

Search for Low-Mass WIMPs with SuperCDMS

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We report a first search for weakly interacting massive particles (WIMPs) using the background rejection capabilities of SuperCDMS. An exposure of 577 kg-days was analyzed for WIMPs with mass $< 30 \text{ GeV}/c^2$, with the signal region blinded. Eleven events were observed after unblinding. We set an upper limit on the spin-independent WIMP-nucleon cross section of $1.2 \times 10^{-42} \text{ cm}^2$ at $8 \text{ GeV}/c^2$. This result is in tension with WIMP interpretations of recent experiments and probes new parameter space for WIMP-nucleon scattering for WIMP masses $< 6 \text{ GeV}/c^2$.

Evidence on galactic and cosmological scales strongly indicates that $\sim 80\%$ of the matter density of the Universe consists of non-luminous, non-baryonic dark matter, whose particle nature remains unknown [1]. Weakly interacting massive particles (WIMPs) are one class of theoretically well-motivated candidates for dark matter and may be detectable by searching for keV-scale nuclear recoils in terrestrial detectors [2]. Recent excesses of events reported by CDMS II (Si) [3], CoGeNT [4],

CRESST-II [5], DAMA [6], and possible indirect evidence from gamma rays from the galactic center [7], may have been caused by a light WIMP with mass in the 6–30 GeV/c^2 range. A variety of theoretical models also favor light WIMPs in this mass range [8–15].

Since light WIMPs produce only low-energy nuclear recoils, experiments optimized for masses $\gtrsim 30 \text{ GeV}/c^2$ have searched for light WIMPs by lowering their analysis energy thresholds [16–19]. This additional sensitiv-

ity comes with higher background rates because resolution effects degrade particle discrimination at low energies. Following this approach, we analyzed low-energy recoils in the range 1.6–10 keV_{nr} (nuclear-recoil equivalent energy) from the SuperCDMS experiment at the Soudan Underground Laboratory (SUL) [20, 21]. Although background discrimination gradually degrades with decreasing event energy, some discrimination can still be achieved using the relative signals measured by the different readout channels on each detector.

SuperCDMS at Soudan is an upgrade to the Cryogenic Dark Matter Search (CDMS II) [22] with new detector hardware, and is operating in the same location with the same low-radioactivity setup [23]. The target consists of fifteen 0.6-kg cylindrical germanium crystals stacked in groups of three to form five towers. These detectors, known as iZIPs, are instrumented with interleaved ionization and phonon sensors on their flat faces. From the measured ionization and phonon energy, we derive the recoil energy and the “ionization yield,” the ratio between ionization and recoil energy. Nuclear recoils, expected from WIMPs, exhibit a reduced ionization yield compared to electron recoils, which are expected from most backgrounds. The iZIP sensor layout improves the ability to define a fiducial volume in the bulk (fiducialization) compared to the CDMS II design [24]. The fraction of the total phonon or ionization energy measured by the guard sensors provides radial fiducialization through the “radial partition” parameter, and the fraction measured by the sensors on each face provides z fiducialization through the “ z partition” parameter. Such fiducialization rejects events in the peripheral regions of the detectors. These “surface events” often suffer from reduced ionization signal, thus polluting the WIMP signal region.

The SuperCDMS payload has been operating in SUL since March 2012. The data presented here, recorded between October 2012 and June 2013, are a subset of the ongoing exposure. The seven detectors with the lowest trigger thresholds are used for this search. The remaining detectors were used to reject events with energy deposition in more than one detector. Consistency tests are used to remove periods of abnormal detector behavior and elevated noise. After accounting for these losses, the exposure is 577 kg-days. To prevent bias when defining the event-selection criteria, all single-detector hits with recoil energies in the range 1.6–10 keV_{nr} and ionization energy consistent with nuclear recoils were removed from the sample, i.e. blinded. An exception was made for periods following ²⁵²Cf calibrations, when background rates were higher because of neutron activation of the detectors and their copper housings. This “open” dataset constitutes 97 kg-days of exposure that is distinct from the 577 kg-days of data analyzed for WIMPs, and was not used in the final limit calculation or to optimize selection criteria.

