



# Overview of the MicroBooNE LArTPC Detector Calibration

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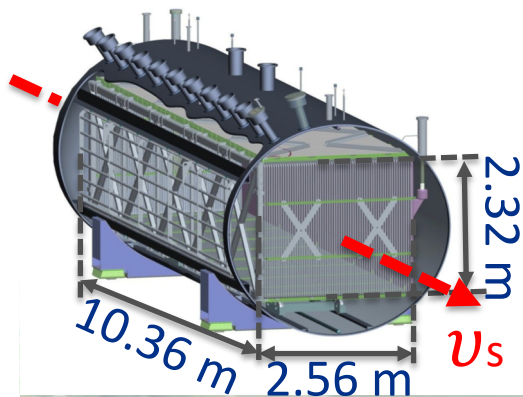
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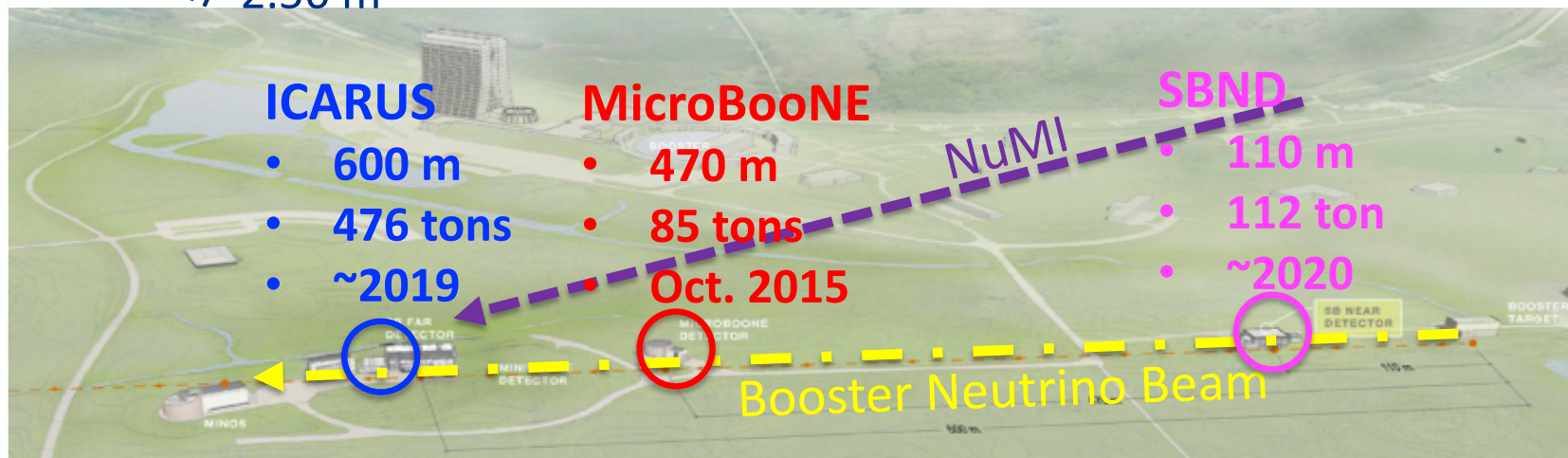
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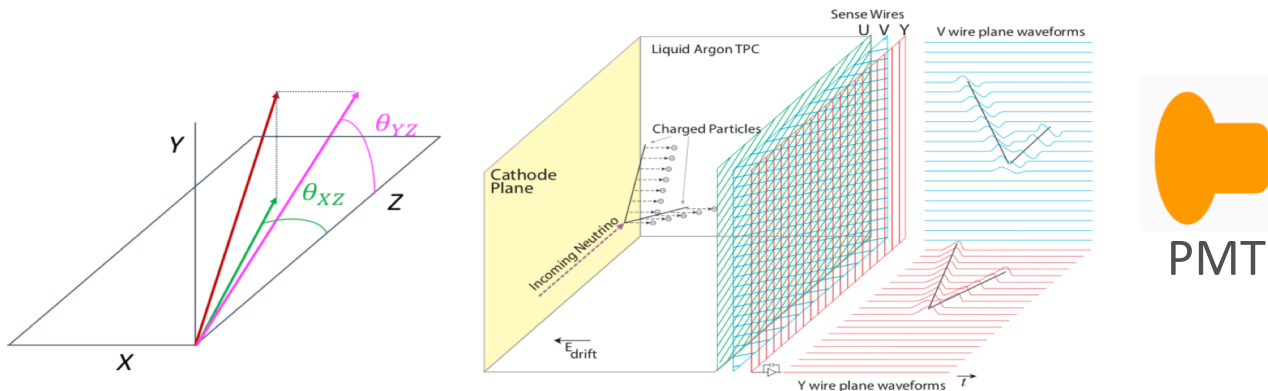
# The MicroBooNE Experiment



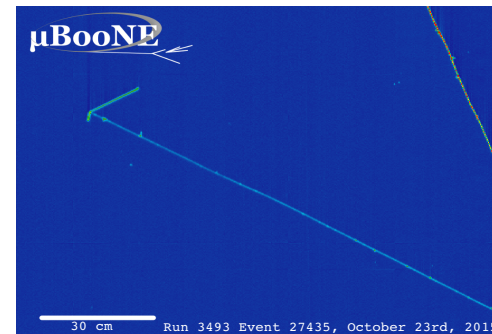
- Liquid Argon Time Projection Chamber (LArTPC)
- Investigate low-energy excess (LEE) of electron-like neutrino events observed by MiniBooNE in 2009
- Measure the low energy neutrino cross sections



