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Function Uncertainties**

Walter T. Giele and Stephane Keller

*Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510*

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Implications of Hadron Collider Observables on Parton Distribution Function Uncertainties

WALTER T. GIELE AND STEPHANE KELLER

*Fermilab, MS 106
Batavia, IL 60510, USA*

Abstract

Standard parton distribution function sets do not have rigorously quantified uncertainties. In recent years it has become apparent that these uncertainties play an important role in the interpretation of hadron collider data. In this paper, using the framework of statistical inference, we illustrate a technique that can be used to efficiently propagate the uncertainties to new observables, assess the compatibility of new data with an initial fit, and, in case the compatibility is good, include the new data in the fit.

1 Introduction

Current standard sets of Parton Distribution Function (PDF) do not include uncertainties [1]. In practice, as long as the PDF's are used to calculate observables that themselves have large experimental uncertainties this shortcoming is obviously not a problem. In the past the precision of the hadron collider data was such that there was no ostensible need for the PDF uncertainties, as was testified by the good agreement between the theory and measurements. However, the need for PDF uncertainties became apparent with the measurement of the one jet inclusive transverse energy at the Tevatron [2]. At large transverse jet energies the data was significantly above the theoretical prediction, a possible signal for new physics. The deviation was ultimately "fixed" by changing the PDF's in such a manner that they still were consistent with the observables used to determine the PDF [3]. This is a reflection of the significant PDF uncertainties for this observable. Knowing the uncertainties on the PDF's would have cleared the situation immediately. Note that once the data is used in the PDF fit, it can not be used for other purposes. Specifically, setting limits on possible physics beyond the Standard Model. In that case, one should fit the PDF's and the new physics simultaneously. The technique presented in this paper is well suited for this sort of problem.

The spread between different sets of PDF's is often associated with PDF uncertainties. Currently, this is what is used for the determination of the PDF uncertainty on the W -boson mass at the Tevatron. It is not possible to argue that this spread is an accurate representation of all experimental and theoretical PDF uncertainties. For the next planned high luminosity run at Fermilab, assuming an integrated luminosity of $2 fb^{-1}$, the expected 40 MeV uncertainty on the W -boson mass is dominated by a 30 MeV production model uncertainty. The latter uncertainty itself is dominated by the PDF uncertainty, estimated to be 25 MeV [4]. This determination of the PDF uncertainty is currently nothing more than an educated guess. It is made by ruling out existing PDF's using the lepton charge asymmetry in W -boson decay events. The spread of the remaining PDF's determines the uncertainty on the extracted W -boson mass. Because the PDF uncertainty seems to be the dominant source of uncertainty in the determination of the W -boson mass, such a procedure must be replaced by a more rigorous quantitative approach. The method described in this paper is well suited for this purpose.

In this paper, using the framework of statistical inference [5, 6], we illustrate a method that can be used for many purposes. First of all, it is easy to propagate the PDF uncertainties to a new observable without the need to calculate the derivative of the observable with respect to the different PDF parameters. Secondly, it is straightforward to assess the compatibility of new data with the current fit and determine whether the new data should be included in the fit. Finally, the new data can be included in the fit without redoing the whole fit.

This method is significantly different from more traditional approaches to fit the PDF's to the data. It is very flexible and beside solving the problems already mentioned, it offers additional advantages. First, the experimental uncertainties and the probability density distributions for the fitted parameters do not have to be Gaussian distributed. However, such a generalization would require a significant increase in computer resources. Second, once a fit has been made to all the data sets, a specific data set can be easily excluded from the fit. Such an option is important in order to be able to investigate the effect of the different data sets. This is particularly useful in the case of incompatible new data. In that case one can easily investigate the origin of the incompatibility. Finally, because it is not necessary to redo a global fit in order to include a new data set, experimenters can include their own new data into the PDF's during the analysis phase.

The outline for the rest of the paper is as follows. In Sec. 2, we describe the inference method. The flexibility and simplicity of the method is illustrated in Sec. 3, by applying it to the CDF one jet inclusive transverse jet energy distribution [2] and the CDF lepton charge asymmetry data [7]. In Sec. 4 we draw our conclusions and outline future improvements and extensions to our method.

2 The Method of Inference

Statistical inference requires an initial probability density distribution for the PDF parameters. This initial distribution can be rather arbitrary, in particular it can be solely based on theoretical considerations. Once enough experimental data are used to constrain the probability density distribution of

the parameters the initial choices become irrelevant ¹. Obviously, the initial choice does play a role at intermediate stages. The initial distribution can also be the result of a former fit to other data. The data that we will use later in this paper do not constrain the PDF's enough by themselves to consider using an initial distribution based only on theory. The final answer would depend too much on our initial guess. We therefore decided to use the results of Ref. [8]. In this work the probability density distribution was assumed to be Gaussian distributed and was constrained using Deep Inelastic Scattering (DIS) data. All the experimental uncertainties, including correlations, were included in the fit, but no theoretical uncertainties were considered. The fact that no Tevatron data were used allows us to illustrate the method with Tevatron data ². We briefly summarize Ref. [8] in the appendix.

