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First Observation of the Decay $B_s^0 \to D_s^- D_s^+$ and Measurement of Its Branching Ratio

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We report the first measurement of the exclusive decay $B_s^0 \to D_s^- D_s^+$ using 355 pb^{-1} of data collected by the CDF II detector in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV at the Fermilab Tevatron. We measure the relative branching ratio $\mathcal{B}(B_s^0 \to D_s^- D_s^+)/\mathcal{B}(B^0 \to D^- D_s^+) = 1.44^{+0.48}_{-0.44}$. Using the world average value for $\mathcal{B}(B^0 \to D^- D_s^+)$, we find $\mathcal{B}(B_s^0 \to D_s^- D_s^+) = (9.4^{+4.4}_{-4.2}) \times 10^{-3}$. This provides a lower bound $\Delta\Gamma_s^{CP}/\Gamma_s \geq 2\mathcal{B}(B_s^0 \to D_s^- D_s^+) > 1.2 \times 10^{-2}$ at 95% C.L.

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The $B_s^0 - \bar{B}_s^0$ system exhibits mixing, with two distinct mass eigenstates B_H and B_L having a mass difference $\Delta m_s = m_s^H - m_s^L$, which has recently been measured [1]. In the standard model, these two states have decay widths Γ_s^L and Γ_s^H, with difference $\Delta\Gamma_s = \Gamma_s^L - \Gamma_s^H$, and average $\Gamma_s = (\Gamma_s^L + \Gamma_s^H)/2$. To good approximation, the two mass eigenstates are expected to be eigenstates of CP: B_s^{even} and B_s^{odd}, so that $\Delta\Gamma_s \approx \Delta\Gamma_s^{CP}$, where $\Delta\Gamma_s^{CP} = \Gamma(B_s^{\text{even}}) - \Gamma(B_s^{\text{odd}})$.Measuring the B_s^0 decay rate to $D_s^{(*)-} D_s^{(*)+}$, where $D_s^{(*)\pm}$ stands for either $D_s^{\pm}$ or $D_s^{*\pm}$, determines $\Delta\Gamma_s^{CP}/\Gamma_s$ assuming the $b \to c\bar{c}s$ transitions are dominated by these decays, and neglecting small CP-odd components [2]:

$$\frac{\Delta\Gamma_s^{CP}}{\Gamma_s} = 2\mathcal{B}(B_s^0 \to D_s^{(*)-} D_s^{(*)+}). \tag{1}$$

The inclusive measurement of $B_s^0 \to D_s^{(*)-} D_s^{(*)+}$ decay rate has been reported previously [3] using $B_s^0 \to \phi\phi X$ correlations. In this letter we present the first observation of the exclusive decay $B_s^0 \to D_s^- D_s^+$ [4], measure the ratio of its branching fraction with respect to that for $B^0 \to D^- D_s^+$, and set a lower bound on $\Delta\Gamma_s^{CP}/\Gamma_s$. We use CDF II detector data corresponding to 355 pb^{-1} of integrated luminosity of $p\overline{p}$ collisions at $\sqrt{s} = 1.96$ TeV, at the Fermilab Tevatron [5].

This analysis depends primarily on the charged particle tracking systems. Charged particle tracks are reconstructed using the hits in the silicon microstrip detector system and the central outer tracker (COT) in the pseudorapidity range $|\eta| \leq 1.0$, where η is defined as $-\ln\tan(\theta/2)$, and θ represents the angle between the particle and the proton beam direction [6]. Both detectors are inside a 1.4 T uniform magnetic field. The silicon detector is composed of L00 (single layer of silicon microstrip sensors close to the beam pipe), SVX II (five cylindrical layers of double-sided sensors), and ISL (intermediate silicon layers), providing up to 8 coordinate measurements in the r-ϕ view [7]. Surrounding the SVX is the COT, a cylindrical drift chamber with 96 layers of sense wires [8].

