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## **SIMULATION AND VERIFICATION OF DPA IN MATERIALS\*<sup>†</sup>**

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### **Abstract**

A recent implementation and verification of consistent modeling of displacements per atom (DPA) in the MARS15 code are described for high-energy particles and heavy ions.

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## SIMULATION AND VERIFICATION OF DPA IN MATERIALS

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### 1. Introduction

Radiation damage is displacement of atoms from their equilibrium position in a crystalline lattice due to irradiation with formation of interstitial atoms and vacancies in the lattice. Resulting deterioration of material (critical) properties is measured – in the most universal way – as a function of displacements per target atom (DPA). DPA is a strong function of projectile type, energy and charge as well as material properties including its temperature. The phenomenon becomes very serious for high-intensity beams especially for high-charge heavy ions ( $\sim z^2$ ), being identified, for example at FRIB and FAIR, as one of the critical issues, limiting lifetime of targets to as low as a few weeks. A recent implementation of consistent DPA modeling into the MARS15 code [1] and its verification are described in this paper.

### 2. DPA Modeling in MARS15 Code

A model used in the MARS15 code for DPA calculations in electromagnetic processes [2] has been extended to an arbitrary projectile of energy ranging from 1 keV to 10 TeV. A primary knock-on atom (PKA) created in nuclear collisions can generate a cascade of atomic displacements. This is taken into account via damage function  $\nu(T)$ . DPA is expressed in terms of damage cross section  $\sigma_d$ :

$$\sigma_d(E) = \int_{T_d}^{T_{\max}} \frac{d\sigma(E, T)}{dT} \nu(T) dT,$$

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where  $E$  is kinetic energy of the projectile,  $T$  is kinetic energy transferred to the recoil atom,  $T_d$  is the displacement energy, and  $T_{\max}$  is the highest recoil energy according to kinematics. In a modified Kinchin-Pease model [3],  $v(T)$  is zero at  $T < T_d$ , unity at  $T_d < T < 2.5T_d$ , and  $k(T)E_d/2T_d$  at  $2.5T_d < T$ , where  $E_d$  is “damage” energy available to generate atomic displacements by elastic collisions (Fig. 1).  $T_d$  is an irregular function of atomic number ( $\sim 40$  eV). The displacement efficiency,  $k(T)$ , introduced as a result of simulation studies on evolution of atomic displacement cascades [4], drops from 1.4 to 0.3 once the PKA energy is increased from 0.1 to 100 keV, and exhibits a weak dependence on target material and temperature.

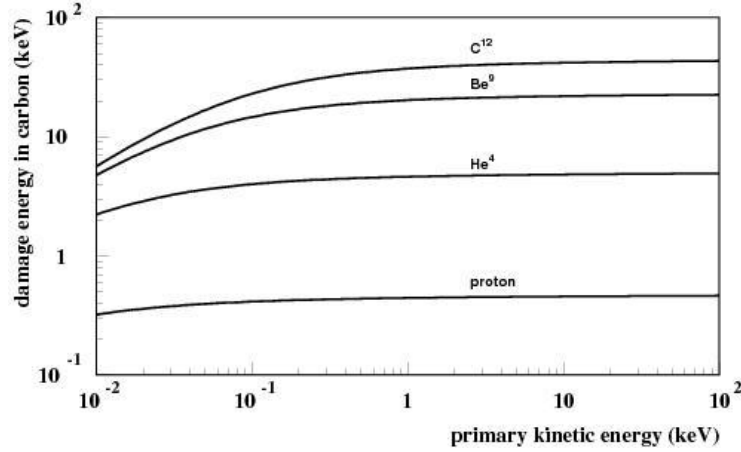


Figure 1. Damage energy  $E_d$  versus PKA energy.

For electromagnetic elastic (Coulomb) scattering, the Rutherford cross-section with Mott corrections and nuclear form-factors (a factor of two effect) are used in our model. Resulting displacement cross-sections due to Coulomb scattering are shown in Fig. 2 for various projectiles on silicon and carbon targets. For elementary particles, energy dependence of  $\sigma_d$  disappears above 2-3 GeV, while it continues to higher energies for heavy ions. For projectiles heavier than a proton,  $\sigma_d$  grows with a projectile charge  $z$  as  $z^2/\beta^2$  at  $\gamma\beta > 0.01$ , where  $\beta$  is a projectile velocity. All products of elastic and inelastic nuclear interactions as well as Coulomb elastic scattering of transported charged particles (hadrons, electrons, muons and heavy ions) from 1 keV to 10 TeV contribute to DPA in the MARS15 model.