In Sec. 2.1 we explain the propagation of the uncertainty to new observables. Sec. 2.2 shows how the compatibility of new data with the PDF can be estimated. Finally, in Sec. 2.3 we demonstrate how the effect of new data can be included in the PDF's by updating the probability density distribution of the PDF parameters.

2.1 Propagation of the uncertainty

We now assume that the PDF's are parametrized at an initial factorization scale Q_0 , with N_{par} parameters, $\{\lambda\} \equiv \lambda_1, \lambda_2, \dots, \lambda_{N_{par}}$ and that the probability density distribution is given by $P_{init}(\lambda)$. Note that $P_{init}(\lambda)$ does not have to be a Gaussian distribution.

By definition $P_{init}(\lambda)$ is normalized to unity,

$$\int_V P_{init}(\lambda) d\lambda = 1 , \quad (1)$$

where the integration is performed over the full multi-dimensional parameter space and $d\lambda \equiv \prod_{i=1}^{N_{par}} d\lambda_i$. To calculate the parameter space integrals we use a Monte-Carlo (MC) integration approach with importance sampling. We generate N_{pdf} random sets of parameters $\{\lambda\}$ distributed according to $P_{init}(\lambda)$. This choice should minimize the MC uncertainty for most of the

¹The standard PDF sets of Ref. [1] basically assume that the initial probability density distribution for the parameters is uniform.

²Recent PDF sets have also included the Tevatron data that we will use, but none of these sets included uncertainties.

integrals we are interested in. For reference we also generate one set at the central values of the $\{\lambda\}$, the $\mu_{\{\lambda\}}$. The number of parameter sets to be used depends on the quality of the data. The smaller the experimental uncertainty is compared to the PDF uncertainty, the more PDF's we need. We must ensure a sufficient fraction of PDF's span the region of interest (i.e. close to the data). For the purposes of this paper, we found that $N_{pdf} = 100$ is adequate. Clearly, to each of the N_{pdf} sets of parameters $\{\lambda\}$ correspond a set of unique PDF's. Each of these PDF sets have to be evolved using the Altarelli-Parisi evolution equations. We used the CTEQ package to do this evolution [9].

We now can evaluate any integral I over the parameter space as a finite sum [6]

$$\begin{aligned}
 I &= \int_V f(\lambda) P_{init}(\lambda) d\lambda \\
 &\approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} f(\lambda^j) \\
 &\equiv \langle f \rangle,
 \end{aligned}
 \tag{2}$$

with λ^j is the j -th random set of $\{\lambda\}$. The function f represents an integrable function of the PDF parameters. The uncertainty on the integral I due to the MC integration is given by,

$$\delta I = \sqrt{\frac{\langle f^2 \rangle - \langle f \rangle^2}{N_{pdf}}}.
 \tag{3}$$

For any quantity, $x(\lambda)$, that depends on the PDF parameters $\{\lambda\}$ (for example an observable, one of the flavor PDF's or for that matter one of the parameter itself), the theory prediction is given by its average value, μ_x , and its uncertainty, σ_x ³:

$$\begin{aligned}
 \mu_x &= \int_V x(\lambda) P_{init}(\lambda) d\lambda \approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} x(\lambda^j) \\
 \sigma_x^2 &= \int_V (x(\lambda) - \mu_x)^2 P_{init}(\lambda) d\lambda \approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} (x(\lambda^j) - \mu_x)^2.
 \end{aligned}
 \tag{4}$$

³If the uncertainty distribution is not Gaussian the average and the standard deviation might not properly quantify the distribution.

Note that μ_x is not necessarily equal to the value of $x(\lambda)$ evaluated at the central value of the $\{\lambda\}$. However, this is how observables are evaluated if one has only access to PDF's without uncertainties.

Given $y(\lambda)$, another quantity calculable from the $\{\lambda\}$, the covariance of $x(\lambda)$ and $y(\lambda)$ is given by the usual expression:

$$\begin{aligned} C_{xy} &= \int (x(\lambda) - \mu_x)(y(\lambda) - \mu_y) P_{init}(\lambda) d\lambda \\ &\approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} (x(\lambda^j) - \mu_x)(y(\lambda^j) - \mu_y) . \end{aligned} \quad (5)$$

The correlation between $x(\lambda)$ and $y(\lambda)$ is given by $cor_{xy} = C_{xy}/(\sigma_x\sigma_y)$. For example, this can be used to calculate the correlation between two experimental observables, between an observable and one of the PDF parameters, or between an observable and a specific flavor PDF at a fixed Bjorken- x .

Using Eq. 3, the MC uncertainty on the average and (co)variance is given by

$$\begin{aligned} \delta\mu_x &= \frac{\sigma_x}{\sqrt{N_{pdf}}} \\ \delta\sigma_x^2 &= \sigma_x^2 \sqrt{\frac{2}{N_{pdf}}} \\ \delta C_{xy} &= C_{xy} \sqrt{\frac{2}{N_{pdf}}} . \end{aligned} \quad (6)$$

The MC technique presented in this sub-section, gives a simple way to propagate uncertainties to a new observable, without the need for calculating the derivatives of the observable with respect to the parameters.

