

Electric Field Effects on the Muon Anomalous Precession Frequency in the Fermilab Muon $g-2$ Experiment

Wanwei Wu

Department of Physics and Astronomy

University of Mississippi

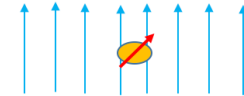
(on behalf of the Muon $g-2$ Collaboration)

APS April Meeting, Columbus, OH

Session X08, Tuesday, April 17, 2018

A particle with spin has a magnetic moment ($\vec{\mu}$) aligned with its spin (\vec{S}):

$$\vec{\mu} = g \frac{q}{2m} \vec{S}$$



This differs from (1) by the two extra terms

Dirac, 1928

$$\frac{eh}{c} (\boldsymbol{\sigma}, \mathbf{H}) + \frac{ieh}{c} \rho_1 (\boldsymbol{\sigma}, \mathbf{E})$$

in F. These two terms, when divided by the factor $2m$, can be regarded as the additional potential energy of the electron due to its new degree of freedom. The electron will therefore behave as though it has a magnetic moment $eh/2mc \cdot \boldsymbol{\sigma}$ and an electric moment $ieh/2mc \cdot \rho_1 \boldsymbol{\sigma}$. This magnetic moment is just that assumed in the spinning electron model. The electric moment, being a pure

Dirac Theory predicts that

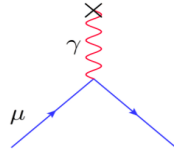
$$g = 2.$$

However, experiments showed that

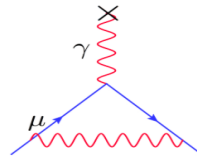
$$g \neq 2.$$

➔ Anomalous Magnetic Dipole Moment: $a = \frac{g-2}{2}.$

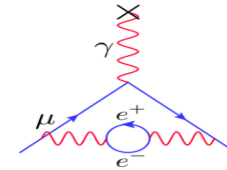
QED:



Dirac term



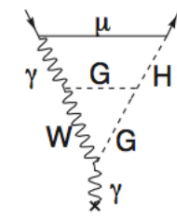
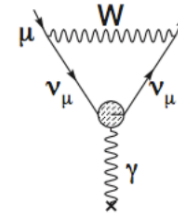
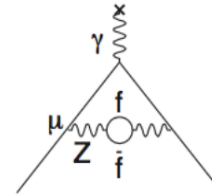
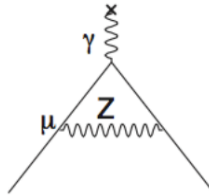
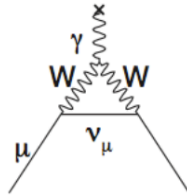
Schwinger term



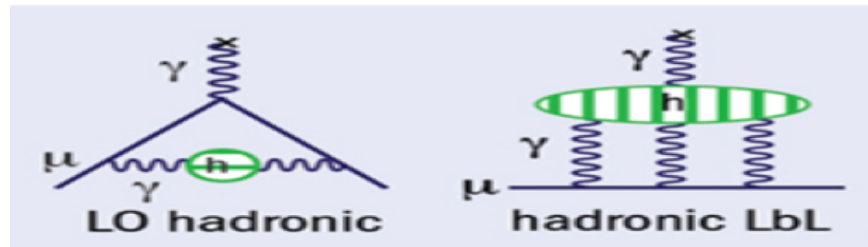
Vacuum polarization

EW:

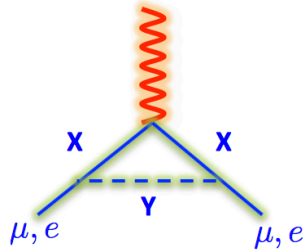
Electroweak Contributions (known to 0.01 ppm vs 0.5 ppm experiment)



Hadron:



$$a_{\mu}^{SM} = a_{\mu}^{QED} + a_{\mu}^{EW} + a_{\mu}^{Had} = 116591828(50) \times 10^{-11}$$

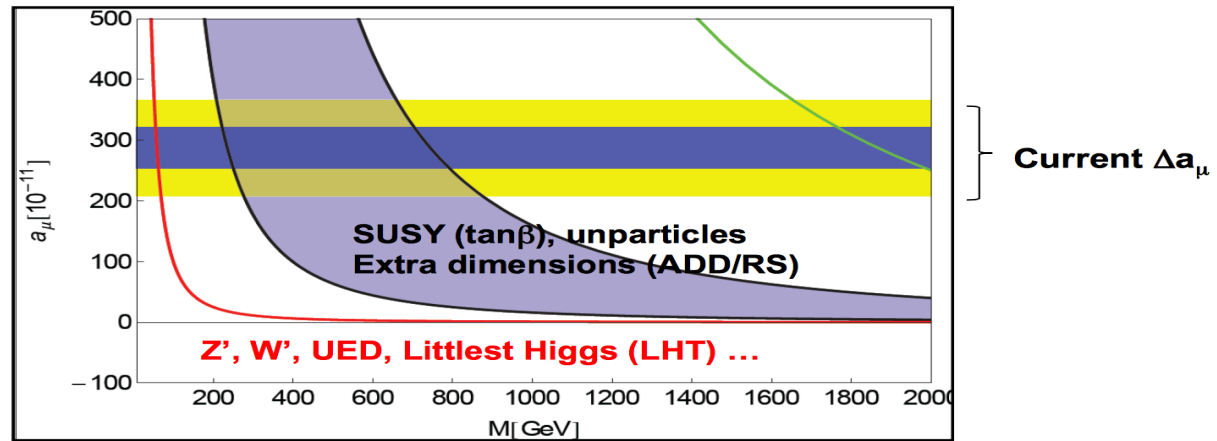


New Physics
proportional to:

$$\left(\frac{m_\ell}{M_{\text{NEW}}} \right)^2$$

a_μ represents a sum over all physics, it is sensitive to a wide range of potential new physics.

radiative muon mass generation



$$a_\mu^{NP} = a_\mu^{Exp} - a_\mu^{SM}$$

SM: $116591828(50) \times 10^{-11}$

BNL g-2: $116592080(63)_{tot} \times 10^{-11}$

$\sim 3.3\sigma$

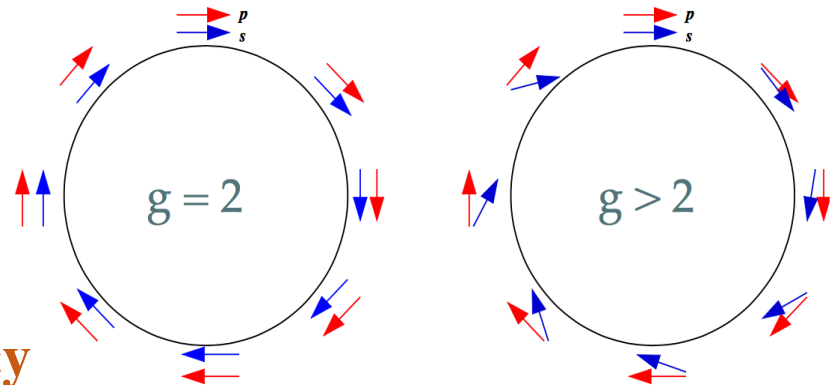
Spin rotation of a muon in a magnetic field

- Spin precession frequency

$$\vec{\omega}_S = -\frac{qg\vec{B}}{2m} - \frac{q\vec{B}}{\gamma m}(1 - \gamma)$$

- Cyclotron rotation frequency

$$\vec{\omega}_C = -\frac{q\vec{B}}{m\gamma}.$$

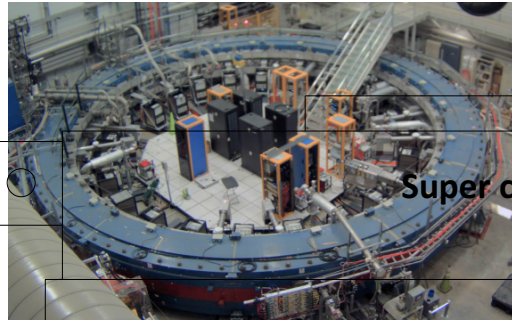


Muon anomalous precession frequency:

