

On the Origin of Neutrino Masses

Pavel Fileviez Perez

University of Wisconsin-Madison, USA

NuFact09, IIT, Chicago



Collaborators

In the US:

V. Barger (Madison)
S. Blanchet (Maryland)
L. Everett (Madison)
T. Han (Madison)
G. Y. Huan (UC-Davis)
S. Spinner (Madison)
M. B. Wise (Caltech)

In Asia:

K. Wang (IPMU, Japan)
T. Li (Pekin Univ., China)
H. Iminniyaz (Pekin Univ., China)

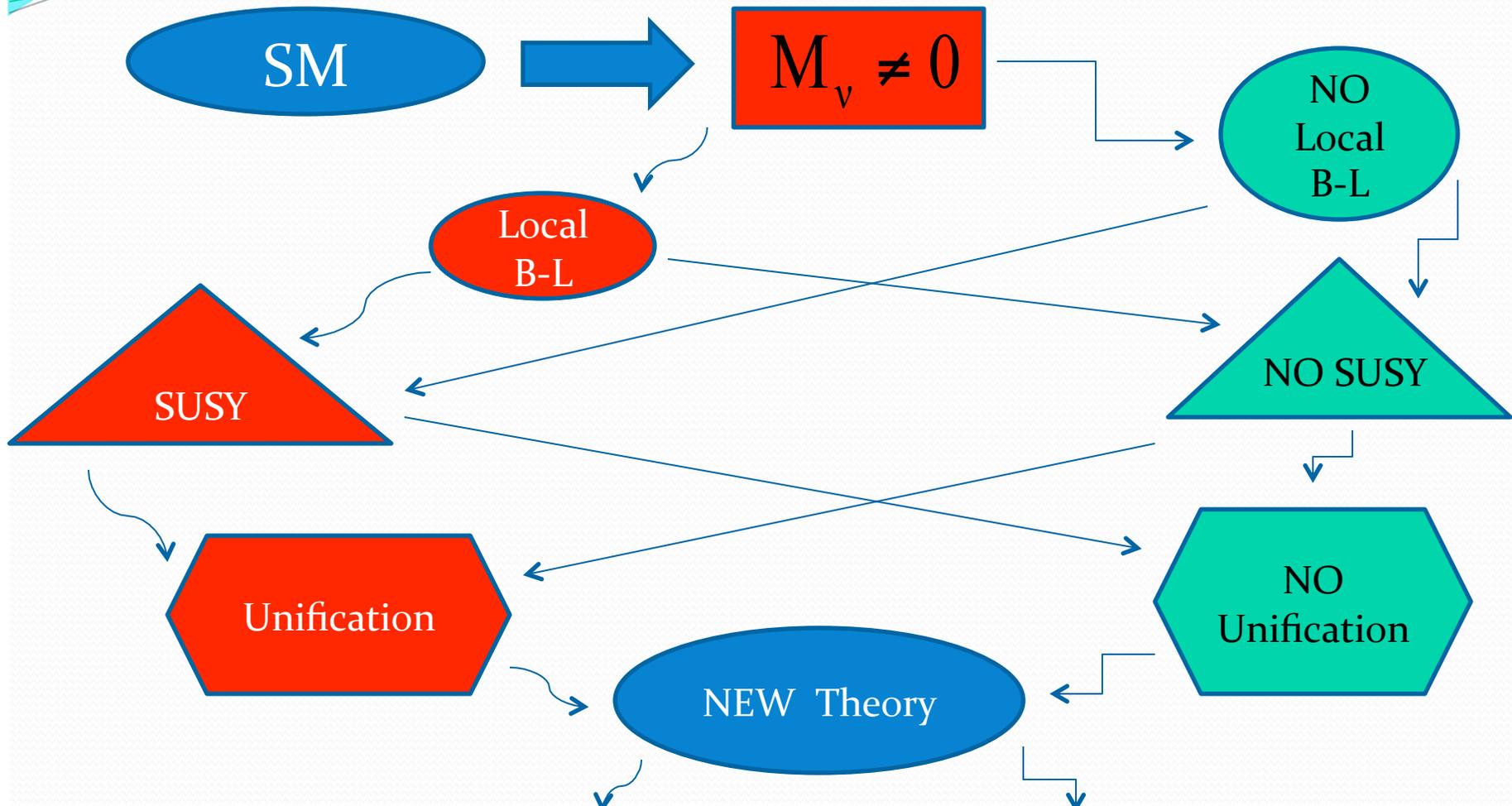
In Europe:

I. Dorsner (Italy)
R. Gonzalez-Felipe (Portugal)
G. Rodrigo (Spain)
G. Senjanovic (Italy)

In Latinamerica:

