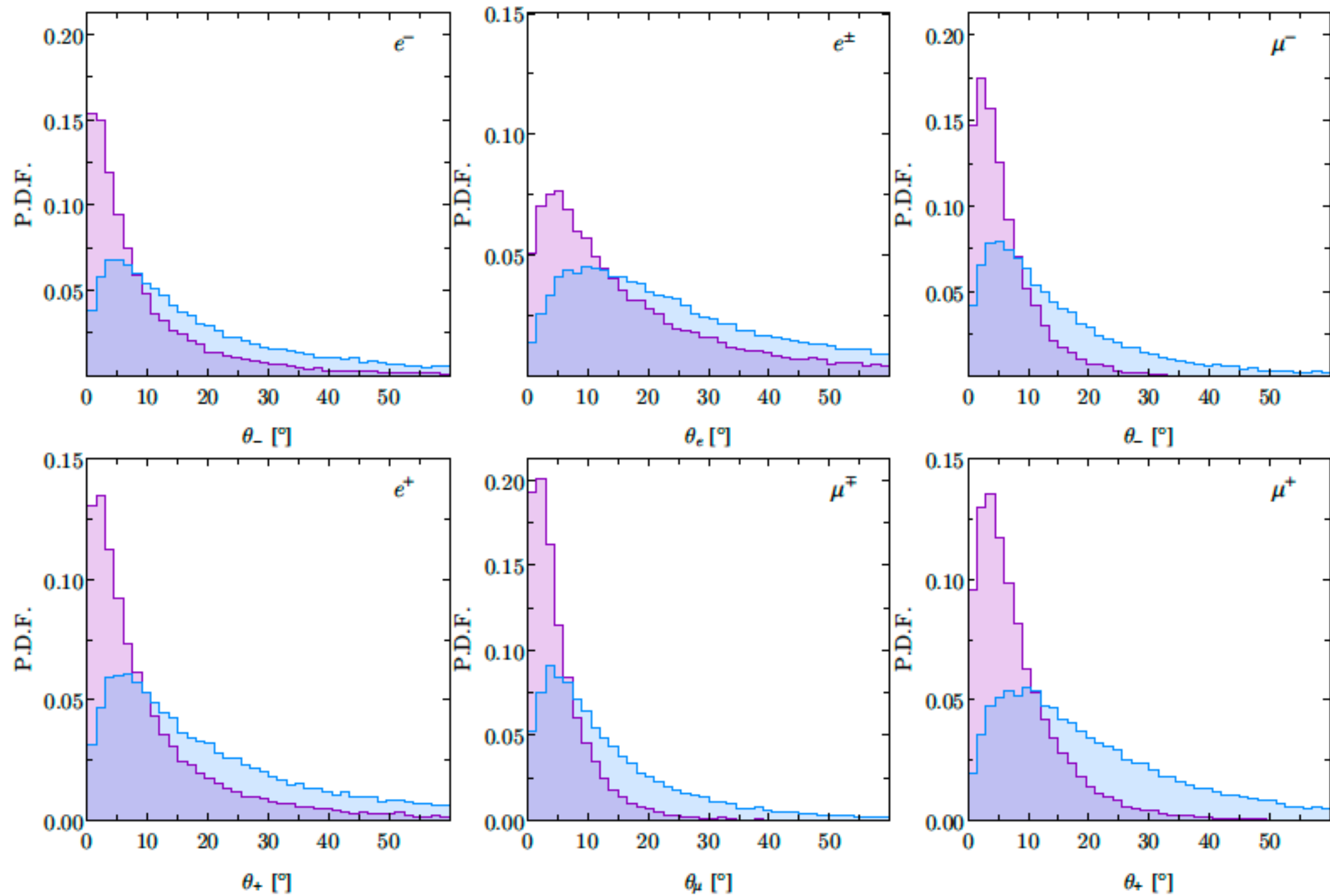
Kinematical Distributions



Kinematical Distributions



Kinematical Distributions



Events at other ND facilities

Rates

$$N_X^\psi = \text{Norm} \times \int dE_\nu \, \sigma_{\nu X}(E_\nu) \frac{d\phi_\nu(E_\nu)}{dE_\nu} \epsilon(E_\nu)$$