For the detectors analyzed, the standard deviation of

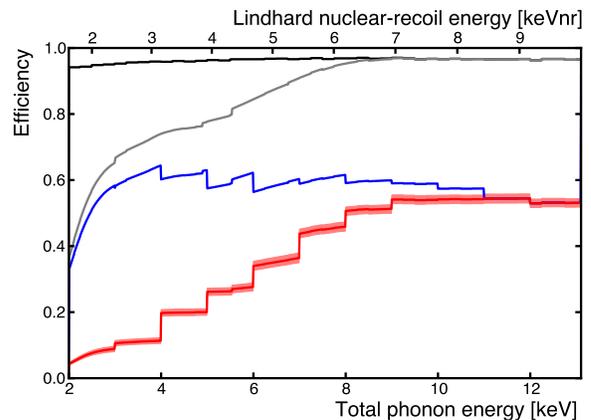


FIG. 1. Cumulative efficiencies after sequential application of each stage of event selection. From top to bottom, these are data-quality criteria, trigger and analysis threshold efficiencies, preselection criteria, and BDT discrimination with 68% C.L. (stat. + syst.) uncertainty band. Steps are due to the combination of smooth fits of the trigger efficiency and binned measurements. For illustrative purposes, an approximate nuclear-recoil energy scale is provided.

the baseline noise is $\lesssim 260$ eV for summed phonon channels and $\lesssim 460$ eV_{ee} (electron-recoil equivalent) for individual ionization channels. The electron- and nuclear-recoil energy scales are calibrated in a fashion similar to the CDMS II light-WIMP search [17] using ¹³³Ba and ²⁵²Cf sources respectively. A small ($\lesssim 10\%$) variation of the phonon signal gain with the cryostat base temperature, which varied over the range 54–62 mK, is taken into account by the phonon calibration. In each detector, the mean ionization energy of nuclear recoils as a function of total phonon energy, as determined from ²⁵²Cf calibration data, is consistent with, or slightly below, the prediction of Lindhard [25, 26]. A nuclear-recoil band was constructed by accepting events within 3σ of the mean ionization energy. Nuclear-recoil equivalent energies are reconstructed from the total phonon energy by subtracting the contribution of Luke-Neganov phonons [27, 28] corresponding to the mean nuclear-recoil ionization response for the respective total phonon energy.

Hardware trigger thresholds for each detector were adjusted several times during the WIMP search. For each period of constant trigger threshold, the trigger efficiencies as functions of total phonon energy were measured using ¹³³Ba calibration data. The fit results were found to be consistent with, and more precise than, ones obtained using ²⁵²Cf and multiple-hit WIMP-search data. Analysis thresholds are set to be 1σ below the energy at which the detector trigger efficiency is 50%. For some time intervals, analysis thresholds are raised further according to baseline noise levels. The combined efficiency is an exposure-weighted sum of the measured efficiency for each detector and period, shown in Fig. 1.

153 To be selected as WIMP candidates, triggered events
 154 had to pass three levels of data-selection criteria: data
 155 quality, preselection, and event discrimination. Figure 1
 156 shows the cumulative efficiency after applying each level
 157 of selection criteria and the analysis thresholds. The
 158 first level of criteria (data quality) rejects poorly recon-
 159 structed and noise-induced events. Periods of abnormal
 160 noise are removed by requiring that the pre-trigger base-
 161 line noise of each event be consistent with normal periods.
 162 Spurious triggers caused by electronic glitches and low-
 163 frequency noise in the phonon channels, which populate
 164 the low-energy region, were rejected using a pulse-shape
 165 discrimination method. Using a Monte Carlo pulse simu-
 166 lation that added experimental noise to template pulses
 167 to account for variation in the noise environment, the
 168 WIMP acceptance of this data-quality selection was de-
 169 termined to be $\gtrsim 95\%$.

170 The second level of event-selection criteria (prese-
 171 lection) removes event configurations inconsistent with
 172 WIMPs. Events coincident with the muon veto are re-
 173 jected (98.7% acceptance). A single-scatter requirement
 174 removes events with energy depositions in multiple de-
 175 tectors, a common signature for background interactions
 176 but not expected for a WIMP-nucleon scatter ($>99\%$ ac-
 177 ceptance, with losses due to noise fluctuations). We also
 178 require events to lie within the 3σ nuclear-recoil band
 179 and to have phonon partitions consistent with bulk nu-
 180 clear recoils. A loose fiducial volume constructed from
 181 the ionization partitions further restricts events to be
 182 consistent with bulk nuclear recoils. In the radial direc-
 183 tion, events near the detectors' sidewalls are rejected by
 184 requiring the guard electrodes on both faces to be within
 185 2σ from the mean of the baseline noise. For one detector
 186 (T5Z3) that has a malfunctioning guard electrode on one
 187 side, this requirement is applied on only the functioning
 188 face. A second detector (T5Z2) suffered sporadic excess
 189 noise on one guard, so only the guard on the functioning
 190 face was used for part of the dataset. In the z direction,
 191 events on the flat faces are excluded by requiring that the
 192 inner electrodes on each side measure similar ionization
 193 energies [24].