A sample rich in charm and beauty hadrons is selected by a three-level displaced track trigger. At level 1, tracks are reconstructed in the COT by the track trigger processor (XFT) [9]. The trigger requires two tracks with transverse momenta $p_T > 2$ GeV/c and the scalar sum $p_{T1} + p_{T2} > 4.0$ GeV/c. The level 2 silicon vertex trigger (SVT) [10] associates SVX II r-ϕ position measurements with XFT tracks and provides a precise measurement of the track impact parameter (d_0), the distance of closest approach of the track helix to the beam axis in the transverse plane. Decays of heavy flavor particles are identified by requiring two tracks with $0.12\ \text{mm} \leq d_0 \leq 1$ mm and an opening angle in the transverse plane $2^\circ \leq |\Delta\phi| \leq 90^\circ$. A requirement $L_{xy} > 0.2$ mm is also applied, where L_{xy} is defined as the distance in the transverse plane from the beam line to the two-track vertex projected onto the two-track momentum vector. The level 3 trigger applies the level 1 and level 2 selection requirements after a full event reconstruction.

We measure the branching fraction ratio $\mathcal{B}(B_s^0 \to D_s^- D_s^+)/\mathcal{B}(B^0 \to D^- D_s^+)$ in which the D_s^+ meson decay rates and part of the systematic uncertainties cancel. In searching for $B_s^0 \to D_s^- D_s^+$ we require the decay $D_s^- \to \phi\pi^-$. To enhance the search sensitivity we reconstruct D_s^+ meson candidates in the $\phi\pi^+$, $\overline{K}^{*0}K^+$, or $\pi^+\pi^+\pi^-$ decay channels for both the B^0 and B_s^0 signals. The ratio is measured independently for three D_s^+ decay modes, and is calculated using

$$\frac{\mathcal{B}(B_s^0 \to D_s^- D_s^+)}{\mathcal{B}(B^0 \to D^- D_s^+)} = \frac{N_{B_s^0}}{N_{B^0}} \frac{\epsilon_{B^0}}{\epsilon_{B_s^0}} \frac{f_d}{f_s} \frac{\mathcal{B}(D^- \to K^+\pi^-\pi^-)}{\mathcal{B}(D_s^- \to \phi\pi^-)}, \quad (2)$$

where $N_{B_s^0}$ and N_{B^0} are the measured signal yields, $\epsilon_{B^0}/\epsilon_{B_s^0}$ is the ratio of reconstruction and trigger efficiencies extracted from Monte Carlo simulation, f_s/f_d is the ratio of b quark fragmentation fractions into B_s^0 and B^0 mesons, and $\mathcal{B}(D^- \to K^+\pi^-\pi^-)/\mathcal{B}(D_s^- \to \phi\pi^-)$ is the ratio of branching fractions determined by other experiments.

All tracks used in reconstruction must have $p_T > 350$ MeV/c and are assumed to be either pions or kaons depending on the specific reconstruction hypothesis. The reconstruction of $B^0 \to D^-(K^+\pi^-\pi^-)D_s^+(\phi\pi^+)$, for example, begins by searching for $D_s^+ \to \phi\pi^+$ candidates. We require two oppositely charged tracks to form $\phi \to K^+K^-$, and then add a third track to form $D_s^+ \to \phi\pi^+$. The reconstruction of D^- mesons uses the $D^- \to K^+\pi^-\pi^-$ mode. We reconstruct $B^0 \to D^- D_s^+$ candidates by applying a fit to six tracks with constraints on a primary B meson decay vertex, two secondary D meson decay vertices, and the masses of the D mesons.