M. A. Diaz (Chile)

Aim



Predictions for the LHC, Cosmology, Proton Decay and others

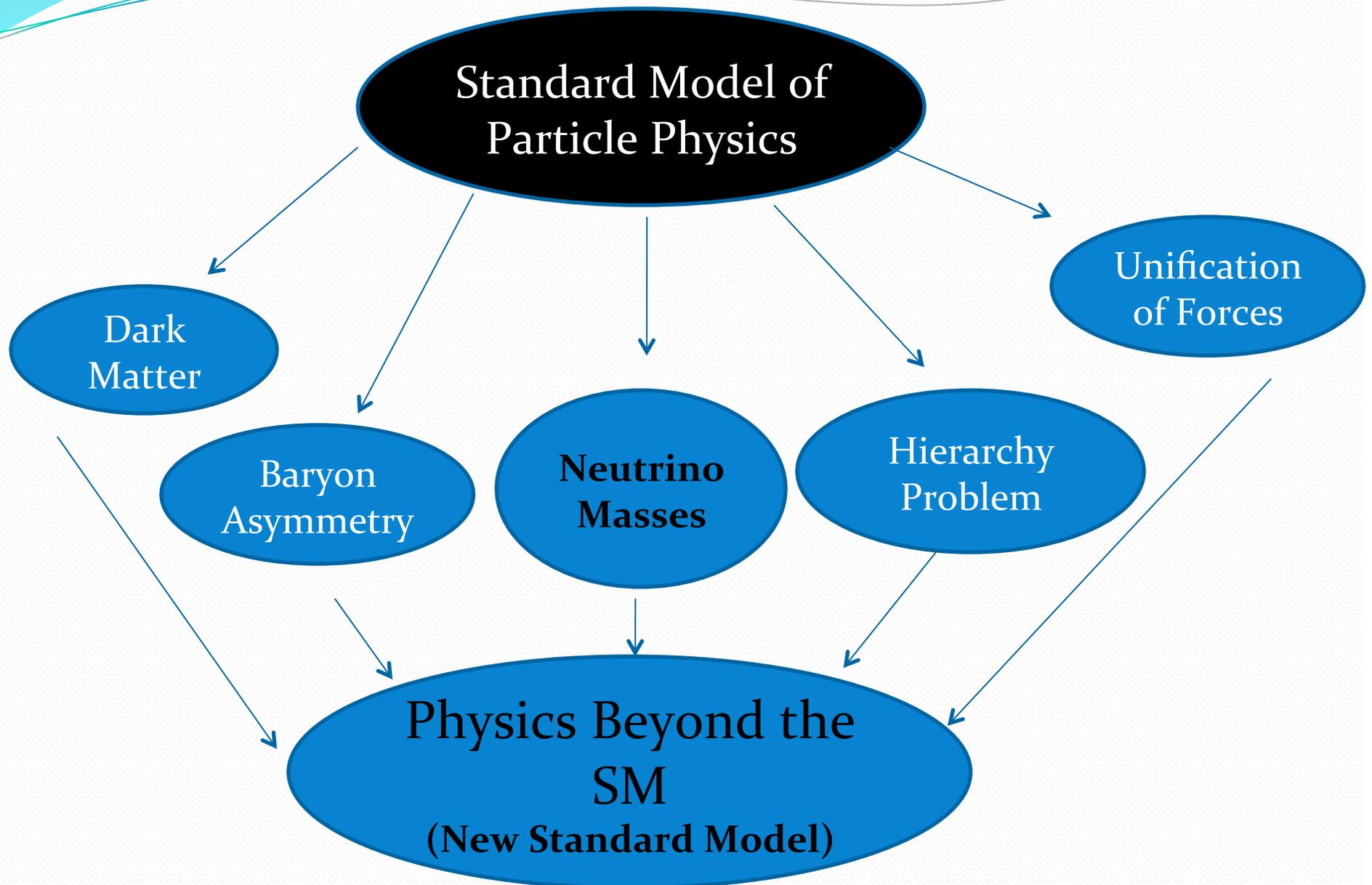


Outline

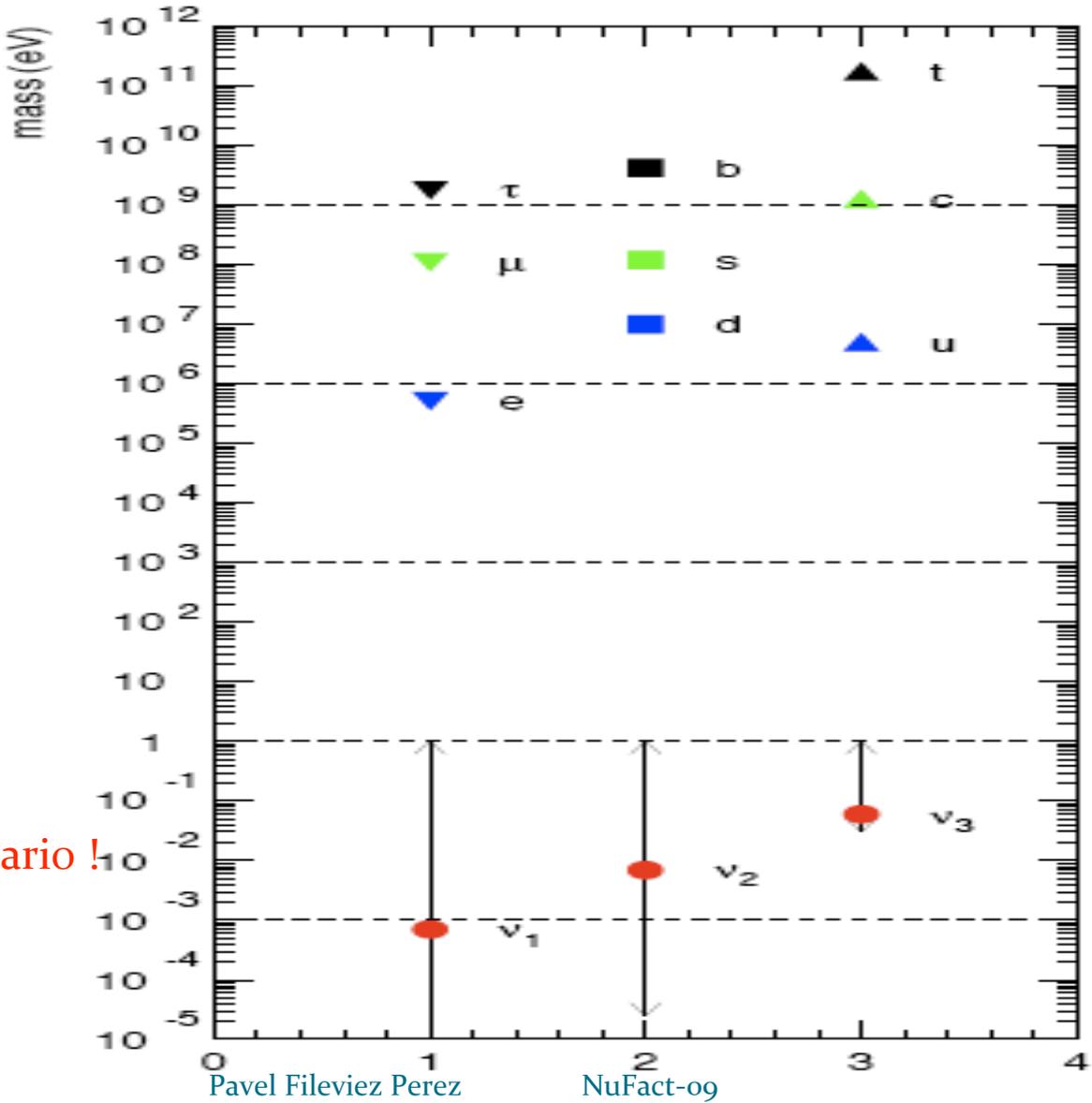
- Introduction
- Mechanisms for neutrino masses
- Grand unification and massive neutrinos
- Summary



INTRODUCTION



FERMIONIC SPECTRUM IN THE SM



A possible scenario !

Neutrino Properties

The leptonic mixing matrix is given by:

$$\begin{pmatrix} c_{12}c_{13} & c_{13}s_{12} & e^{-i\delta}s_{13} \\ -c_{12}s_{13}s_{23}e^{i\delta} - c_{23}s_{12} & c_{12}c_{23} - e^{i\delta}s_{12}s_{13}s_{23} & c_{13}s_{23} \\ s_{12}s_{23} - e^{i\delta}c_{12}c_{23}s_{13} & -c_{23}s_{12}s_{13}e^{i\delta} - c_{12}s_{23} & c_{13}c_{23} \end{pmatrix}$$

$$K_M = \text{diag}(e^{i\Phi_1/2}, 1, e^{i\Phi_2/2})$$

The experimental constraints are:

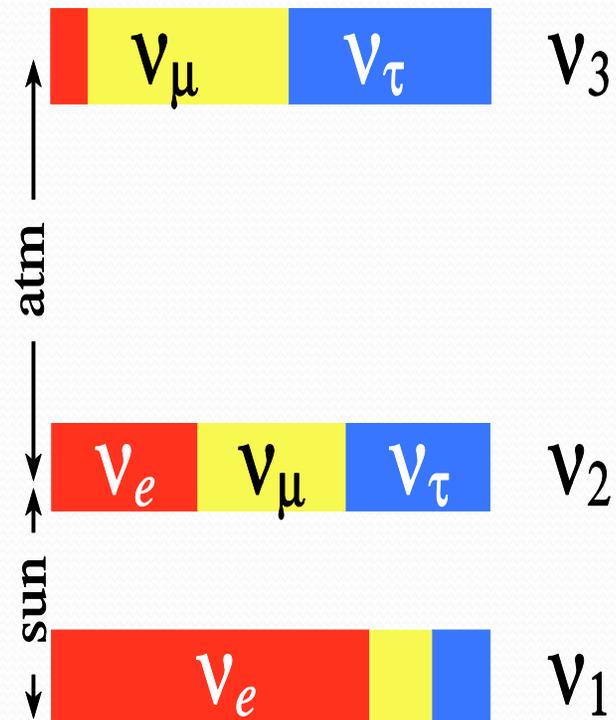
$$\Delta m_{21}^2 = (7.2 - 8.9) \times 10^{-5} \text{eV}^2$$

$$|\Delta m_{23}^2| = (2.1 - 3.1) \times 10^{-3} \text{eV}^2$$

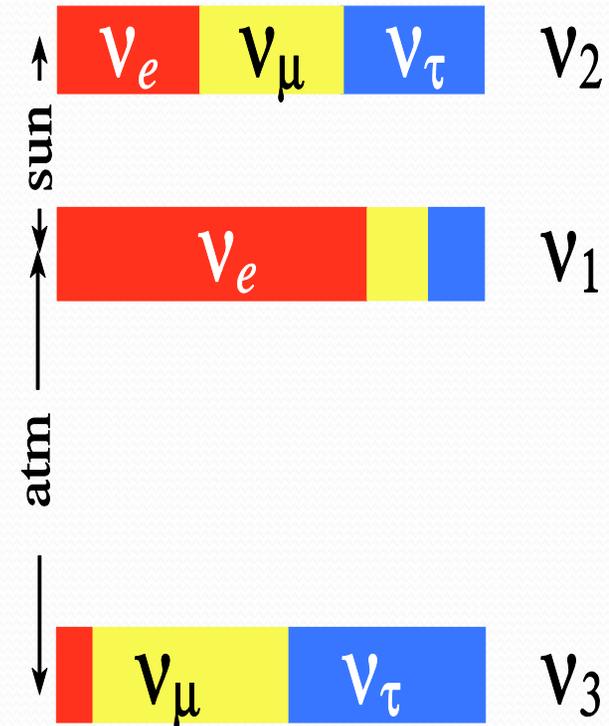
$$30^\circ < \theta_{12} < 38^\circ, \quad 36^\circ < \theta_{23} < 54^\circ, \quad \theta_{13} < 10^\circ$$

Neutrino Spectrum

Normal Hierarchy



Inverted Hierarchy



Neutrinos

In the SM the lepton number (or B-L) is an accidental symmetry and the neutrinos are massless !

Today we know that the neutrinos are massive and they can be:

Dirac Fermion: $\nu^C \sim (1, 1, 0)$ (L is conserved by hand!)

$$-\mathcal{L}_\nu^D = Y_\nu l H \nu^C + \text{h.c.} \implies M_\nu^D = Y_\nu v / \sqrt{2}$$

If $Y_\nu \sim 10^{-11}$ one gets $M_\nu \sim 1 \text{ eV}$.

Majorana Fermion: (L is broken) Weinberg'79

$$-\mathcal{L}_\nu^M = c_\nu (l H)^2 / \Lambda_\nu + \text{h.c.}$$

$$\implies M_\nu^M = c_\nu v^2 / 2\Lambda_\nu$$

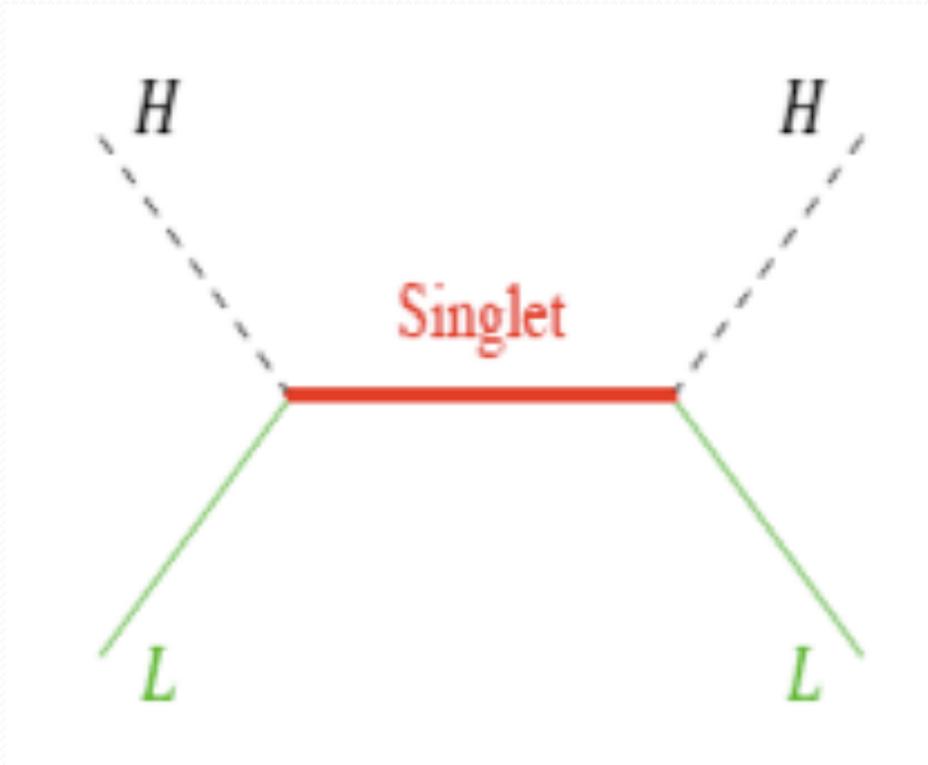
What is the origin of this operator?

