Blois - 16 July '10

# The Physics of the LHC

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# The LHC physics run is going on, .... finally!

3.5 TeV per beam. Goal 1$fb^{-1}$ by the end of 2011

## Top physics priorities at the LHC (ATLAS&CMS):

- Clarify the EW symmetry breaking sector
- Search for new physics at the TeV scale
- Identify the particle(s) that make the Dark Matter in the Universe

## Also:

- LHCb: precision B physics (CKM matrix and CP violation)
- ALICE: Heavy ion collisions & QCD phase diagram

⊕ At this point, fresh input from experiment is badly needed!!

# Particle physics at a glance

The SM is a low energy effective theory
(nobody can believe it is the ultimate theory)

It happens to be renormalizable, hence highly predictive.
And is well supported by the data.

However, we expect corrections from higher energies

not only from the GUT or Planck scales
but also from the TeV scale (LHC!)

In fact even just as a low energy effective theory
the SM it is not satisfactory

QCD + the gauge part of the EW theory are fine,
but the Higgs sector is so far only a conjecture

# The Higgs problem is central in particle physics today

A review: G.A. ArXiv:1003.3180

## The main problems of the SM show up in the Higgs sector

$$V_{Higgs} = V_0 - \mu^2 \phi^\dagger \phi + \lambda \left( \phi^\dagger \phi \right)^2 + [\bar{\psi}_{Li} Y_{ij} \psi_{Rj} \phi + h.c.]$$

Vacuum energy
$V_{0exp} \sim (2.10^{-3}\ eV)^4$

Origin of quadratic
divergences.
Hierarchy problem

Possible instability
depending on $m_H$

The flavour problem:
large unexplained ratios
of $Y_{ij}$ Yukawa constants

# The Standard EW theory: $\mathcal{L} = \mathcal{L}_{\text{symm}} + \mathcal{L}_{\text{Higgs}}$

$$\mathcal{L}_{symm} = -\frac{1}{4}[\partial_\mu W_\nu^A - \partial_\nu W_\mu^A - ig\varepsilon_{ABC} W_\mu^A W_\nu^B]^2 +$$

$$-\frac{1}{4}[\partial_\mu B_\nu - \partial_\nu B_\mu]^2 +$$

$$+\bar{\psi}\gamma^\mu[i\partial_\mu + g W_\mu^A t^A + g' B_\mu \frac{Y}{2}]\psi$$

$$\mathcal{L}_{Higgs} = |[\partial_\mu - ig W_\mu^A t^A - ig' B_\mu \frac{Y}{2}]\phi|^2 +$$

$$+ V[\phi^\dagger\phi] + \bar{\psi}\Gamma\psi\phi + h.c$$

with $V[\phi^\dagger\phi] = \mu^2(\phi^\dagger\phi)^2 + \lambda(\phi^\dagger\phi)^4$

$\mathcal{L}_{\text{symm}}$: well tested (LEP, SLC, Tevatron...), $\mathcal{L}_{\text{Higgs}}$: ~ untested

All we know from experiment about the SM Higgs:

No Higgs seen at LEP2 -> $m_H > 114.4$ GeV (95%cl)

Rad. corr's -> $m_H < 186$ GeV (95%cl, incl. direct search bound)

$v = \langle\phi\rangle = \sim 174$ GeV ; $m_W = m_Z \cos\theta_W$ ⟶ doublet Higgs

That some sort of spontaneous symmetry breaking mechanism is at work has already been established (couplings symmetric, spectrum totally non symmetric)

The question is on the nature of the Higgs mechanism/particle(s)

- One doublet, more doublets, additional singlets?
- SM Higgs or SUSY Higgses
- Fundamental or composite (of fermions, of WW....)
- Pseudo-Goldstone boson of an enlarged symmetry
- A manifestation of extra dimensions (fifth comp. of a gauge boson, an effect of orbifolding or of boundary conditions....)
- Some combination of the above

## Can we do without the Higgs?

Suppose we take the gauge symmetric part of the SM and put masses by hand.

Gauge invariance is broken explicitly. The theory is no more renormalizable. One loses understanding of the observed accurate validity of gauge predictions for couplings.

Still, what is the fatal problem at the LHC scale?

The most immediate disease that needs a solution is the occurrence of unitarity violations in some amplitudes

To avoid this either there is one or more Higgs particles or some new states (e.g. new vector bosons)

Thus something must happen at the few TeV scale!!

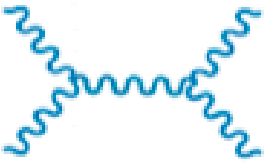
## With no Higgs unitarity violations for $E_{CM}$ ~ 1-3 TeV

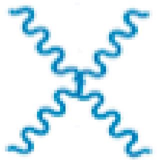
Unitarity implies that scattering amplitudes cannot grow indefinitely with the centre-of-mass energy s

In the SM, the Higgs particle is essential in ensuring that the scattering amplitudes with longitudinal weak bosons ($W_L$ , $Z_L$) satisfy (tree-level) unitarity constraints
[Veltman, 1977; Lee-Quigg-Thacker, 1977; …]

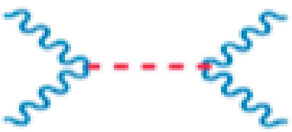
Zwirner

An example: $\mathcal{A}(W_L^+ \, W_L^- \longrightarrow Z_L \, Z_L) \quad (s \gg m_W^2)$





$$i \frac{s}{v^2}$$



$$-i \frac{s^2}{v^2(s - m_h^2)}$$

$$\rightarrow \quad -i \frac{m_h^2}{v^2} \frac{s}{s - m_h^2}$$

If no Higgs then something must happen!

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## A crucial question for the LHC

### What saves unitarity?

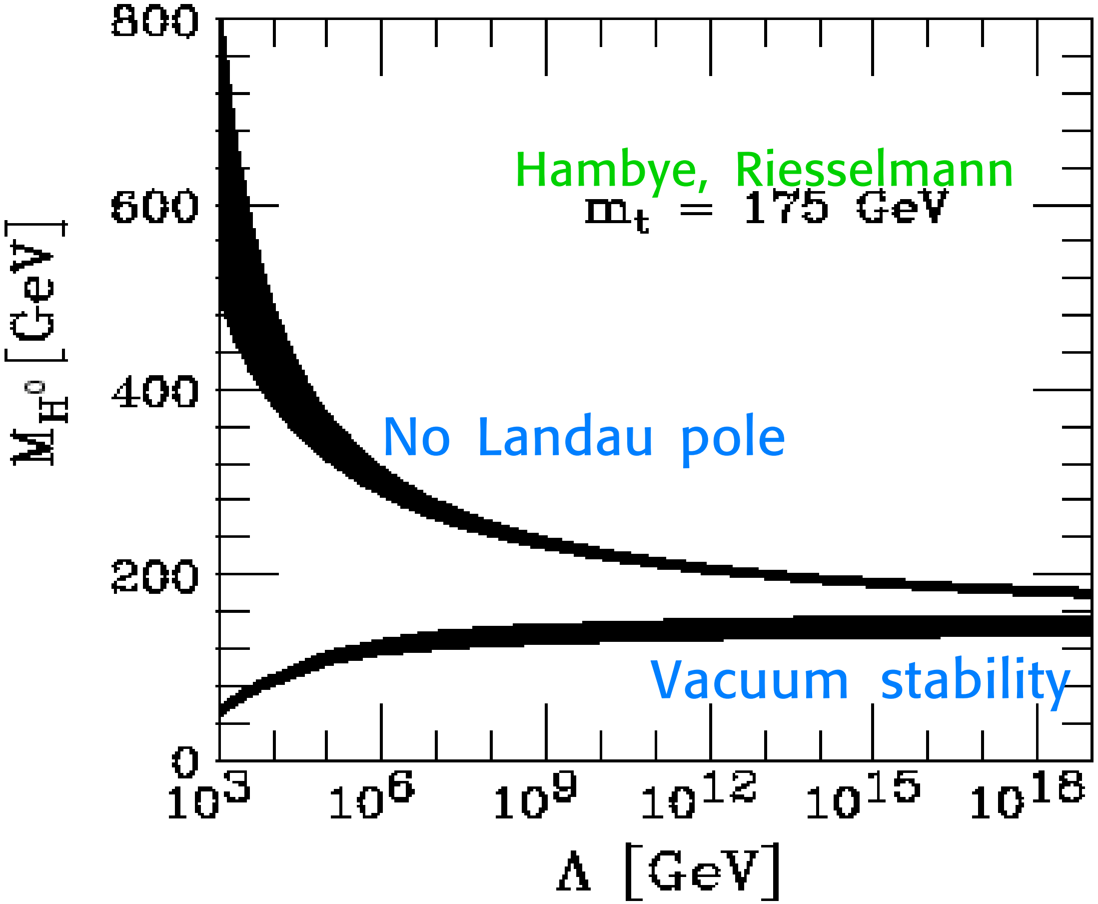
- the Higgs
- some new vector boson
  - W', Z'
  - KK recurrences
  - resonances from a strong sector
  - ......