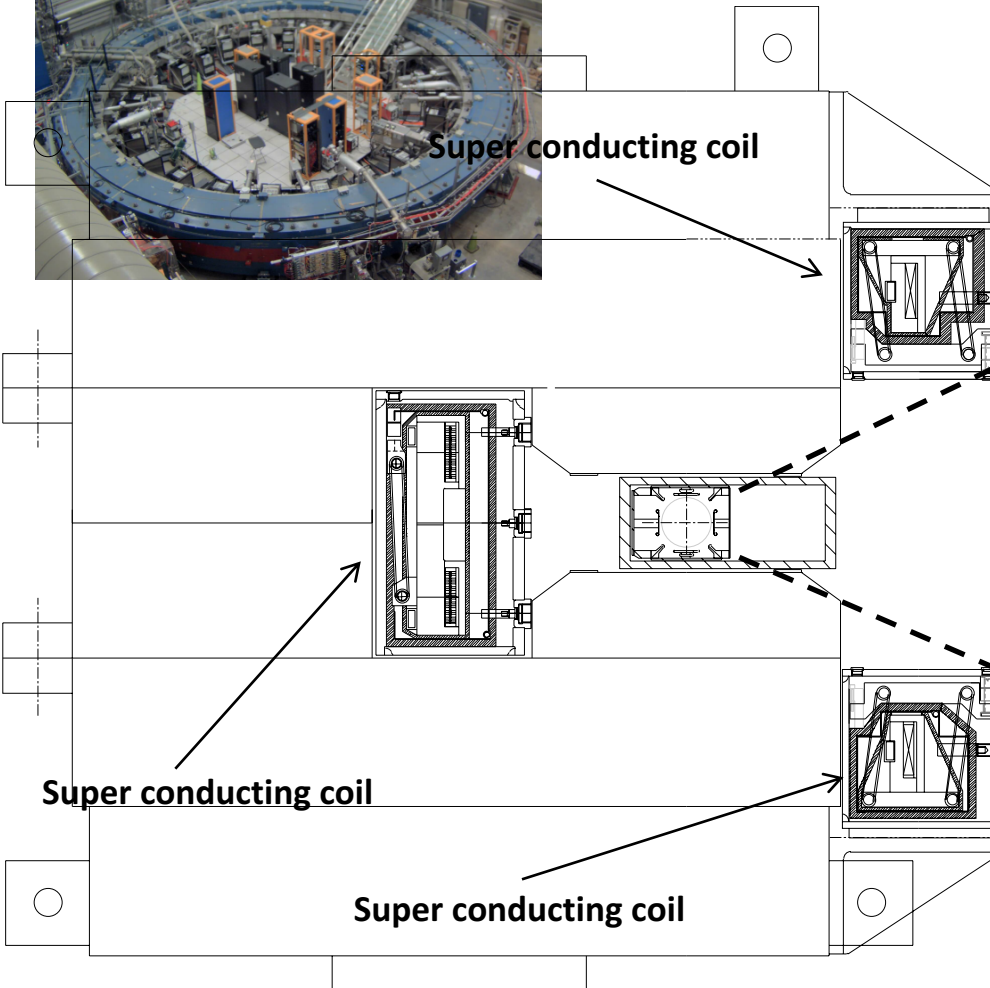


$$\vec{\omega}_a = \vec{\omega}_S - \vec{\omega}_C = -\left(\frac{g-2}{2}\right) \frac{q\vec{B}}{m} = -a_\mu \frac{q\vec{B}}{m}.$$

Weak focusing muon storage ring— Electrostatic quadrupoles provide vertical focusing

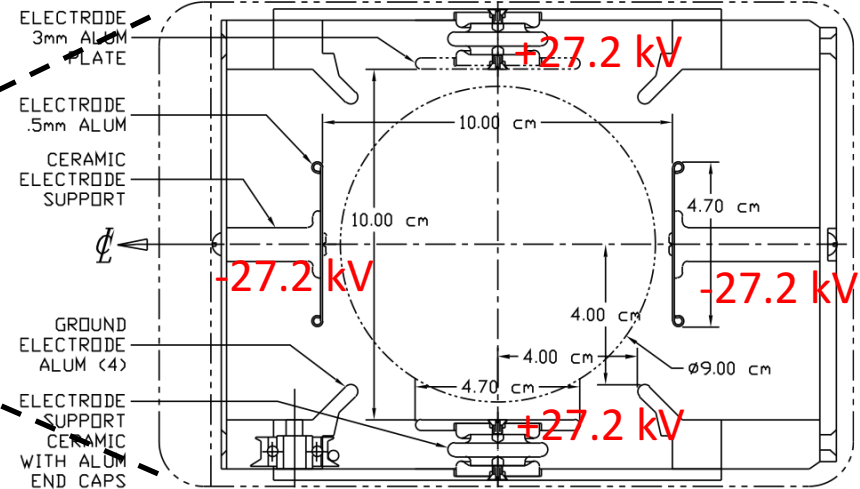


Super conducting coil



Super conducting coil

Super conducting coil

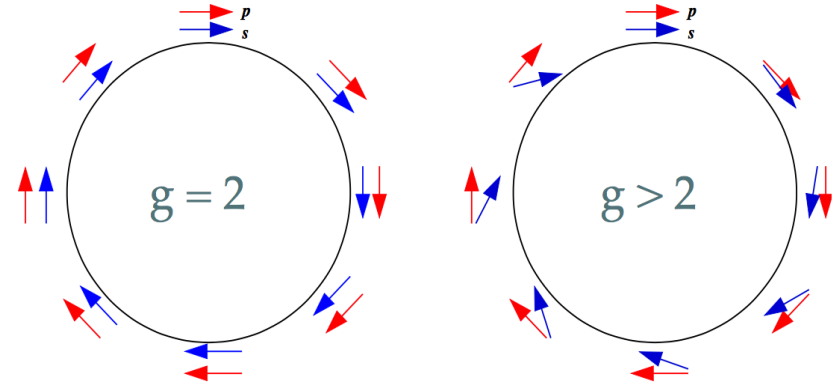


ELECTRODE AND SUPPORT FRAME - END VIEW

Electrostatic focusing

- Cyclotron rotation frequency

$$\vec{\omega}_C = -\frac{q}{m} \left[\frac{\vec{B}}{\gamma} - \frac{\gamma}{\gamma^2 - 1} \left(\frac{\vec{\beta} \times \vec{E}}{c} \right) \right]$$



- Spin precession frequency

$$\vec{\omega}_S = -\frac{q}{m} \left[\left(\frac{g}{2} - 1 + \frac{1}{\gamma} \right) \vec{B} - \left(\frac{g}{2} - 1 \right) \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left(\frac{g}{2} - \frac{\gamma}{\gamma + 1} \right) \left(\frac{\vec{\beta} \times \vec{E}}{c} \right) \right]$$

Muon anomalous precession frequency:

$$\vec{\omega}_a = -\frac{q}{m} \left[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

$$\vec{\omega}_a = -\frac{q}{m} \left[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

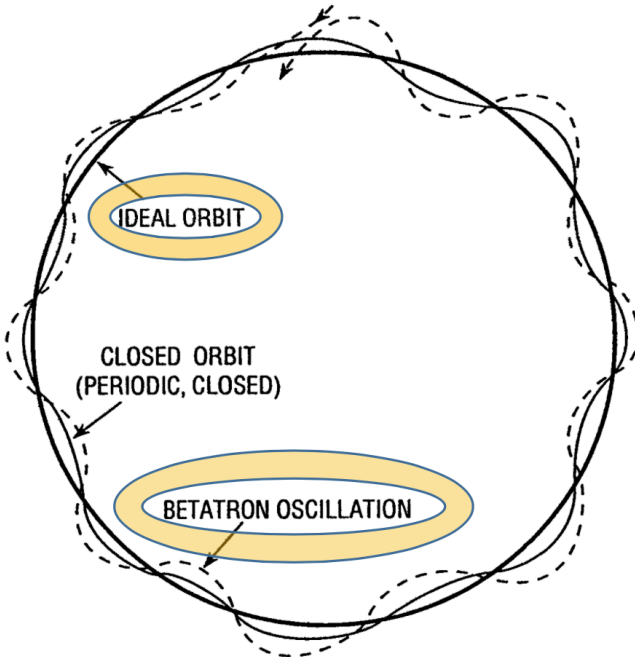
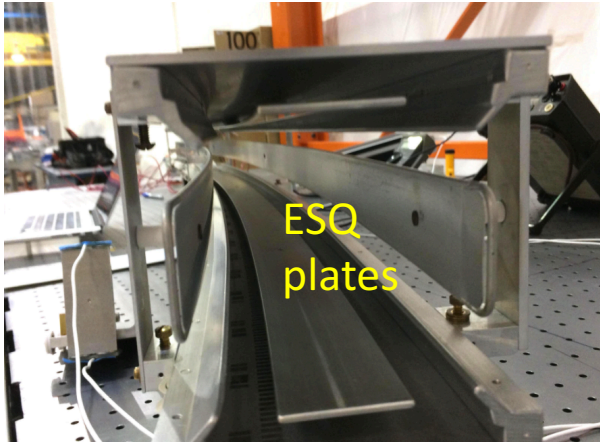
- Assuming $\vec{\beta} \cdot \vec{B} = 0$, the **second term** vanishes.
- By choosing $\gamma = 29.3$, the **third term** vanishes. The corresponding momentum (3.904 GeV/c)/radius (711.2 cm) is then called “magic” momentum/radius.



$$\omega_a = \omega_S - \omega_C = a_\mu \frac{q}{m} B$$

- We measure ω_a by using the decay positron signals and measure B by observing the Larmor frequency of stationary protons ($\omega_p = \frac{2\mu_p B}{\hbar}$) with NMR probes.
- For our final analysis, we can solve for a_μ and rewrite it as (using $\mu_e = g_e e \hbar / 4m_e$):

$$a_\mu = \frac{\omega_a}{\omega_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$