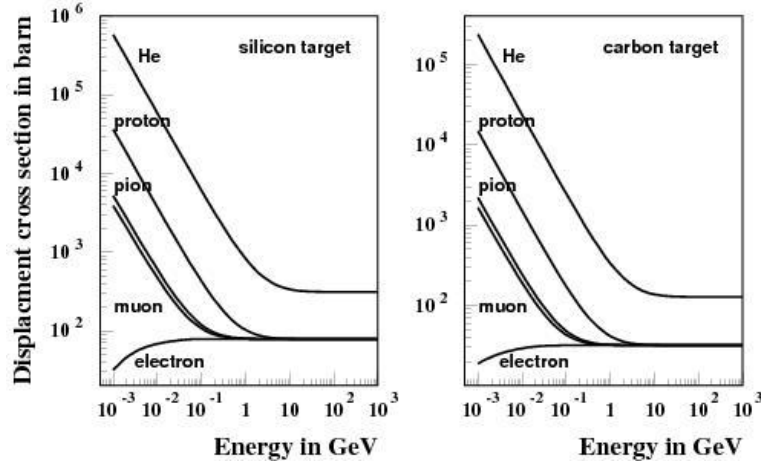


Figure 2. Displacement cross-section in silicon and carbon for various projectiles.

### 3. DPA Calculation Comparison

In this section, results on DPA calculated with the new MARS15 model for three cases are compared with those obtained with other DPA-capable codes, SRIM/TRIM, PHITS and MCNPX. The first case is a 1-GeV proton beam of 1-cm<sup>2</sup> area on a 3-mm thick iron target. SRIM, PHITS and MCNPX results are courtesy of Susana Reyes. As one can see in Table 1, there is a quite substantial difference between the predictions, with SRIM giving a very small value and the MARS15 result being a factor of 2.6 to 2.9 above those by PHITS and MCNPX. Calculated with MARS15 contributions to DPA of physics processes are as follows: 75.5% nuclear inelastic, 16% nuclear elastic, 2.75% electromagnetic elastic, 5.5% low-energy neutrons, and 0.25% electrons. The dominance of nuclear interactions in this case explains the above differences.

Table 1. DPA for 1-GeV protons on 3-mm iron.

| Code    | SRIM     | PHITS    | MCNPX    | MARS15   |
|---------|----------|----------|----------|----------|
| DPA/pot | 1.18e-22 | 2.96e-21 | 3.35e-21 | 8.73e-21 |

The second case is a 0.32-GeV/u Uranium beam of 9-cm<sup>2</sup> area on a 1-mm thick beryllium target. SRIM and PHITS results are again a courtesy of Susana Reyes. Table 2 shows that SRIM and MARS15 results are now very close to each other, while those calculated with PHITS are a factor of 70 lower. Calculated with MARS15 contributions to DPA of physics processes are as

follows: 0.3% nuclear inelastic, 99.06% electromagnetic elastic, 0.02% low-energy neutrons, and 0.62% electrons. The dominant role of Coulomb scattering in this case explains the similarity of the SRIM and MARS15 predictions.

Table 2. DPA for 0.32-GeV/u Uranium on 1-mm beryllium target.

| Code    | SRIM     | PHITS    | MARS15   |
|---------|----------|----------|----------|
| DPA/pot | 2.97e-20 | 5.02e-22 | 2.13e-20 |

The third case is a 0.13-GeV/u Germanium beam of  $0.004\text{-cm}^2$  area on a 1.2-mm thick tungsten target. TRIM and PHITS results are a courtesy of Yosuke Iwamoto. Table 3 gives calculated DPA values in the first hundred microns of the target. The difference between TRIM and MARS15 needs to be understood.

Table 3. Entrance DPA for 0.13-GeV/u Germanium on 1.2-mm tungsten target.

| Code    | TRIM     | PHITS    | MARS15   |
|---------|----------|----------|----------|
| DPA/pot | 8.04e-16 | 1.25e-17 | 1.43e-16 |

#### 4. BLIP Beam Tests for 0.7-MW NuMI/LBNE Target

A majority of data on radiation damage is available for reactor neutrons. Studies with hundred MeV protons [5] have revealed that a threshold of about 0.2 DPA exists for carbon composites and graphite. MARS15 studies helped realize that the BLIP beam tests with 0.165-GeV protons can emulate the NuMI neutrino target situation for a 120-GeV proton beam (Table 4). It turns out that despite a substantial difference in the beam energies in these cases, nuclear interactions and Coulomb scattering contribute about the same way (45-50% each) to the peak DPA in thick graphite targets irradiated at these two facilities.

Table 4. Peak DPA in POCO graphite targets at BLIP and NuMI.

| Target | $E_p$ (GeV) | Beam $\sigma$ (mm) | $N_p$ (1/yr) | DPA (1/yr) |
|--------|-------------|--------------------|--------------|------------|
| NuMI   | 120         | 1.1                | 4.0e20       | 0.45       |
| BLIP   | 0.165       | 4.23               | 1.12e22      | 1.5        |

#### References

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