# Theoretical bounds on the SM Higgs mass

$\Lambda$: scale of new physics beyond the SM

Upper limit: No Landau pole up to $\Lambda$
Lower limit: Vacuum (meta)stability

The LHC was designed to cover the whole range




**If the SM would be valid up to $M_{GUT}$, $M_{Pl}$ then $m_H$ would be limited in a small range**

Lower now because of $m_t$ → 128 GeV $< m_H <$ 180 GeV

## Status of the SM Higgs fit

## Radiative corr's indicate a light H

Summer '09

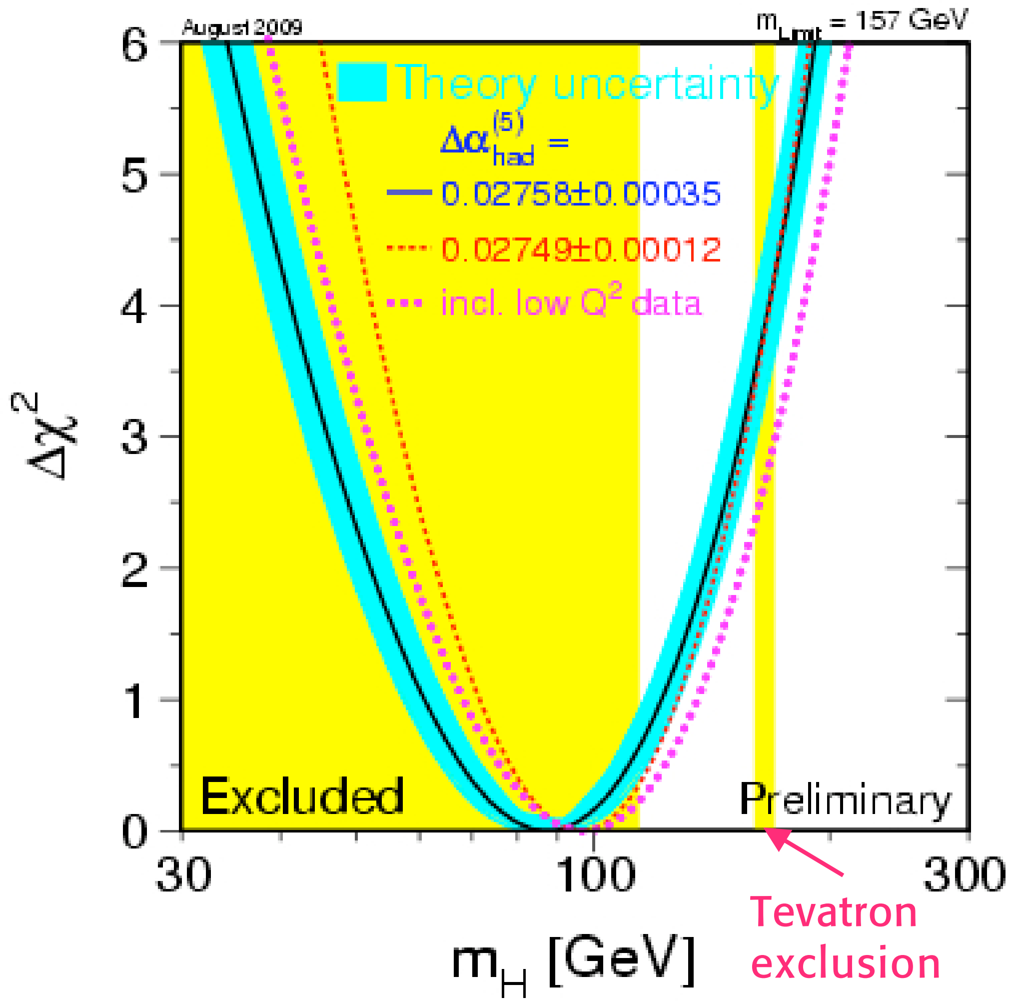
Sensitive to log $m_H$

Rad Corr.s ->
$\log_{10} m_H(\mathrm{GeV}) = 1.94 \pm 0.15$

$m_H = 87 + 35 - 26$ GeV

This is a great triumph for the SM: ~right in the narrow allowed range $\log_{10} m_H \sim 2 - 3$

Direct search: $m_H > 114.4$ GeV




At 95 % cl
$m_H < 157$ GeV (rad corr.'s)
$m_H < 186$ GeV (incl. direct search bound)

# Is it possible that the Higgs is not found at the LHC?

Here "Higgs" means the "the EW symmetry breaking mechanism"

Looks pretty unlikely!!

The LHC discovery range is large enough: $m_H < \sim 1$ TeV
the Higgs should be really heavy!

Rad. corr's indicate a light Higgs (whatever its nature)

A heavy Higgs would make perturbation theory to collapse nearby (violations of unitarity for $m_H > \sim$ TeV)

Such nearby collapse of pert. th. is very difficult to reconcile with EW precision tests plus simulating a light Higgs

The SM good agreement with the data favours forms of new physics that keep at least some Higgs light

# The Standard Model works very well

So, why not find the Higgs and declare particle physics solved?

First, you have to find it!

⟶ LHC

Because of both:

Conceptual problems

- Quantum gravity
- The hierarchy problem
- The flavour puzzle
- •••••

and experimental clues:

- Neutrino masses
- Coupling unification
- Dark matter
- Baryogenesis
- Vacuum energy
- •••••

Some of these problems point at new physics at the weak scale: eg
Hierarchy
Dark matter (perhaps)

# Dark Matter

WMAP, SDSS, 2dFGRS….

Most of the Universe is not made up of atoms: $\Omega_{tot}\sim 1$, $\Omega_b \sim 0.045$, $\Omega_m \sim 0.27$
Most is Dark Matter and Dark Energy

Most Dark Matter is Cold (non relativistic at freeze out)
Significant Hot Dark matter is disfavoured
Neutrinos are not much cosmo-relevant: $\Omega_\nu < 0.015$

SUSY has excellent DM candidates: eg Neutralinos (--> LHC)
Also Axions are still viable (introduced to solve strong CPV)
(in a mass window around m $\sim 10^{-4}$ eV and $f_a \sim 10^{11}$ GeV
but these values are simply a-posteriori)

Identification of Dark Matter is a task of enormous importance for particle physics and cosmology

LHC?

LHC has good chances because it can reach any kind of WIMP:

WIMP: Weakly Interacting Massive Particle
with m ~ $10^1$-$10^3$ GeV

For WIMP's in thermal equilibrium after inflation the density is:

$$\Omega_\chi h^2 \simeq const. \cdot \frac{T_0^3}{M_{\rm Pl}^3 \langle \sigma_A v \rangle} \simeq \frac{0.1 \text{ pb} \cdot c}{\langle \sigma_A v \rangle}$$

can work for typical weak cross-sections!!!

This "coincidence" is a good indication in favour of a WIMP explanation of Dark Matter

# A crucial question for the LHC

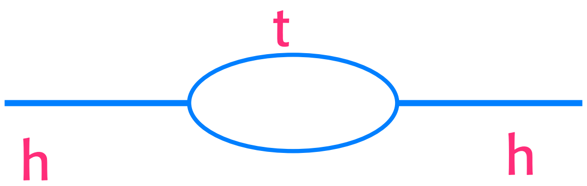
Is Dark Matter a WIMP?

LHC will tell yes or no to WIMPS

# The "little hierarchy" problem

e.g. the top loop (the most pressing):

$m_h^2 = m^2_{bare} + \delta m_h^2$




$$\delta m^2_{h|top} = -\frac{3G_F}{2\sqrt{2}\pi^2} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$$

$\Lambda \sim o(1\text{TeV})$

This hierarchy problem demands new physics near the weak scale

$\Lambda$: scale of new physics beyond the SM

- $\Lambda >> m_Z$: the SM is so good at LEP
- $\Lambda \sim$ few times $G_F^{-1/2} \sim o(1\text{TeV})$ for a natural explanation of $m_h$ or $m_W$

Barbieri, Strumia

The LEP Paradox: $m_h$ light, new physics must be close but its effects were not visible at LEP2

The B-factory Paradox: and not visible in flavour physics

## A crucial question for the LHC

What damps the top loop $\Lambda^2$ dependence?

- the s-top (SUSY)
- some new fermion

    t′ (Little Higgs)
    KK recurrences of the top (Extra dim.)
    ……

# Solutions to the hierarchy problem

- Supersymmetry: boson-fermion symm.

  exact (unrealistic): cancellation of $\Lambda^2$ in $\delta m_h{}^2$

  approximate (possible): $\Lambda \sim m_{SUSY}-m_{ord}$ → top loop $\Lambda \sim m_{stop}$

  The most widely accepted

- The Higgs is a $\bar{\psi}\psi$ condensate. No fund. scalars. But needs new very strong binding force: $\Lambda_{new} \sim 10^3 \Lambda_{QCD}$ (technicolor).

  Strongly disfavoured by LEP. Coming back in new forms

- Models where extra symmetries allow $m_h$ only at 2 loops and non pert. regime starts at $\Lambda \sim 10$ TeV

  "Little Higgs" models. Some extra trick needed to solve problems with EW precision tests

- Extra spacetime dim's that "bring" $M_{Pl}$ down to o(1TeV)

Exciting. Many facets. Rich potentiality. No baseline model emerged so far

- ⊕ Ignore the problem: invoke the anthropic principle

# The anthropic route

The scale of the cosmological constant is a big mystery.

$\Omega_\Lambda \sim 0.75$ → $\rho_\Lambda \sim (2\ 10^{-3}\ \mathrm{eV})^4 \sim (0.1\mathrm{mm})^{-4}$ → Similar to $m_\nu$!?