If $c_\nu \sim 1$ and $\Lambda_\nu \sim 10^{14-15}$ GeV one gets $M_\nu \sim 1$ eV.



Mechanisms for Neutrino Masses

TYPE I SEESAW MECHANISM



Minkowski'77;
Yanagida'79; Gell-Mann, Ramond, Slansky'79;
Glashow'80; Mohapatra, Senjanovic'80

TYPE I SEESAW

Extra Fermions: $\nu^C \sim (1, 1, 0)$

$$-\mathcal{L}^I = Y_\nu l H \nu^C + \frac{1}{2} M \nu^C \nu^C + \text{h.c.}$$



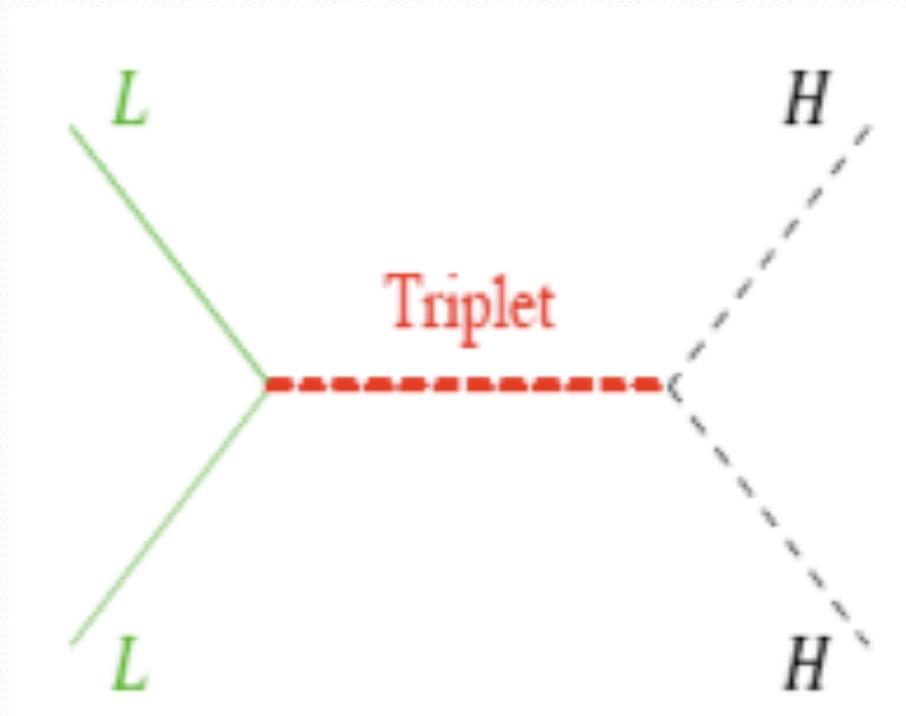
$$\mathcal{M}_\nu^I = \frac{v^2}{2} Y_\nu M^{-1} Y_\nu^T$$

$Y_\nu? \quad M?$

if $Y_\nu \sim 1$ and $M \sim 10^{14-15}$ GeV one has $\mathcal{M}_\nu^I \sim 1$ eV

M is defined by the $U(1)_{B-L}$ breaking scale

TYPE II SEESAW MECHANISM



Type II Seesaw

Konetschny, Kummer'77
Chen, Li'80; Lazarides, Shafi,
Wetterich'81; Schechter, Valle'80;
Mohapatra, Senjanovic'81

Extra Scalar: $\Delta \sim (1, 3, 1)$ $\Delta = \begin{pmatrix} \delta^+/\sqrt{2} & \delta^{++} \\ \delta^0 & -\delta^+/\sqrt{2} \end{pmatrix}$

$$\mathcal{L}_\nu^{II} = -Y_\nu l \Delta l + \mu H \Delta^\dagger H + \text{h.c.} + \dots$$

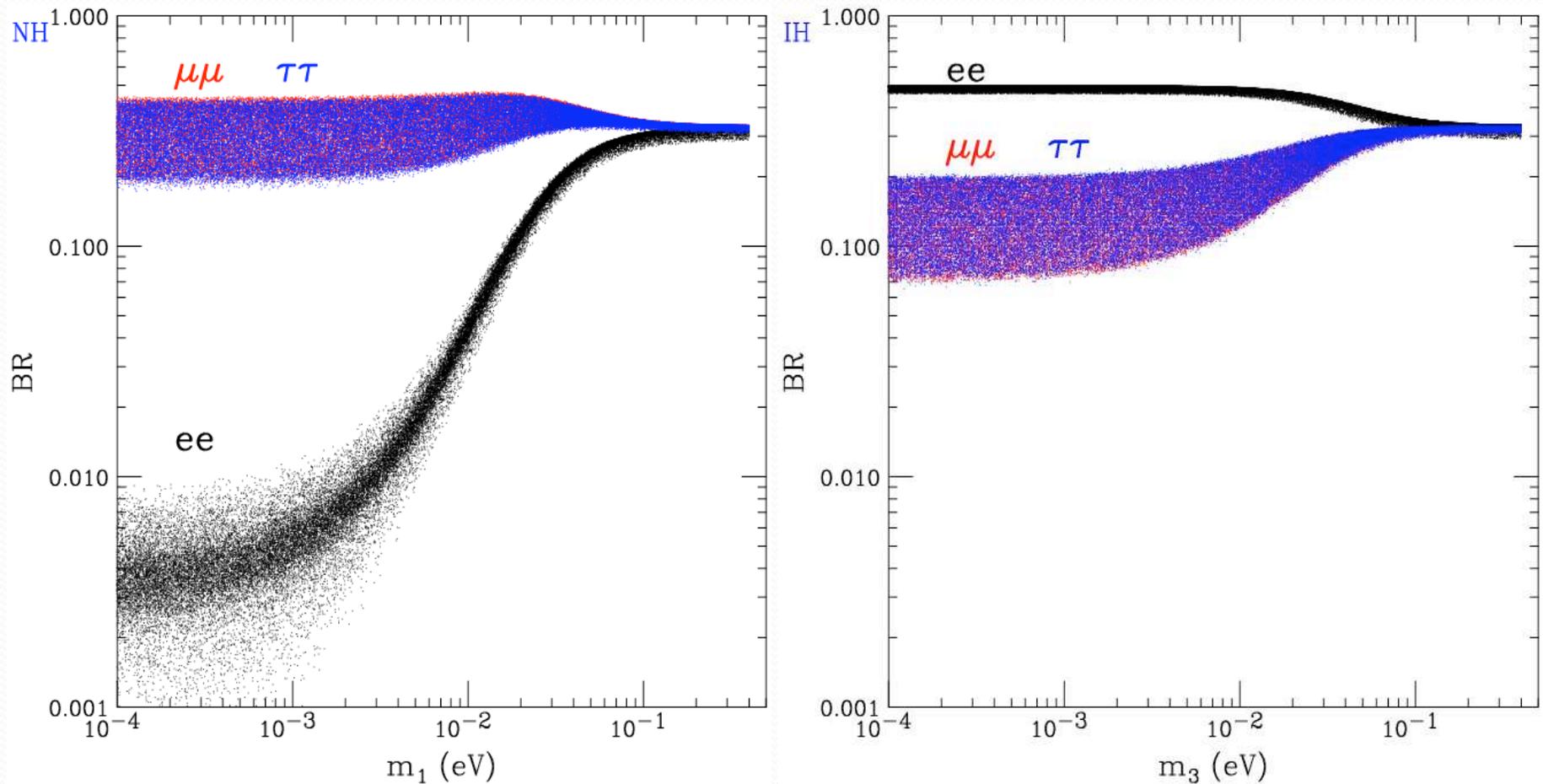
 $\mathcal{M}_\nu^{II} = Y_\nu \mu v^2 / M_\Delta^2$