if electrostatic quadrupole (ESQ)
plates misaligned



introduce E-field multipoles



closed-orbit distortion

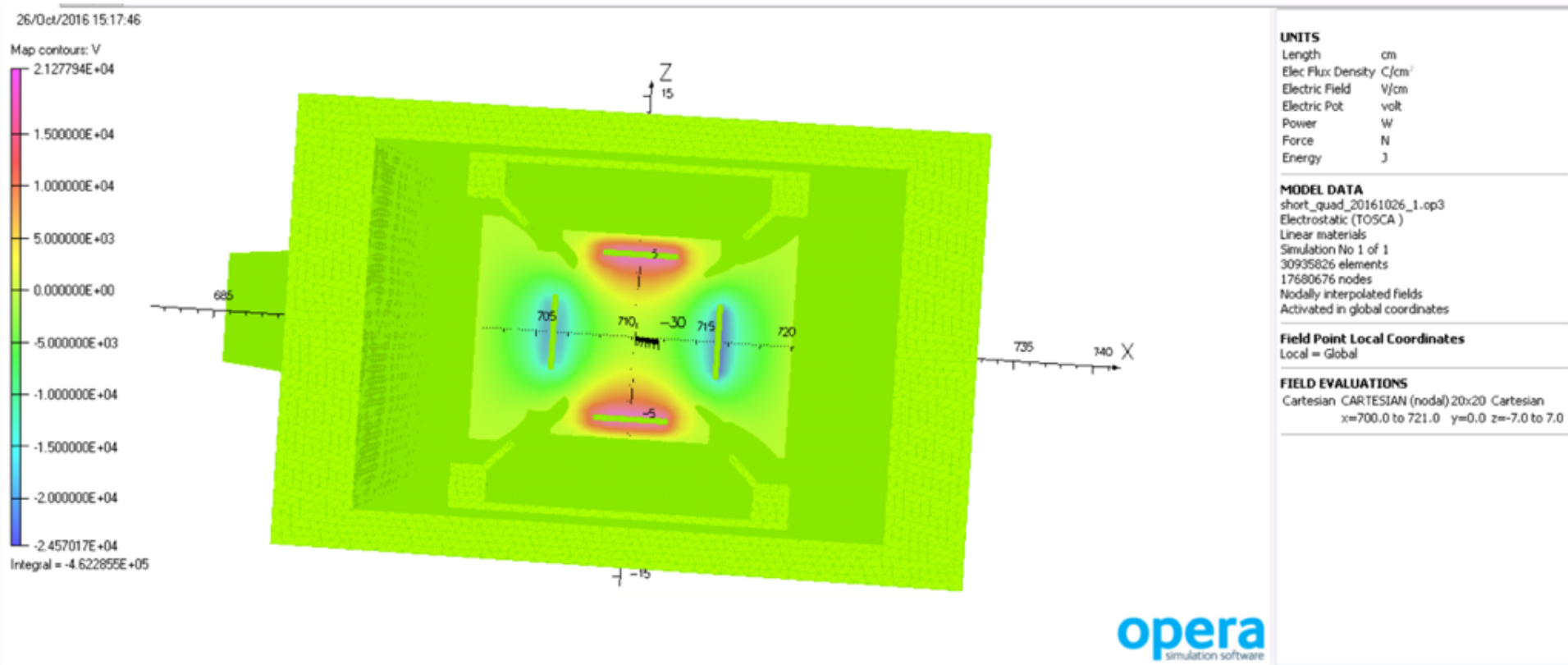


muon loss



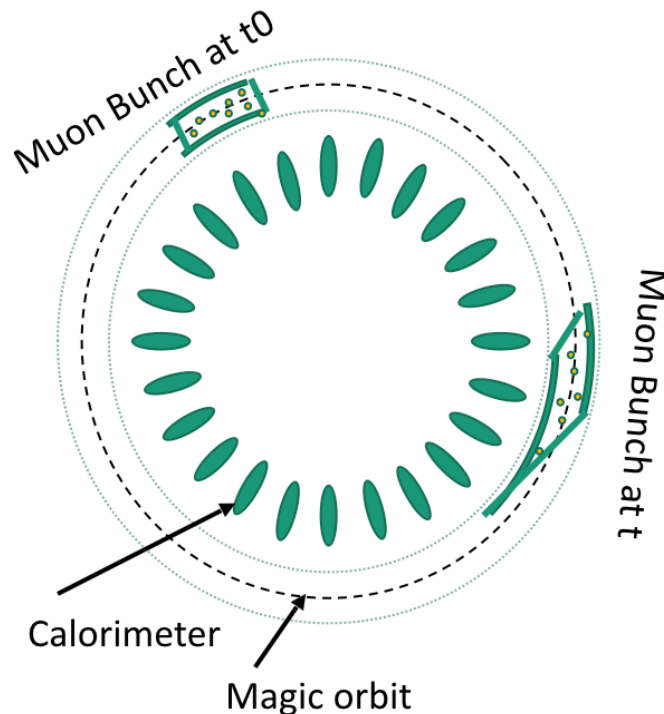
shifts in E-field and pitch ($\vec{\beta} \cdot \vec{B}$) corrections





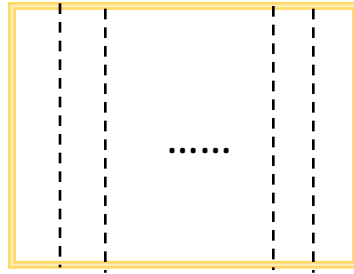
Electric Potential at the end of plates (zone map)
V=27.2kV, from Downstream $\theta=0^\circ$

$$\vec{\omega}_a = -\frac{q}{m}[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1}\right)(\vec{\beta} \cdot \vec{B})\vec{\beta} - \boxed{\left(a_\mu - \frac{1}{\gamma^2 - 1}\right)\frac{\vec{\beta} \times \vec{E}}{c}}]$$

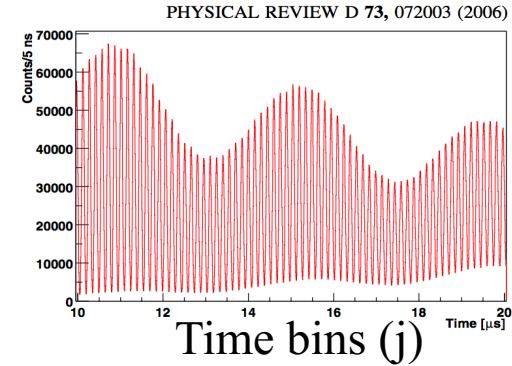


- Muons are injected into the storage ring as a bunch—radial distribution
- Muons at inner equilibrium radii will go steadily ahead of those at outer equilibrium radii ->**debunching**
- Modulation of decay positron count (fast rotation signals) can be used to study the debunching
- Fast rotation analysis: use a model of the time evolution of the bunch structure to obtain the momentum (radial) distribution of decayed muons

Two bin sets:



Radial bins (i) ($L_x=90$ mm)
(i.e., 50 bins w/width=1.8mm)

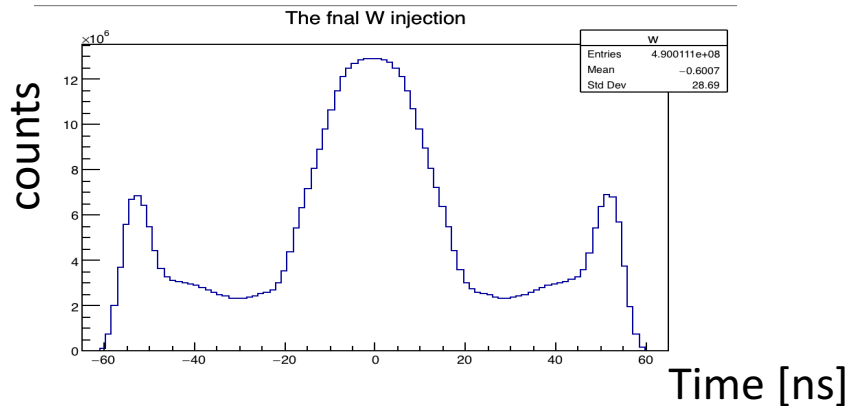


(positron count histogram)

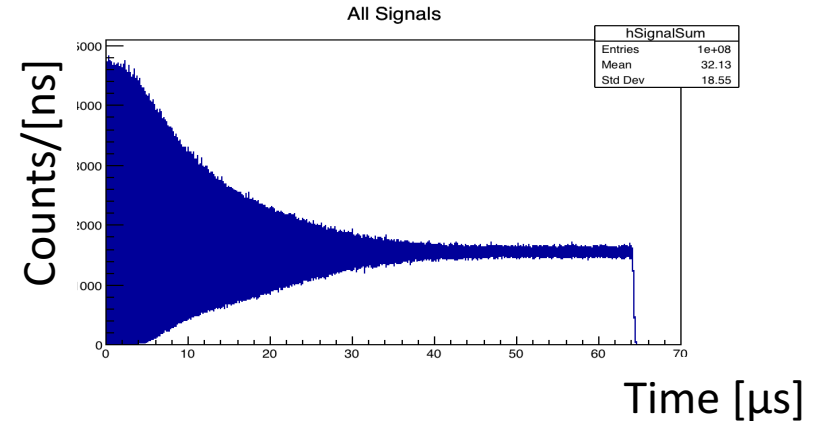
- f_i : the content of the radial bin i , fraction of the beam oscillating around radial bin i
- N_j : ($N(j)_{obs}$) counts in time bin j
- β_{ij} : contribution from radial bin i to the counts in time bin j
- C_j : ($N(j)_{exp}$) expected counts in time bin j
- Z_j : weighting factor which should equal to C_j

$$\chi^2 = \sum_j \frac{(N_j - C_j)^2}{Z_j} = \sum_j \frac{(N_j - \sum_i f_i \beta_{ij})^2}{Z_j}$$

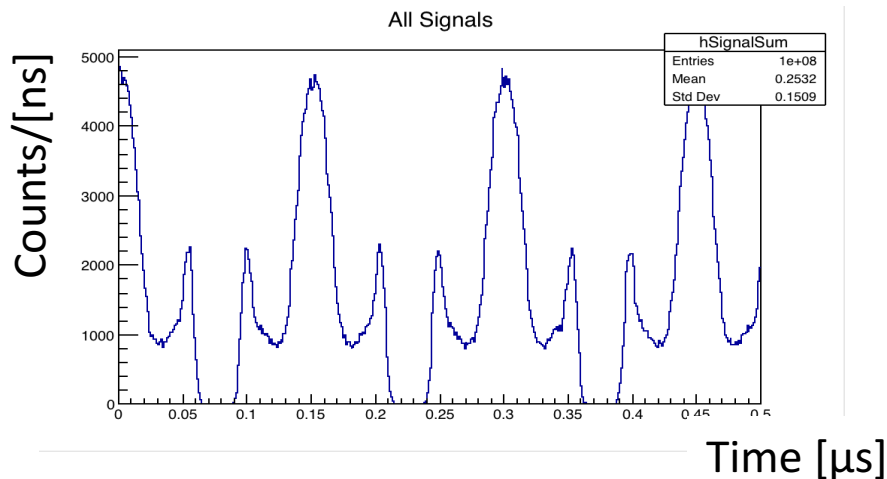
Note: The geometry factors β_{ij} are known functions of ring geometry and the apparent time structure of the injected bunch (injection zero time and beam revolution time).