In Quantum Field Theory: $\rho_\Lambda \sim (\Lambda_{cutoff})^4$

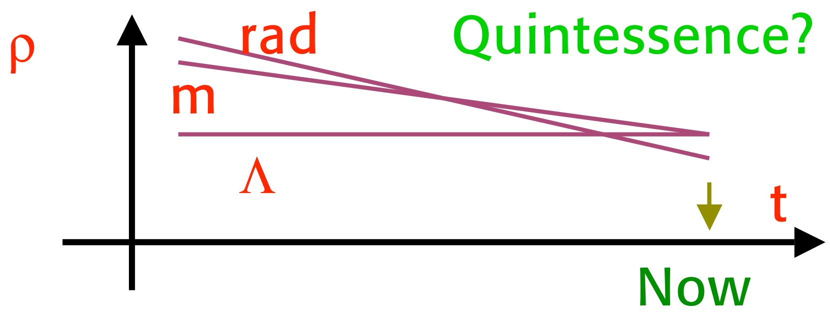
If $\Lambda_{cutoff} \sim M_{Pl}$ → $\rho_\Lambda \sim 10^{123} \rho_{obs}$

Exact SUSY would solve the problem: $\rho_\Lambda = 0$

But SUSY is broken: $\rho_\Lambda \sim (\Lambda_{SUSY})^4 \sim 10^{59} \rho_{obs}$

It is interesting that the correct order is $(\rho_\Lambda)^{1/4} \sim (\Lambda_{EW})^2/M_{Pl}$

Other problem:
"Why now"?




"Quintessence"
$\Lambda$ as a vev of a field $\phi$?

Coupled to gauge singlet matter, eg $\nu_R$, to solve magnitude and why now?

# Is naturalness relevant?

Speculative physics reasons to doubt:

- The empirical value of the cosmological constant $\Lambda$ poses a tremendous, unsolved naturalness problem

  yet the value of $\Lambda$ is close to the Weinberg upper bound for galaxy formation

- Possibly our Universe is just one of infinitely many continuously created from the vacuum by quantum fluctuations

- Different physics in different Universes according to the multitude of string theory solutions ($\sim 10^{500}$)

Perhaps we live in a very unlikely Universe but one that allows our existence

I find applying the anthropic principle to the SM hierarchy problem not appropriate

After all we can find plenty of models that reduce the fine tuning from $10^{14}$ to $10^{2}$:
so why make our Universe so terribly unlikely?

The case of the cosmological constant is a lot different: the context is not as fully specified as the for the SM (quantum gravity, string cosmology, branes in extra dims., wormholes thru different Universes....)

# SUSY: boson fermion symmetry

The hierarchy problem: $$\delta m^2_{h|top} = -\frac{3G_F}{2\sqrt{2}\pi^2} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$$

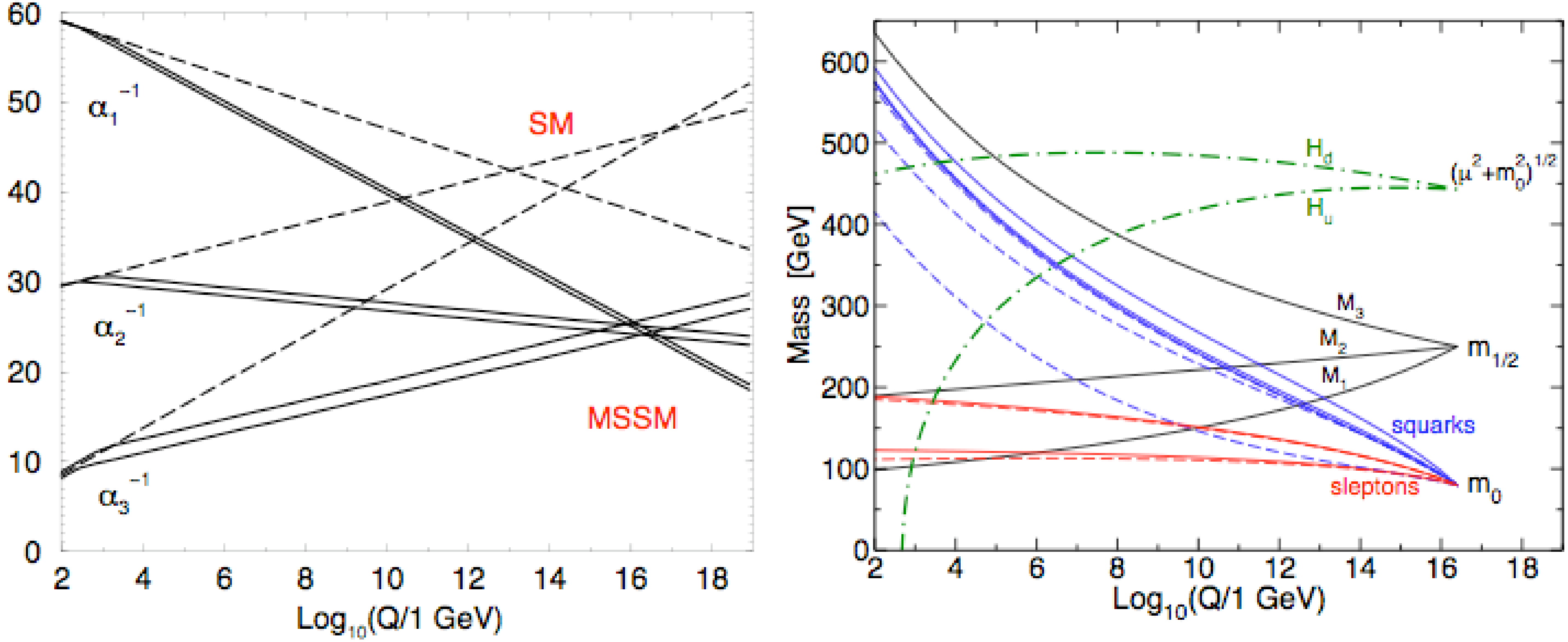
In broken SUSY $\Lambda^2$ is replaced by $(m_{stop}^2 - m_t^2)\log\Lambda$

$m_H > 114.4$ GeV, $m_{\chi+} > 100$ GeV, EW precision tests, success of CKM, absence of FCNC, all together, impose sizable Fine Tuning (FT) particularly on minimal realizations (MSSM, CMSSM…).

Yet SUSY is a completely specified, consistent, computable model, perturbative up to $M_{Pl}$ quantitatively in agreement with coupling unification (GUT's)
(unique among NP models)
and has a good DM candidate: the neutralino (actually more than one).

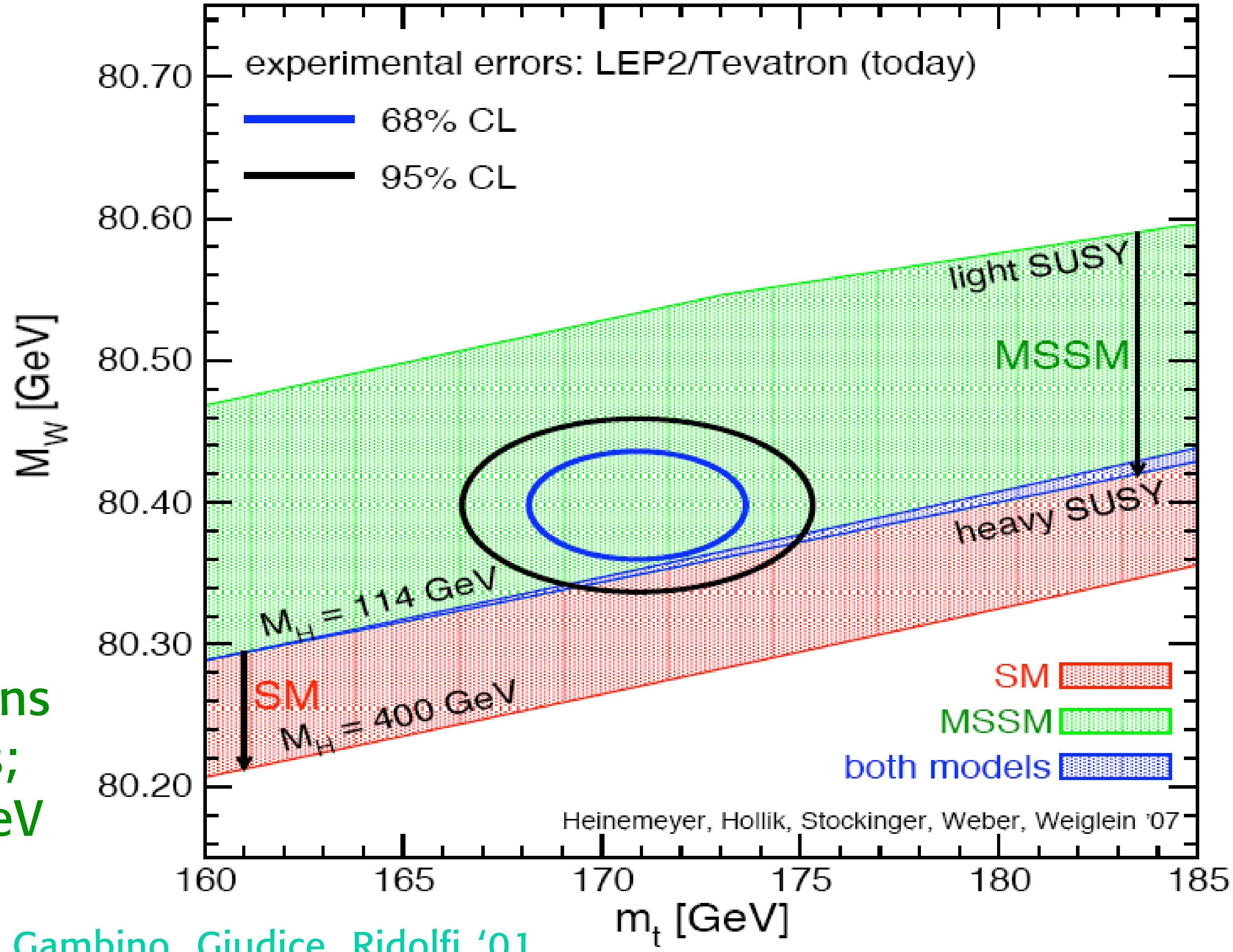
Remains the reference model for NP

## SUSY is unique in providing a weakly interacting theory up to the GUT/Planck scale




Other BSM models (little Higgs, composite Higgs, Higgsless….) all become strongly interacting at a multi-TeV scale