Monte Carlo simulations are used to optimize the selection requirements, to derive fitting functions for signal and background, and to determine the trigger and reconstruction efficiencies. Single B hadrons are generated without fragmentation products of underlying event particles, and their decays are simulated using EvtGen[11]. The detector response, including the trigger, is modeled using the CDF simulation package [12]. The selection requirements are optimized by maximizing the significance of the Monte Carlo simulated signal, scaled to the expected yield, relative to the combinatorial background using a method valid for low statistics [13]. Combinatorial background is fitted in the interval $[5.4, 6.0]$ GeV/c^2 and extrapolated into a 60 MeV/c^2 wide signal region centered around the appropriate B-meson mass. Selection requirements are made on the minimum p_T of the tracks, the impact parameter of the B meson, the χ^2 masses of ϕ and K^{*0} candidates. We also make requirements on the significance of the L_{xy} measurement, $L_{xy}/\sigma(L_{xy})$, where $\sigma(L_{xy})$ is the L_{xy} uncertainty, of B and D meson vertices, and of the displacement of the D meson vertices with respect to the B meson vertex. For decays involving resonant states, we require 1010 MeV/c^2 $< m(K^+K^-) <$ 1029 MeV/c^2 for ϕ candidates, and 840 MeV/c^2 $< m(K^-\pi^+) <$ 940 MeV/c^2, for K^{*0} candidates. The background from $\overline{B^0} \to D^+(K^-\pi^+\pi^+)D_s^-(\phi\pi^-)$ is removed from the $B_s^0 \to D_s^-(\phi\pi^-)D_s^+(\overline{K}^{*0}K^+)$ signal by reconstructing $D_s^+ \to \overline{K}^{*0}K^+$ as $D^+ \to K^-\pi^+\pi^+$ and removing events with the D^+ candidate mass in the range 1845 MeV/c^2 $< m(D^+) <$ 1893 MeV/c^2.

Figure 1 shows the reconstructed mass spectra for $B_s^0 \to D_s^- D_s^+$ and $B^0 \to D^- D_s^+$ decays. The signal yields $N_{B_s^0}$ and N_{B^0} are extracted from a binned likelihood fit of these spectra. The fitting functions for all the modes have terms describing the signal, combinatorial background, partially reconstructed B hadrons, and contributions from B decays to different D_s^+ decay modes. The combinatorial background is represented by the sum of a constant plus an exponential. The signal and partially reconstructed modes are fitted with templates that