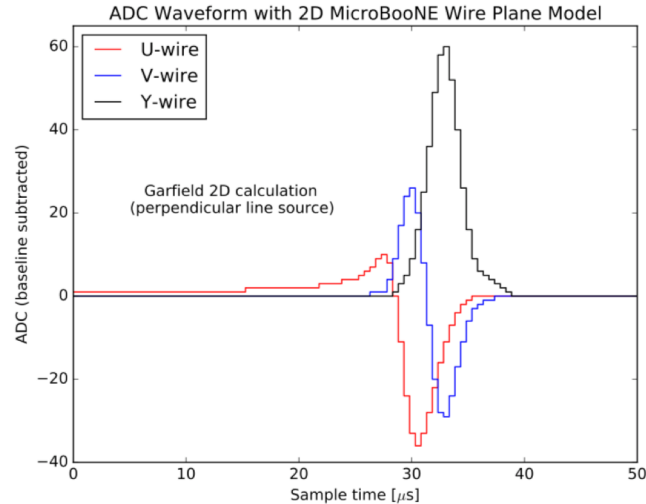
# The MicroBooNE Experiment



- 3 wire planes:
  - separated by 3 mm, with 3 mm wire spacing for each
  - to reconstruct the event, tracking and calorimetry
- 32 PMTs:
  - to detect scintillation light, mainly for trigger and event selection



# Example of Simulated Digitized Signals



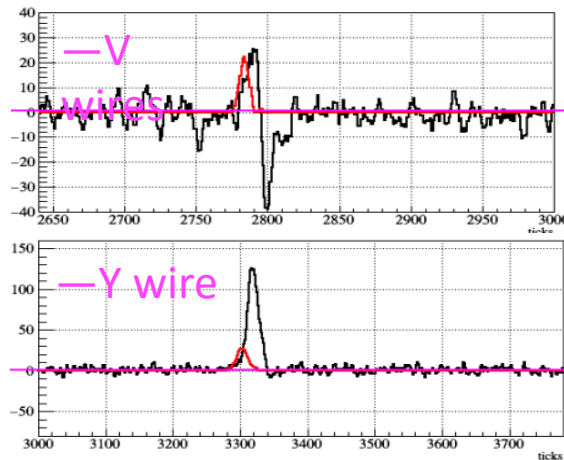
(ADC: Analog-to-digital converter)

The digitized signals from a central wire from each plane that are induced by an ideal MIP track in a 2D model of the MicroBooNE TPC.

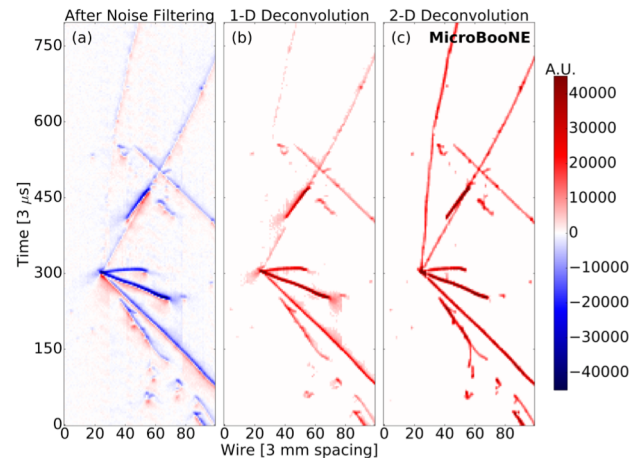
----Refer: [JINST 12, P08003 \(2017\)](#)