$Y_\nu? \quad M_\Delta? \quad \mu?$

if $Y_\nu \sim 1$ and $v_\Delta \sim 1$ eV one gets the natural neutrino scale

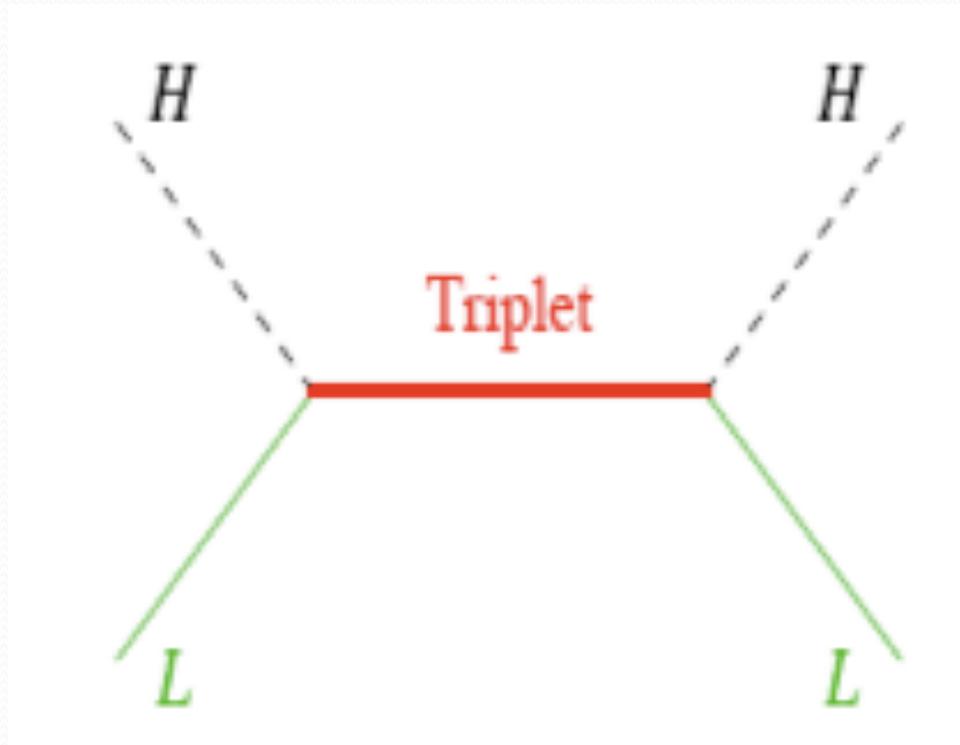
HIGGS DECAYS AND NEUTRINO MASSES

$$H^{++} \rightarrow e_i^+ e_j^+$$



See Talk by F. de Aguilá. See also Garayoa, Schwetz; F. de Aguilá et al, and others

TYPE III SEESAW MECHANISM



Type III Seesaw

Extra Fermions: $\rho \sim (1, 3, 0)$ $\rho = \begin{pmatrix} \rho^0/\sqrt{2} & \rho^+ \\ \rho^- & -\rho^0/\sqrt{2} \end{pmatrix}$

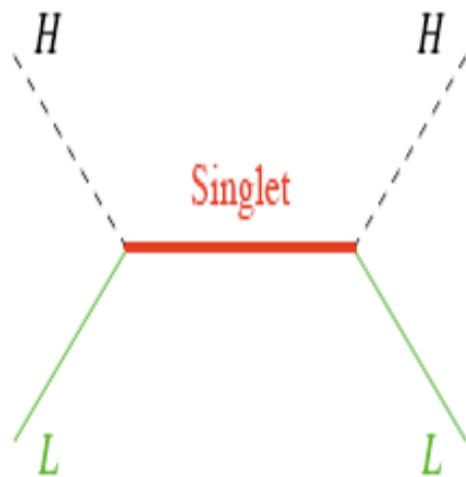
$$-\mathcal{L}^{III} = Y_\nu l \rho H + M_\rho \text{Tr} \rho^2 + \text{h.c.}$$

 $\mathcal{M}_\nu^{III} = v^2 Y_\nu M_\rho^{-1} Y_\nu^T / 2$

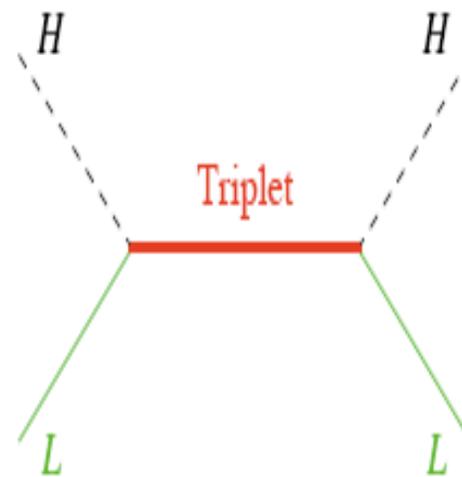
$Y_\nu? \quad M_\rho?$

If $Y_\nu \sim 1$ and $M_\rho \sim 10^{14-15}$ GeV one has $M_\nu^{III} \sim 1$ eV.

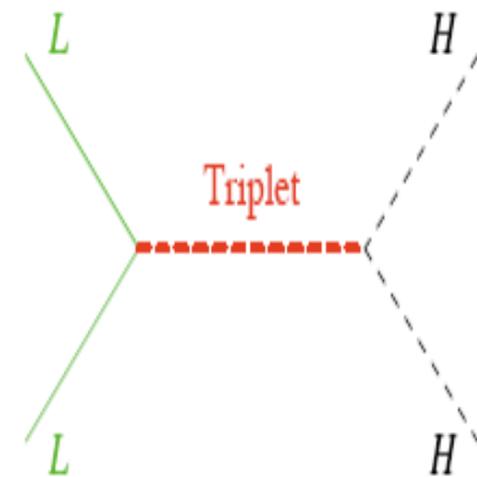
Seesaw Mechanisms at Tree Level



Type I Seesaw



Type III Seesaw



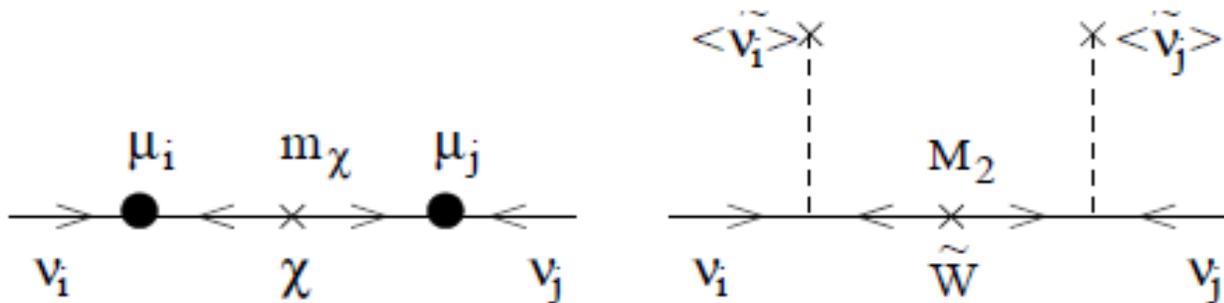
Type II Seesaw

R-Parity Violation in SUSY

In the MSSM:

$$R = (-1)^{3(B-L)+2S}$$

$$W_{RpV} = \epsilon_i \hat{L}_i \hat{H}_u + \lambda_{ijk} \hat{L}_i \hat{L}_j \hat{E}_k^C + \lambda'_{ijk} \hat{Q}_i \hat{L}_j \hat{D}_k^C + \lambda''_{ijk} \hat{U}_i^C \hat{D}_j^C \hat{D}_k^C$$



We need a mechanism for spontaneous R-parity violation !!!!

V. Barger, **P.F.P.**, S. Spinner, *Phys. Rev. Lett.* 102:181802, 2009

Local B-L and Spontaneous R-Parity Violation

$$SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \otimes U(1)_{B-L}$$

Matter: \hat{Q} \hat{L} \hat{U}^C \hat{D}^C \hat{E}^C

$$\hat{N}^C \sim (1, 1, 0, 1) \quad (\text{for anomaly cancellation})$$

-Some R-Parity Violating couplings

$$\frac{1}{2} g_{BL} v_R \nu^C \tilde{B}'$$

$$\frac{1}{2} g_2 v_L \nu \tilde{W}^0$$

$$\frac{1}{2} g_1 v_L \nu \tilde{B}$$

Neutrino-Neutralino Mass Matrix

$$(\nu, \nu^c, \tilde{\chi}^0) \longrightarrow \mathcal{M}_N = \begin{pmatrix} 0 & M_\nu^D & \Gamma \\ (M_\nu^D)^T & 0 & G \\ \Gamma^T & G^T & M_{\tilde{\chi}^0} \end{pmatrix}$$

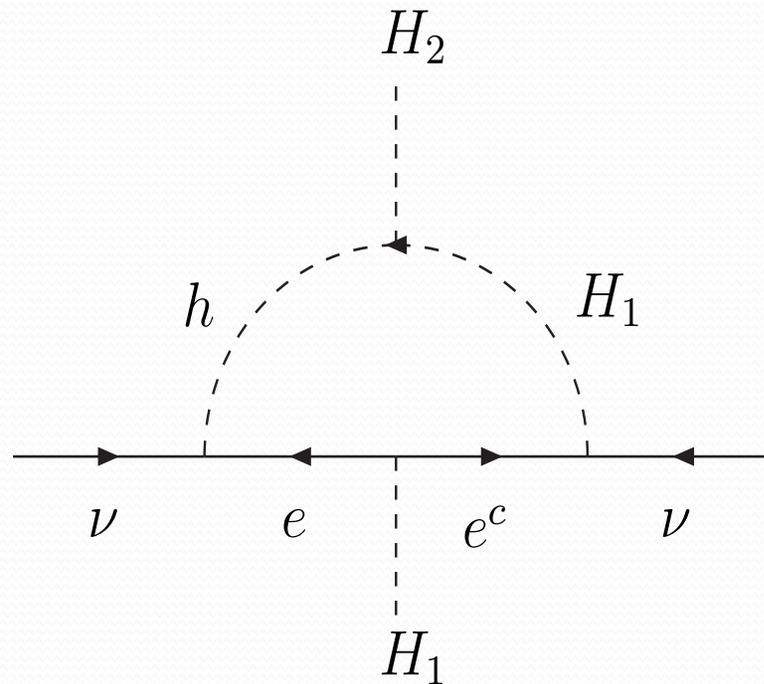
Type I Seesaw and RpV:
$$M_\nu = (\tilde{M}_\nu^D) M_{\nu^c}^{-1} (\tilde{M}_\nu^D)^T + \Gamma M_{\tilde{\chi}^0}^{-1} \Gamma^T$$



MECHANISMS AT ONE-LOOP LEVEL

Zee Model

A. Zee'80



The Zee-Wolfenstein model is ruled out !