2.2 Compatibility of New Data

We will assume that one or several new experiments, not used in the determination of the initial probability density distribution, have measured a set of N_{obs} observables $\{x^e\} = x_1^e, x_2^e, \dots, x_{N_{obs}}^e$. The experimental uncertainties, including the systematic uncertainties, are summarized by the $N_{obs} \times N_{obs}$ experimental covariance matrix C^{exp} . Note that the correlations between experiments are easily incorporated. Here however, we have to assume that the

new experiments are not correlated with any of the experiments used in the determination of P_{init} . The probability density distribution of $\{x^e\}$ is given by

$$P(x^e) = \int_V P(x^e|\lambda)P_{init}(\lambda)d\lambda \quad (7)$$

$$\approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} P(x^e|\lambda^j) ,$$

where $P(x^e|\lambda)$ is the conditional probability density distribution (often referred to as likelihood function). This distribution quantifies the probability of measuring the specific set of experimental values $\{x^e\}$ given the set of PDF parameters $\{\lambda\}$. In PDF sets without uncertainties, $P_{init}(\lambda)$ is a delta function and $P(x^e) = P(x^e|\lambda)$.

Instead of dealing with the probability density distribution of Eq. 7, one often quotes the confidence level to determine the agreement between the data and the model. The confidence level is defined as the probability that a repeat of the given experiment(s) would observe a worse agreement with the model. The confidence level of $\{x^e\}$ is given by

$$CL(x^e) = \int_V CL(x^e|\lambda)P_{init}(\lambda)d\lambda \quad (8)$$

$$\approx \frac{1}{N_{pdf}} \sum_{j=1}^{N_{pdf}} CL(x^e|\lambda) ,$$

where $CL(x^e|\lambda)$ is the confidence level of $\{x^e\}$ given $\{\lambda\}$. If $CL(x^e)$ is larger than an agreed value, the data are considered consistent with the PDF and can be included in the fit. If it is smaller, the data are inconsistent and we have to determine the source of discrepancy.

For non-Gaussian uncertainties the calculation of the confidence level might be ambiguous. In this paper we assume that the uncertainties are Gaussian. The conditional probability density distribution and confidence level are then given by

$$P(x^e|\lambda) = P(\chi_{new}^2) = \frac{e^{-\frac{1}{2}\chi_{new}^2(\lambda)}}{\sqrt{(2\pi)^{N_{obs}}|C^{tot}|}} \quad (9)$$

$$CL(x^e|\lambda) = CL(\chi_{new}^2) = \int_{\chi_{new}^2}^{\infty} P(\chi^2) d\chi^2 , \quad (10)$$

where

$$\chi_{new}^2(\lambda) = \sum_{k,l}^{N_{obs}} \left(x_k^e - x_k^t(\lambda) \right) M_{kl}^{tot} \left(x_l^e - x_l^t(\lambda) \right) , \quad (11)$$

is the chi-squared of the new data. The theory prediction for the k -th experimental observable, $x_k^t(\lambda)$, is calculated using the PDF set given by the parameters $\{\lambda\}$. The matrix M^{tot} is the inverse of the total covariance matrix, C^{tot} , which in turn is given by the sum of the experimental, C^{exp} , and theoretical, C^{theor} , covariance matrix. We assume that there is no correlation between the experimental and theoretical uncertainties. We will use a minimal value of 0.27% on the confidence level, corresponding to a three sigma deviation, as a measure of compatibility of the data with the theory. If the new data are consistent with the theory prediction then the maximum of the distribution of the χ_{new}^2 should be close to N_{obs} (within the expected $\sqrt{2N_{obs}}$ uncertainty). The standard deviation of χ_{new}^2 , $\sigma_{\chi_{new}^2}$, tells us something about the relative size of the PDF uncertainty compared to the size of the data uncertainty. The larger the value of $\sigma_{\chi_{new}^2}$ is compared to $\sqrt{2N_{obs}}$, the more the data will be useful in constraining the PDF's.

Note that if there are several uncorrelated experiments, the total χ_{new}^2 is equal to the sum of the χ_{new}^2 of the individual experiments and the conditional probability is equal to the product of the individual conditional probabilities.

2.3 Effect of new data on the PDF's

Once we have decided that the new data are compatible with the initial PDF's, we can constrain the PDF's further. We do this within the formalism of statistical inference, using Bayes theorem. The idea is to update the probability density distribution taking into account the new data. This new probability density distribution is in fact the conditional probability density distribution for the $\{\lambda\}$ considering the new data $\{x^e\}$ and is given directly by Bayes theorem

$$P_{new}(\lambda) = P(\lambda|x^e) = \frac{P(x^e|\lambda) P_{init}(\lambda)}{P(x^e)} , \quad (12)$$

where $P(x^e)$, defined in Eq. 7, acts as a normalization factor such that $P(\lambda|x^e)$ is normalized to one. Because $P_{new}(\lambda)$ is normalized to unity, we

can replace $P(x^e|\lambda)$ in Eq. 12 simply by $e^{-\frac{\chi_{new}^2(\lambda)}{2}}$. This factor acts as a new weight on each of the PDF's.