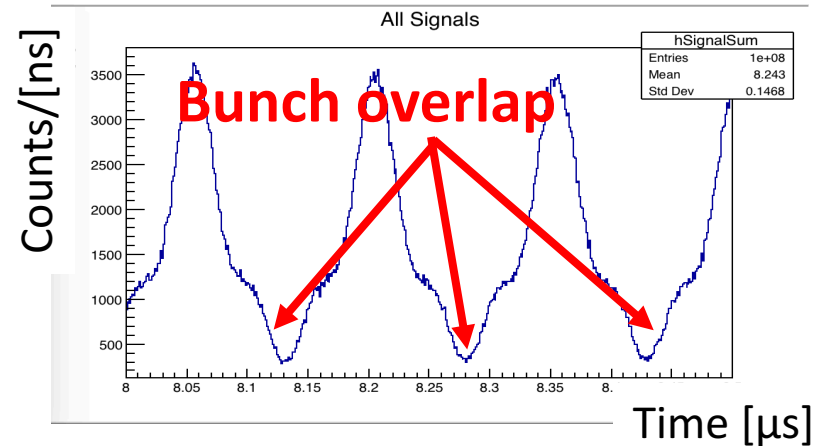
Injection Pulse



Signals seen by Detector



Signals seen by Detector (*early time*)



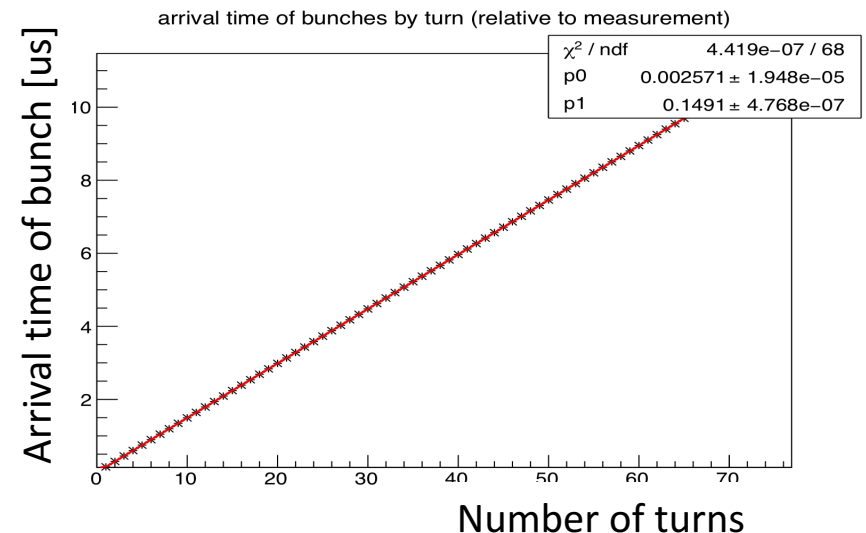
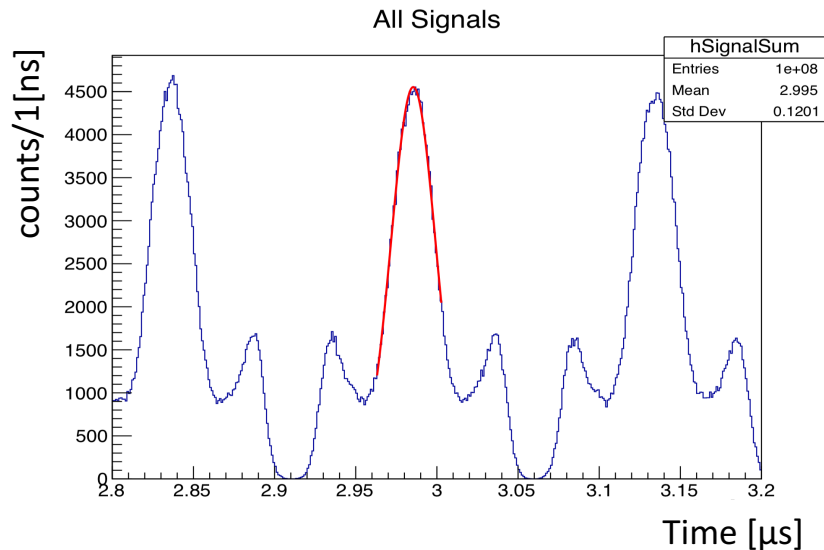
Signals seen by Detector (*late time*)

$$\chi^2 = \sum_j \frac{(N_j - C_j)^2}{Z_j} = \sum_j \frac{(N_j - \sum_i f_i \beta_{ij})^2}{Z_j}$$

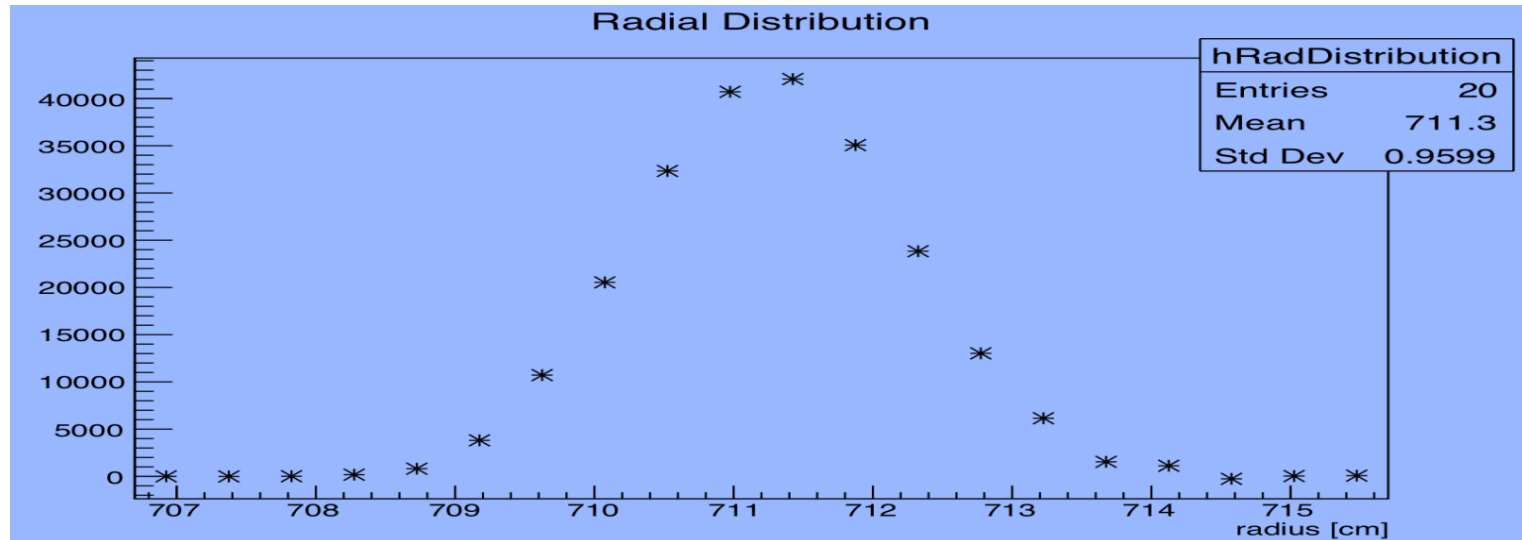
Fit the arrival time of bunches by number of turns to find out the injection zero time and beam revolution time

$$t_{peak} = nT_C + t_0$$

Arrival time of bunch by turns



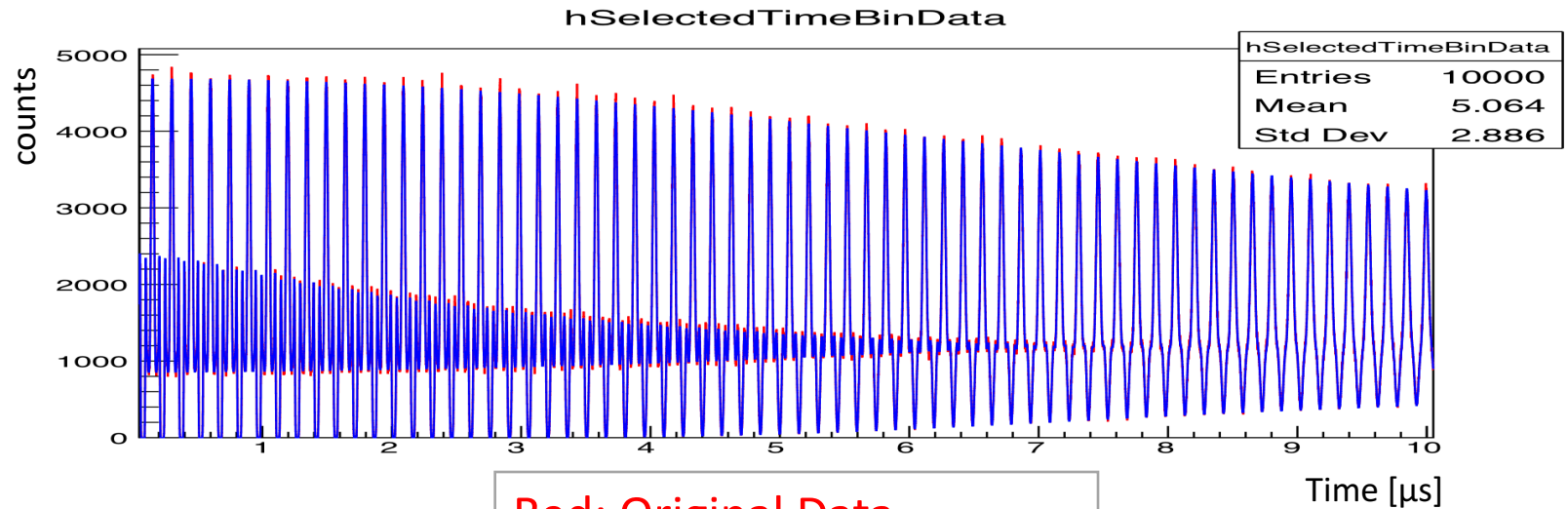
Apply the χ^2 minimization analysis and solve for the radial (momentum) bin contents:



With this distribution, we are able to evaluate the electric field corrections to g-2.