# Reconstruction of Drifted Electrons



----Image from W. Gu

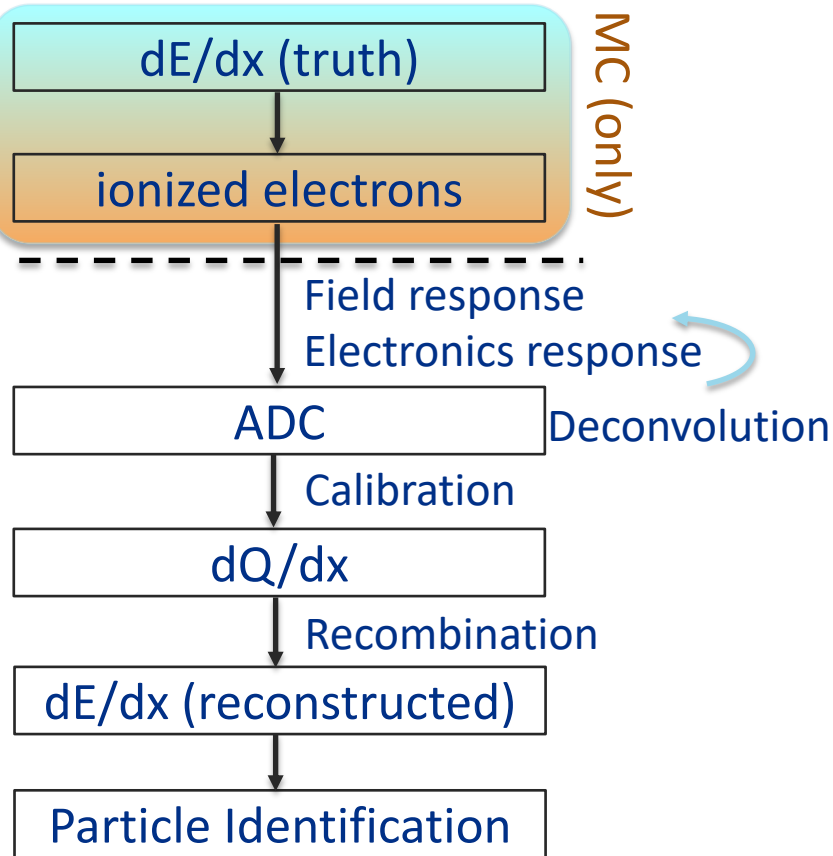


----Refer: [JINST 13, P07006 \(2018\)](#)

- We use deconvolution to extract original signals:
  - Noise filtering & Signal processing
  - Deconvolution is a mathematical technique to extract the original signal  $S(t)$  from the measured signal  $M(t')$ .

$$M(t') = \int_{-\infty}^{\infty} R(t, t') \cdot S(t) dt$$

$R(t, t')$ : detector response function



Many other effects need to be calibrated:

## I. Detector uniformity calibration ( $dQ/dx$ )

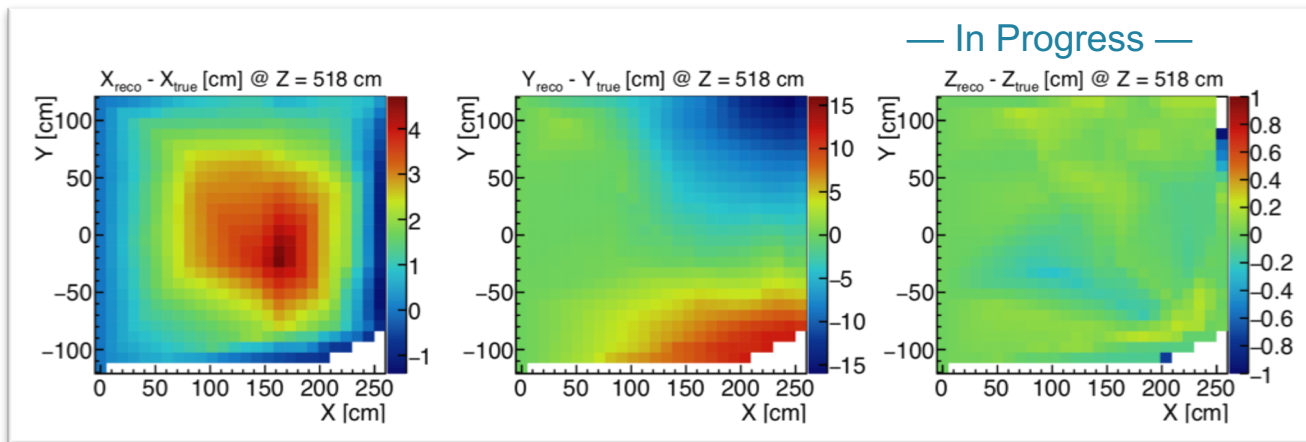
- Variation in electronics gains
- Electron attenuation ( $e^-$  lifetime)
- Nonuniform E field due to Space Charge Effects (SCE), diffusion, disconnected wires, TPC edges, etc.
- Temporal variations: i.e., change of run condition, temperature

## II. Energy scale determination ( $dE/dx$ )

- Use the well-known energy loss profile

# Space Charge Effects

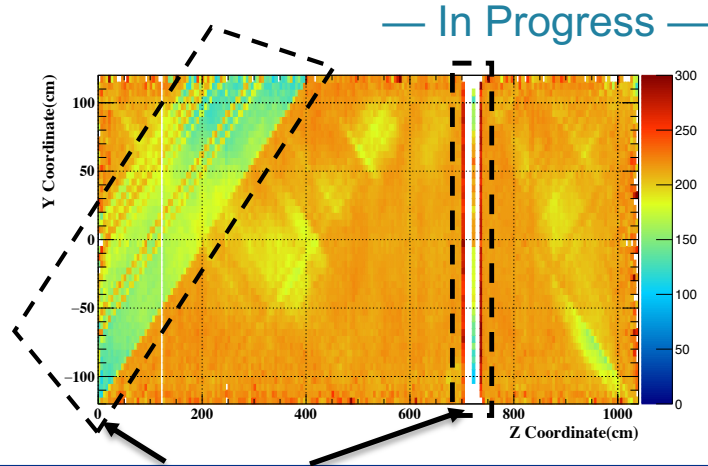
- MicroBooNE detector is on the surface (~20-30 cosmic muons per 4.8 ms readout window)
- Distortion of E field and ionization drift trajectories due to accumulation of slow-moving argon ions produced from cosmic muons impinging TPC
- Calibration on the E-field can be done by using UV laser system and the cosmic muon tracks