We can now replace $P_{init}(\lambda)$ by $P_{new}(\lambda)$ in the expression for the average, standard deviation and covariance given in Sec. 2.1 and obtain predictions that include the effect of the new data. With the MC integration technique described before, these quantities can be estimated by weighted sums over the N_{pdf} PDF sets

$$\begin{aligned}\mu_x &\approx \sum_{k=1}^{N_{pdf}} w_k x(\lambda^{(k)}) \\ \sigma_x^2 &\approx \sum_{k=1}^{N_{pdf}} w_k (x(\lambda^{(k)}) - \mu_x)^2 \\ C_{xy} &\approx \sum_{k=1}^{N_{pdf}} w_k (x(\lambda^{(k)}) - \mu_x) (y(\lambda^{(k)}) - \mu_y) ,\end{aligned}\tag{13}$$

where the weights are given by

$$w_k = \frac{e^{-\frac{1}{2}\chi_{new}^2(\lambda^k)}}{\sum_{l=1}^{N_{pdf}} e^{-\frac{1}{2}\chi_{new}^2(\lambda^l)}} .\tag{14}$$

Note that for the calculation of the Monte-Carlo uncertainty of the weighted sums, the correlation between the numerator and denominator in Eq. 13 has to be taken into account properly.

Our strategy is very flexible. Once the theory predictions $x_i^t(\lambda)$ using the N_{pdf} PDF sets are known for each of the experiments, it is trivial to include or exclude the effect of one of the experiments on the probability density distribution. If the different experiments are uncorrelated then all what is needed is the χ_{new}^2 of each individual experiments for all the PDF sets. In that case, each experiment is compressed into N_{pdf} χ_{new}^2 values.

One other advantage is that all the needed $x_i^t(\lambda)$ can be calculated beforehand in a systematic manner, whereas standard chi-squared or maximum likelihood fits require many evaluations of $x_k^t(\lambda)$ during the fit as the parameters are changed in order to find the extremum. These methods are not very flexible, as a new fit is required each time an experiment is added or removed.

The new probability density distribution of the PDF parameters is Gaussian if the following three conditions are met. First, the initial probability density distribution, $P_{init}(\lambda)$, must be Gaussian. Second, all the uncertainties on the data points must be Gaussian distributed (that includes systematic and theoretical uncertainties). Finally, the theory predictions, $x_i^t(\lambda)$, must be linear in $\{\lambda\}$ in the region of interest. This last requirement is fulfilled once the PDF uncertainties are small enough. For the studies in this paper all three requirements are fulfilled. The new probability density distribution can therefore be characterized by the average value of the parameters and their covariance matrix, which can be calculated, together with their MC integration uncertainty, using Eq. 13. Once the new values of the average and the covariance matrix have been calculated, a new set of PDF parameters can be generated according to the new distribution and used to make further prediction instead of using the initial set of PDF with the weights.

An alternative way to generate a PDF set distributed according to $P_{new}(\lambda)$ is to unweight the now weighted initial PDF set. The simplest way to unweight the PDF sets is to use a rejection algorithm. That is, define w_{max} as the largest of the N_{pdf} weights given in Eq. 14. Next generate for each PDF set a uniform stochastic number, r_k , between zero and one. If the weight w_k is larger or equal to $r_k \times w_{max}$ we keep PDF set k , otherwise it is discarded. The surviving PDF's are now distributed according to $P_{new}(\lambda)$. The number of surviving PDF's is on average given by $N_{pdf}^{new} = 1/w_{max}$. We can now apply all the techniques of the previous sub-sections, using the new unweighted PDF set. The MC integration uncertainties are easily estimated using the expected number of surviving PDF's. In the extreme case that w_{max} is close to one and only a few PDF survive the unweighting procedure, the number of initial PDF's must be increased. The other extreme occurs when all the weights are approximately equal, i.e. $w_k \sim 1/N_{pdf}$. In that case the new data puts hardly any additional constraints on the PDF.

The χ_{new}^2 is only used to calculate the weight of a particular PDF, so that the new probability density distribution of the PDF parameters can be determined. We do not perform a chi-squared fit. However, if the new probability density distribution of the parameters is Gaussian distributed then our method is equivalent to a chi-squared fit. In that case the average value of the parameters correspond to the maximum of the probability density distribution. The minimum chi-squared can be estimated (with MC uncertainties) from the average χ_{new}^2 calculated with the new probability density

distribution. Indeed, by definition this average must be equal to the minimum chi-squared, χ_{min}^2 , plus the known number of parameter. Note that the variance of the χ_{new}^2 must itself be equal to twice the number of parameters. To obtain the overall minimum chi-squared, the value of the minimum chi-squared of the initial fit must be added to χ_{min}^2 . As long as the confidence level of the new data that were included in the fit is sufficiently high, the overall minimum chi-squared obtained is guaranteed to be in accordance with expectations ⁴.

3 Expanding the PDF sets

The viability of the method described in Sec. 2 is studied using two CDF measurements. In Sec. 3.1 the one jet inclusive transverse energy distribution is considered, while the lepton charge asymmetry in W -boson decay is examined in Sec. 3.2. The statistical, systematic, and theoretical uncertainties on the observables will be taken into account.

