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# Evidence for Parton $k_T$ Effects in High- $p_T$ Particle Production

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# Abstract

Inclusive  $\pi^0$  and direct-photon cross sections in the kinematic range 3.5<  $p_T < 12~{\rm GeV/}c$  with central rapidities  $(y_{\rm cm})$  are presented for 530 and 800  ${\rm GeV/}c$  proton beams and a 515  ${\rm GeV/}c$   $\pi^-$  beam incident on Be targets. Current Next-to-Leading-Order perturbative QCD calculations fail to adequately describe the data for conventional choices of scales. Kinematic distributions from these hard scattering events provide evidence that the interacting partons carry significant initial-state parton transverse momentum  $(k_T)$ . Incorporating these  $k_T$  effects phenomenologically greatly improves the agreement between calculations and the measured cross sections.

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In recent years, perturbative QCD (PQCD) has been tested in a wide variety of processes involving strong interactions at short distances, and increasing attention is now being directed towards areas that may be sensitive to shortcomings in the current theoretical description [1]. The high statistics samples of hard scattering data accumulated by Fermilab fixed-target experiment E706 provide an opportunity to probe such issues. This paper presents comparisons of PQCD calculations to our data on the production of direct photons and  $\pi^0$ 's with large transverse momenta  $(p_T)$ . Direct-photon data have long been expected to provide an accurate determination of the distributions of gluons in hadrons, especially at large longitudinal momentum fraction (x), where information has proven difficult to obtain from other measurements. Inclusive meson production at large  $p_T$  probes a different mix of hard scattering processes and provides insights into parton fragmentation. For conventional choices of parameters, our data are not described satisfactorily by Next-to-Leading-Order (NLO) PQCD calculations [2]. Resolving the observed discrepancies is important for improving the understanding of both parton distribution functions (PDF) and parton fragmentation functions (FF).

Several interesting aspects of QCD contributions beyond Leading-Order (LO) can be investigated experimentally through studies of processes sensitive to transverse motion of the partons prior to the hard scatter. This  $k_T$  is presumably due to effects of hadron size (primordial  $k_T$ ) as well as initial-state gluon radiation. Measurements of Drell-Yan pair production [3] and direct di-photon production [4] have demonstrated the presence of substantial effective  $k_T$ , (significantly larger than can be attributed to primordial  $k_T$ ), and have revealed a significant  $\sqrt{s}$  dependence of  $\langle k_T \rangle$ . A resummation of soft gluon emissions has recently been used to reproduce the size of the effect observed in the WA70 direct di-photon data [5]. Other data also suggest  $\langle k_T \rangle$  values larger than those expected from NLO PQCD calculations. Recent comparisons of  $p_T$  spectra from charm-particle hadroproduction to NLO PQCD results provide evidence that supplemental  $k_T$  may be required to properly describe the data [6]. Likewise, it has been suggested that the observed pattern of discrepancies between data from various direct-photon experiments and results from NLO PQCD calculations could be related to  $k_T$  effects [7].

E706 is designed to measure the production of direct photons, neutral mesons, and associated particles at high  $p_T$ . The apparatus features a large lead and liquid argon electromagnetic calorimeter and a charged particle spectrometer [8]. The experiment accumulated  $\approx 10$  events/pb of  $\pi^-$  beam data at 515 GeV/c,  $\approx 9$  events/pb of proton beam data at 530 GeV/c, and  $\approx 11$  events/pb of proton beam data at 800 GeV/c on Be, Cu, and H targets (primarily Be). A variety of event selection triggers sensitive to high- $p_T$  electromagnetic showers were employed (using different prescale factors) to accumulate data over a broad range of  $p_T$ .

The steep  $p_T$  dependences of neutral meson and direct-photon production make the measured cross sections very sensitive to uncertainties in the energy calibration. Therefore, achieving a precise and accurate calibration of the response of the electromagnetic calorimeter was essential to the success of E706. As a result of detailed studies of the data acquired, the uncertainty in the calibration of the energy response of the calorimeter was reduced to less than 0.5% [9].