194 The final level of event selection (discrimination) uses
 195 a boosted decision tree (BDT) [29]. The discriminators
 196 used by the BDT are the total phonon energy, ioniza-
 197 tion energy, phonon radial partition and phonon z par-
 198 tition. Near threshold, the latter two variables provide
 199 identification of surface events superior to the ionization
 200 partitions, while the two energy quantities together opti-
 201 mize the discrimination at low energy where the electron-
 202 and nuclear-recoil bands overlap. A BDT was trained
 203 for each detector using simulated background events (de-
 204 scribed below) and nuclear recoils from ^{252}Cf calibration
 205 weighted to mimic a WIMP energy spectrum, accounting
 206 for the selection criteria acceptance. The BDT discrim-
 207 ination thresholds for individual detectors were chosen
 208 simultaneously to minimize the expected 90% confidence

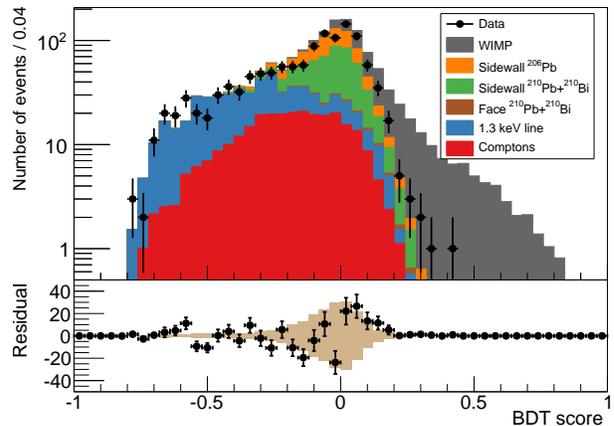


FIG. 2. Top: Stacked histogram showing the components of the background model passing the preselection criteria, summed over all detectors (neutron backgrounds are negligible and not included). For comparison, a 10 GeV/ c^2 WIMP with cross section 6×10^{-42} cm 2 is shown on top of the total background. Events passing preselection criteria are overlaid (markers with statistical errors). A p-value statistic comparing the data to background model is 14% for this selection. Bottom: Difference between the data and the background expectation. Tan bars indicate the systematic uncertainty (68% C.L.) on the background estimate. The background spectrum was computed prior to unblinding and was not fit or rescaled to match the data.

level (C.L.) Poisson upper limit of the rate of passing events per WIMP exposure. The BDT was trained and optimized separately for 5, 7, 10, and 15 GeV/ c^2 WIMPs. Events that pass any of the four WIMP-mass optimizations are accepted into the signal region as candidates. When a limit is set using the optimum interval method [30, 31], this acceptance technique provides sensitivity to a range of masses, but incurs only modest sensitivity loss compared to an analysis optimized at every WIMP mass. In addition to the BDT, two other discrimination methods were developed and similarly optimized for WIMP masses between 5 and 15 GeV/ c^2 . The BDT was chosen as the primary discrimination method before unblinding because of its better expected sensitivity on the background simulation data.

The acceptance of the preselection criteria and the BDT was evaluated using the fraction of ^{252}Cf nuclear recoils passing as a function of the energy. Unlike WIMPs, the ^{252}Cf neutrons can multiply scatter within a single detector, which necessitates correcting the acceptance upwards by $\sim 25\%$ above ~ 5 keV $_{\text{nr}}$ based on a Geant4 [32] neutron simulation, which includes constraints on the resolution effects and the size of the fiducial volume. The uncertainty of the total acceptance is dominated by systematic uncertainty on the size of the fiducial volume and is shown in Fig. 1.

A background model was developed that includes

Compton recoils from the gamma-ray background; 1.1–292
 1.3 keV X-rays and Auger electrons from L-shell electron-293
 capture (EC) decay of ^{65}Zn , ^{68}Ga , ^{68}Ge and ^{71}Ge ; and294
 decay products from ^{210}Pb contamination on the detec-295
 tors and their copper housings. We normalize the flat296
 Compton background to the observed rate of electron297
 recoils in the range 2.6–5.1 keV_{ee}. The average rate of298
 L-shell EC events is estimated by scaling the observed299
 rate in the open dataset by the ratio of the K-shell event300
 rates in the WIMP-search and open datasets. We use301
 Geant4 to simulate the implantation and decay of ^{222}Rn 302
 daughters starting from ^{214}Po as described in [24]. Back-303
 ground components from ^{210}Pb decay products (betas,304
 conversion electrons, X-rays), ^{210}Bi betas, and ^{206}Pb nu-305
 clei from ^{210}Po decays are considered, with rates nor-306
 malized to the alpha and ^{206}Pb decay products of ^{210}Po 307
 under the assumption of secular equilibrium. 308