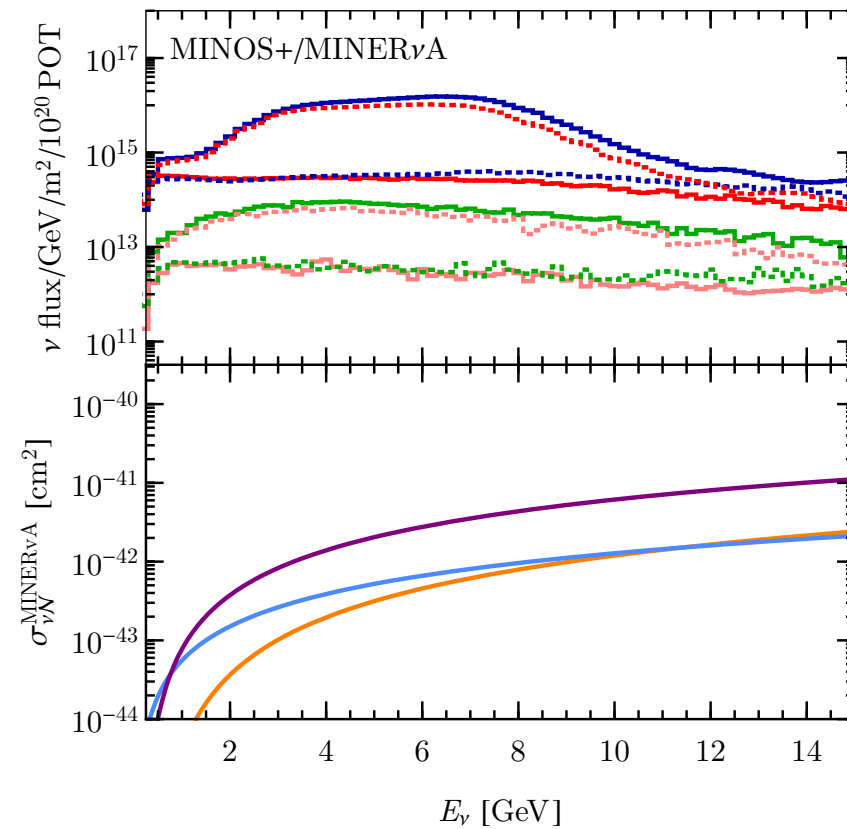
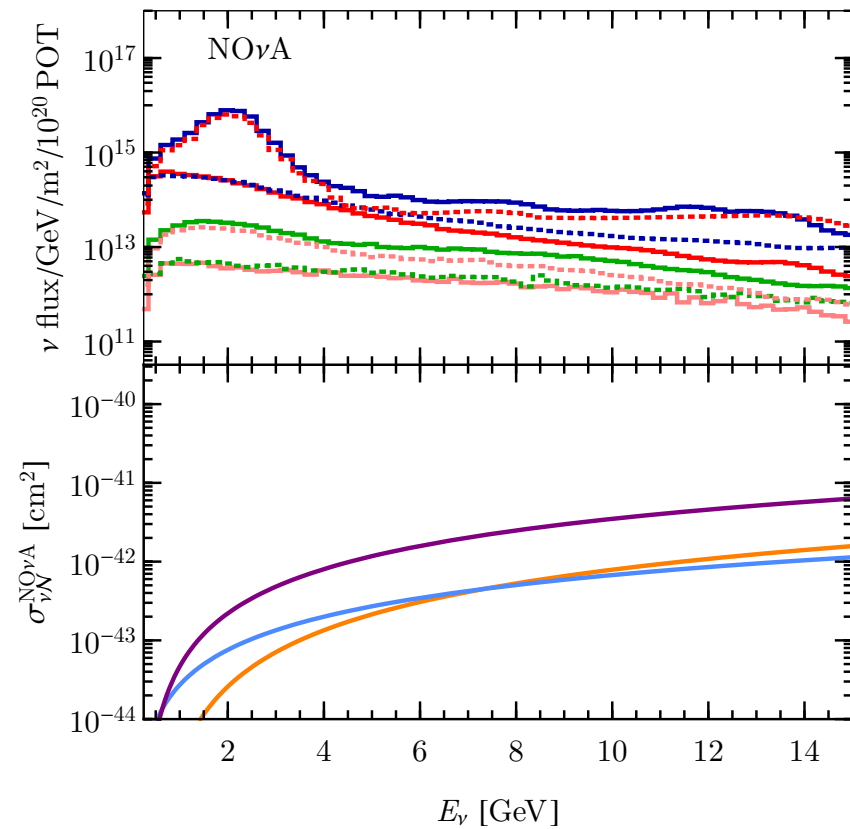