The general model still alive. See He'03

Zee Model

A. Zee'80

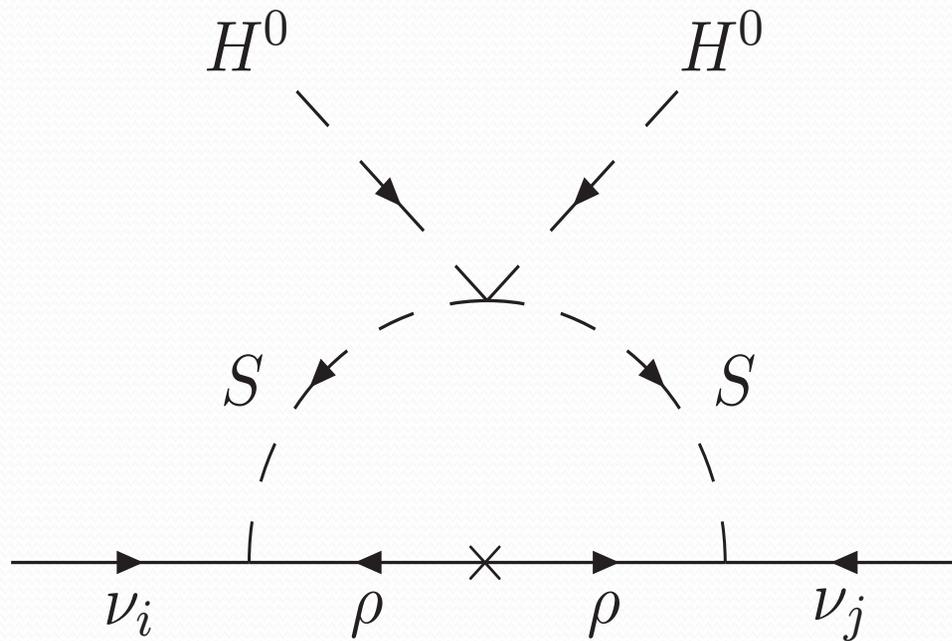
Extra Higgses: $h \sim (1, 1, 1)$ $H_2 \sim (1, 2, 1/2)$

$$-\mathcal{L}_{Zee} = Y l l h + \mu H_1 H_2 h^\dagger + Y_1 e^C H_1^\dagger l + Y_2 e^C H_2^\dagger l + \text{h.c.}$$

$$Y = -Y^T$$

$$\mathcal{O}_5^\nu = c_{11} l l H_1^2 + c_{12} l l H_1 H_2 + c_{22} l l H_2^2$$

COLORED MECHANISM



COLORED MECHANISM

$$-\mathcal{L} = Y_2 l S_1 \rho_1 + M_{\rho_1} \text{Tr} \rho_1^2 + \lambda_2 \text{Tr} \left(S_1^\dagger H \right)^2 + \text{h.c.}$$

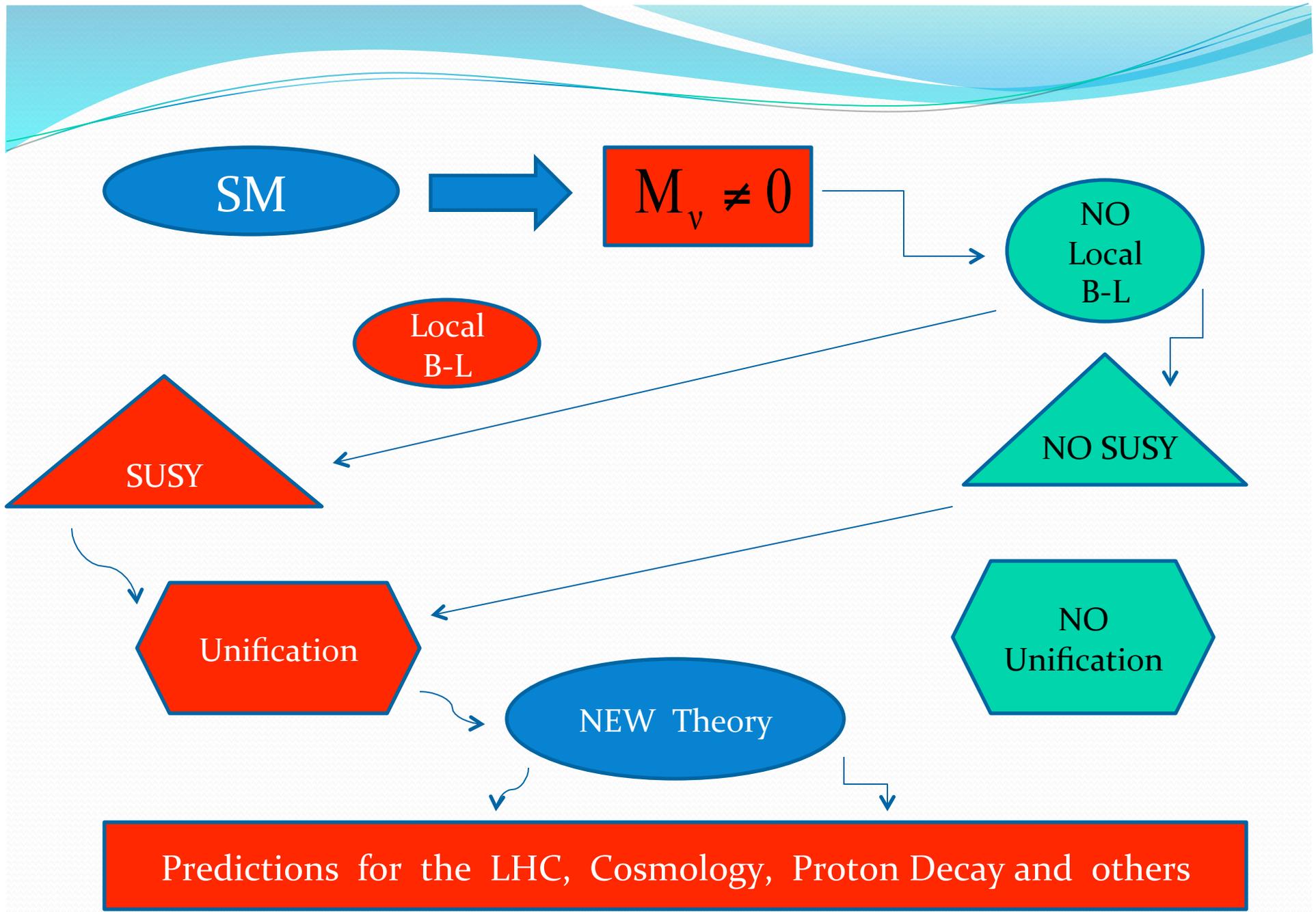
$$1) \rho_1 \sim (8, 1, 0), \quad S_1 \sim (8, 2, 1/2)$$

$$\mathcal{M}_\nu^{ij} = Y_2^{i\alpha} Y_2^{j\alpha} \frac{\lambda_2}{4\pi^2} v^2 \frac{M_{\rho_1^\alpha}}{M_{S_1}^2}.$$

$$M_{\rho_1} = 200 \text{ GeV}, v = 246 \text{ GeV} \text{ and } M_{S_1} = 2 \text{ TeV} \quad \Rightarrow \quad Y_2^2 \lambda_2 \sim 10^{-8}$$



Grand Unification and Massive Neutrinos





Neutrino Masses and Renormalizable SU(5)

Georgi-Glashow Model

Symmetry: $SU(5)$

Gauge Bosons: $A_\mu = (G_\mu, W_\mu, B_\mu, V_\mu, \bar{V}_\mu)$

Fermions:

$$10 = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & u_3^C & -u_2^C & u_1 & d_1 \\ -u_3^C & 0 & u_1^C & u_2 & d_2 \\ u_2^C & -u_1^C & 0 & u_3 & d_3 \\ -u_1 & -u_2 & -u_3 & 0 & e^C \\ -d_1 & -d_2 & -d_3 & -e^C & 0 \end{pmatrix}, \bar{5} = \begin{pmatrix} d_1^C \\ d_2^C \\ d_3^C \\ e \\ -\nu \end{pmatrix}$$

Higgs Sector:

$$5_{\text{H}} = \begin{pmatrix} T_1 \\ T_2 \\ T_3 \\ H^+ \\ H^0 \end{pmatrix}, \quad 24_{\text{H}} = \begin{pmatrix} \Sigma_8 & \Sigma_{(3,2)} \\ \Sigma_{(\bar{3},2)} & \Sigma_3 \end{pmatrix} + \frac{1}{2\sqrt{15}} \begin{pmatrix} 2 & 0 \\ 0 & -3 \end{pmatrix} \Sigma_{24}$$

Ruled out by unification and/or fermion masses !

Why the Georgi-Glashow model is ruled out ?