# SUSY effects could modify the SM fit




"light SUSY"=
= light s-leptons
and charginos;
s-quarks ~1 TeV

G.A, Caravaglios, Gambino, Giudice, Ridolfi '01

## A recent study indicates that $m_h$ goes up in CMSSM when b->s$\gamma$, $a_\mu$, $\Omega_{DM}$ are added

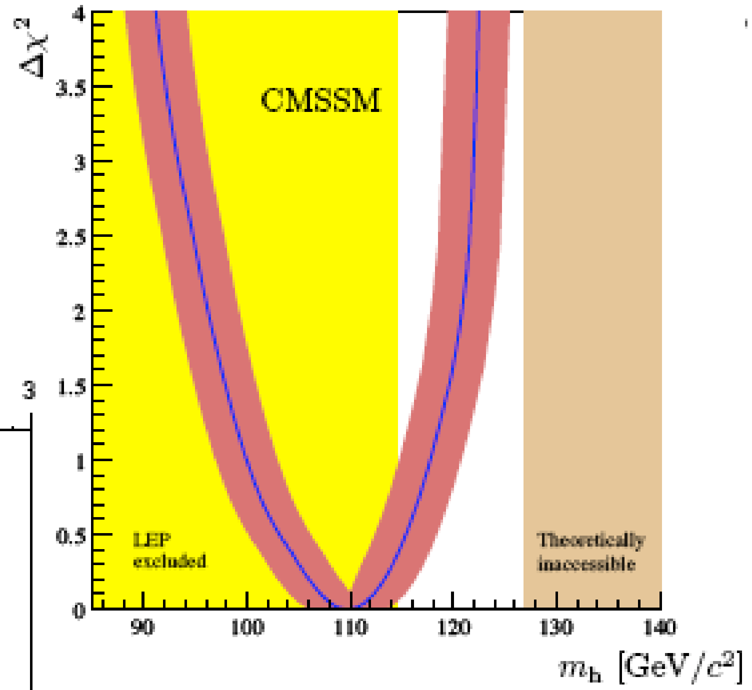
O. Buchmuller et al '07, '08 [0808.4128]

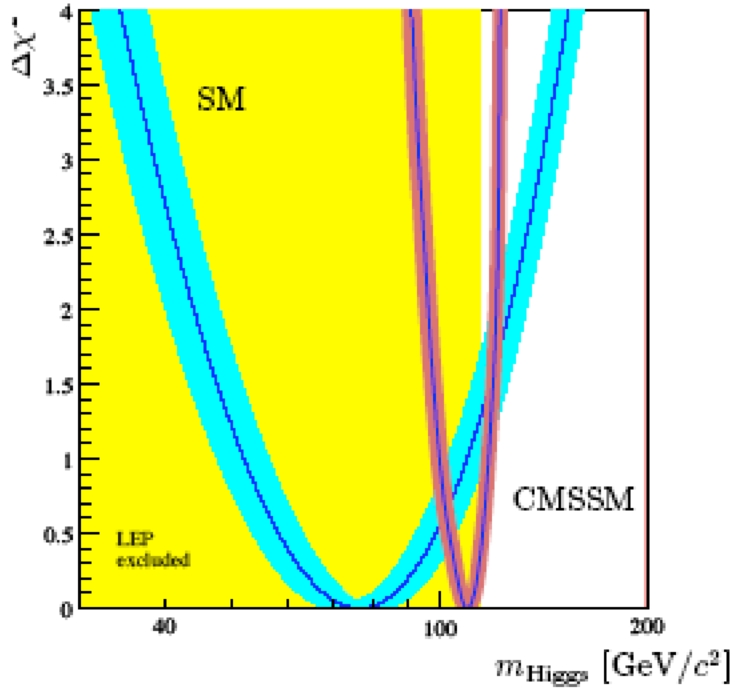
also:
J. Ellis et al '07

CMSSM

| Variable | Measurement | Fit |
|---|---|---|
| $\Delta\alpha^{(5)}_{had}(m_Z)$ | 0.02758± 0.00035 | 0.02774 |
| $m_Z$ [GeV] | 91.1875± 0.0021 | 91.1873 |
| $\Gamma_Z$ [GeV] | 2.4952± 0.0023 | 2.4952 |
| $\sigma^0_{had}$ [nb] | 41.540± 0.037 | 41.486 |
| $R_l$ | 20.767± 0.025 | 20.744 |
| $A^{0,l}_{fb}$ | 0.01714± 0.00095 | 0.01641 |
| $A_l(P_\tau)$ | 0.1465± 0.0032 | 0.1479 |
| $R_b$ | 0.21629± 0.00066 | 0.21613 |
| $R_c$ | 0.1721± 0.0030 | 0.1722 |
| $A^{0,b}_{fb}$ | 0.0992± 0.0016 | 0.1037 |
| $A^{0,c}_{fb}$ | 0.0707± 0.0035 | 0.0741 |
| $A_b$ | 0.923± 0.020 | 0.935 |
| $A_c$ | 0.670± 0.027 | 0.668 |
| $A_l$(SLD) | 0.1513± 0.0021 | 0.1479 |
| $\sin^2\theta^{lept}_{eff}(Q_{fb})$ | 0.2324± 0.0012 | 0.2314 |
| $m_W$ [GeV] | 80.398± 0.025 | 80.382 |
| $m_t$ [GeV] | 170.9± 1.8 | 170.8 |
| $R(b\to s\gamma)$ | 1.13± 0.12 | 1.12 |
| $B_s\to\mu\mu$ [$\times10^{-8}$] | < 8.00 | 0.33 |
| $\Delta a_\mu$ [$\times10^{-9}$] | 2.95± 0.87 | 2.95 |
| $\Omega h^2$ | 0.113± 0.009 | 0.113 |

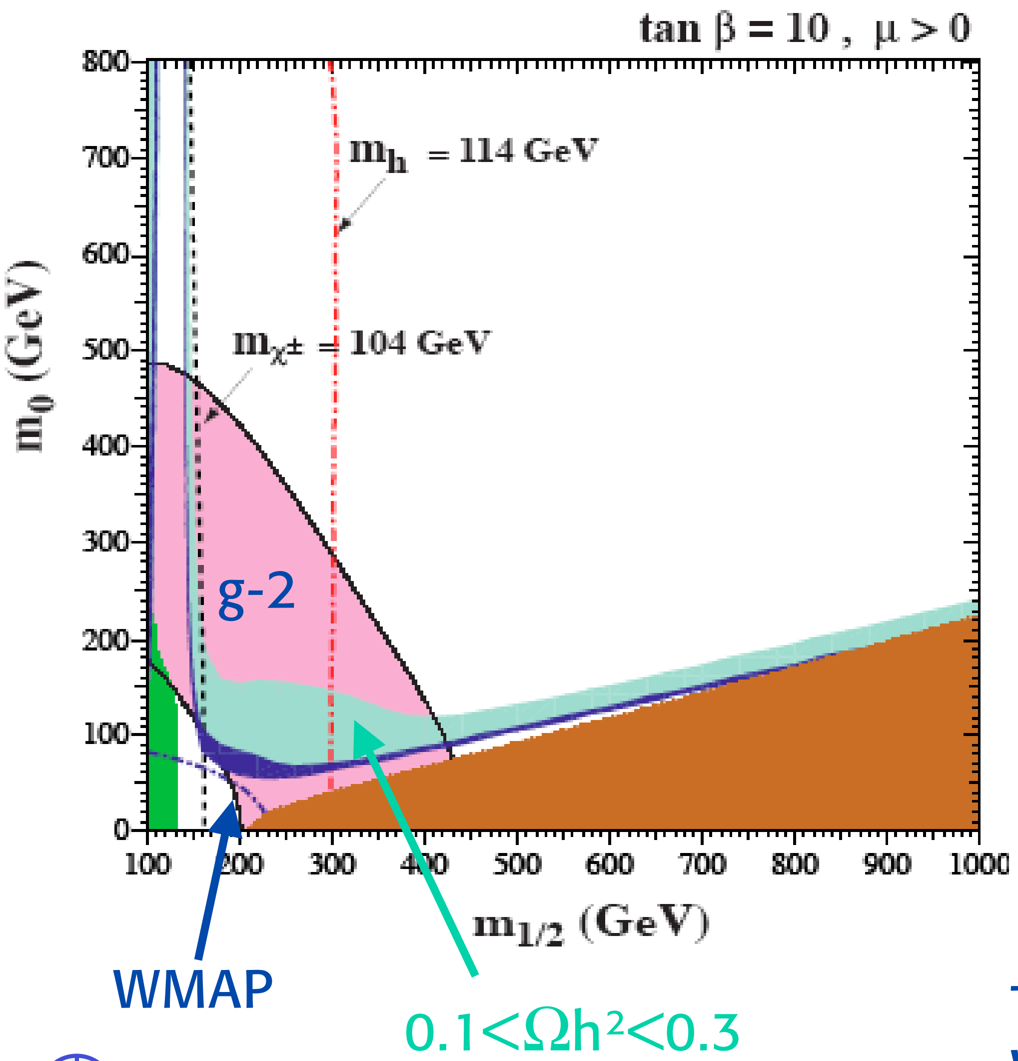
$|O^{meas}-O^{fit}|/\sigma^{meas}$ (0, 1, 2, 3); $B_s\to\mu\mu$: N/A (upper limit)