The single-photon sample is composed of those photons not identified as elements of reconstructed two photon decays of  $\pi^0$  or  $\eta$  meson candidates. The direct-photon signal

is extracted from the single-photon sample via statistical subtraction of the background contributions. These backgrounds are primarily due to photons from unreconstructed decays of neutral mesons. Failure to correctly identify a photon as originating from a  $\pi^0$  or  $\eta$  decay occurs when the other photon from that decay converts in the target region, escapes the fiducial volume of the calorimeter, or is otherwise not reconstructed. Sources of direct-photon background have been modeled using the HERWIG event generator [10] and a detailed GEANT simulation [11] of the spectrometer response. These Monte Carlo generated events have been weighted to accurately represent our measured neutral meson production spectra.

The  $p_T$  dependences of inclusive  $\pi^0$  and direct-photon cross sections are shown in Figs. 1, 2, and 3. The results of NLO PQCD calculations (using BKK FF for the  $\pi^0$  [12]) are compared with the data [13]. For simplicity, all QCD scales (renormalization, factorization, and, where appropriate, fragmentation) have been set equal. The broken curves in Fig. 1 represent the results of NLO PQCD calculations using conventional choices of scales (and GRV PDF [14]). The calculations are quite sensitive to the scales (an indication of the importance of still higher order contributions), but even for rather small scales, the NLO calculations fail to describe our data. Using other recent PDF [15,16] in the calculations also does not adequately account for the discrepancy between the NLO PQCD results and our data (broken curves in Fig. 2) [17]. Differences between LO [18] and NLO calculations are likewise not large when compared to the difference observed between either and the measured cross sections (Fig. 3). However, PQCD at NLO may not adequately account for soft gluon radiation that imparts an effective transverse momentum to the incident partons.

Kinematic distributions for high-mass pairs of particles directly probe the transverse momentum of incident partons in hard scattering events. The  $p_T$  distribution of  $\pi^0$  pairs produced in 515 GeV/c  $\pi^-$ -nucleon collisions is shown in the inset of Fig. 4. The results of LO PQCD calculations in which the incident partons have Gaussian transverse momentum distributions with  $\langle k_T \rangle = 1.3$  GeV/c (dashed curve) and  $\langle k_T \rangle = 1.7$  GeV/c (solid curve) are shown in Fig. 4 [18]; the dotted curve does not include any incident parton  $k_T$ . For these high-mass  $\pi^0$  pairs, Fig. 4 also shows the distribution of the angle between the two  $\pi^0$ 's in the transverse plane  $(\Delta\phi)$ , and provides a comparison to LO PQCD, both with and without supplemental  $k_T$ . Fragmentation itself generates significant pair- $p_T$  and contributes to the width of the  $\Delta\phi$  distributions, but supplemental  $\langle k_T \rangle > 1$  GeV/c provides a much better description of the data.

Another kinematic variable that is sensitive to  $k_T$  is the out-of-plane momentum,  $p_{\text{out}}$  (the component of the momentum of one high- $p_T$  particle, perpendicular to the plane defined by the incident beam direction and the direction of the other high- $p_T$  particle). The  $p_{\text{out}}$  distributions for  $\gamma\pi^0$  pairs produced in proton-nucleon interactions are shown in the upper part of Fig. 5, compared to LO results without  $k_T$ , and with  $\langle k_T \rangle$  values chosen to bracket the data. The  $p_{\text{out}}$  distributions for  $\pi^0$  pairs are also shown in Fig. 5 (lower plots). These distributions (pair- $p_T$ ,  $\Delta\phi$ , and  $p_{\text{out}}$ ) show clear evidence for the presence of significant  $k_T$  (> 1 GeV/c) in the hard scattering interactions. The corresponding distributions for our other data samples also support this conclusion [19].

Our preliminary analyses of the kinematic distributions of pairs of direct photons, as well as studies of the distribution of the fractional momentum carried by individual charged particles in jets recoiling against isolated photons, also show evidence of substantial  $k_T$ , as do our comparisons of the measured high- $p_T$  charged-D cross section to NLO PQCD results [8].