The background model is implemented using events309
 from high-energy sidebands and calibration data as310
 templates for low-energy backgrounds. Ionization and311
 phonon pulses are scaled to lower energies, injected with312
 noise from randomly triggered events throughout the313
 data, and reconstructed as actual data. ^{133}Ba calibra-314
 tion data and K-shell EC events are used as templates315
 for the Compton recoils and L-shell EC events, respec-316
 tively. Templates for ^{210}Pb daughters are sampled from317
 high-energy betas and ^{206}Pb recoils. 318

Figure 2 shows the individual components of the back-319
 ground model as a function of the 10 GeV/ c^2 BDT dis-320
 crimination parameter after applying the preselection cri-321
 teria. This background model was finalized prior to un-322
 blinding and predicted $6.1_{-0.8}^{+1.1}$ (stat.+syst.) events pass-323
 ing the BDT selection. Simulations of radiogenic and324
 cosmogenic neutrons, as described in [22], predict an ad-325
 ditional 0.098 ± 0.015 (stat.) events. These estimates326
 included only known systematic effects. Because the ac-
 curacy in background modeling required for a full like-
 lihood analysis is difficult to achieve in a blind analysis
 of this type, the decision was made before unblinding to
 report an upper limit on the WIMP-nucleon cross section

Upon unblinding, eleven candidates were observed as
 indicated in Fig. 3. The events were found to be of
 high quality and occurring during good periods of experi-
 mental operation, except for the lowest-energy candidate,
 which has an abnormal pulse shape and is suspected to be
 noise. As seen in Table I, the observed number of events
 is consistent with the background prediction for most de-
 tectors. However, the three high-energy events in detec-
 tor T5Z3 strongly disagree with the background predic-
 tion. The probability to observe at least this many back-
 ground events on this detector is 4×10^{-4} . These events
 are observed on the only detector in this dataset that
 has an ionization guard electrode shorted to ground. Al-
 though the background model was developed to account
 for the shorted channel, we realized after unblinding that
 the altered electric field may have affected the selection

of background model templates, potentially making the
 background estimate on this detector inaccurate.

The background model is compared to unblinded
 events passing all preselection criteria in Fig. 2. The
 systematic uncertainty, shown with tan fill, is dominated
 by the uncertainty of the expected ionization of sidewall
 events originating from ^{210}Pb and ^{210}Bi . P-value statis-
 tics comparing the data passing the preselection criteria
 with the blind background model prediction range from
 8–26% for the BDTs trained to each of the four masses.
 This reasonable compatibility, based on the sum over all
 detectors, suggests that the background model correctly
 reproduces most features of the true background.

A 90% C.L. upper limit on the spin-independent
 WIMP-nucleon cross section was calculated using the op-
 timum interval method without background subtraction.
 The calculation used standard halo assumptions as dis-
 cussed in [33]. The result is shown in Fig. 4. Statistical
 and systematic uncertainties in the fiducial-volume effi-
 ciency, the nuclear-recoil energy scale, and the trigger
 efficiency were propagated into the limit by Monte Carlo
 and are represented by the narrow gray band around the
 limit. The limit is consistent with the expected sensitiv-
 ity for masses below 10 GeV/ c^2 as shown by the green
 band in Fig. 4. The discrepancy above 10 GeV/ c^2 is due
 to the three high-energy events in T5Z3, which are in
 tension with the background expectation.

This work represents the first search for WIMPs with
 the background rejection capability of SuperCDMS de-
 tectors. A physically motivated background model gener-
 ally agrees with the data, except for the detector with a
 shorted ionization guard. This analysis strongly disfavors
 a WIMP-nucleon scattering interpretation of the excess
 reported by CoGeNT, which also uses a germanium tar-
 get. Similar tension exists with WIMP interpretations

Detector	Candidate energies [keV _{nr}]	Expected background
T1Z1	—	$0.03_{-0.01}^{+0.01}$
T2Z1	1.7, 1.8	$1.4_{-0.2}^{+0.2}$
T2Z2	1.9, 2.7	$1.8_{-0.3}^{+0.4}$
T4Z2	—	$0.04_{-0.02}^{+0.02}$
T4Z3	—	$1.7_{-0.3}^{+0.4}$
T5Z2	1.9, 2.3, 3.0, 5.8	$1.1_{-0.3}^{+0.3}$
T5Z3	7.0, 7.8, 9.4	$0.13_{-0.04}^{+0.06}$

TABLE I. Energies of candidate events in each detector, labeled by tower (first number) and position within tower from top to bottom (second number). Expected background is based on the model used to train the BDT and includes the estimated systematic uncertainty. Differences in expected background across detectors reflect different trigger thresholds and background event rates. Event energies are calculated using the measured mean ionization energy for nuclear recoils.