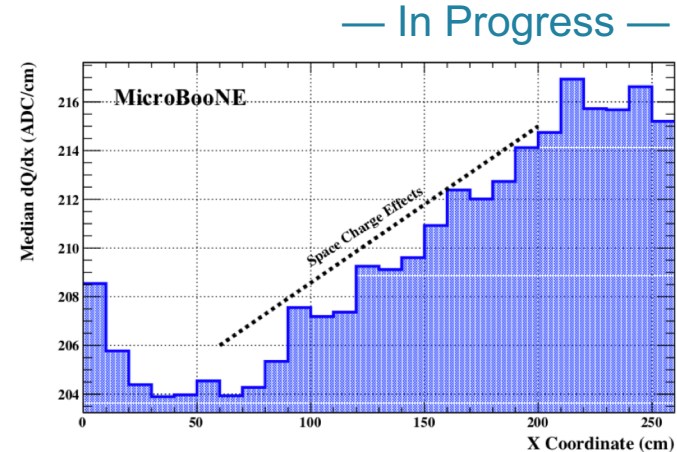
At central Z the spatial distortion of  $dX = X_{reco} - X_{true}$  (left),  $dY = Y_{reco} - Y_{true}$  and  $dZ = Z_{reco} - Z_{true}$  from laser data.

# Detector Uniformity Calibration–dQ/dx

- **dQ/dx calibration** (dQ/dx: ionization charge per unit length)
  - make the detector response to ionization charge uniformly throughout the detector and in time: *YZ plane, X direction, T*
  - use anode-cathode crossing cosmic muons to cover whole drift distance
- Variation of detector response (dQ/dx) in YZ plane and in the drift direction X



Misconfigured/cross-connected TPC channels

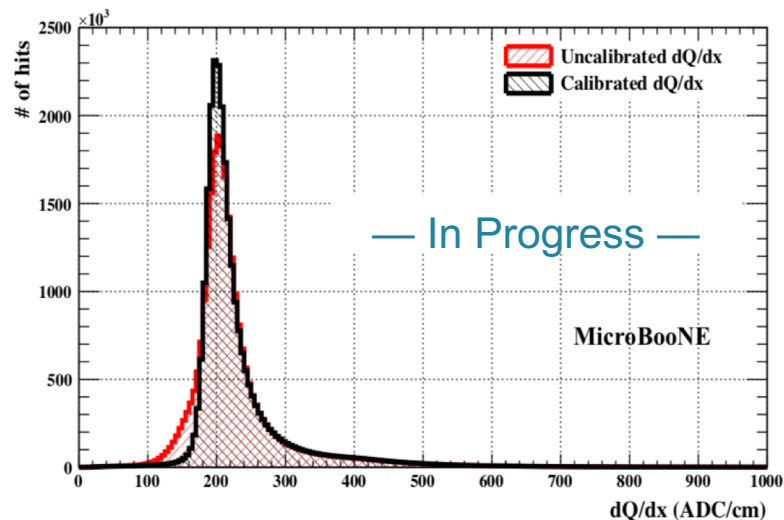
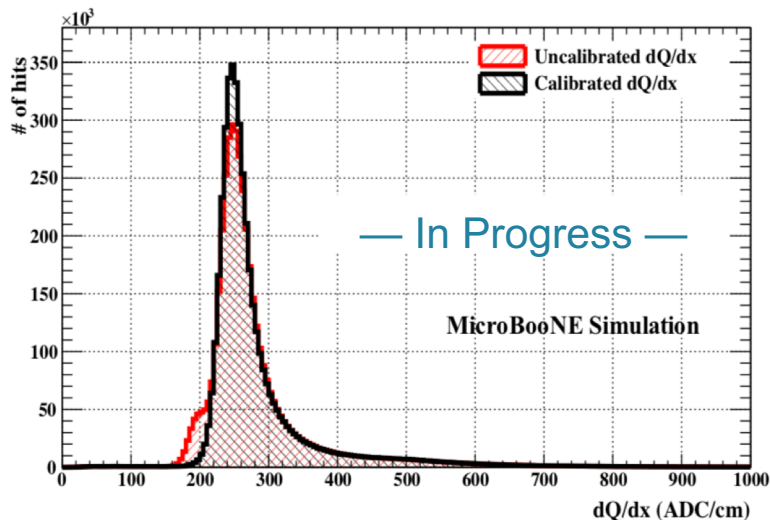


----Refer: MICROBOONE-NOTE-1048-PUB



# Detector Uniformity Calibration–dQ/dx

- Results (MC does not have time dependence):
  - Data:  $(dQ/dx)_{corr}^{data} = C(t) \cdot C(x) \cdot C(y, z) \cdot (dQ/dx)_{uncorr.}$
  - Monte Carlo (MC):  $(dQ/dx)_{corr}^{MC} = C(x) \cdot C(y, z) \cdot (dQ/dx)_{uncorr.}$
- Comparisons between calibrated and uncalibrated dQ/dx for both MC and data