All these results suggest a supplemental  $\langle k_T \rangle$  of order 1 GeV/c.

Since the inclusive spectra fall rapidly with increasing  $p_T$ , the introduction of  $k_T$ -smearing has a significant effect on predicted cross sections. To approximate the effect of supplemental  $k_T$  smearing on the inclusive NLO PQCD calculations for direct-photon (and  $\pi^0$ ) production, we calculated  $k_T$  factors (as functions of  $p_T$ ) for different values of  $\langle k_T \rangle$ , by computing ratios of results from LO PQCD calculations [18] for different  $\langle k_T \rangle$  values compared to results without  $k_T$  [20]. These same  $k_T$  factors were then applied to the results of NLO PQCD calculations. As indicated by the solid curves in Figs. 1, 2, and 3, reasonable representations of both the direct-photon and  $\pi^0$  results are obtained using  $\langle k_T \rangle$  values > 1 GeV/c [21]. The kinematic distributions exhibit a pattern consistent with increasing  $\langle k_T \rangle$  as  $\sqrt{s}$  increases, a trend reflected in the choices of  $\langle k_T \rangle$  factors employed in the theory curves (solid curves) shown in the inclusive cross section plots.

As an illustration of the sensitivity of our data to the gluon distribution, Fig. 6 compares our direct-photon cross sections to NLO PQCD calculations using CTEQ4M and CTEQ4HJ PDF [16]. Once soft-gluon effects are satisfactorily taken into account, either approximately as in this paper or in a more theoretically rigorous manner, our data can be used to help discriminate between PDF that otherwise provide acceptable descriptions of the data sets used in Ref. [16].

In conclusion, we have measured the inclusive production of high- $p_T$  neutral mesons and direct photons by 530 GeV/c and 800 GeV/c proton and 515 GeV/c  $\pi^-$  beams. Current NLO PQCD calculations (which exhibit substantial dependences on QCD scales) fail to account for the measured cross sections using conventional choices of scales. Significant  $k_T$  effects (> 1 GeV/c) have been observed in the kinematic distributions of high-mass pairs of  $\pi^0$ 's, as well as high-mass  $\gamma\pi^0$  pairs. A simple implementation of supplemental parton  $k_T$  in PQCD calculations, using  $k_T$  values consistent with observations, provides a reasonable description of the inclusive cross sections. Our high statistics direct-photon data samples are directly sensitive to the gluon distribution at large x values. An improved theoretical understanding of soft-gluon effects in inclusive direct-photon production will facilitate the global determination of the gluon distribution function.

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- [21] The  $\langle k_T \rangle$  values that provide a corresponding level of agreement between LO direct-photon calculations (not shown) and our inclusive cross sections are somewhat larger ( $\approx 10\%$ ) than the corresponding values for the NLO comparisons.

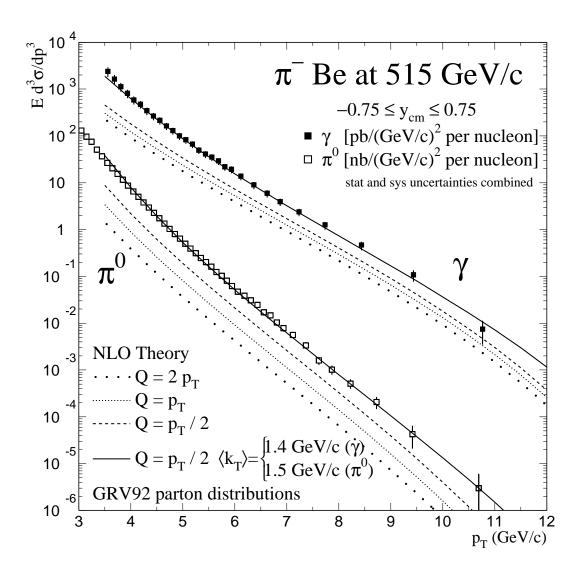


FIG. 1. The  $\pi^0$  and direct-photon inclusive cross sections as functions of  $p_T$  for 515 GeV/c  $\pi^-$ -nucleon interactions compared to NLO PQCD results for several choices of scales. The solid curves show the NLO PQCD results for  $Q=p_T/2$  scales adjusted for supplemental  $\langle k_T \rangle$ . (Note that the units for the  $\pi^0$  and  $\gamma$  results differ by a factor of 1000.)