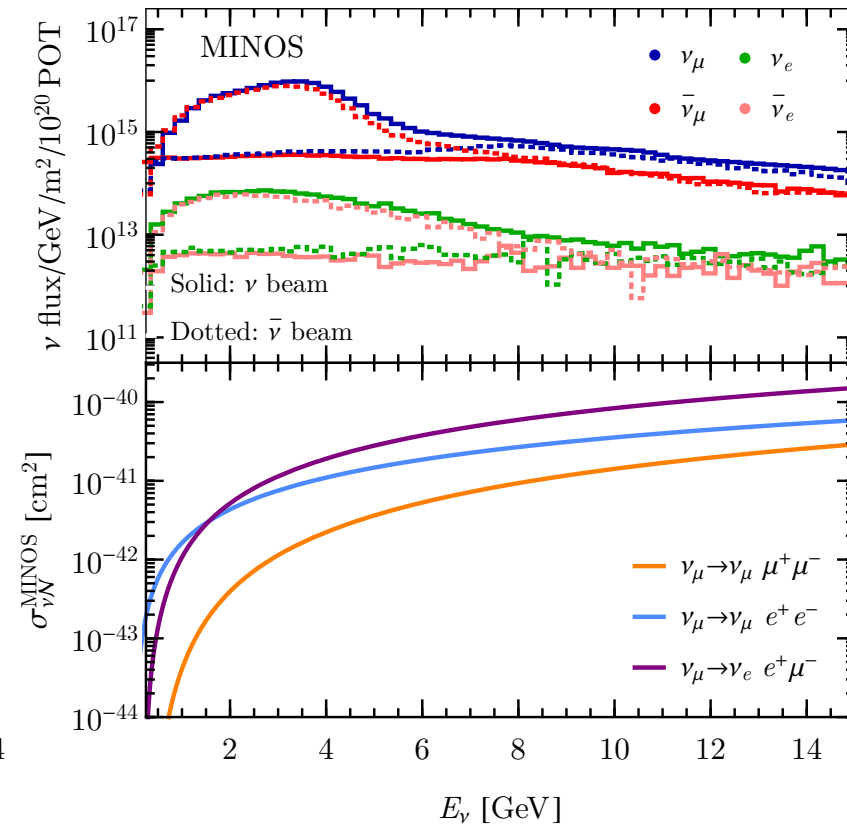
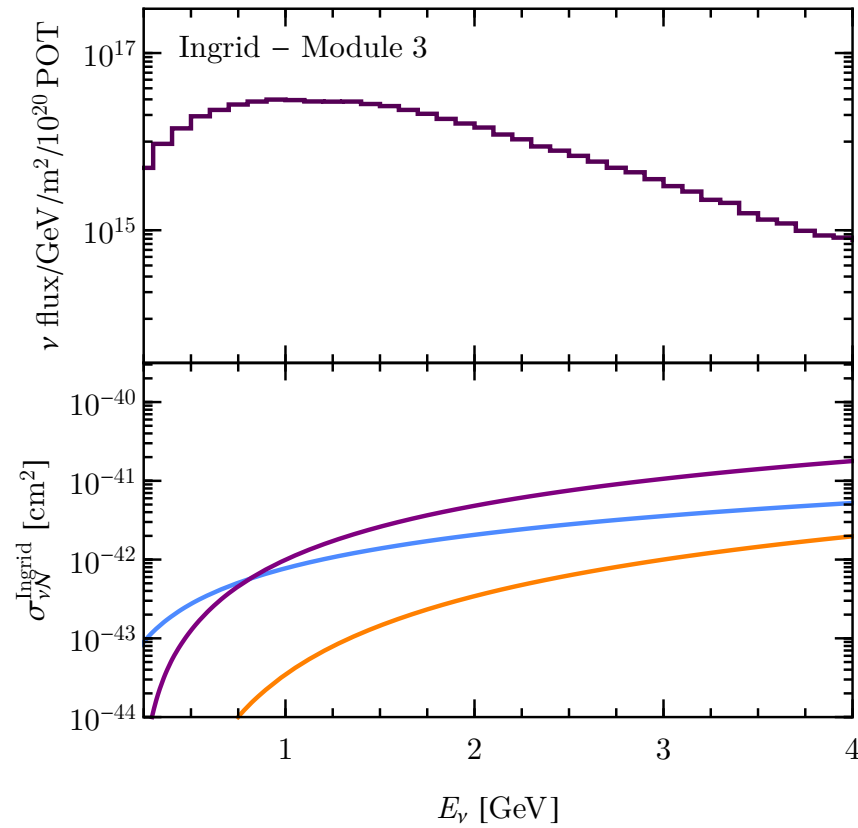
Fraction