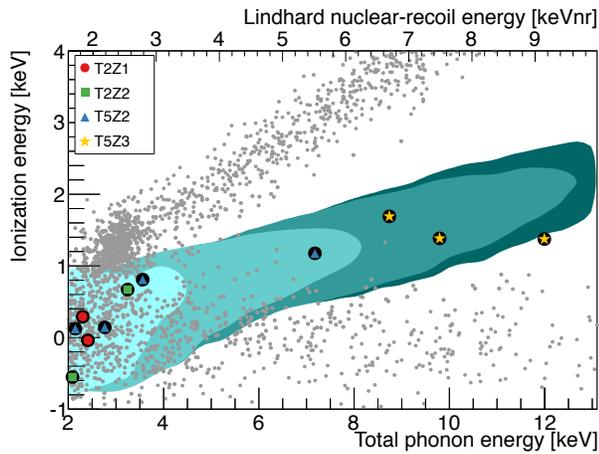


FIG. 3. Small gray dots are all veto-anticoincident single-scatter events within the ionization-partition fiducial volume that pass the data-quality selection criteria. Large encircled shapes are the 11 candidate events. Overlapping shaded regions (from light to dark) are the 95% confidence contours expected for 5, 7, 10 and 15 GeV/c^2 WIMPs, after application of all selection criteria. The three highest-energy events occur on detector T5Z3, which has a shorted ionization guard. The band of events above the expected signal contours corresponds to bulk electron recoils, including the 1.3 keV activation line at a total phonon energy of ~ 3 keV. High-radius events near the detector sidewalls form the wide band of events with near-zero ionization energy. For illustrative purposes, an approximate nuclear-recoil energy scale is provided.

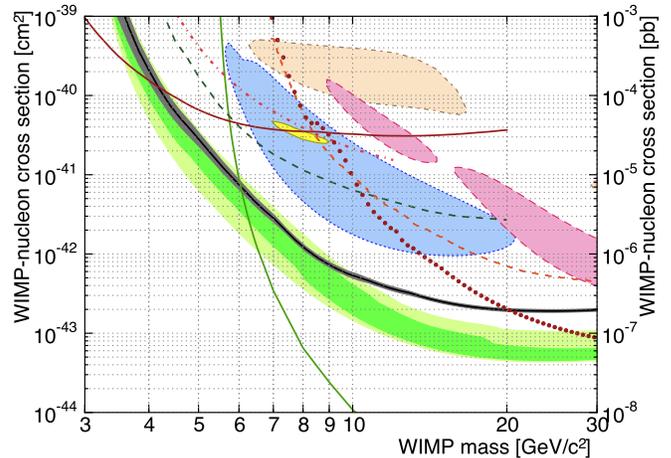


FIG. 4. The 90% confidence upper limit (solid black) based on all observed events is shown with 95% C.L. systematic uncertainty band (gray). The pre-unblinding expected sensitivity in the absence of a signal is shown as 68% (dark green) and 95% (light green) C.L. bands. The disagreement between the limit and sensitivity at high WIMP mass is due to the events in T5Z3. Closed contours shown are CDMS II Si [3] (dotted blue, 90% C.L.), CoGeNT [4] (yellow, 90% C.L.), CRESST-II [5] (dashed pink, 95% C.L.), and DAMA/LIBRA [34] (dash-dotted tan, 90% C.L.). 90% C.L. exclusion limits shown are CDMS II Ge [22] (dotted dark red), CDMS II Ge low-threshold [17] (dashed-dotted red), CDMSlite [20] (solid dark red), LUX [35] (solid green), XENON10 S2-only [19, 36] (dashed dark green), and EDELWEISS low-threshold [18] (dashed orange).

of several other experiments, including CDMS II (Si), assuming spin-independent interactions and a standard halo model. New regions of WIMP-nucleon scattering for WIMP masses below $6 \text{ GeV}/c^2$ are excluded.

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