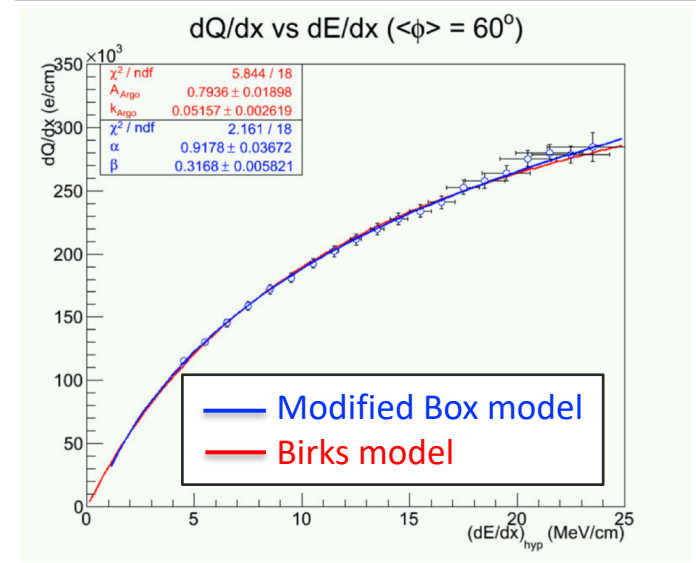


# Energy Scale Determination—dE/dx



- Determine the calibration constant  $C$  (unit: ADC/e), which translates the corrected dQ/dx (ADC/cm) to dQ/dx (e/cm).
- Stopping muons from neutrino interactions or cosmic rays are used to study the measured and predicted Most Probable dE/dx Value (MPV):

$$\left(\frac{dE}{dx}\right)_{corr.} = \frac{\exp\left(\frac{\left(\frac{dQ}{dx}\right)_{corr.}}{C} \cdot \frac{\beta' W_{ion}}{\rho \xi}\right) - \alpha}{\frac{\beta'}{\rho \xi}}$$

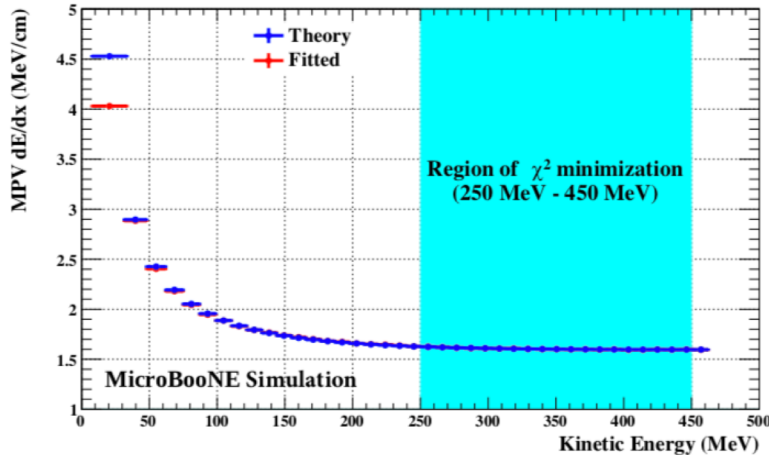


----Refer: [ArgoNeuT JINST 8, P08005 \(2013\)](#)

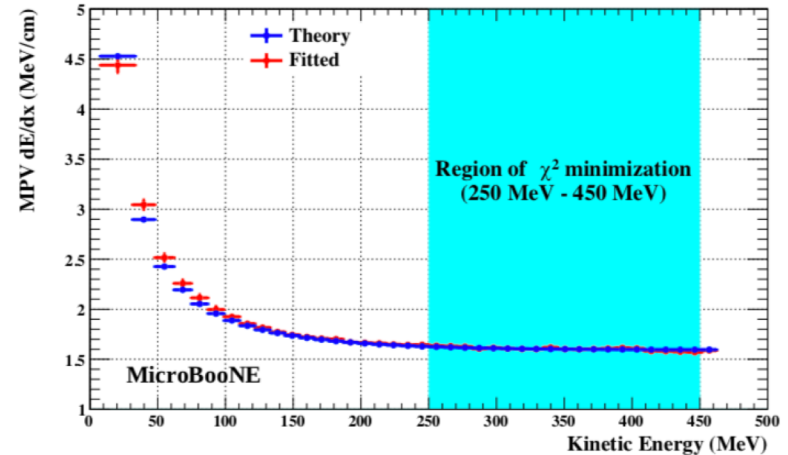
# Energy Scale Determination—dE/dx

- Once we find the calibration constants, we would expect the calibrated dE/dx matches expectation.

— In Progress —



— In Progress —



- Kinetic Energy (K.E.) =  $\sum \frac{dE}{dx} dx$ , where  $dx$  is the track pitch.

- In this talk, we present a briefly overview of the MicroBooNE LArTPC detector calibration.
- MicroBooNE detector performs at a high signal-to-noise ratio.
- Deconvolution is used to extract the charge signal after removing noise.
- Many detector effects will affect the measurement and need to be addressed appropriately.
- The detector calibration can correct most of the detector effects and improve the energy reconstruction and particle identification.
- The calibration strategy carried out here can be used by other LArTPC experiments to calibration their detectors.