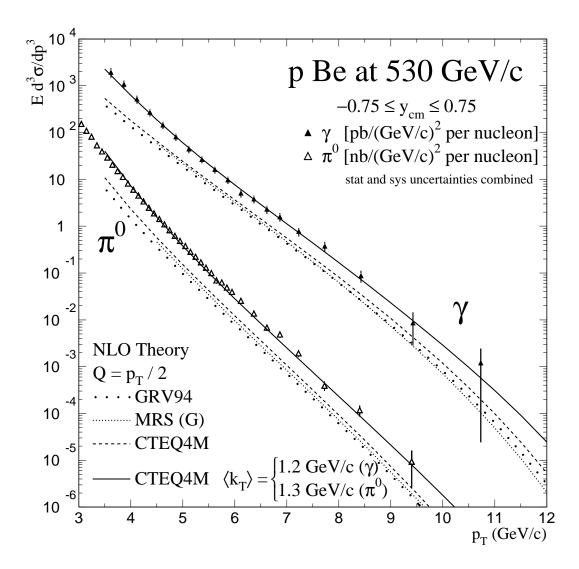


FIG. 2. The  $\pi^0$  and direct-photon inclusive cross sections as functions of  $p_T$  for 530 GeV/c proton-nucleon interactions compared to NLO PQCD results for several choices of PDF. The solid curves show the NLO result (using the CTEQ4M PDF) adjusted for supplemental  $\langle k_T \rangle$ . (Note that the units for the  $\pi^0$  and  $\gamma$  results differ by a factor of 1000.)

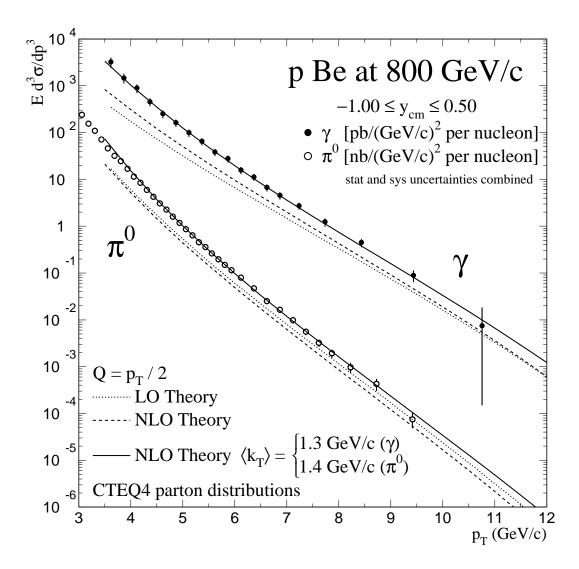


FIG. 3. The  $\pi^0$  and direct-photon cross sections as functions of  $p_T$  for 800 GeV/c proton-nucleon interactions compared to LO and NLO PQCD results. The solid curves show NLO results adjusted for supplemental  $\langle k_T \rangle$ . (Note that the units for the  $\pi^0$  and  $\gamma$  results differ by a factor of 1000.)