3.1 The one jet inclusive measurement

The CDF results on the one jet inclusive transverse energy distribution [2] demonstrated the weakness of the current standard PDF sets due to the absence of uncertainties on the PDF parameters.

The observables are the inclusive jet cross section at different transverse energies ⁵, E_T^i

$$x_i = \frac{d\sigma}{dE_T}(E_T^i). \quad (15)$$

We first have to construct the experimental covariance matrix, C_{ij}^{exp} , using the information contained in Ref. [2]. The paper lists the statistical uncertainty at the different experimental points, $\Delta_0(E_T^i)$, together with eight independent sources of systematic uncertainties, $\Delta_k(E_T^i)$. Hence, the experimental

⁴We are assuming that the initial χ_{min}^2 was within expectations.

⁵To be more precise, the inclusive jet cross section in different bins of transverse energy. In the numerical results presented here we take the finite binning effects into account.

measurements, x_i^e are given by

$$x_i^e = x_i^t(\lambda) + \eta_0^i \Delta_0(E_T^i) + \sum_{k=1}^8 \eta_k \Delta_k(E_T^i), \quad (16)$$

where as before, $x_i^t(\lambda)$ is the theoretical prediction for the observable calculated with the set of parameters $\{\lambda\}$. The η_0^i and η_k are independent random variables normally distributed with zero average and unit standard deviation. Note that some of the systematic uncertainties given in Ref. [2] are asymmetric. In those cases we symmetrized the uncertainty using the average deviation from zero. From Eq. 16 we can construct the experimental covariance matrix

$$C_{ij}^{exp} = \left(\Delta_0(E_T^i)\right)^2 \delta_{ij} + \sum_{k=1}^8 \Delta_k(E_T^i) \Delta_k(E_T^j). \quad (17)$$

We also need to estimate the theoretical uncertainty. In Eq. 16 no theoretical uncertainties were taken into account. We consider two types of uncertainties: the uncertainty due to the numerical Monte Carlo integration over the final state particle phase space, $\Delta_{MC}(E_T^i)$, and the renormalization/factorization scale, μ , uncertainty, $\Delta_\mu(E_T^i)$. The theoretical prediction in Eq. 16 must then be replaced by

$$x_i^t(\lambda) \rightarrow \frac{d\sigma^{NLO}}{dE_T}(E_T^i, \lambda, \mu) + \eta_{MC}^i \Delta_{MC}(E_T^i) + \eta_\mu \Delta_\mu(E_T^i), \quad (18)$$

from which we can derive the theoretical covariance matrix

$$C_{ij}^{theor} = \left(\Delta_{MC}(E_T^i)\right)^2 \delta_{ij} + \Delta_\mu(E_T^i) \Delta_\mu(E_T^j). \quad (19)$$

Here we assume that there is no bin to bin correlation in the MC uncertainty. On the other hand, we take the correlation of the scale uncertainty fully into account. Both Δ_{MC} and Δ_μ are evaluated at the central values of the PDF parameters, assuming that the variation is small.

We evaluate the scale uncertainty in a very straightforward manner. As the central prediction the renormalization and factorization scale are taken to be equal to half the transverse energy of the leading jet in the event, $\mu = \frac{1}{2} E_T^{max}$. To estimate the uncertainty we make another theoretical prediction

now choosing as a scale $\mu = E_T^{max}$. The “one-sigma” uncertainty is defined as

$$\Delta_\mu(E_T) = \frac{d\sigma^{NLO}}{dE_T}(E_T, \mu_\lambda, \mu = \frac{1}{2}E_T^{max}) - \frac{d\sigma^{NLO}}{dE_T}(E_T, \mu_\lambda, \mu = E_T^{max}) . \quad (20)$$

As we will see later in this section the theoretical uncertainties are small compared to the other uncertainties. Therefore this crude estimate suffices for the purposes of this paper. In the future a more detailed study of the theoretical uncertainty is required. The scale uncertainty is often associated with the theoretical uncertainty due to the truncation of the perturbative series. However, it is important to realize this is only a part of the full theoretical uncertainty.

In Fig. 1^a we present results for the single inclusive jet cross section as a function of the jet transverse energy. Both data and theoretical predictions are divided by the average prediction of the initial PDF’s. The NLO predictions are calculated using the JETRAD prediction [10]. The inner (outer) error bar on the experimental points represent the diagonal part of the experimental (total) covariance matrix. The dotted lines represent the initial one-sigma PDF uncertainties. The solid lines are the theory predictions calculated with the new PDF’s (i.e., the new probability density distribution). The plot is somewhat misleading because of the large point-to-point correlation of the uncertainties. The confidence level of 50% is very high, indicating a good agreement between the prediction and the data.