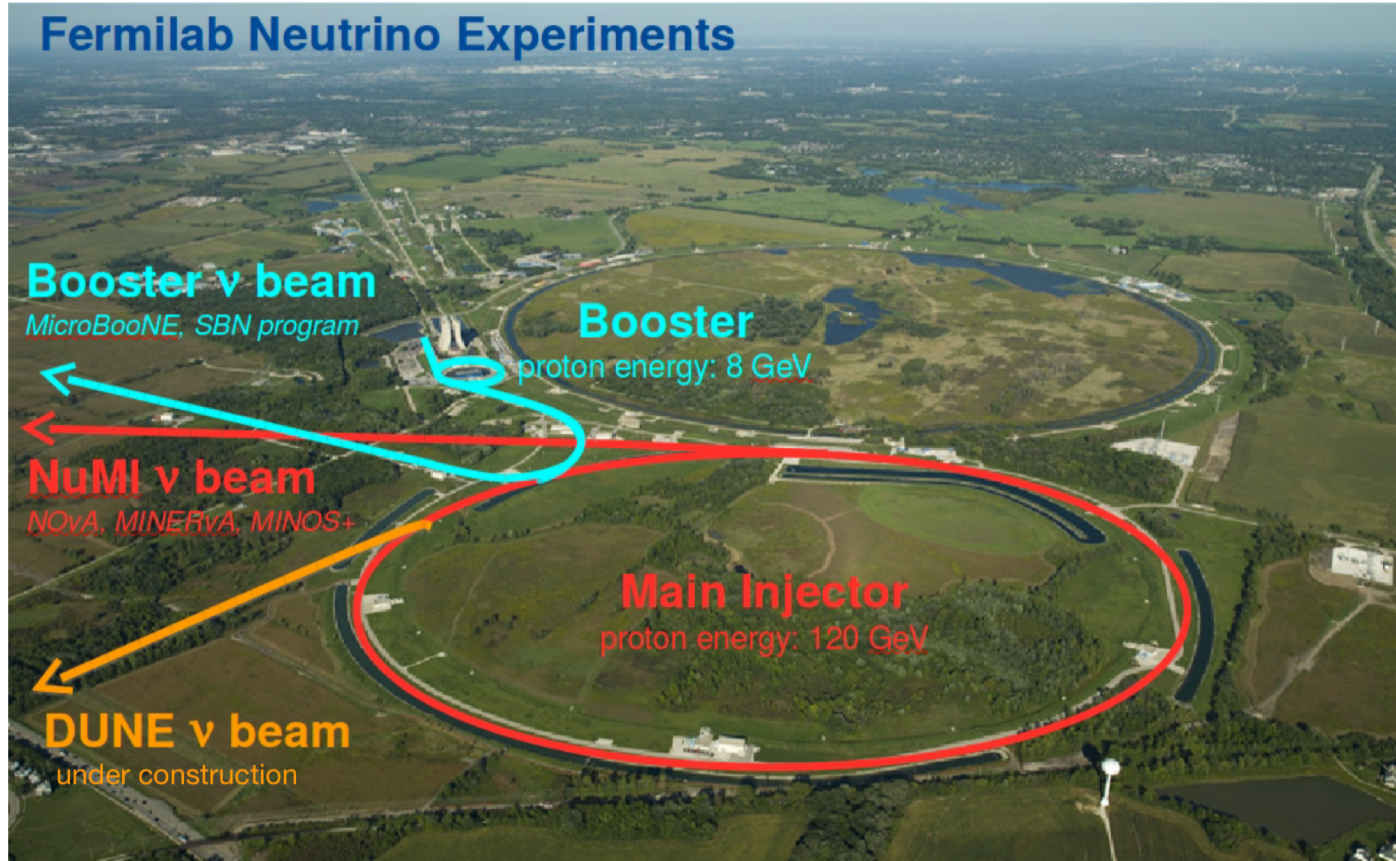
# Acknowledgement



- The calibration method is developed by the MicroBooNE Calibration Group. I would like to express my gratitude and appreciation to the team, especially *Tingjun Yang, Glenn Horton-Smith, Mike Mooney, Yifan Chen, Varuna Meddage* and *Marianette Wospakrik*. Without their effort and help, this talk would not have been possible.
- MicroBooNE is supported by the following: the U.S. Department of Energy, Office of Science, Offices of High Energy Physics and Nuclear Physics; the U.S. National Science Foundation; the Swiss National Science Foundation; the Science and Technology Facilities Council of the United Kingdom; and The Royal Society (United Kingdom).

—Thank you!

# Backup—Fermilab Neutrino Programs



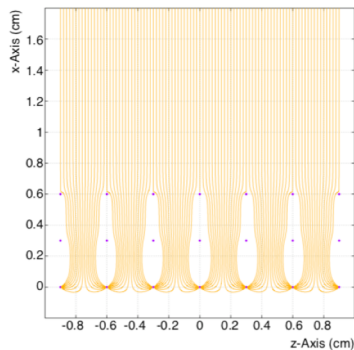
# Backup—PMT and Trigger Systems

- 32 TPB-coated PMTs behind wires
  - Two gains: 1.8% and 18%
  - Shaping: 60 ns rise time
  - 64 MHz ADC (ADS5272): accurate determination of event  $t_0$
- 
- Level-1 trigger: accelerator gates (BNB and NuMI) and random triggers (for cosmics)
  - Level-2 trigger: PMT information in the beam window (in software)

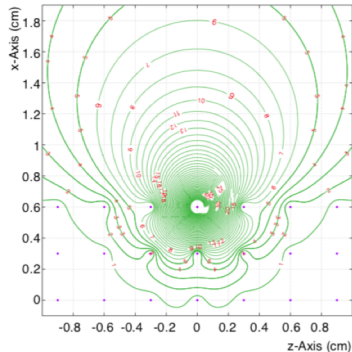


MicroBooNE: JINST 10, T06001 (2015); JINST 8, T07005 (2013)