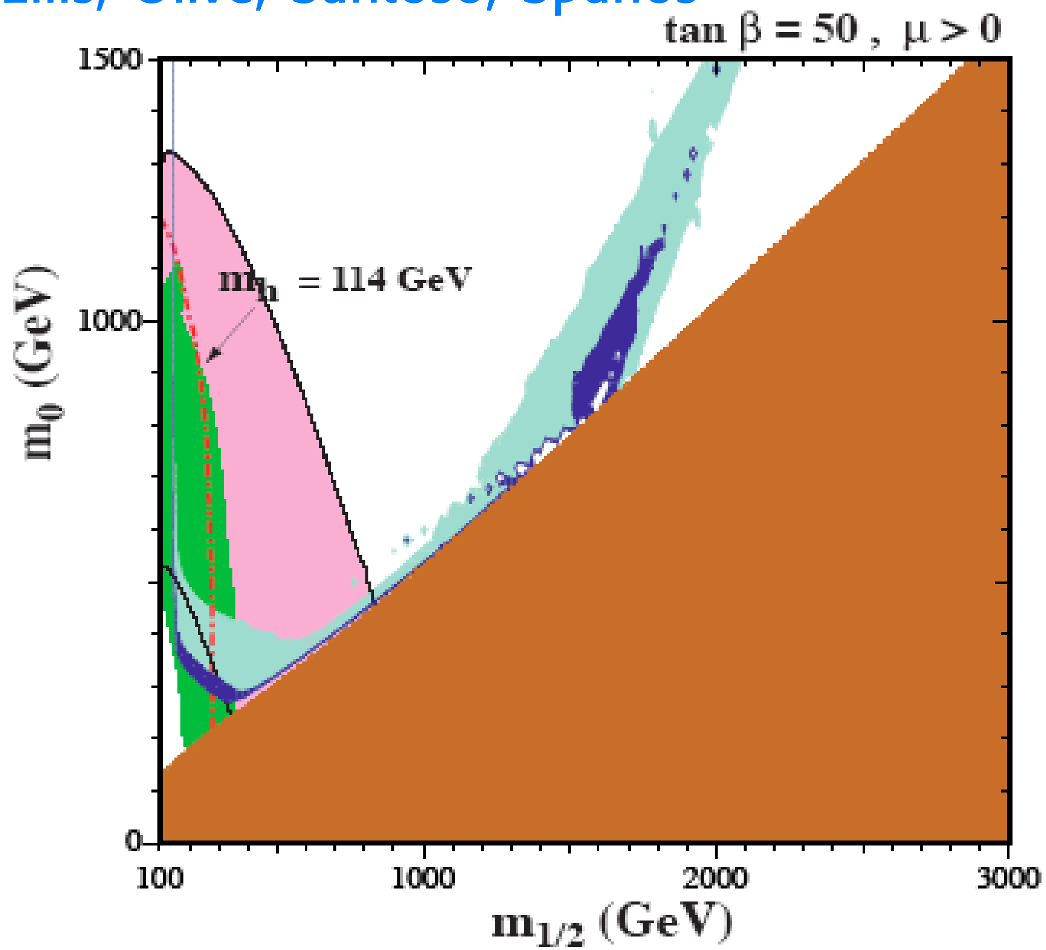


# SUSY Dark Matter: best candidate the neutralino
## [in SUSY the gravitino is a non-WIMP alternative]

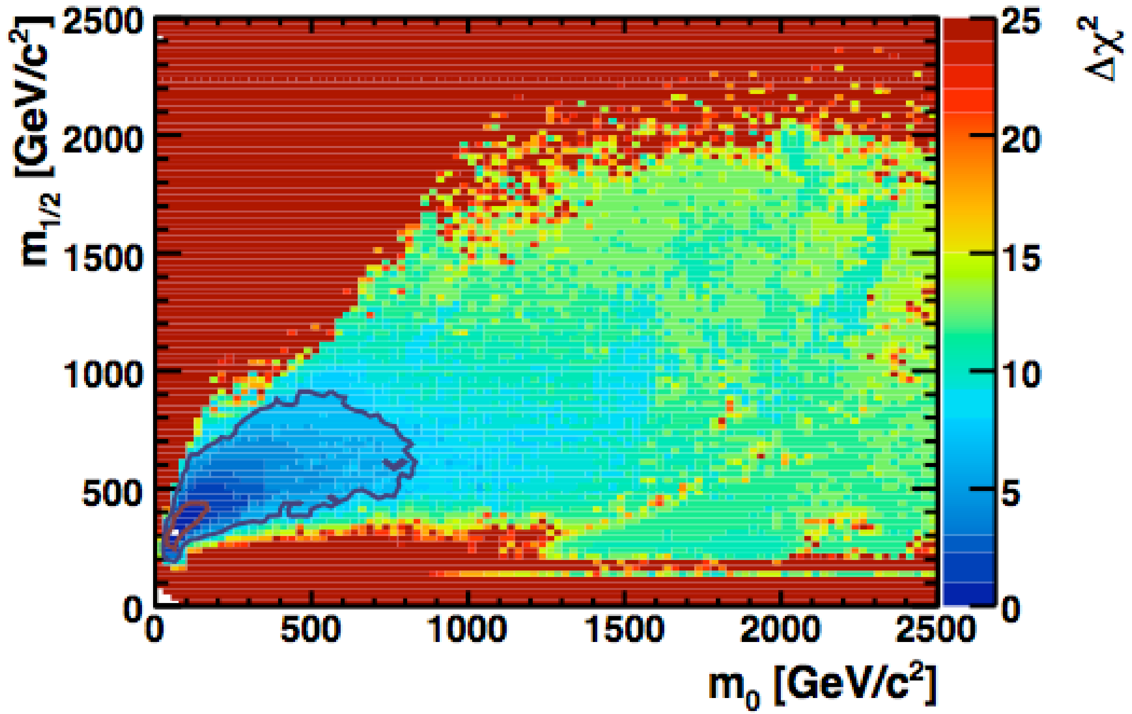
Ellis, Olive, Santoso, Spanos

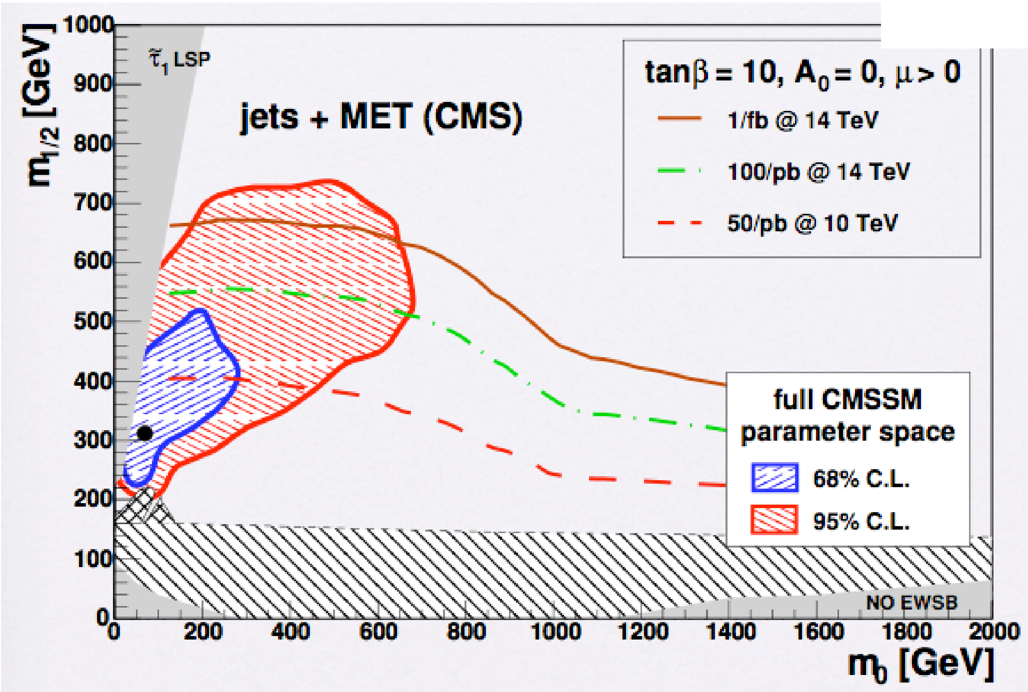



This is for the CMSSM
With less constraints, more space

relatively light SUSY
is indicated




accesssible at the LHC

## Muon g-2 and SUSY

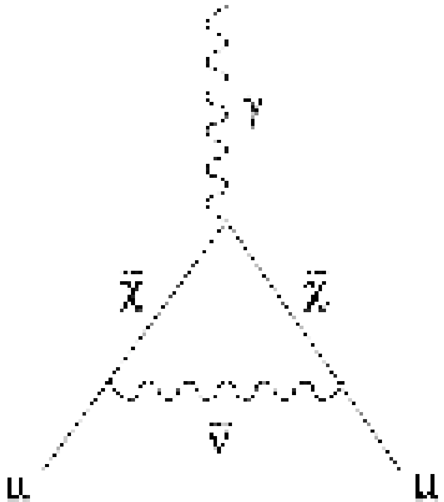
Observed Difference with Experiment:

$$a_\mu^{exp} - a_\mu^{SM} = (27.5 \pm 8.4) \times 10^{-10}$$

➡ 3.3 "standard deviations"