This leads us to the conclusion that the one jet inclusive transverse energy distribution is statistically in agreement with the NLO theoretical expectation based on the initial probability density distribution of the PDF parameters. No indication of new physics is present. Note that the prediction using the initial PDF differs quite a bit from the more traditional fits such as MRSD0, see the dashed line in Fig. 1^a. Having no uncertainties on the traditional fits it is hard to draw any quantitative conclusion from this observation. The larger value of the jet cross section calculated using the initial PDF set at high transverse energies compared to MRSD0 was anticipated in Ref. [8] and can probably be traced back to the larger d and u quark distribution at the reference scale Q_0 and moderate $x \sim 0.2$. This difference in turn was partially attributed to the different way of treating target mass and Fermi motion corrections.

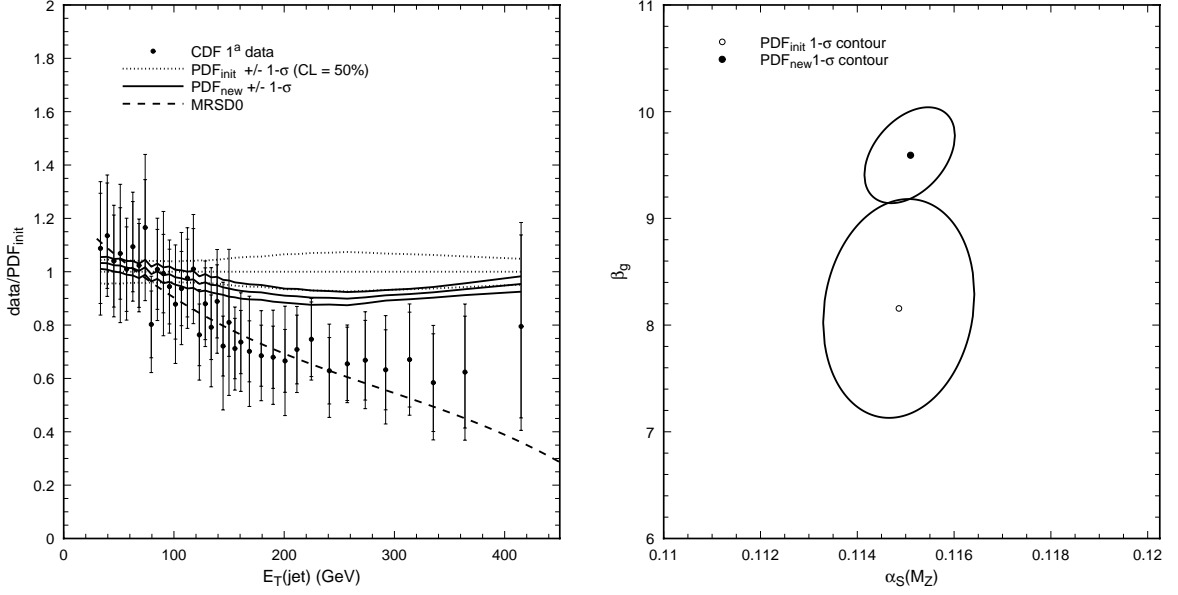


Figure 1: (a) Single inclusive jet cross section as a function of the jet transverse energy. The results are divided by the average prediction calculated with the initial PDF's. The data points are the CDF run 1^a results. The dotted lines represent the initial one-sigma PDF uncertainties. The solid lines are the theory predictions calculated with the new PDF's. The inner (outer) error bars on the data points are the diagonal entries of the experimental (total) covariance matrix. The dashed line is the prediction obtained with the MRSD0 PDF set. (b) The one-sigma correlation contour between the strong coupling constant $\alpha_S(M_Z)$ and the β -parameter in the gluon PDF ($\simeq x^\alpha(1-x)^\beta$ at the initial factorization scale) calculated for both the initial and new PDF's.

Given the confidence level of 50% the one jet inclusive data can be included in the fit. Using Eq. 11 we calculate for each PDF set k the corresponding $\chi_{new}^2(\lambda^k)$. This gives us the 100 weights w_k (conditional probabilities) defined in Eq. 14. Using Eq. 13, we can calculate the effects of including the CDF data into the fit. The results are shown in Figs. 1^a and 1^b. As can be seen in Fig. 1^a the effect is that the central value is pulled closer to the data and the PDF uncertainty is reduced substantially. Two of the fourteen PDF parameters are affected the most. As expected these are the strong coupling constant $\alpha_s(M_Z)$ and the gluon PDF coefficient β , which controls the high x behavior (the gluon PDF is proportional to $x^\alpha(1-x)^\beta$ at the initial scale). In Fig. 1^b we show the correlation between these two parameters before and after the inclusion of the CDF data. As can be seen the impact on β is very significant. Similarly, the uncertainty on α_s is reduced substantially and the correlation between the two parameters is also changed. This indicates that the one jet inclusive transverse energy distribution in itself has a major impact on the uncertainty of α_s and the determination of the gluon PDF. Note that we do not address the issue of the parametrization uncertainty. Other choices of how to parameterize the initial PDF's will change the results. To obtain a value and uncertainty of $\alpha_s(M_Z)$ which is on the same footing as the one obtained from e^+e^- -colliders, one needs to address this issue.