$$\sigma_{\nu X}^{\text{FAC}} = \sum_i f_i \sigma_{\nu X}^i$$

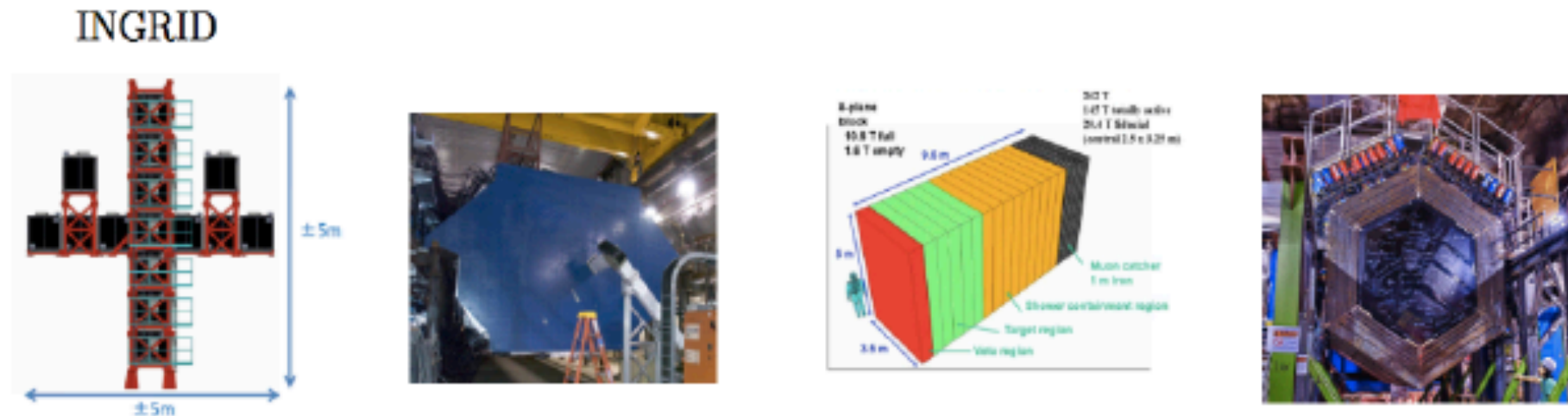
Cross section for the component

Experiment	Material	Baseline (m)	Exposure (POT)	Fiducial Mass (t)	E_ν (GeV)
INGRID	Fe	280	3.9×10^{21} [10 ²²] T2K-I [T2K-II]	99.4	0 – 4
MINOS[+]	Fe and C	1040	$10.56(3.36)[9.69] \times 10^{20}$	28.6	0 – 20
NO ν A	C ₂ H ₃ Cl and CH ₂	1000	$8.85(6.9)$ [36(36)] $\times 10^{20}$ [NO ν A-II]	231	0 – 20
MINER ν A	CH, H ₂ O, Fe, Pb, C	1035	$12(12) \times 10^{20}$	7.98	0 – 20

Rates



Rates



Channel	T2K-I	T2K-II	MINOS	MINOS+	NO ν A-I	NO ν A-II	MINER ν A
Total $e^\pm\mu^\mp$	563	1444	222 (56)	730	83 (72)	340 (374)	149 (102)
	96	246	46 (11)	151	25 (22)	102 (114)	56 (39)
Total e^+e^-	277	711	61 (15)	62	29 (22)	119 (114)	39 (27)
	24	62	9 (2)	8	4 (4)	16 (21)	10 (7)
Total $\mu^+\mu^-$	30	76	26 (6)	86	9 (9)	37 (47)	18 (13)
	21	54	15 (3)	49	8 (8)	34 (36)	18 (13)