Could be new physics
eg light SUSY

$$\delta a_\mu = 13 \cdot 10^{-10} \left( \frac{100 GeV}{M_{SUSY}} \right)^2 tg\beta$$




$a_\mu$ is a plausible
location for a
new physics signal!!

# Little Higgs Models

Georgi (moose)/Arkani-Hamed et al/Low, Skiba, Smith/Kaplan, Schmaltz/Chang,Wacker/Gregoire et al

$$G \supset [SU(2) \otimes U(1)]^2 \supset SU(2) \otimes U(1)$$

global gauged SM

H is (pseudo)-Goldstone boson of G: takes mass only at 2-loops (needs breaking of 2 subgroups or 2 couplings)

recall: $\delta m^2_{h|top} = -\frac{3G_F}{2\sqrt{2}\pi^2} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$ $G_F \sim g^2 \rightarrow g^4$

cutoff $\Lambda$ ~10 TeV

$\Lambda^2$ divergences canceled by:

| | | |
|---|---|---|
| $\delta m^2_{H\|top}$ | new coloured fermion $\chi$ with Q=2/3 | ~1 TeV |
| $\delta m^2_{H\|gauge}$ | W', Z', $\gamma'$ | |
| $\delta m^2_{H\|Higgs}$ | new scalars | |

$\oplus$ 2 Higgs doublets ~0.2 TeV

With some tension, Little Higgs models technically can work.

T parity interchanges the two SU(2)xU(1) groups  Cheng, Low

Standard gauge bosons are T even, heavy ones are T odd

Lightest T-odd particle stable --> Dark Matter

Technically sophisticated. But the main drawback is:
Little Higgs provides just a postponement:
UV completion beyond ~10 TeV? GUT's?

Still important as it offers well specified signals and signatures for searching at the LHC:
a light Higgs, a new top-like fermion $\chi$ to damp the top loop, new W′, Z′ for the W, Z loops,.....

# Extra Dimensions (ED)

String Theory ---> ED at $M_{Pl}$

Perhaps ED have a direct impact on physics below $M_{Pl}$

Exciting possibilities (a large domain of contemporary BSM)

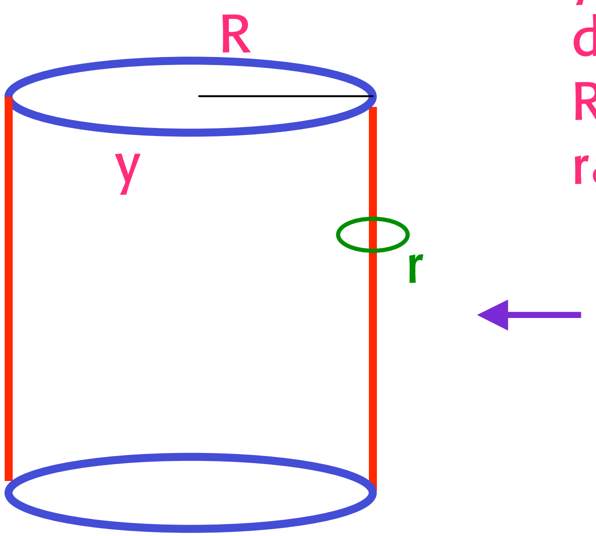
- GUT's in ED ($M_{GUT}$)
- ED as (part of the) solution of the hierarchy probem ($M_{EW}$)
- EW symmetry breaking from ED ($M_{EW}$)

## Early formulation

Solve the hierachy problem by bringing gravity down from $M_{Pl}$ to o(1TeV)

Arkani-Hamed, Dimopoulos/ Dvali+Antoniadis

- Large compactified extra dimensions: 1/R <~ 1 TeV
- SM fields are on a brane
- Gravity propagates in the whole bulk




y: extra dimension
R: compact'n radius

y=0 "our" brane (possibly with thickness r)

$G_N \sim 1/M^2_{Pl}$: Newton const.

$M_{Pl}$ large as $G_N$ weak

The idea is that gravity appears weak as a lot of lines of force escape in extra dimensions

## Generic feature of extra dim. models:

compact dim. → Kaluza-Klein (KK) modes

$p=n/R$ $m^2=n^2/R^2$ (quantization in a box)

Many possibilities:

- SM fields on a brane or in bulk

cfr: • Gravity always on bulk

- Factorized metric:

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu + h_{ij}(y) dy^i dy^j$$

- Warped metric: Randall-Sundrum (R-S)

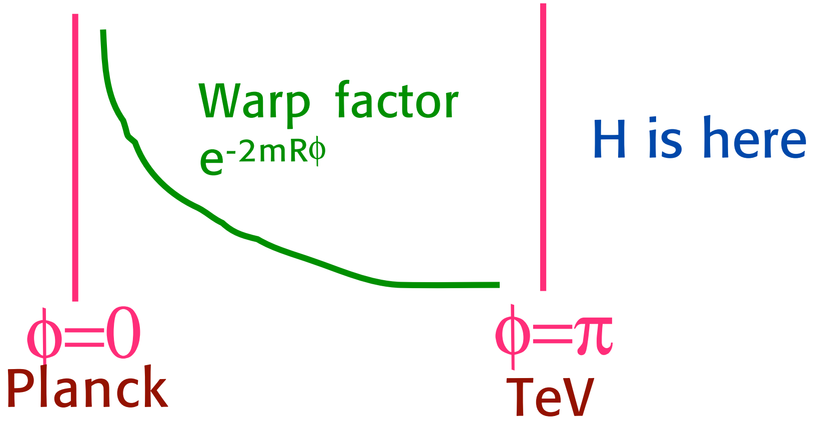
$$ds^2 = e^{-2mR|\varphi|} \eta_{\mu\nu} dx^\mu dx^\nu - R^2 \varphi^2$$

$m_{weak}=M_{Pl}\exp(-mR\pi)$ → $Rm \sim 12$

emerges as the most promising formulation

## A more promising scheme (warped ED)

$$ds^2 = e^{-2mR|\varphi|}\eta_{\mu\nu}dx^\mu dx^\nu - R^2\varphi^2$$




Randall-Sundrum:

This non-fact.ble metric is solution of Einstein eq.s with 2 branes at $\phi=0,\pi$ and specified 5-dim cosmological term

All SM particles in bulk except the H

$m \sim M_{Pl}$ for all mR: $m^2 \sim M_{Pl}^2(1-e^{-2mR\phi})$

All 4-dim masses $m_4$ are scaled down with respect to 5-dim masses $m_5 \sim M_{Pl}$ by the warp factor: $m_4 = M_{Pl}e^{-mR\pi}$

The hierarchy problem demands that $mR \sim 12$: not too large!!

R not large in this case!

Stabilization of mR at a compatible value can be assured by a scalar field in the bulk with a suitable potential

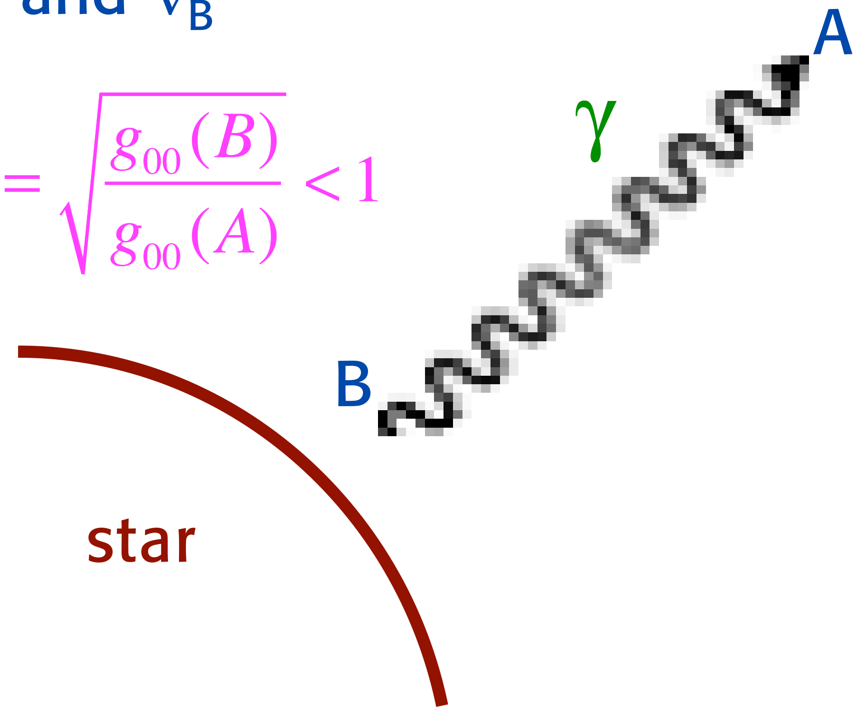
→ "radion"