- Unification of gauge couplings in disagreement with the values of α_s , $\sin \theta_W$ and α_{em} at the electroweak scale
- $M_E = M_D^T$ at the GUT scale
- $M_\nu = 0$

$$M_D = M_E^T$$



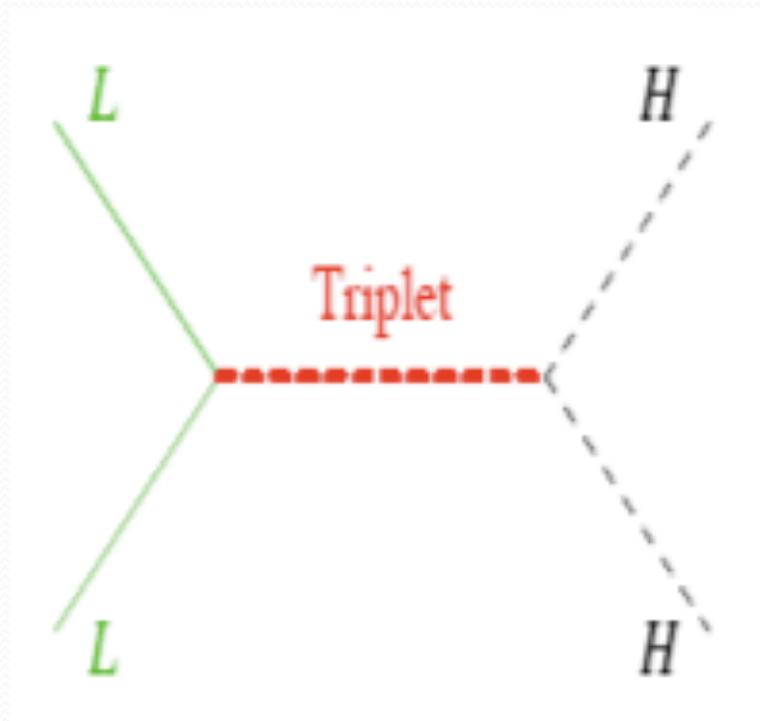
- 45_H
- h.d.ops

Neutrino Masses:

$$M_\nu \neq 0$$

- Type I seesaw: $1_i, i=1,2,\dots$
- Type II seesaw: $\Delta \subset 15_H \quad \left(15_H \oplus \overline{15}_H \right)$
- Type III seesaw: $\rho \subset 24$

Type II Seesaw



$$\Delta \subset 15_H$$

Type II – SU(5)

Higgs Sector: $5_H, 24_H, 15_H$

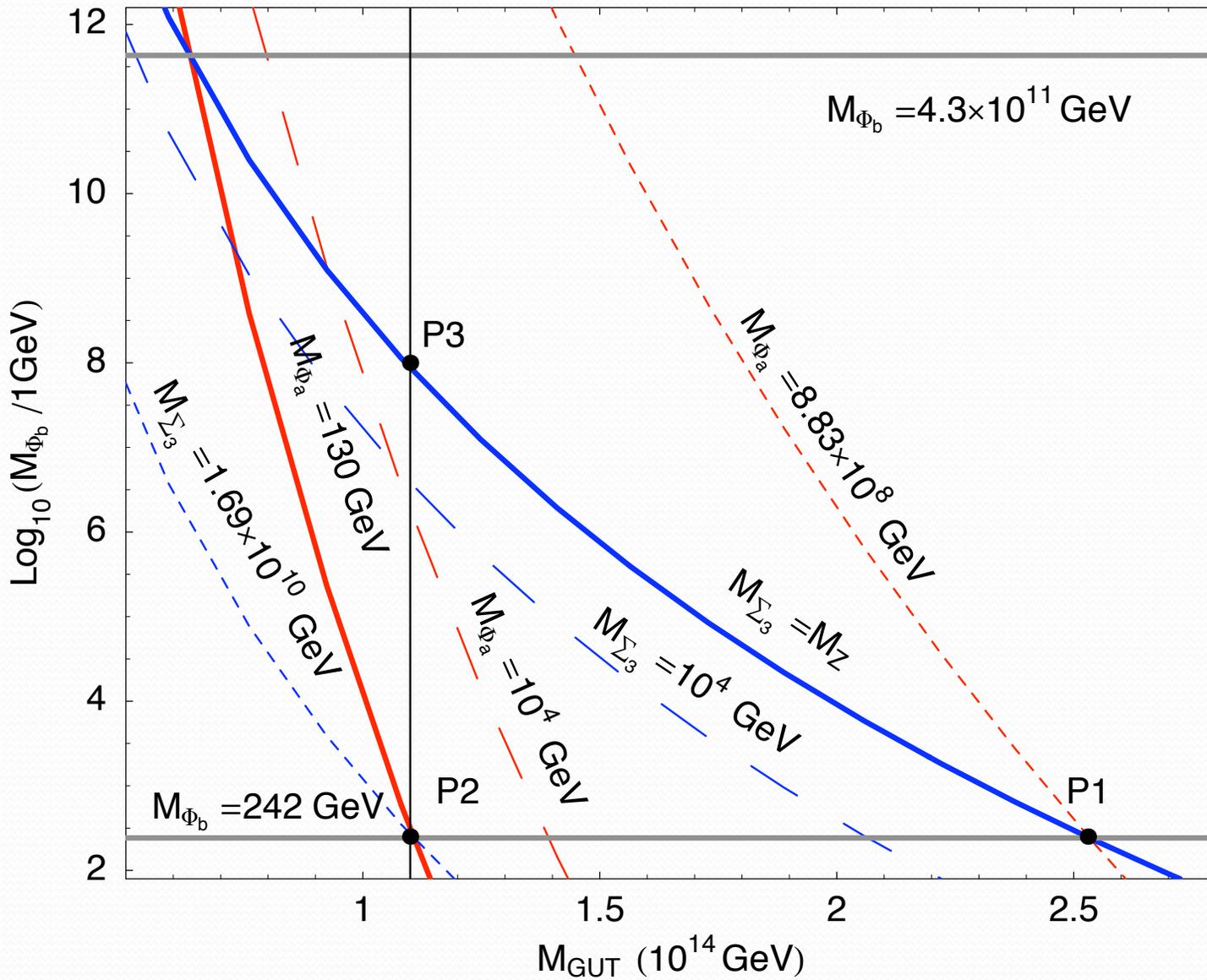
$$15_H = \underbrace{(1, 3, 1)}_{\Phi_a} \oplus \underbrace{(3, 2, 1/6)}_{\Phi_b} \oplus \underbrace{(6, 1, -2/3)}_{\Phi_c}$$

$$\Phi_a = i\sigma_2 \Delta$$

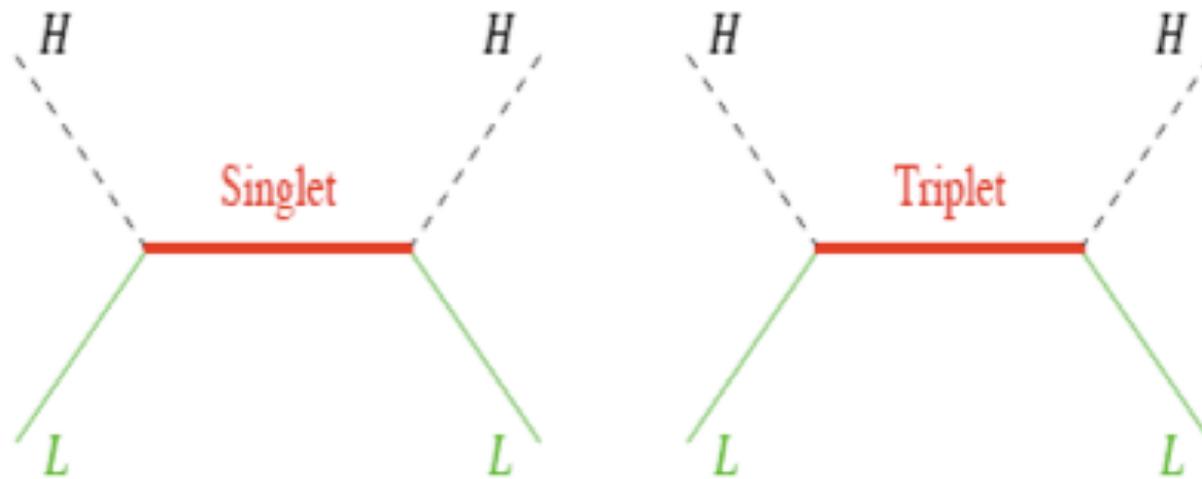
Neutrino Masses: Type II seesaw mechanism

$$V_\nu = Y_\nu \bar{5} \bar{5} 15_H + \mu 5_H^* 5_H^* 15_H + \text{h.c.}$$

$$M_\nu = Y_\nu \mu v^2 / M_\Delta^2$$



Type I plus Type III Seesaw



Type I Seesaw

Type III Seesaw

Adjoint SU(5)

P. Fileviez Pérez, PLB 654 (2007) 189.

Matter: $\bar{5} = (d^C, e, \nu)$, $10 = (u^C, Q, e^C)$, 24

Higgs Sector: $5_H, 24_H, 45_H$

$$24 = \underbrace{(8, 1)}_{\rho 8} \oplus \underbrace{(1, 3)}_{\rho 3} \oplus \underbrace{(3, 2)}_{\rho(3,2)} \oplus \underbrace{(\bar{3}, 2)}_{\rho(\bar{3},2)} \oplus \underbrace{(1, 1)}_{\rho 0}$$

Neutrino Mass: Type I and Type III seesaw

$$V_\nu = \alpha_i \bar{5}_i 24 5_H + \rho_i \bar{5}_i 24 45_H$$

See also: P. Fileviez Pérez, PRD 76 (2007) 071701.