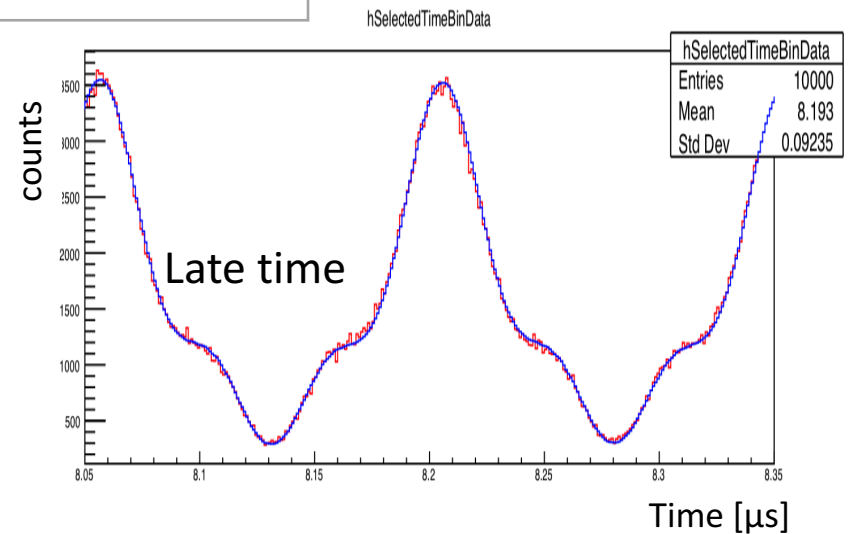
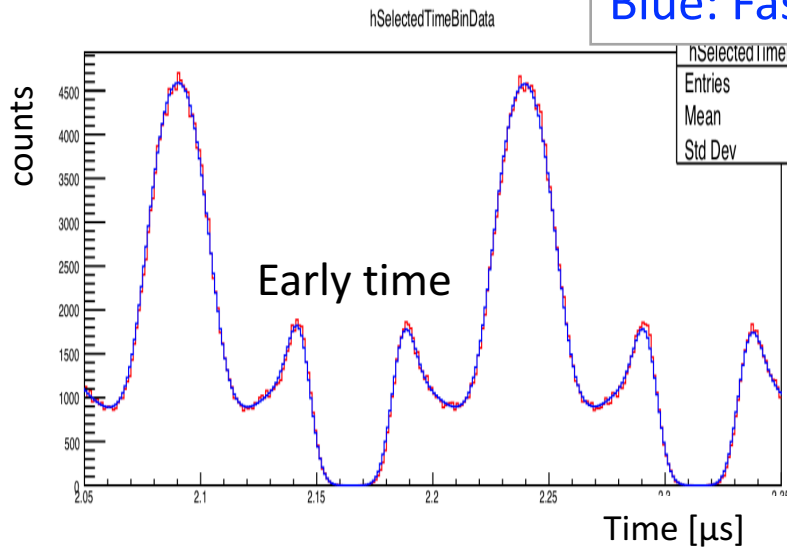
$$\Rightarrow \left\langle \frac{\delta \omega_a}{\omega_a} \right\rangle = -2\beta^2 n(1 - n) \left\langle \left(\frac{x_e}{R_o} \right)^2 \right\rangle$$

$$(x_e = R_{mean} - R_0, R_0 \text{ is the "magic" radius; } n = \frac{\rho}{v_z B_0} \frac{\partial E_R}{\partial R})$$



Red: Original Data

Blue: Fast Rotation Results



Conclusion

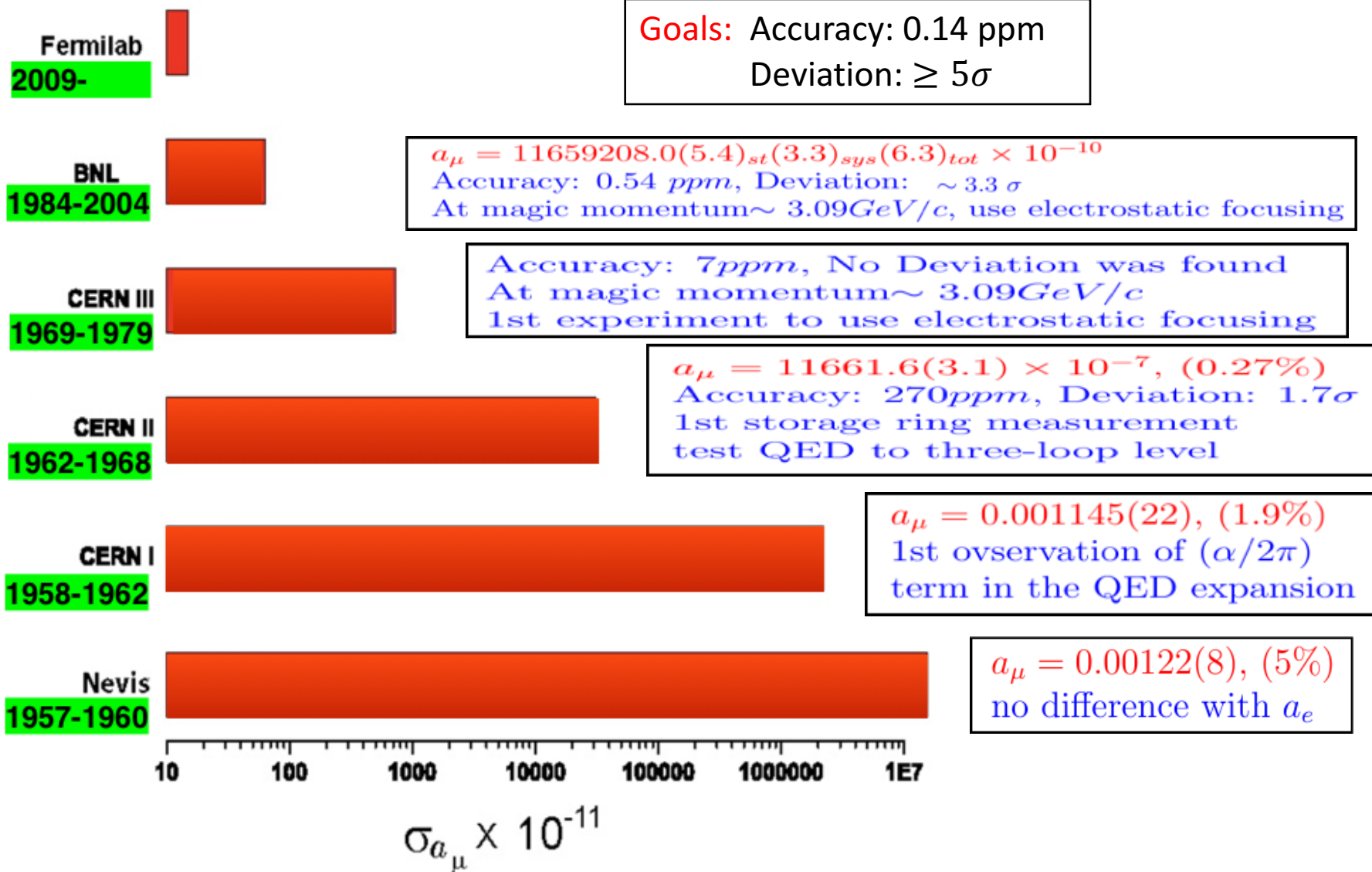
$$\left\langle \frac{\delta \omega_a}{\omega_a} \right\rangle = -2\beta^2 n(1 - n) \left\langle \left(\frac{x_e}{R_o} \right)^2 \right\rangle$$