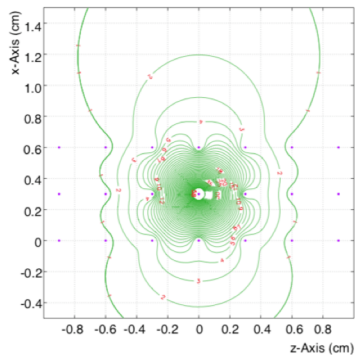
# Backup—Drift Electron Path



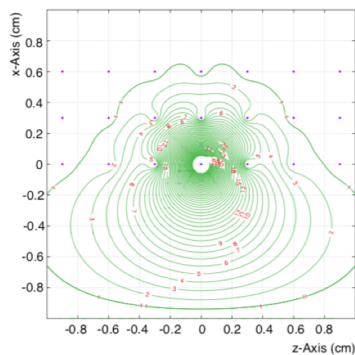
(a) Electron drift paths.



(b) Weighting potential on a U wire.



(c) Weighting potential on a V wire.



(d) Weighting potential on a Y wire.

----Refer: [JINST 13, P07006 \(2018\)](#)



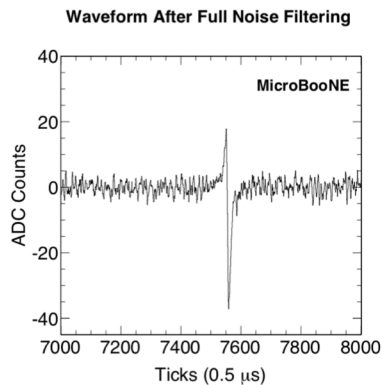
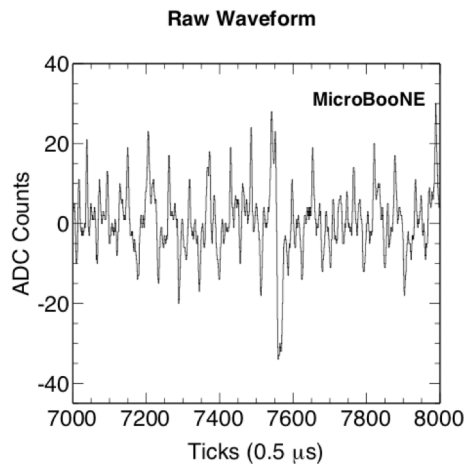
The principal method to extract charge is deconvolution. This procedure in its one-dimensional (1D) form has been used in the data analysis of previous liquid argon experiments [29, 30]. This technique has the advantages of being robust and fast. It is an essential step in the overall drifted-charge profiling process.

Deconvolution is a mathematical technique to extract the *original signal*  $S(t)$  from the *measured signal*  $M(t')$ . The measured signal is modeled as a convolution integral over the original signal  $S(t)$  and a given detector *response function*  $R(t, t')$ , which gives the instantaneous portion of the measured signal at some time  $t'$  due to an element of original signal at time  $t$ :

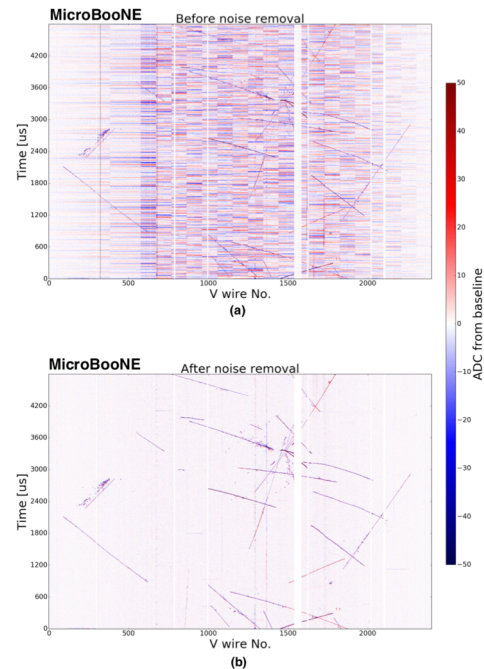
$$M(t') = \int_{-\infty}^{\infty} R(t, t') \cdot S(t) dt. \quad (3.1)$$

----Refer: [JINST 13, P07006 \(2018\)](#)

# Backup—Noise Filter



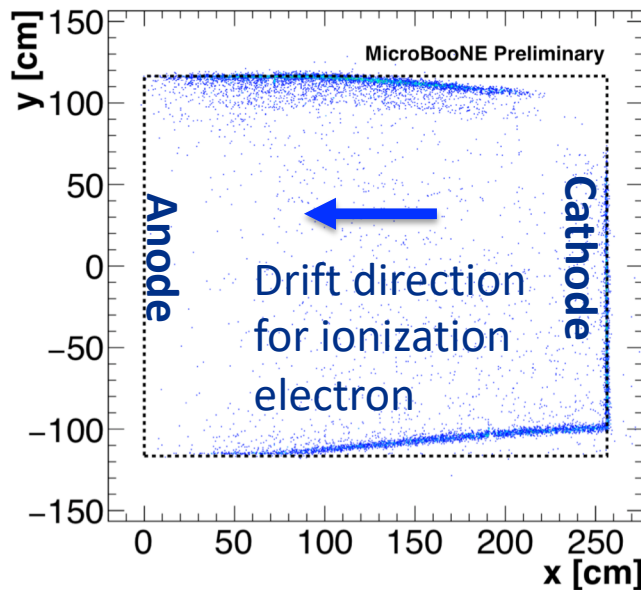
----Refer: [JINST 12, P08003 \(2017\)](#)



**Figure 14.** 2-D event display of the V plane from run 3493 event 41075 showing the raw signal (a) before and (b) after offline noise filtering. A clean event signature is recovered once all the identified noise sources are subtracted.