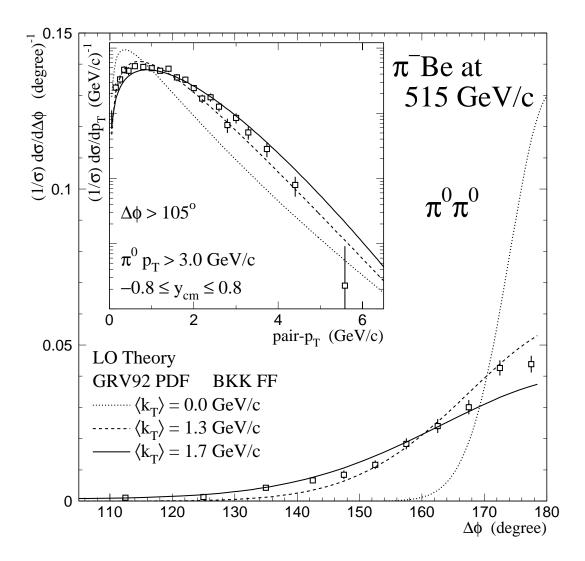


FIG. 4. The  $\Delta\phi$  distribution for high-mass  $\pi^0$  pairs produced in 515 GeV/c  $\pi^-$ -nucleon interactions compared to curves showing LO PQCD results using various  $\langle k_T \rangle$  values. The inset shows the pair- $p_T$  distribution for such pairs and the corresponding results of LO PQCD calculations.

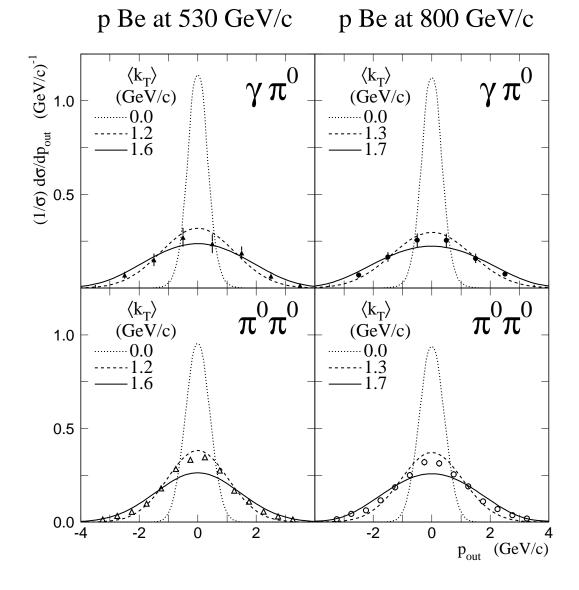


FIG. 5. The out-of-plane momentum distributions for high-mass pairs produced in proton-nucleon interactions at 530 and 800 GeV/c compared to results of LO PQCD calculations (using CTEQ4L PDF) for several  $\langle k_T \rangle$  values.

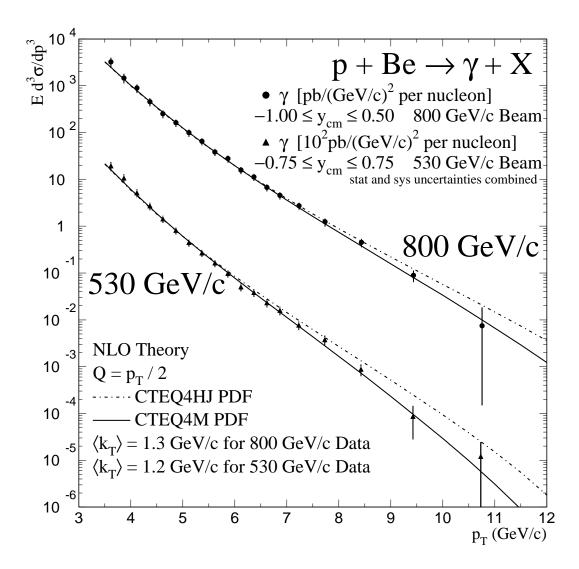


FIG. 6. Direct-photon inclusive cross sections as functions of  $p_T$  for 530 and 800 GeV/c proton-nucleon interactions compared to results of NLO PQCD calculations using CTEQ4HJ (dot-dashed curve) and CTEQ4M (solid curve) PDF. Factors for supplemental  $\langle k_T \rangle$  are included. (Note that the units for the 530 and 800 GeV/c results differ by a factor of 100.)