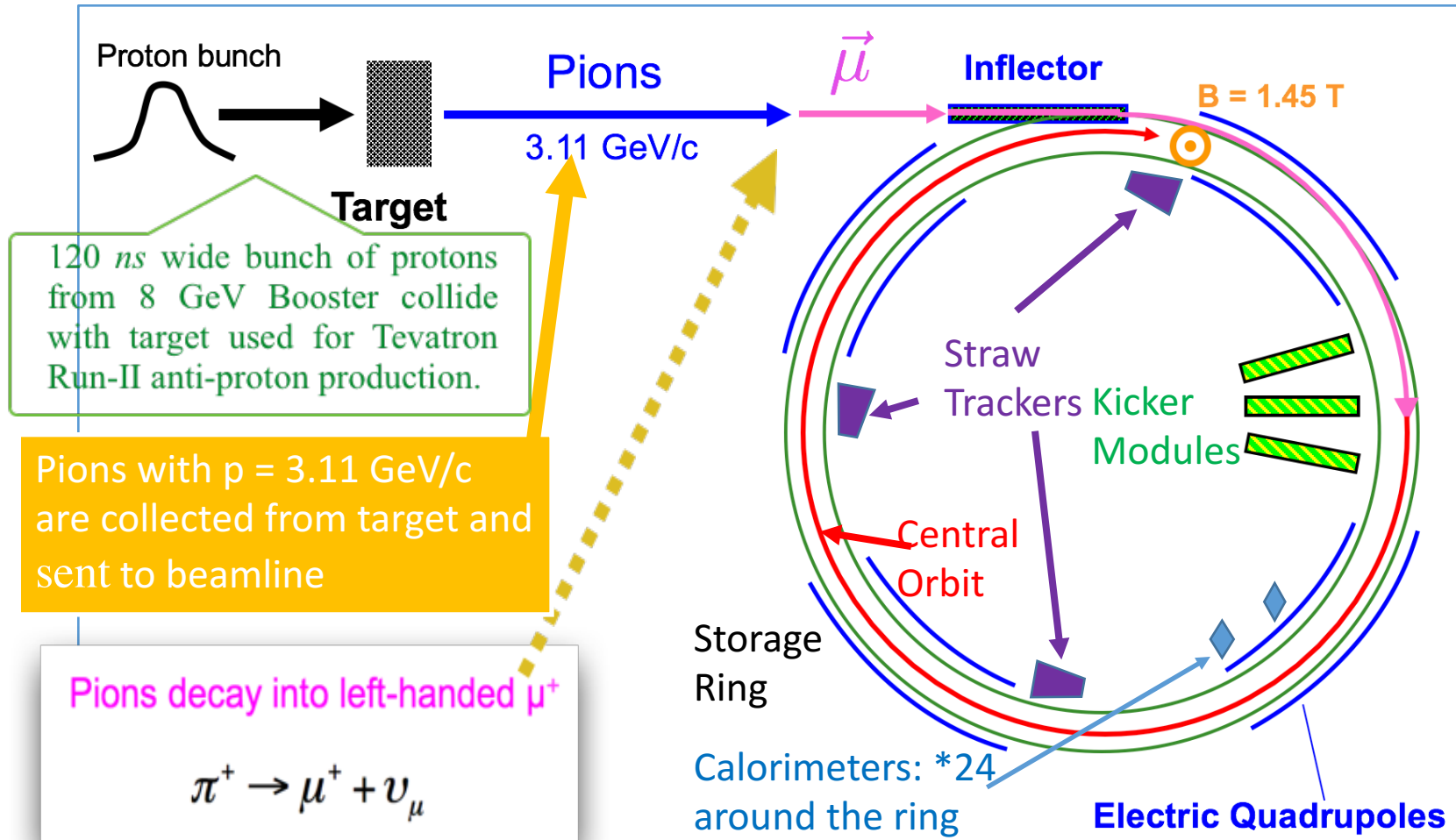
$$(x_e = R_{mean} - R_0, R_0 \text{ is the "magic" radius; } n = \frac{\rho}{v_z B_0} \frac{\partial E_R}{\partial R})$$

- The electric field plays a very important role in the Fermilab Muon g-2 Experiment in many ways, i.e., muon orbits, muon losses, E-field correction.
- The electric field has a significant effect on the muon anomalous precession frequency:
 - We need to align the electrostatic quadrupole plates very carefully;
 - We need to know the 3D electric field map;
 - We need to know the muon momentum/radius distribution through the so-called fast rotation analysis.
- The Experiment is running and it will be a great time to study muon anomaly including its electric field corrections.

—Thank You!

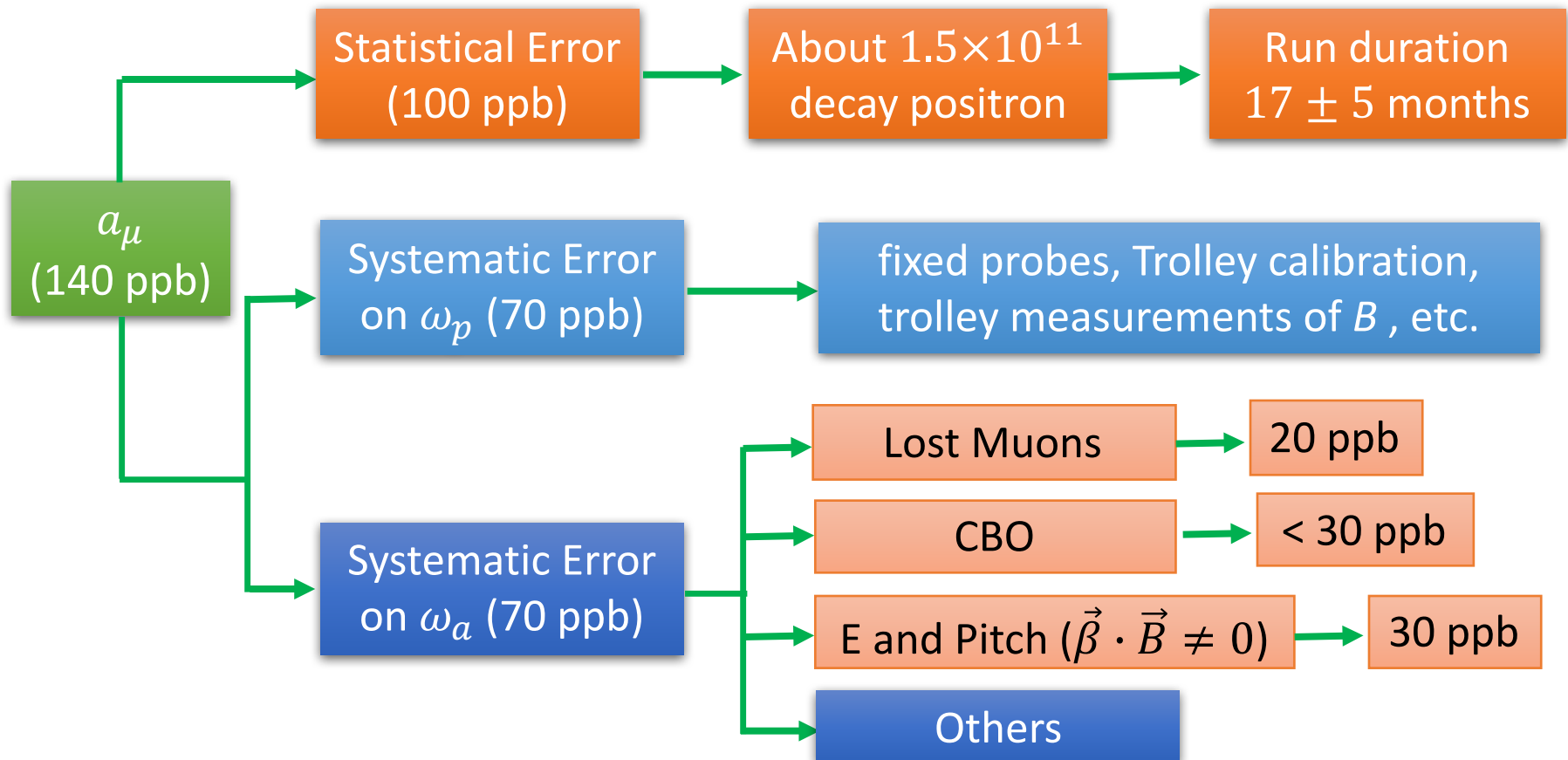
Backup



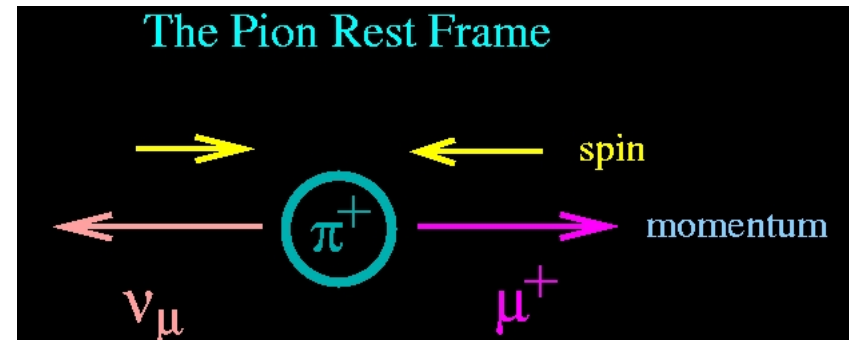


- For our final analysis, we can solve for a_μ and rewrite it as (using $\mu_e = g_e e \hbar / 4 m_e$):

$$a_\mu = \frac{\omega_a}{\omega_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$