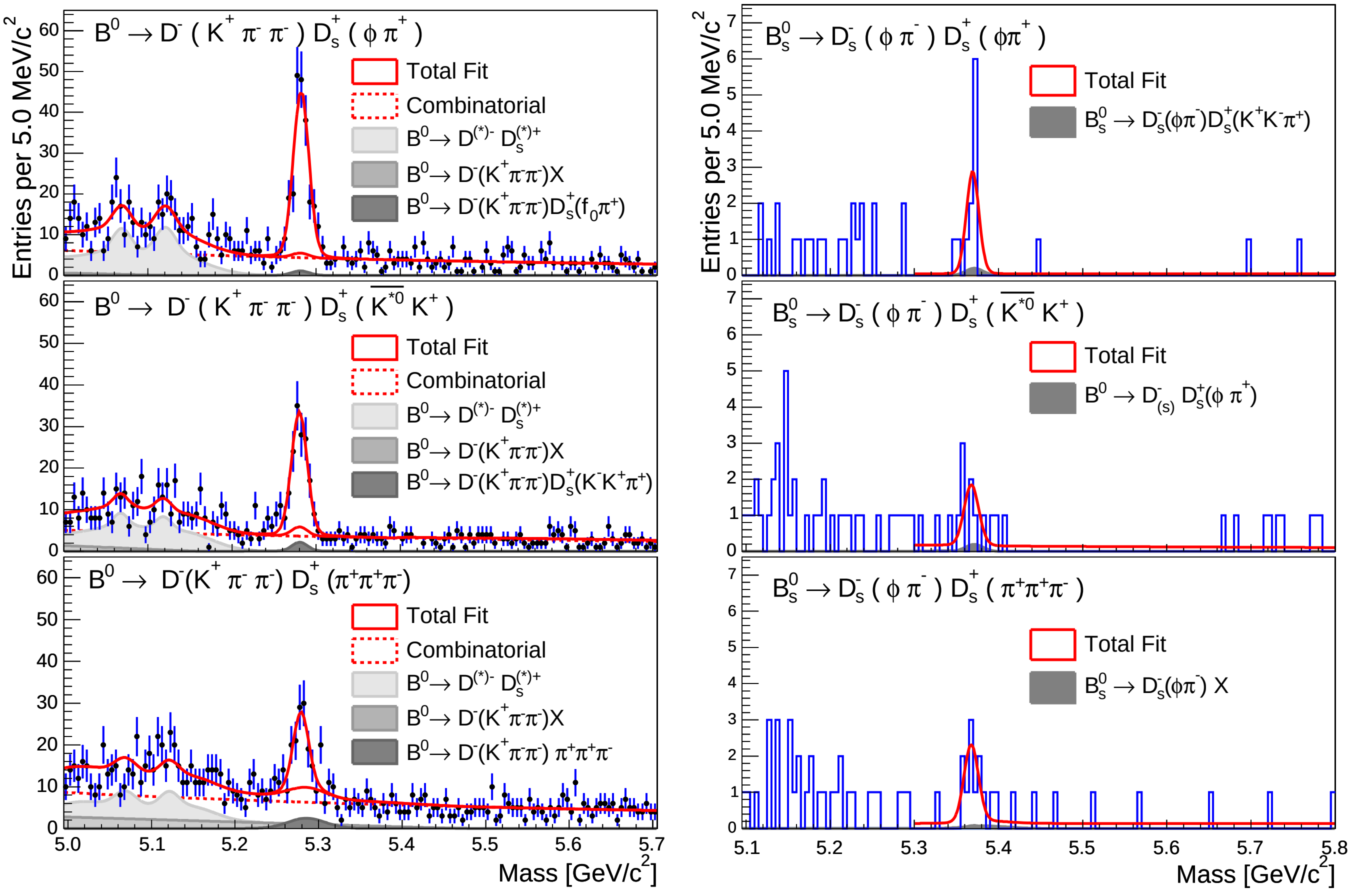

FIG. 1: Mass spectra for $B^0 \to D^- D_s^+$ (left) and $B_s^0 \to D_s^- D_s^+$ (right) where $D_s^+ \to \phi\pi^+$ (top), $D_s^+ \to \overline{K}^{*0} K^+$ (middle), $D_s^+ \to \pi^+\pi^+\pi^-$ (bottom). The decompositon of the background into combinatorial background and several backgrounds from B hadron decays is shown.

have fixed shapes derived from simulation and floating normalizations. Each signal template is parametrized by two Gaussians with different widths and a common mean. In fitting $B_s^0 \to D_s^- D_s^+$ distributions, we fix the signal masses to PDG values [14] and we fix the Gaussian signal widths, dominated by detector resolution, to the values obtained from Monte Carlo simulation.

We treat the background from the decay mode $B^0 \to D^-\pi^+\pi^+\pi^-$ to the signal for $B^0 \to D^-(K^+\pi^-\pi^-)D_s^+(\overline{K}^{*0}K^+)$ and $B^0 \to D^-(K^+\pi^-\pi^-)D_s^+(\pi^+\pi^+\pi^-)$ modes by introducing templates normalized to the B^0 yield. Similarly we introduce the template normalized to the B_s^0 yield to model $B_s^0 \to D_s^-(\phi\pi^-)\pi^+\pi^+\pi^-$ background under the signal for $B_s^0 \to D_s^-(\phi\pi^-)D_s^+(\pi^+\pi^+\pi^-)$ mode. These corrections lead to less than 2% change in the signal.

Monte Carlo studies show that B meson signal, reconstructed in a specific D_s^+ decay mode, have contributions from misreconstructed B meson candidates decaying through other D_s channels. The decay $B^0 \to D^- D_s^+$, followed by $D_s^+ \to f_0(980)(K^+K^-)\pi^+$, contributes to the reconstructed $B^0 \to D^-(K^+\pi^-\pi^-)D_s^+(\phi\pi^+)$. Similarly B meson decays followed by a non-resonant $D_s^+ \to K^+K^-\pi^+$ contribute to the B meson signal reconstructed with $D_s^+ \to \overline{K}^{*0}K^+$. The $D_s^+ \to K^+K^-\pi^+$ decay model takes into account the measured branching fractions of its resonant substructure [15]. The forementioned effects are taken into account by introducing correction templates with relative normalizations derived from Monte Carlo simulations and result in a 4% correction to the signal yield. Other b hadron backgrounds are described with templates derived from semi-generic simulations ($B \to D_{(s)}^- X$), where one of the $D_{(s)}^-$ mesons is forced to decay in the signal channel and the rest of the decay chain (X), follows the best available measurement of branching fractions.