Goldberger, Wise

2 identical atoms in A and B emit light with frequencies $\nu_A$ and $\nu_B$

seen from A the B frequency is smaller: as if the photon kinetic energy lost by climbing out of grav. field

$$\frac{\nu_B}{\nu_A} = \sqrt{\frac{g_{00}(B)}{g_{00}(A)}} < 1$$




Similarly in RS $mc^2$ is smaller by the corresponding factor $g_{00}{}^{1/2}$--> $m_4 = M_{Pl} e^{-mR\pi}$

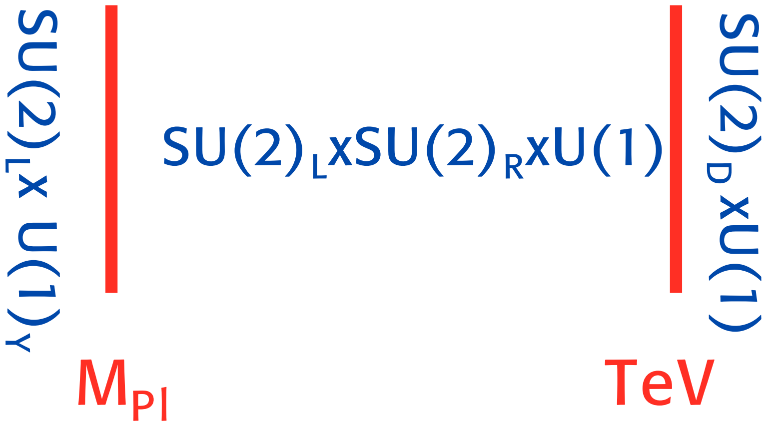
Good tutorials:
R. Sundrum '04 TASI lectures
R. Rattazzi '05 Cargese Lectures
Csaki et al '05
Gherghetta'06

# Applications

- Gauge Symmetry Breaking (Higgsless theories)

Csaki et al/Grojean/Papucci/Nomura/Davoudiasl et al/Barbieri, Pomarol, Rattazzi;....

The only models were no Higgs would be found at LHC.
But signals of new physics would be observed




Warped R-S background

Symmetries broken by Boundary Conditions (BC) on the branes

Altogether only $U(1)_Q$ unbroken

- Unitarity breaking (no Higgs) delayed by KK recurrences
- Dirac fermions on the bulk (L and R doublets). Only one chirality has a zero mode on the brane

With no Higgs unitarity violations, eg:

$$A\left(W_L^+ W_L^- \rightarrow Z_L Z_L\right) = \frac{G_F E^2}{8\sqrt{2}\pi}$$

At E ~ 1.2 TeV unitarity is violated

In Higgsless models unitarity breaking is delayed by the exchange of KK recurrences

$Z_k = k_{th}$ KK

Cancellation guaranteed by sum rules implied by 5-dim symmetry

$$g^2_{WWWW} - e^2 - \sum_k g^2_{WWZ_k} = 0\ ;$$

$$4M^2_W\, g^2_{WWWW} - 3\sum_k g^2_{WWZ_k} M_{Z^2_k} = 0\ .$$

The small W, Z mass implies a small KK gap --> W', Z' at the LHC

---

⊕ Higgsless models can also be formulated in 4 dimensions (pioneered by Casalbuoni, De Curtis, Dominici, Gatto '85)

No convincing, realistic Higgsless model for EW symmetry breaking emerged so far

Serious problems with EW precision tests

e.g. Chivukula et al ‘02, Barbieri, Pomarol, Rattazzi '03;

also with Z->bb

$m_W$ fixes the KK gap and it is not sufficiently large

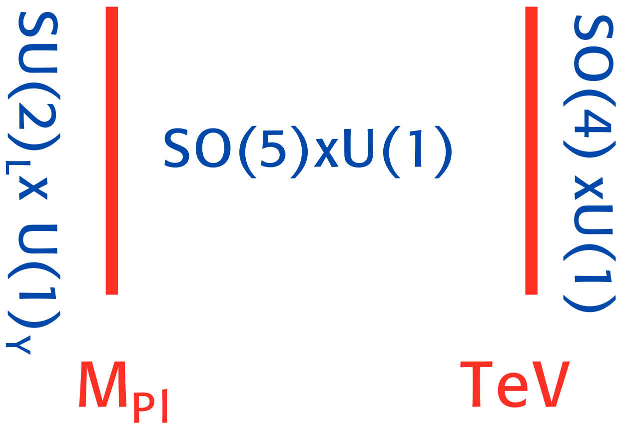
Substantial fine tuning required

Best try: Cacciapaglia et al '06

However be alerted of possible signals at the LHC: no Higgs but KK recurrences of W, Z and/or additional gauge bosons

• Composite Higgs in a 5-dim holographic theory

Agashe, Contino, Pomarol......




A new way to look at walking technicolor using AdS/CFT corresp.

All SM fields in the bulk (but the Higgs is localised on the TeV brane)

Warped R-S background

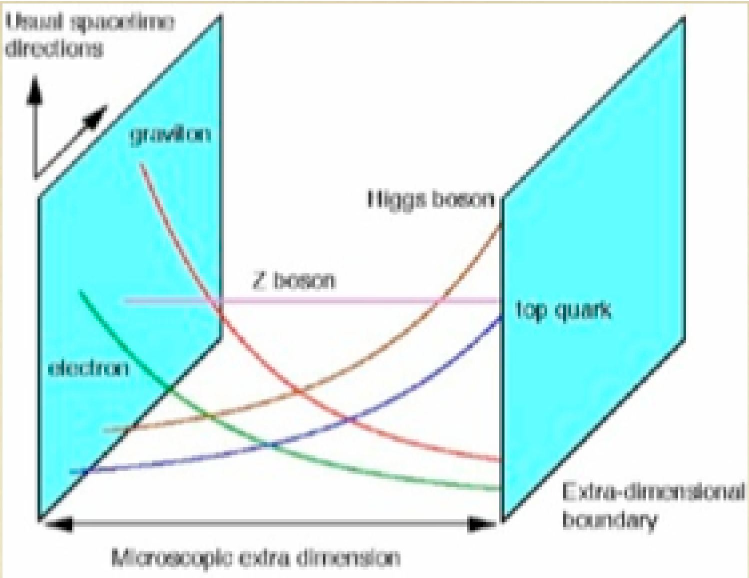
As in Little Higgs models

The Higgs is a PGB and EW symmetry breaking is triggered by bulk effects (in 4-dim the bulk appears as a strong sector).

The 5-dim theory is weakly coupled so that the Higgs potential and EW observables can be computed

The Higgs is rather light: $m_H < 185$ GeV

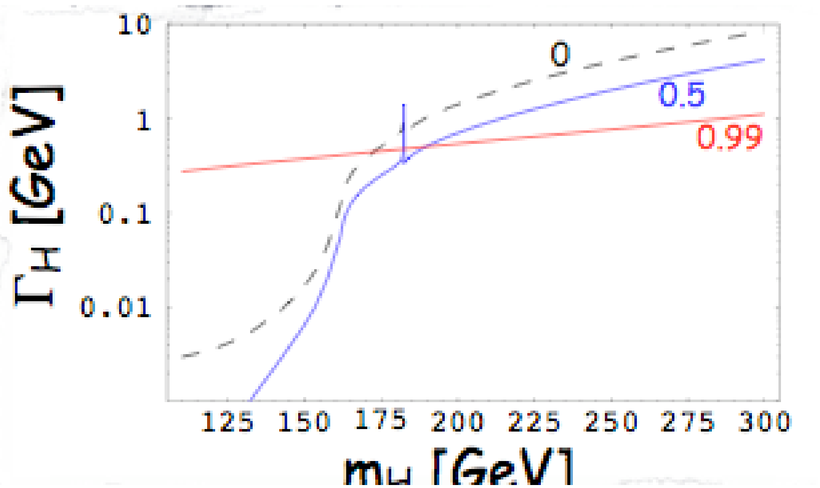
⊕ Also in these models a sizable fine-tuning is required