3.2 The lepton charge asymmetry measurement

Our second example is the lepton charge asymmetry in W -boson decay at the Tevatron. As already explained, this observable is important for the reduction of the PDF uncertainties in the W -boson mass extraction at hadron colliders. The asymmetry is given by

$$A(\eta_e) = \frac{(N^+(\eta_e) - N^-(\eta_e))}{(N^+(\eta_e) + N^-(\eta_e))}, \quad (21)$$

where N^+ and N^- are respectively the number of positrons and electrons at the pseudo-rapidity η_e .

In Fig. 2^a, we show the preliminary CDF data of run 1^b (solid points) for the asymmetry, along with the NLO predictions (dotted lines) including the PDF uncertainties, relative to the theory average prediction using the initial PDF's. For the NLO calculations the DYRAD prediction [10] was used. The inner error bars on the experimental points are the statistical uncertainties;

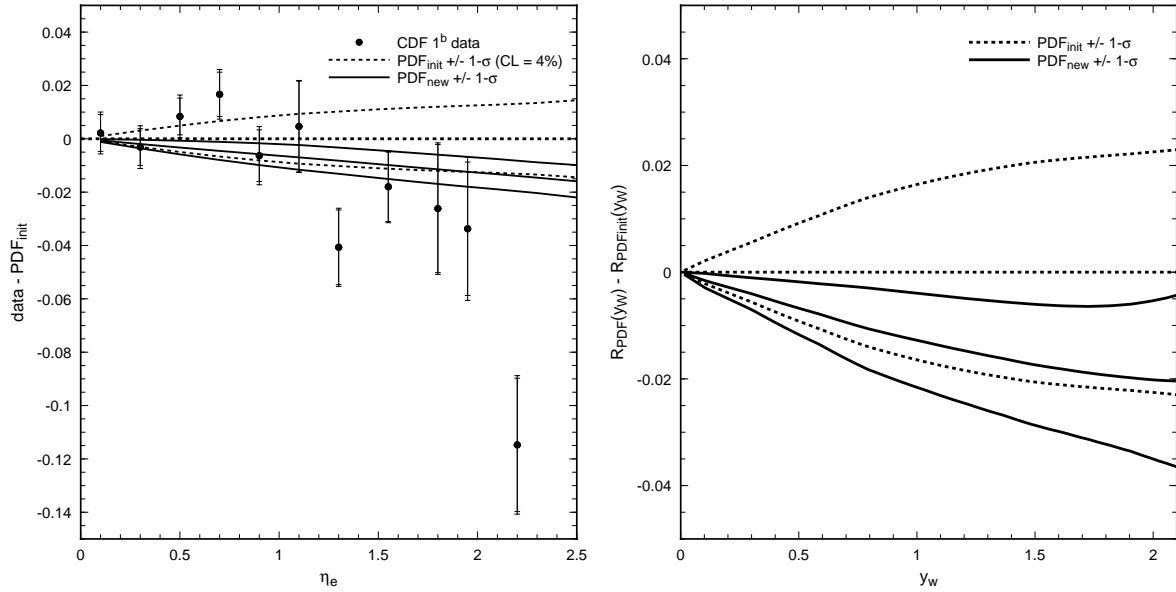


Figure 2: (a) The lepton charge asymmetry as a function of the lepton pseudo-rapidity. The results are normalized to the theory prediction using the average value of the initial PDF's. The data are the CDF run 1^b preliminary results. The error bars, dotted and solid lines have the same definition as in Fig. 1. (b) The ratio $R(y_W)$ normalized as in (a) as a function of the W -boson rapidity. The dotted and solid lines are defined as in Fig. 1.

the systematic uncertainties are small and we can safely neglect them. The outer error bars are the diagonal of the total covariance matrix. In this case, the theoretical uncertainty is dominated by the phase space Monte Carlo integration uncertainty; we took its bin to bin correlation into account. Similar to the one jet inclusive transverse energy case, the scale uncertainty is defined by the difference between the theoretical prediction calculated using two scales, $\mu = M_W$ and $\mu = 2 \times M_W$.

As is clear from Fig. 2^a, there is a good agreement between the data and the NLO prediction, except for the last experimental point at the highest pseudo-rapidity. The confidence level including the last point is well below our threshold of 0.27%. In order to be able to include the data into the PDF fit we decided to simply exclude this data point from our analysis. Without the highest pseudo-rapidity point we obtain a reasonable confidence level of 4%. It is not as good as in the single inclusive jet case even though the plots appear to indicate otherwise. The reason for this is the absence of significant point-to-point correlation for the charge asymmetry uncertainties.

We can now include the lepton charge asymmetry data into the fit by updating the probability density distribution with Bayes theorem, as described in the previous section. In Fig. 2^a the prediction obtained with the new probability density distribution are shown by the solid lines. As expected, the data are pulling the theory down and reducing the PDF uncertainties.

It is difficult to correlate the change in the asymmetry to a change in a particular PDF parameter. On the other hand, it is well known that the lepton asymmetry can be approximately related to the following asymmetry of the ratio of up quark (u) and down quark (d) distribution function

$$R(y_W) = \frac{\frac{u(x_1)}{d(x_1)} - \frac{u(x_2)}{d(x_2)}}{\frac{u(x_1)}{d(x_1)} + \frac{u(x_2)}{d(x_2)}}. \quad (22)$$

The Bjorken- x are given by

$$x_{1,2} = \frac{M_W}{\sqrt{s}} e^{\pm y_W}. \quad (23)$$

where M_W is the mass of the W -boson, \sqrt{s} the center of mass of the collider, and y_W the W -boson rapidity. The PDF's were evaluated with the factorization scale equal to M_W . The ratio $R(y_W)$ is approximately the W -boson asymmetry and obviously is sensitive to the slope of the u/d ratio.