Yields and ratios of reconstruction efficiencies extracted from signal simulations are summarized in Tab. I. Using Eq. 2 and the latest PDG [14] values $f_s/f_d = 0.259 \pm 0.038$, and $\mathcal{B}(D_s^- \to \phi\pi^-) \times \mathcal{B}(\phi \to K^+K^-) = (2.16 \pm 0.28) \times 10^{-2}$, we calculate the ratio of branching fractions $\mathcal{B}(B_s^0 \to D_s^- D_s^+)/\mathcal{B}(B^0 \to D^- D_s^+)$ for the three D_s^+ modes shown in Tab. I along with corresponding statistical uncertainties, systematic uncertainties dis-

	$\phi\pi^+$	$\overline{K}^{*0}K^+$	$\pi^+\pi^+\pi^-$
$\mathrm{N}(B^0 \to D^- D_s^+)$	183 ± 15	128 ± 13	84 ± 13
$\mathrm{N}(B_s^0 \to D_s^- D_s^+)$	$9.2^{+3.5}_{-2.9}$	$6.0^{+3.4}_{-2.7}$	$8.3^{+3.5}_{-2.8}$
$\epsilon(B_s^0 \to D_s^- D_s^+)/\epsilon(B^0 \to D^- D_s^+)$	0.88 ± 0.03	0.53 ± 0.02	0.63 ± 0.02
$\mathcal{B}(B_s^0 \to D_s^- D_s^+)/\mathcal{B}(B^0 \to D^- D_s^+)$	$0.98^{+0.38}_{-0.32}$	$1.51^{+0.87}_{-0.70}$	$2.67^{+1.20}_{-0.99}$
systematic uncertainty	$+0.06/-0.08$	$+0.15/-0.25$	$+0.27/-0.29$
(f_s/f_d) uncertainty	± 0.14	± 0.22	± 0.39
$\mathcal{B}(D_s^- \to \phi\pi^-)/\mathcal{B}(D^- \to K^+\pi^-\pi^-)$ uncertainty	± 0.13	± 0.20	± 0.36

TABLE I: Summary of event yields, efficiencies, measured ratios of branching fractions, and corresponding uncertainties for three D_s^+ decay modes.

cussed below, and the uncertainties from the measurements of f_s/f_d and $\mathcal{B}(D_s^- \to \phi\pi^-)/\mathcal{B}(D^- \to K^+\pi^-\pi^-)$.

The systematic uncertainties, summarized in Tab. II, are evaluated from the change in the ratio of the branching fractions $\mathcal{B}(B_s^0 \to D_s^- D_s^+)/\mathcal{B}(B^0 \to D^- D_s^+)$ for each effect under consideration. Fit systematic uncertainties are estimated by varying the fit window, binning, and template parameters. In the simulations, the known branching fractions are varied within their PDG [14] uncertainties. The effect of the B meson p_T spectrum is evaluated by comparing the efficiency determined from simulation based on next-to-leading order calculations [16] and on the measured B hadron spectrum [17]. The effect of meson lifetimes is studied by varying the world average B_s^0 and B^0 lifetimes within their uncertainties in the simulations. Trigger related systematic uncertainties are estimated from simulations. The effects due to the limited knowledge of the $D_s^+ \to \pi^+\pi^+\pi^-$ composition are studied by varying the relative branching fractions of the components of the decay within their PDG [14] uncertainties. Finally, using the combinatorial background from data to optimize the selection may tune on a fluctuation, because only few events pass the selection. This effect has been estimated using simulation based on the expected combinatorial background distribution.