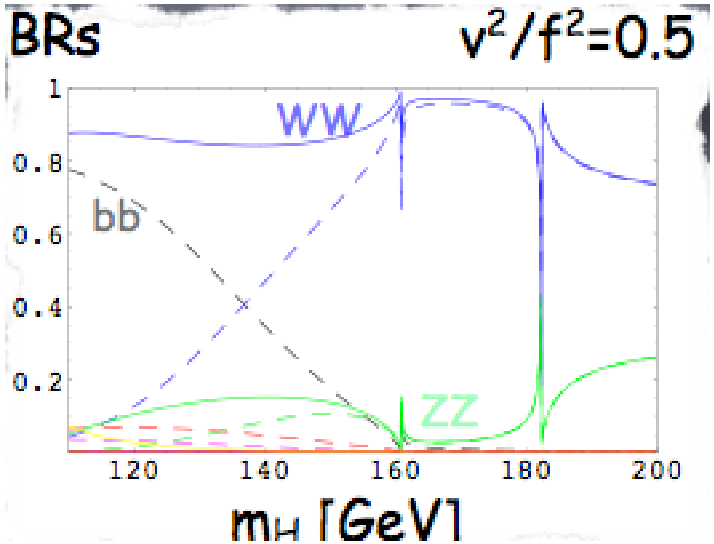



A qualitative description of flavour

The fermion mass hierarchies explained by exp warp factors with o(1) exponents

Higgs couplings modified



# Composite Higgs: a more model independent approach

Georgi, Kaplan '84

The light Higgs is a bound state of a strongly interacting sector.
Pseudo-Goldstone boson of an enlarged symmetry.
eg. SO(5)/SO(4)

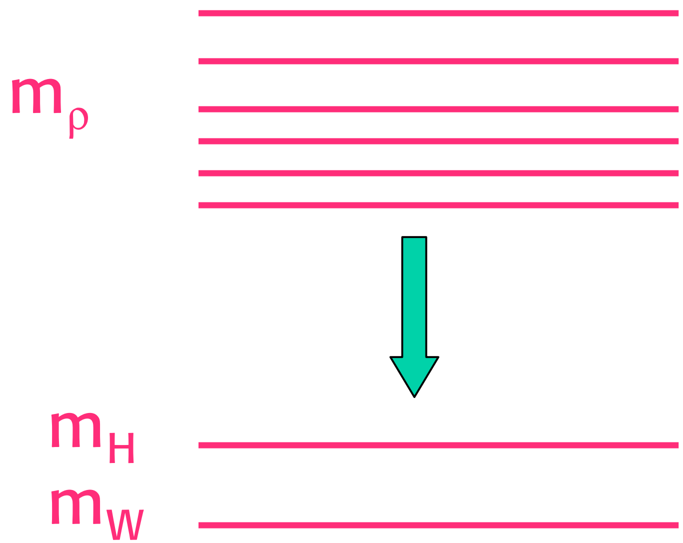
Agashe/ Contino/Pomarol/Sundrum/ Grojean/Rattazzi....

$v \sim$ EW scale $\qquad f \sim$ SI scale

$\sim f < m_\rho <\sim 4\pi f$

$\xi = (v/f)^2$

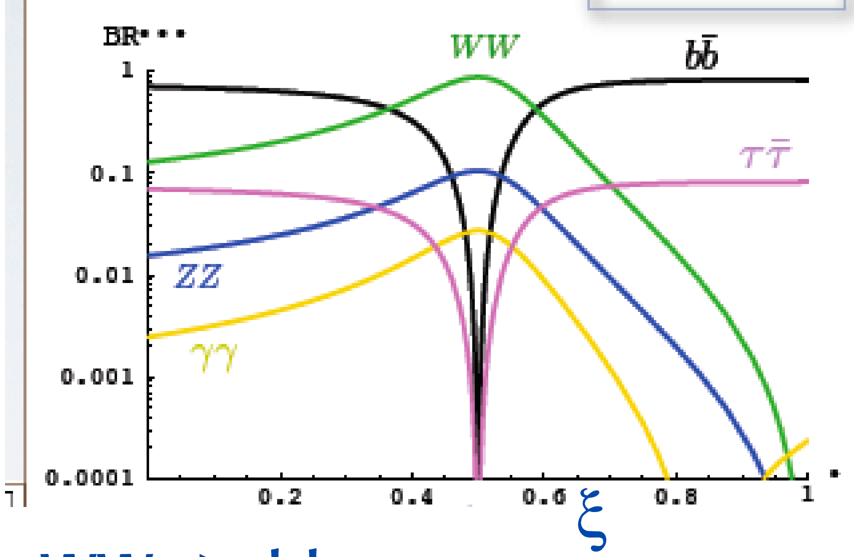
$\xi$ interpolates between SM [$\xi \sim 0$]
and some degree of
compositeness
[$\xi \sim o(1)$ limited by precision EW tests]



# Detectable ξ effects at the LHC

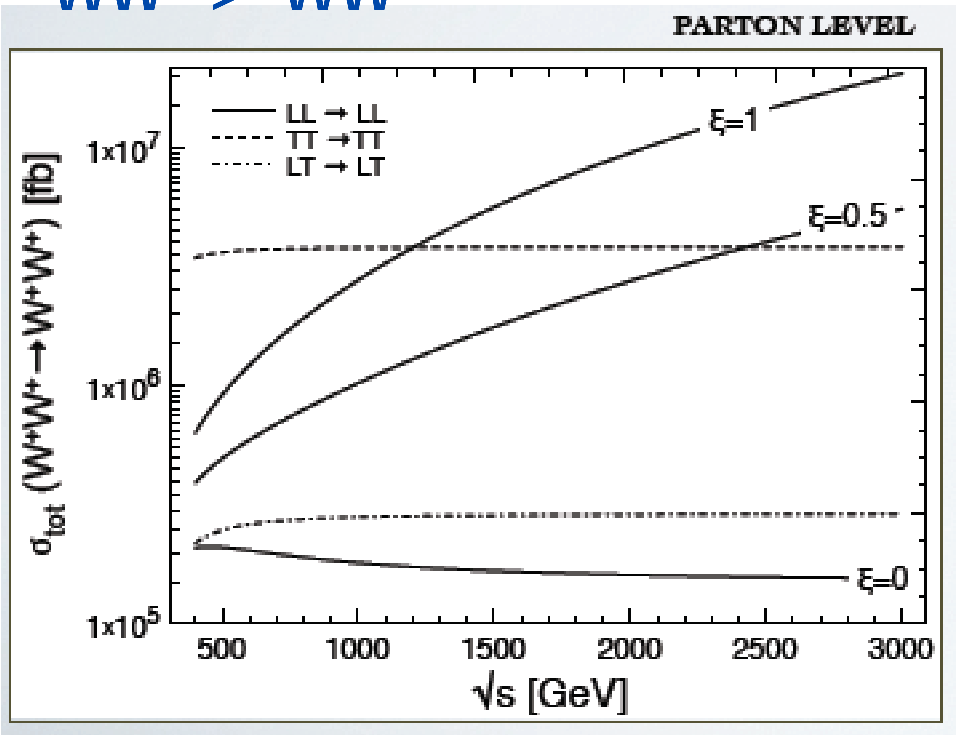
- Higgs couplings
- WW scattering
- 2-Higgs Production

## H Br Ratios

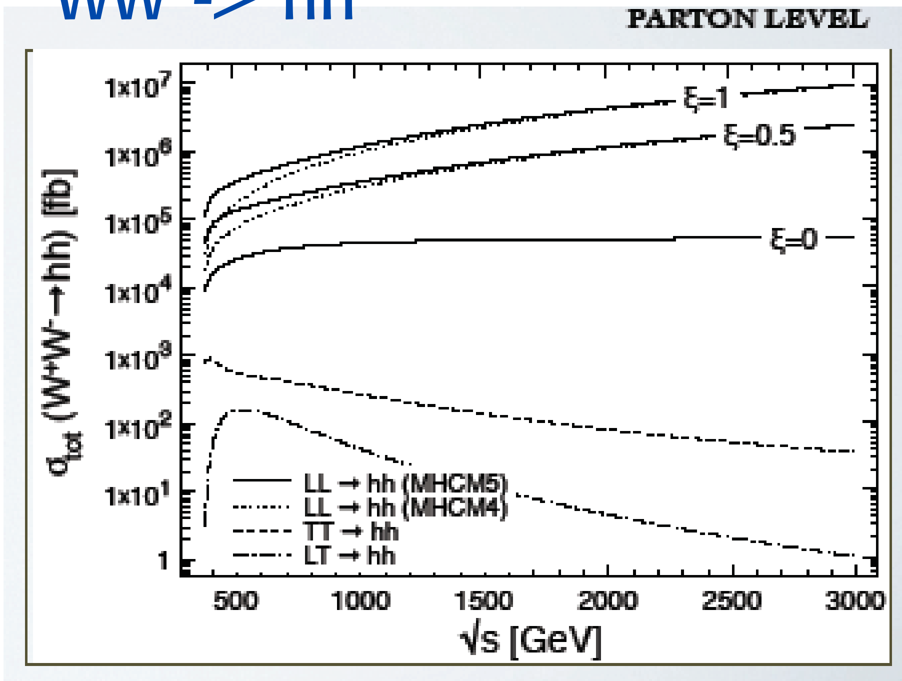



Contino et al

## WW -> WW




## WW -> hh



## Lessons from model building

In all the new physics models we mentioned

there is a light Higgs ($< 200$ GeV)

[except in Higgsless models (if any) but new light new vector bosons exist in this case]

there is at least a % fine tuning

Fine tuning appears to be imposed on us by the data

# In conclusion

Is it possible that the LHC does not find the Higgs particle?

Yes, it is possible, but then something else must be found

Is it possible that the LHC finds the Higgs particle but no other new physics (pure and simple SM)?

Yes, it is technically possible but it is not natural

Is it possible that the LHC finds neither the Higgs nor new physics?

No, it is "approximately impossible"