- Challenges for MicroBooNE LArTPC detector:
  - **Misconfigured or cross-connected TPC channels:** change gain of electronic channels, distort field between wire planes
  - **Space Charge Effects (SCE):** cosmic rays, accumulation of slow moving  $\text{Ar}^+$ , causing distortion the magnitude and direction of drift E field
  - **Electron attenuation:** electronegative contaminants such as  $\text{H}_2\text{O}$  and  $\text{O}_2$  can capture some of the drifting electrons
  - **Diffusion:** drifted electrons may get smeared out
  - **Recombination:** ionization electrons may not completely liberate from their parent argon ions and recombine back to form neutral argon atoms again, causing underestimation of particle energy loss
  - **Temporal variations:** change of temperature, run conditions...
  - ...
- We want to measure the charge and position of ionization signal precisely and improve both the measurement of total deposited energy and particle identification (PID).

# Backup—Space Charge Effects



Start/end points of reconstructed cosmic muon tracks tagged by an external muon counter in the x-y plane for off-beam (cosmic) events.

----Refer: [MICROBOONE-NOTE-1018-PUB](#)



➤  $(\frac{dQ}{dx})_{uncorr.} \rightarrow YZ_{corr.} \rightarrow DriftDirection_{corr.} \rightarrow Time_{corr.} \rightarrow (\frac{dQ}{dx})_{corr.}$

- Consists of 3 steps:

1. YZ plane correction: remove effects from SCE, misconfigured or cross-connected TPC channels and transverse diffusion
2. Drift direction correction: remove effects coming from electron attenuation, SCE and longitudinal diffusion
3. Time correction: remove any temporal variations in detector response (for data only, MC does not have time dependence)

➤  $(\frac{dQ}{dx})_{uncorr.} \rightarrow YZ_{corr.} \rightarrow Lifetime_{corr.} \rightarrow (\frac{dQ}{dx})_{corr.}$

- Consists of 2 steps:

1. YZ plane correction: same as above
2. Lifetime correction:

- remove effects coming from electron attenuation, SCE and longitudinal diffusion, temporal variations
- $Q = Q' \cdot e^{-\frac{t}{\tau}}$  (t: drift time;  $\tau$ : electron lifetime, determined by purity or TPC (laser system, muon count system )

In progress

The YZ plane is segmented into 5 cm by 5 cm cells. The correction factor is defined by:

$$C(y, z) = \frac{(dQ/dx)_{Global}}{(dQ/dx)_{Local}}$$

$(dQ/dx)_{Global}$  – Global median  $dQ/dx$  in the collection plane

$(dQ/dx)_{Local}$  – Local median  $dQ/dx$  for a given YZ plane cell in the collection plane

After YZ plane calibration:

$$(dQ/dx)_{corr.} = C(y, z) \cdot (dQ/dx)_{uncorr.}$$

The drift direction is segmented into 10 cm in data (5 cm in MC). The correction factor is defined by:

$$C(x) = \frac{(dQ/dx)'_{Global}}{(dQ/dx)'_{Local}}$$

$(dQ/dx)'_{Global}$  – Global median  $dQ/dx$  value after correcting for YZ plane irregularities

$(dQ/dx)'_{Local}$  – Local median  $dQ/dx$  value after correcting for YZ plane irregularities

After YZ plane calibration and drift direction calibration:

$$(dQ/dx)_{corr.} = C(x) \cdot C(y, z) \cdot (dQ/dx)_{uncorr.}$$

The time dependent correction is determined for each day:

$$C(t) = \frac{(dQ/dx)_{Reference}}{(dQ/dx)_{Global}}$$

$(dQ/dx)_{Reference}$  – Selected reference  $dQ/dx$  value

$(dQ/dx)_{Global}$  – Global  $dQ/dx$  value in the anode wire plane after correcting for YZ plane and drift direction irregularities

$$\left(\frac{dE}{dx}\right)_{corr.} = \frac{\exp\left(\frac{\left(\frac{dQ}{dx}\right)_{corr.} \cdot \frac{\beta' W_{ion}}{\rho \xi}}{\beta' \rho \xi}\right) - \alpha}{\beta' \rho \xi}$$

- $C$ : Calibration constant to convert ADC values to number of electrons
  - $W_{ion}$ :  $23.6 \times 10^{-6}$  MeV/electron (work function of argon)
  - $\xi$ : 0.273 kV/cm (MicroBooNE drift electric field)
  - $\rho$ : 0.212 (kV/cm)(g/cm<sup>2</sup>)/MeV
  - $\alpha$ : 0.93
- ArgoNeuT: JINST 8, P08005 (2013)

$$\chi^2 = \Sigma\left(\frac{(MPV(dE/dx)_{prediction} - MPV(dE/dx)_{Measured})^2}{\sigma^2}\right)$$