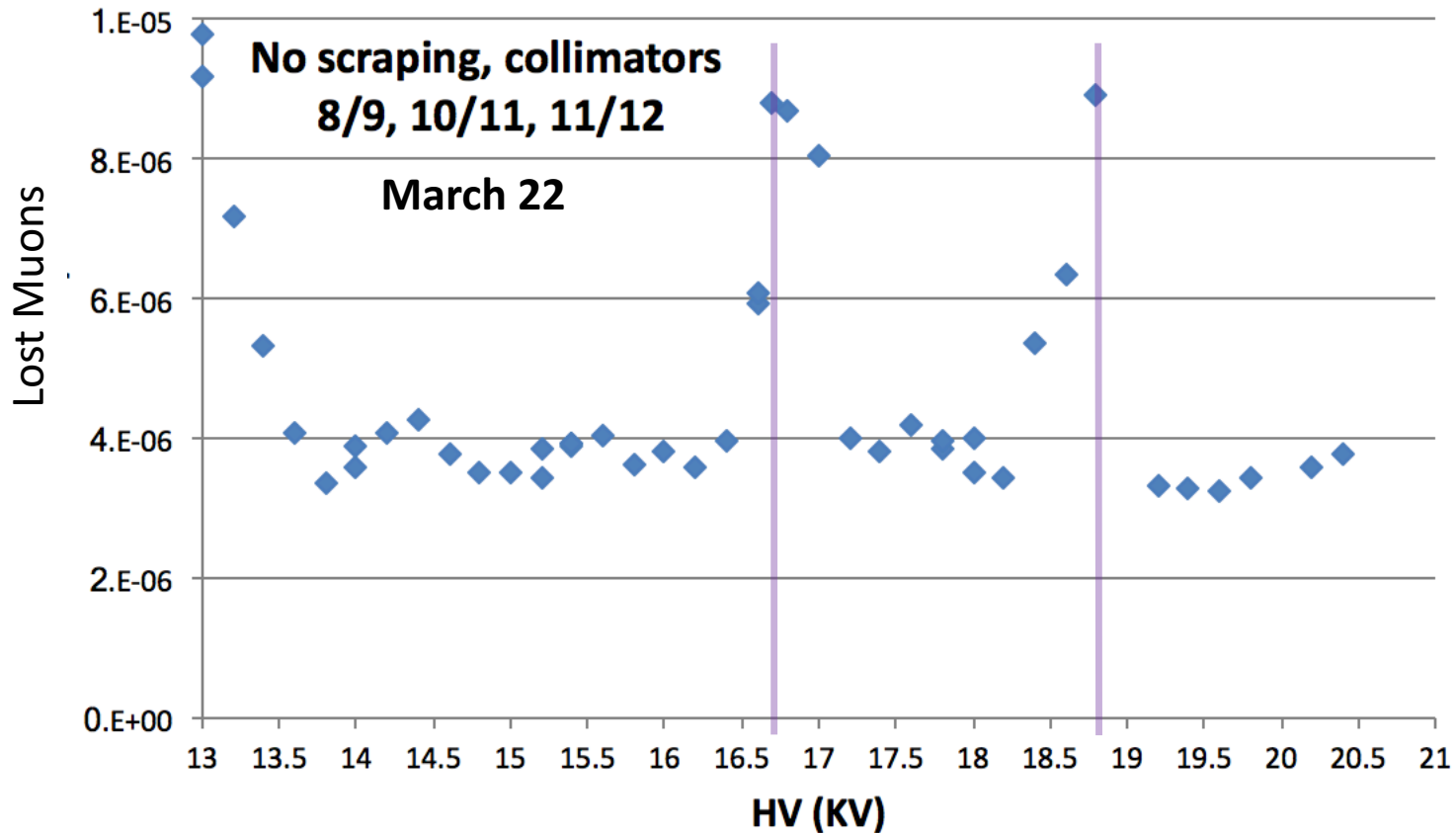
- **Muon is produced polarized:**
in-flight decay, both “forward” and “backward” muons are highly polarized



Pion decay modes:

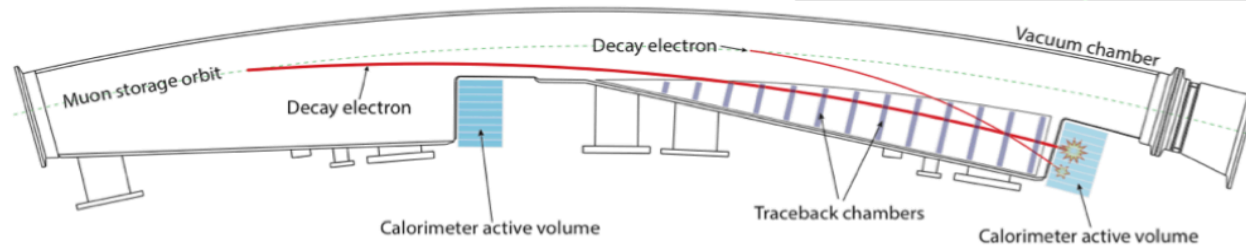
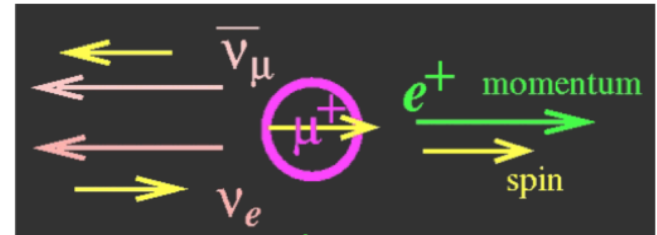
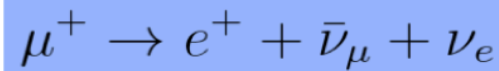
	Mode	Fraction (Γ_i/Γ)	Confidence level
Γ_1	$\mu^+ \nu_\mu$	[a] $(99.98770 \pm 0.00004) \%$	
Γ_2	$\mu^+ \nu_\mu \gamma$	[b] $(2.00 \pm 0.25) \times 10^{-4}$	
Γ_3	$e^+ \nu_e$	[a] $(1.230 \pm 0.004) \times 10^{-4}$	
Γ_4	$e^+ \nu_e \gamma$	[b] $(7.39 \pm 0.05) \times 10^{-7}$	
Γ_5	$e^+ \nu_e \pi^0$	$(1.036 \pm 0.006) \times 10^{-8}$	
Γ_6	$e^+ \nu_e e^+ e^-$	$(3.2 \pm 0.5) \times 10^{-9}$	
Γ_7	$e^+ \nu_e \nu \bar{\nu}$	$< 5 \times 10^{-6}$	90%

-- <http://pdg.lbl.gov/2014/listings/rpp2014-list-pi-plus-minus.pdf>



Quad scan for beam resonance study—Lost Muons

Muon decay:



$$N_e(t) = N_0 e^{-t/\tau} [1 - A \cos(\omega_a t + \phi_a)]$$

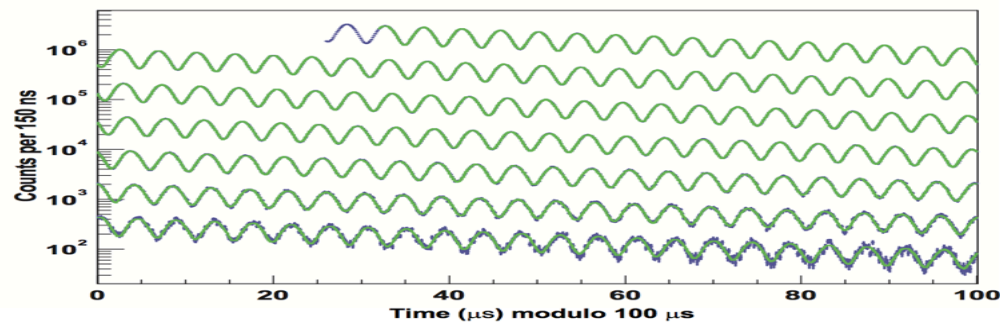
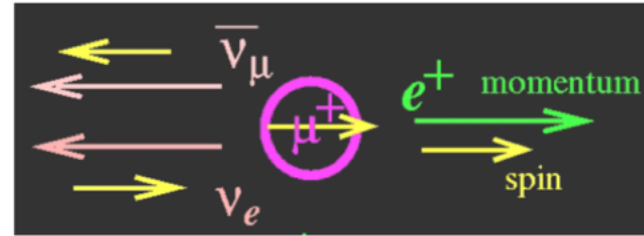
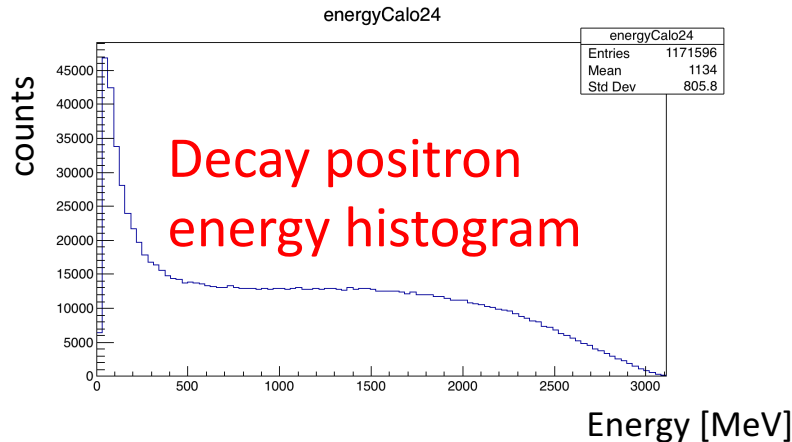
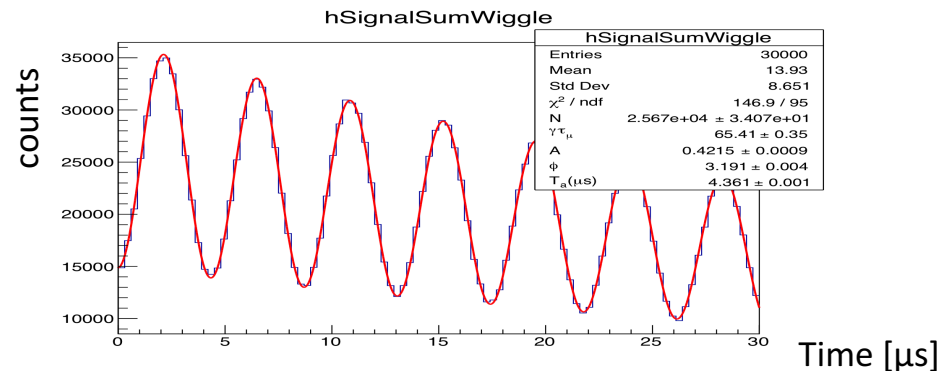
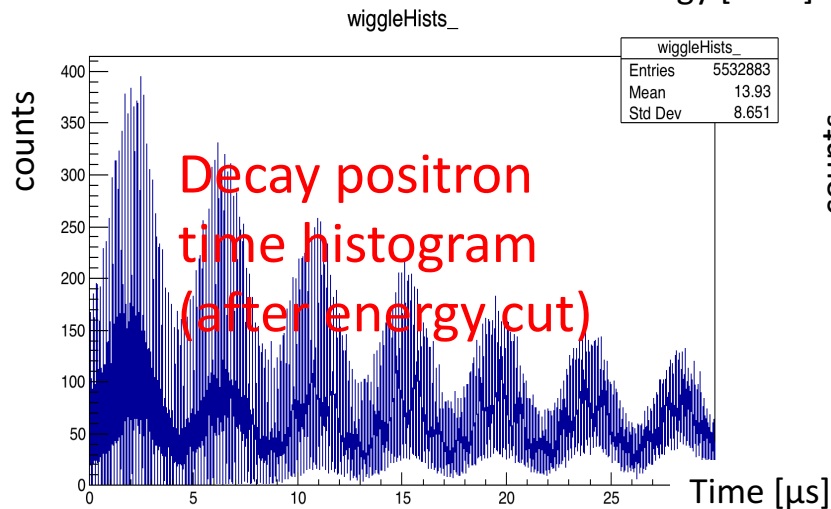