Neutrino Masses in Adjoint SU(5)

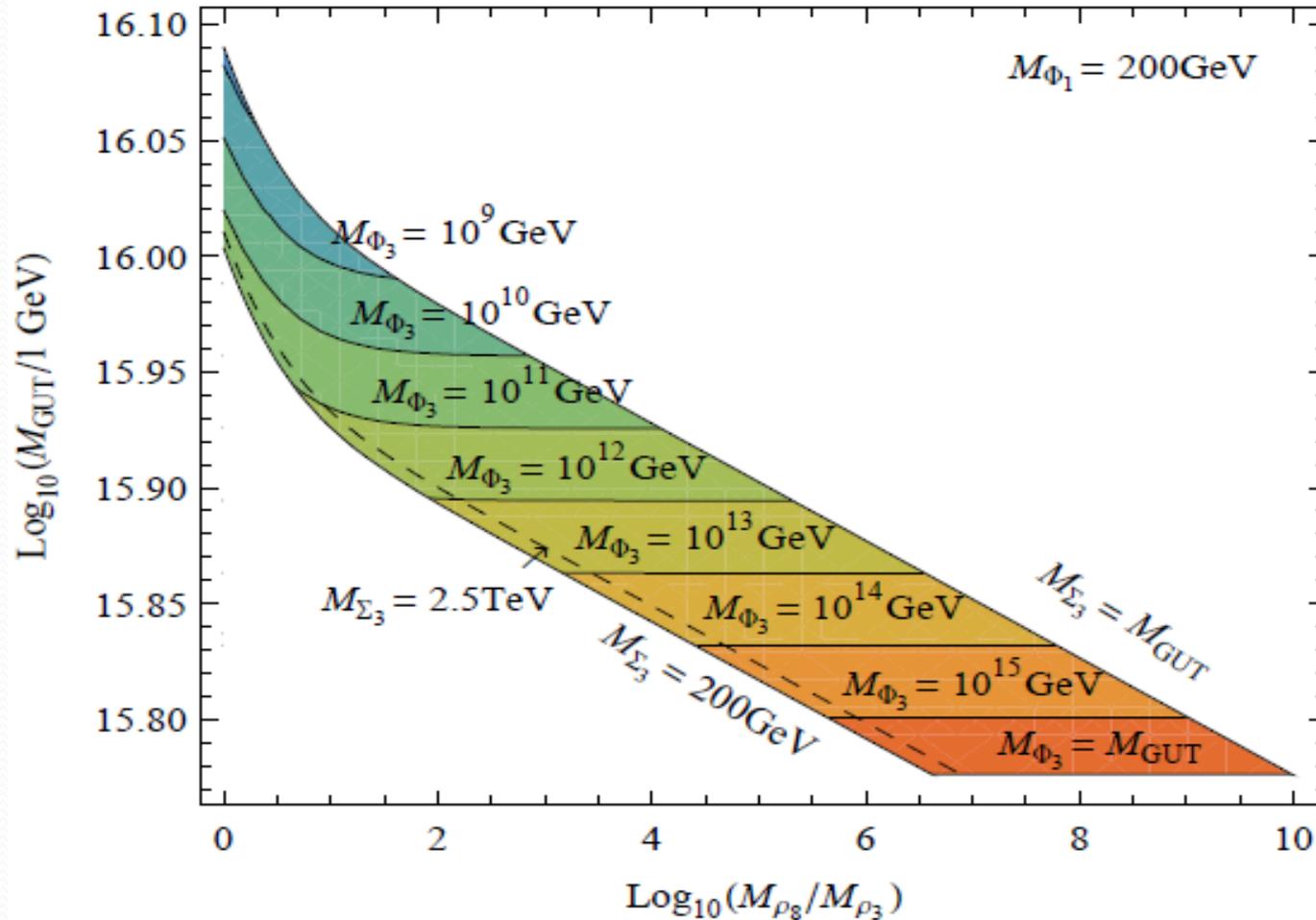
$$V = c_\alpha \bar{5}_\alpha 24 5_H + p_\alpha \bar{5}_\alpha 24 45_H + M \text{Tr} 24^2 + \lambda \text{Tr} (24^2 24_H) + \text{h.c.}$$

$$\rho_3 = \frac{1}{2} \begin{pmatrix} T^0 & \sqrt{2}T^+ \\ \sqrt{2}T^- & -T^0 \end{pmatrix}.$$

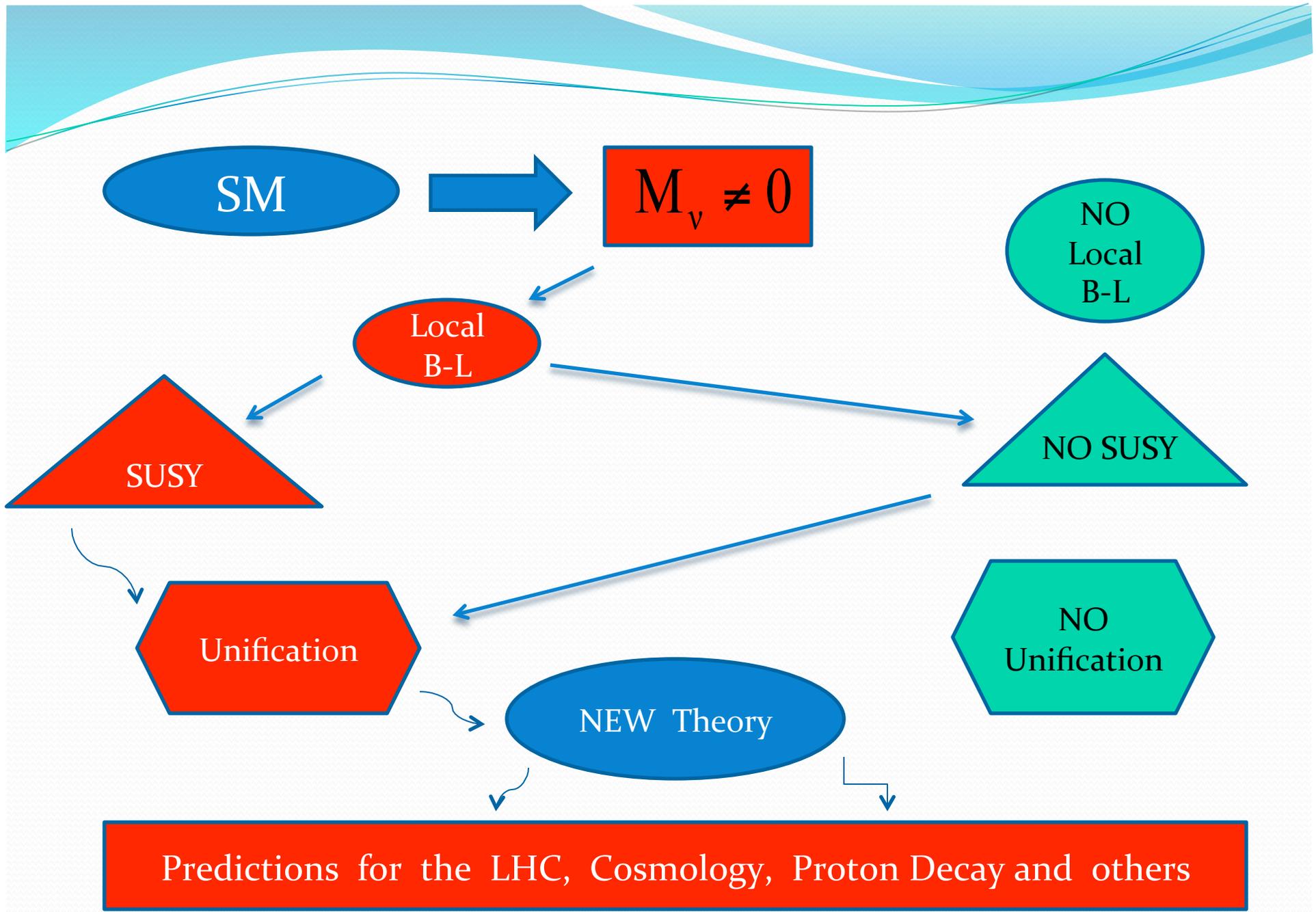
$$M_{\alpha\beta}^\nu = \left(\frac{h_{\alpha 1} h_{\beta 1}}{M_{\rho_3}} + \frac{h_{\alpha 2} h_{\beta 2}}{M_{\rho_0}} \right) v_0^2.$$

One Massless Neutrino !

See also Bajc, Senjanovic'07



Unification at the one-loop level.





Neutrino Masses and Renormalizable $SO(10)$

SO(10) Unification

Gauge Symmetry: $SO(10)$

Matter: $16 = (Q, L, u^C, d^C, e^C, \nu^C)$

Higgs Sector: $10_H = 5_H \oplus \bar{5}_H = (H, T) \oplus (\bar{H}, \bar{T})$

$45_H, \dots$

Naïve $SO(10)$

$$-\mathcal{L}_Y = Y_{10} 16 16 10_H + \text{h.c.}$$

$$M_U = M_\nu^D = v_{10}^u Y_{10} \quad (\textit{wrong!})$$

$$M_D = M_E = v_{10}^d Y_{10} \quad (\textit{wrong!})$$

$$Y_{10} = Y_{10}^T$$

Realistic Renormalizable SO(10)

$$-\mathcal{L}_Y = Y_{10} 16 16 10_H + Y_{126} 16 16 \overline{126}_H + \text{h.c.}$$

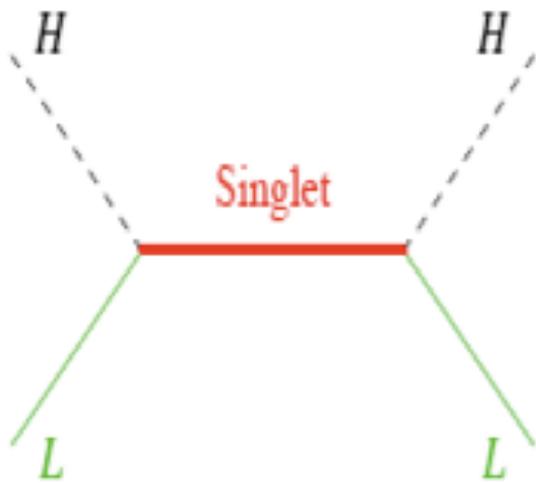
$$M_U = v_{10}^u Y_{10} + v_{126}^u Y_{126}$$

$$M_\nu^D = v_{10}^u Y_{10} - 3 v_{126}^u Y_{126} \quad M_{\nu_R} = Y_{126} v_R$$