In Fig. 2^b we show the ratio $R(y_W)$ calculated with both the initial and the new probability density distributions. As can be seen, the change is very similar to the change experienced by the lepton charge asymmetry itself. The change in $R(y_W)$ can be traced to a simultaneous increase in the anti-up quark distribution and decrease in the anti-down quark distribution at low x .

4 Conclusions and Outlook

Current standard sets of PDF do not include uncertainties. It is clear that we can not continue to discount them. Already current measurements at the Tevatron have highlighted the importance of these uncertainties for the search of physics beyond the Standard Model. Furthermore, the potential of future hadron colliders to measure $\alpha_s(M_Z)$ and the W -boson mass is impressive, but can not be disentangled from PDF uncertainties. The physics at the LHC will also undoubtedly require a good understanding of the PDF uncertainties. On a more general level, if we want to quantitatively test the framework of perturbative QCD over a very large range of parton collision energies the issue of PDF uncertainties can not be sidestepped.

In this paper we have illustrated a method, based on statistical inference, that can be used to easily propagate uncertainties to new observables, assess the compatibility of new data, and if the latter is good to include the effect of the new data on the PDF without having to redo the whole fit. The method is versatile and modular: an experiment can be included in or excluded from the new PDF fit without any additional work. The statistical and systematic uncertainties with the full point-to-point correlation matrix can be included as well as the theoretical uncertainties. None of the uncertainties are required to be Gaussian distributed.

One remaining problem is the uncertainty associated with the choice of parametrization of the input PDF. This is a difficult problem that does not have a clear answer yet and will require a compromise between the number of parameters and the smoothness of the PDF. We plan to address this question in another paper. The next phase would then be to obtain a large number of initial PDF's sets based on theoretical consideration only, in the spirit of the inference method and Bayes theorem. The DIS and Tevatron data could then be used to constraint the range of these PDF's resulting in a set of

PDF's which would include both experimental and theoretical uncertainties.

A Appendix: Input PDF

For our initial PDF parameter probability density distribution we use the results of Ref. [8]. There a chi-squared fit was performed to DIS data from BCDMS [11], NMC [12], H1 [13] and ZEUS [14]. Both statistical uncertainties and experimental systematic uncertainties with point-to-point correlations were included in the fit, assuming Gaussian distributions. However, *no* theoretical uncertainties were considered. It is important to include the correlation of the systematic uncertainties because the latter usually dominate in DIS data. Simply adding them in quadrature to the statistical uncertainty would result in a overestimation of the uncertainty.

A standard parametrization at $Q_0^2 = 9 \text{ GeV}^2$ is used with 14 ($=N_{par}$) parameters: $x d_v$, $x g$, $x \bar{d}$, $x \bar{u}$, and $x s$ are parametrized using the functional form $x^{\lambda_i} (1-x)^{\lambda_j}$ whereas $x u_v$ is parametrized as $x^{\lambda_i} (1-x)^{\lambda_j} (1 + \lambda_k x)$. Here x is the Bjorken- x . Parton number and momentum conservation constraints are imposed. The full covariance matrix of the parameters, C^{init} , is extracted at the same time as the value of the parameters that minimize the chi-squared. The uncertainties on the parameters were assumed to be Gaussian, such that the fitted values also correspond to the average values of the parameters, μ_{λ_i} . The probability density distribution is then given by

$$P_{init}(\lambda) = \frac{e^{-\frac{\chi_{init}^2(\lambda)}{2}}}{\sqrt{(2\pi)^{N_{par}} |C^{init}|}}, \quad (24)$$

where

$$\chi_{init}^2(\lambda) = \sum_{ij}^{N_{par}} (\lambda_i - \mu_{\lambda_i}) M_{ij}^{init} (\lambda_j - \mu_{\lambda_j}), \quad (25)$$

is the difference between the total chi-squared of the experimental data used in the fit and the minimum chi-squared (1256 for 1061 data points) with the PDF's fixed by the set of parameters $\{\lambda\}$. The matrix M^{init} is the inverse of the covariance matrix C^{init} . The $|C^{init}|$ is the determinant of the covariance matrix. All the calculations were done in the \overline{MS} -scheme.

Comparison with MRS and CTEQ sets showed a good overall agreement with a few exceptions. One example is the difference in the gluon distribution

function at large values of x . The CTEQ and MRS distribution are somewhat above the result of Ref. [8]. This difference was attributed to the fact that prompt photon data were included in the CTEQ and MRS fits. Note that the direct photon data have large scale uncertainty, and it might be misleading to include them in a fit without taking the theoretical uncertainty into account. Also, it is important to keep in mind that it is misleading to compare specific PDF's, as the correlations between different PDF parameters are large.

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