Figure 3.8: Histogram, modulo 100 μs , of the number of detected electrons above 1.8 GeV for the 2001 data set as a function of time, summed over detectors, with a least-squares fit to the spectrum superimposed. Total number of electrons is 3.6×10^9 . The data are in blue, the fit in green.

BNL g-2 PRD 2006



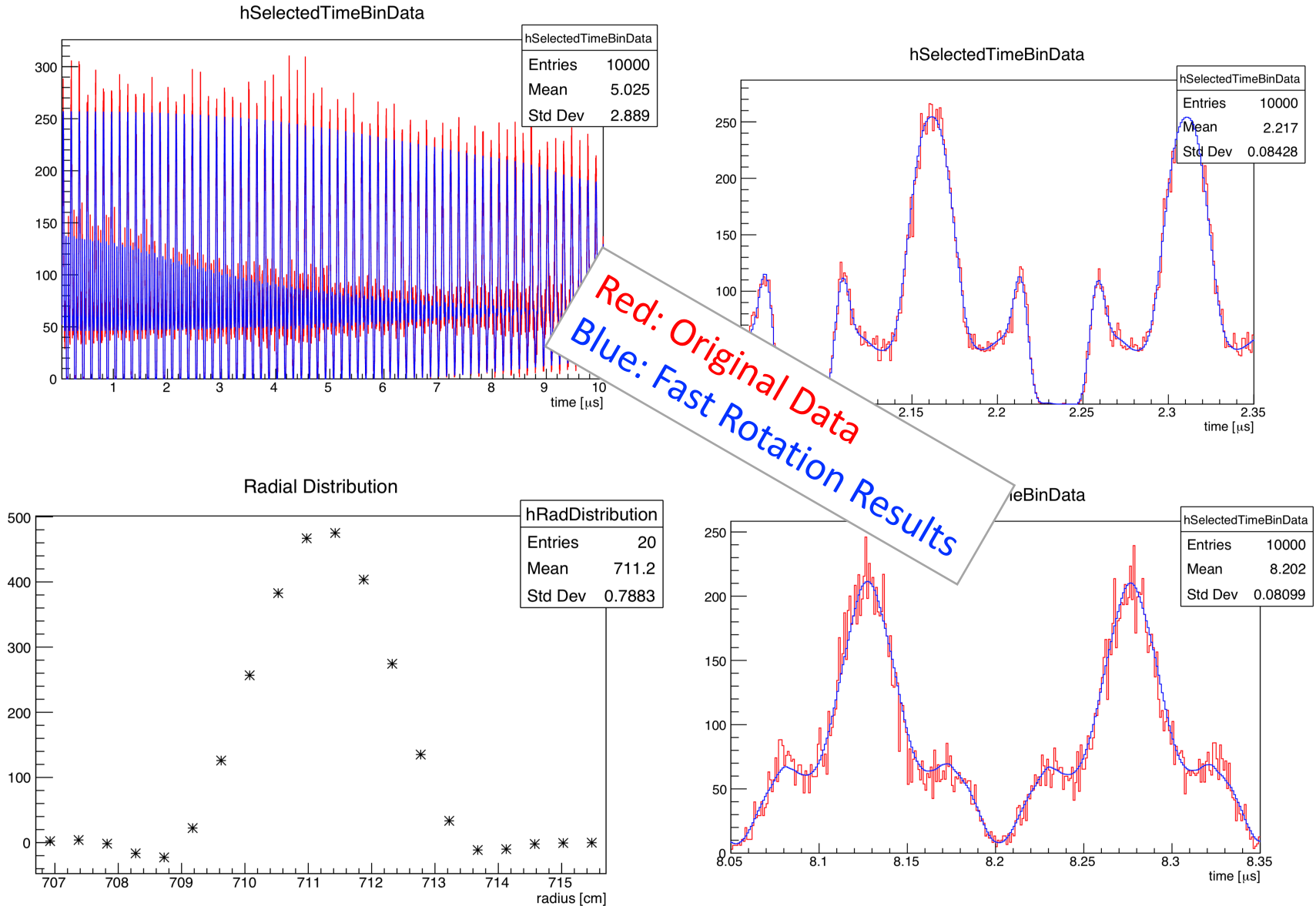
Energy Cut to select the decay positron signal



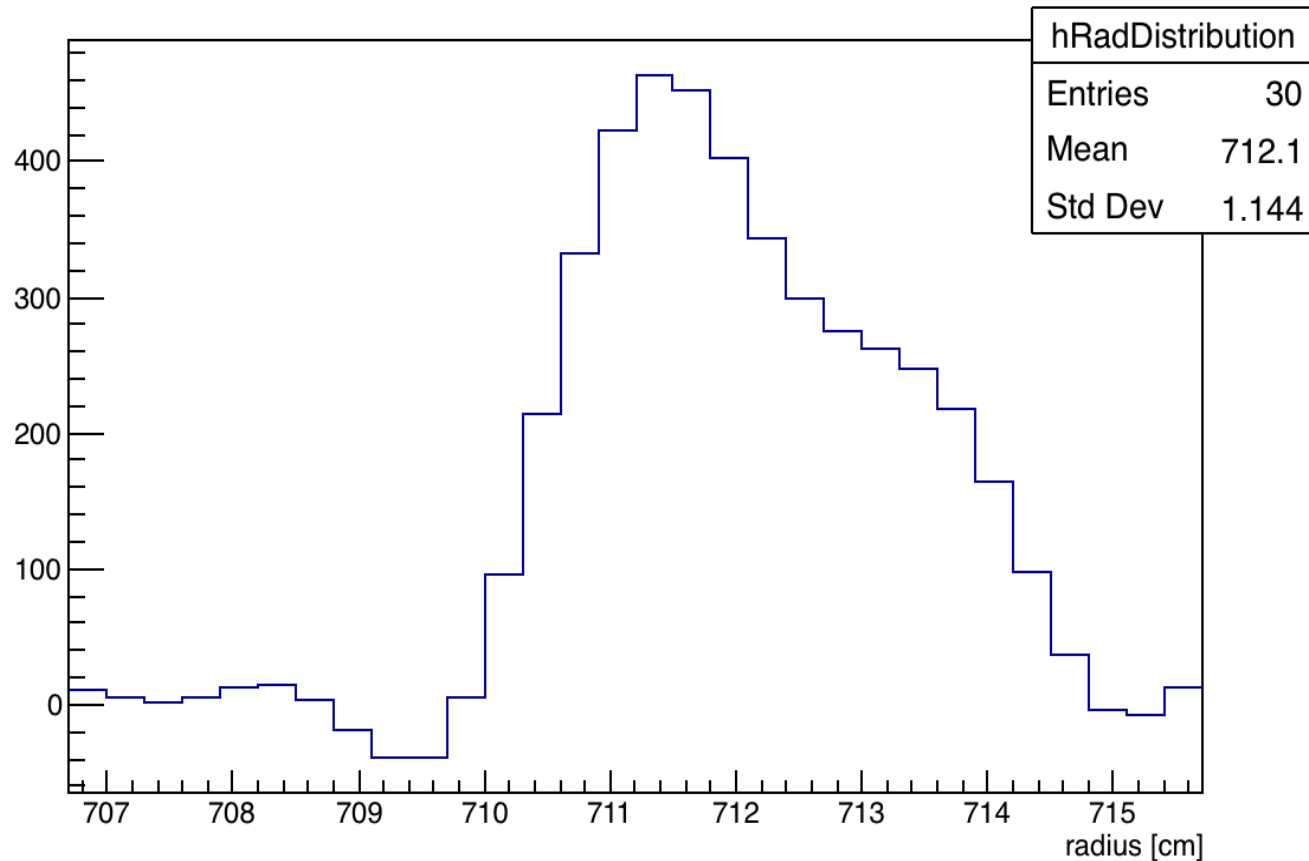
$$N_e(t) = N_0 e^{-t/\gamma\tau} [1 - A \cos(\omega_a t + \phi_a)]$$

Real situations: a lot of backgrounds, i.e., muon lifetime, g-2 frequency, CBO, Muon losses

If we can functionalize the backgrounds, we can remove them in fast rotation analysis, i.e., τ and ω_a



Example (bunch # 0): Radius distribution



Peak at 711.5 cm