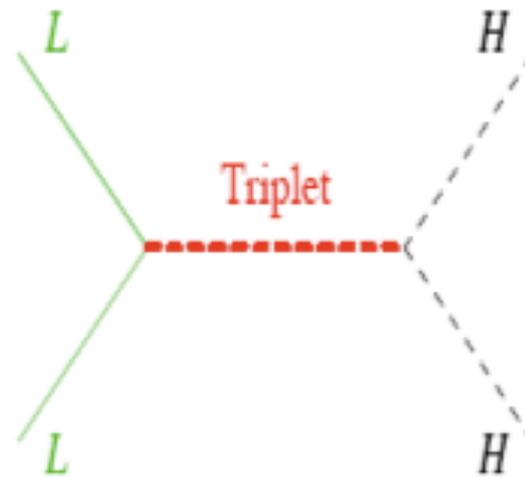
$$M_D = v_{10}^d Y_{10} + v_{126}^d Y_{126}$$

$$M_\nu = -M_\nu^D M_{\nu_R}^{-1} M_\nu^D + Y_{126} v_L$$

Type I plus Type II Seesaw



Type I Seesaw



Type II Seesaw



For the study of neutrino masses in this context see:

Babu, Mohapatra'93

Bajc, Senjanovic, Vissani'02, '04

Goh, Mohapatra, Ng'03

Fukuyama, Okada'03

Bertolini, Frigerio, Malinsky'04

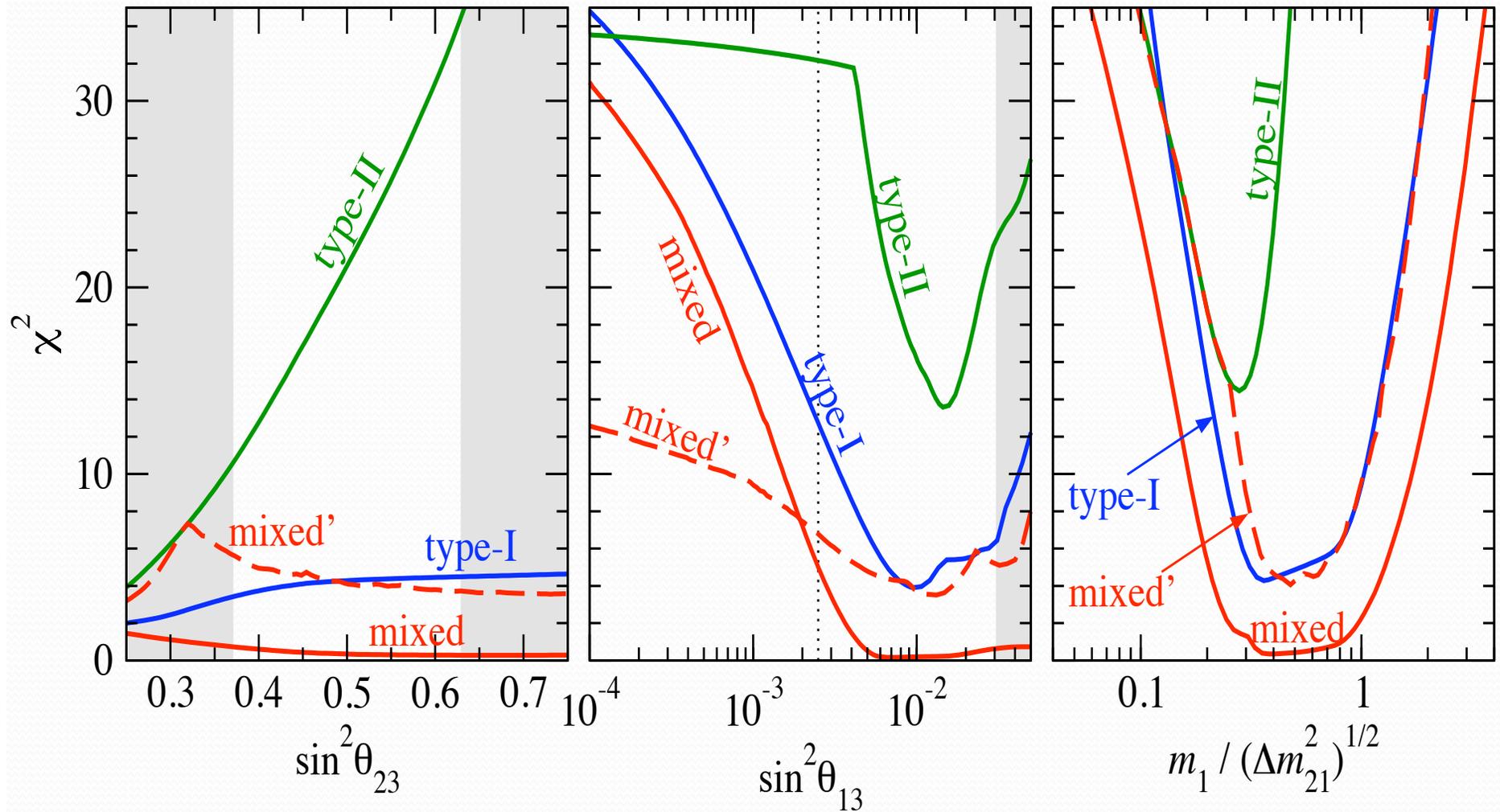
Aulakh, Bajc, Melfo, Senjanovic, Vissani'04,'05'06

Bertolini, Schwetz, Malinsky'06

Bajc, Dorsner, Nemedev'08

	type-II		mixed'		mixed		type-I	
$ \xi $	0		10^{-4}		0.3587		3.59×10^6	
$\arg(\xi)$	-		0.866π		1.018π		1.318π	
$ r $	0.3278		1.9977		0.47896		0.3551	
$\arg(r)$	0.408π		1.849π		0.0013π		0.0057π	
f_u	16.62		11.51		18.77		19.23	
f_ν	1.671×10^{-10}		4.519×10^{-10}		8.732×10^{-10}		3.613×10^{-17}	
observable	pred.	pull	pred.	pull	pred.	pull	pred.	pull
m_d [MeV]	0.7662	- 1.16	0.4956	- 1.82	1.122	-0.29	0.4719	- 1.87
m_s [MeV]	31.33	1.85	22.46	0.15	22.85	0.22	19.99	-0.33
m_b [MeV]	1147	0.61	1096	0.25	1078	0.13	1029	-0.35
m_u [MeV]	0.5543	0.02	0.5576	0.03	0.5512	0.00	0.5538	0.02
m_c [MeV]	213.1	0.17	213.5	0.18	210.6	0.03	213.1	0.16
m_t [MeV]	78030	-0.29	77411	-0.34	81659	-0.05	78117	-0.29
$\sin \phi_{23}^{\text{CKM}}$	0.0345	-0.43	0.0352	0.08	0.0351	0.03	0.0349	-0.13
$\sin \phi_{13}^{\text{CKM}}$	0.00331	0.23	0.00319	-0.02	0.00319	-0.01	0.00323	0.06
$\sin \phi_{12}^{\text{CKM}}$	0.2245	0.11	0.2243	0.02	0.2243	0.01	0.2243	0.01
$\delta_{\text{CKM}} [^\circ]$	79.35	1.38	59.47	-0.04	61.41	0.10	61.11	0.08
$\sin^2 \theta_{23}^{\text{PMNS}}$	0.3586	- 2.17	0.5126	0.19	0.5027	0.04	0.4944	-0.09
$\sin^2 \theta_{13}^{\text{PMNS}}$	0.0145	0.93	0.0106	0.68	0.0066	0.43	0.0095	0.61
$\sin^2 \theta_{12}^{\text{PMNS}}$	0.2829	- 1.08	0.3078	-0.09	0.3094	-0.02	0.3078	-0.09
$\Delta m_{21}^2 [10^{-5} \text{eV}^2]$	7.863	-0.12	7.894	-0.02	7.898	-0.01	7.896	-0.01
$\Delta m_{31}^2 [10^{-3} \text{eV}^2]$	2.385	0.50	2.232	0.09	2.210	0.03	2.223	0.06
$m_1 / \sqrt{\Delta m_{21}^2}$	0.279		0.478		0.382		0.361	
$\delta_{\text{PMNS}} [^\circ]$	-0.70		-59		-0.70		4.9	
$\alpha_1 [^\circ]$	1.1		30		1.8		-2.1	
$\alpha_2 [^\circ]$	91		126		-84		90	
χ^2		14.5		4.1		0.35		4.3

TABLE II: Parameter values and predictions in four example solutions corresponding to different terms dominating the neutrino mass matrix: type-I, type-II, or both contributions of comparable size (mixed and mixed'). In the column "pred." the predicted values P_i for the observables are given, the column "pull" shows the number of standard deviations from the observations, $(P_i - O_i)/\sigma_i$, using the data and errors from Tab. I. Deviations of more than 1σ are highlighted in boldface. The final χ^2 is the sum of the squares of the numbers in the "pull" column. See the text for comments on the values of the leptonic CP phases.



Summary

- We have presented a new mechanism for neutrino masses at one-loop level called “ Colored Mechanism ”.
- It has been shown that if the seesaw scale is low one could learn about the neutrino spectrum at the LHC in the case of Type II seesaw.
- We have discussed how to realize the different seesaw mechanisms in the simplest grand unified theories pointing out the possible predictions.
- In the context of $SO(10)$ theories we have seen the possibility to make some predictions without assuming extra symmetries.
- All the experimental efforts will be of great importance to understand how to prove the mechanism for neutrino masses.



THANKS !!!!