	$\phi\pi^+$	$\overline{K}^{*0}K^+$	$\pi^+\pi^+\pi^-$
$B^0 \to D^- D_s^+$ fit	± 2.3	± 4.2	± 8.4
$B_s^0 \to D_s^- D_s^+$ fit	± 4.4	± 8.2	± 4.6
B meson p_T spectrum	± 3.0	± 3.0	± 3.0
B_s^0 lifetime	± 2.0	± 2.0	± 2.0
trigger	± 1.0	± 1.0	± 1.0
$D_s^+ \to \pi^+\pi^+\pi^-$ composition	-	-	± 3.0
optimization bias	-5.0	-13.0	-4.0
Total	$^{+6.2}_{-8.0}$	$^{+9.9}_{-16.4}$	$^{+10.3}_{-11.0}$

TABLE II: Systematic uncertainties in [%] for the B_s^0/B^0 relative branching fraction measurements for three D_s^+ decay modes.

The significance of the observation of the $B_s^0 \to D_s^- D_s^+$ decay is obtained from the ratio of the likelihood found in the data fit to the likelihood achieved by fitting the same data mass distribution with a background only model. The individual significances of the signal reconstructed with $B_s^0 \to D_s^-(\phi\pi^-)D_s^+(\phi\pi^+)$, $B_s^0 \to D_s^-(\phi\pi^-)D_s^+(\overline{K}^{*0}K^+)$, and $B_s^0 \to D_s^-(\phi\pi^-)D_s^+(\pi^+\pi^+\pi^-)$ decay modes are greater than 5.8σ, 3.4σ, and 4.4σ respectively. From the product of three likelihoods we find the combined result consistent with an observation of $B_s^0 \to D_s^- D_s^+$ at a 7.5σ significance.

When combining the three results, the fit systematic uncertainties are weighed by the measured yields. The rest of the systematic uncertainties, except for the $D_s^+ \to \pi^+\pi^+\pi^-$ composition uncertainty, are considered common for all three modes. We find

$$\frac{\mathcal{B}(B_s^0 \to D_s^- D_s^+)}{\mathcal{B}(B^0 \to D^- D_s^+)} = 1.44^{+0.38}_{-0.31}(\mathrm{stat})^{+0.08}_{-0.12}(\mathrm{syst}) \qquad (3)$$

$$\pm 0.21(f_s/f_d) \pm 0.20(\frac{\mathcal{B}(D^- \to K^+\pi^-\pi^-)}{\mathcal{B}(D_s^- \to \phi\pi^-)}).$$

Using $\mathcal{B}(B^0 \to D^- D_s^+) = (6.5 \pm 2.1) \times 10^{-3}$ [14] and Eq. 2 we determine

$$\mathcal{B}(B_s^0 \to D_s^- D_s^+) = (9.4^{+4.4}_{-4.2}) \times 10^{-3} \qquad (4)$$

from which we derive a lower limit on $\Delta\Gamma_s^{CP}/\Gamma$

$$\begin{aligned} \frac{\Delta\Gamma_s^{CP}}{\Gamma_s} &= 2\mathcal{B}(B_s^0 \to D_s^{(*)-} D_s^{(*)+}) \geq \qquad (5) \\ &\geq 2\mathcal{B}(B_s^0 \to D_s^- D_s^+) \geq 1.2 \times 10^{-2} \text{ at 95\% C.L.} \end{aligned}$$

In the derivation of the lower limit we take into account the Poisson statistical fluctuations of the signal yields and the Gaussian distribution for systematics uncertainties.

We have presented the first observation of the decay $B_s^0 \to D_s^- D_s^+$ and have measured its branching fraction with respect to $B^0 \to D^- D_s^+$. We set a lower bound on $\Delta\Gamma_s^{CP}/\Gamma_s$ which at the 95non-zero decay rate difference and agrees with theoretical predictions [2] and other experimental data [14].

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