

## DETECTION OF HEAVY PARTICLES ...Quarks and Higgs

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### Abstract

Possibilities of detecting heavy particles (quarks and Higgs) by SSC is examined in three different modes: jets only; jets and leptons; and leptons only. The background estimation is not complete, but some hints of suppressing the background by kinematical cuts are obtained.

### Introduction

It is generally taken for granted that lepton identification is essential for detection of new particles. Good momentum resolution up to a few TeV, and pseudo rapidity coverage up to 3 is required. Jet analysis is also very important. One of the new aspect in SSC data analysis is to find out the origin of the jet by measuring all possible physical quantities: particle multiplicity in the event (e.g., Higgs and Z have lower multiplicity), number of jets (as we see later this is a measure of hardness of collision), fatness of the jet, energy flow, energy deposition around a lepton track (lepton isolation), angular correlation, etc. The group of K. Kondo at Univ. of Tsukuba, for example, tries to distinguish quark- and gluon-jets by fitting the moments of particle-momenta with suitable distribution functions. In this case, too, a fast and reliable event generator is essentially important. In this talk, some of the physical quantities which help identify the origin of the jets are investigated, and requirements for the lepton detector are made clear in the process of reconstructing the Higgs particle of  $500 \text{ GeV}/c^2$ .

### SSC-detector

In Fig.1, an example of proposed SSC-detector is shown. There are many other proposals of different types, but essential points are as follows:

#### 1) Tracking.

The tracking is thought to be important in measuring relatively low momenta of particles and matching the tracks which enter the calorimeter or muon chamber. The proposed

momentum resolution is  $10-50\%P_t$ .

2) Super conducting solenoid coil.

The radius is about 2m, field strength is 1.7-2 Tesla.

3) Calorimeter.

Electro-magnetic and hadronic shower calorimeters. There are many possibilities, but proposed detector elements are roughly classified as ionization chamber and scintillation counter. The electro-magnetic calorimeter has energy resolution of  $15-20\%/sqr(E)$ , the hadron calorimeter has  $35-50\%/sqr(E)$ . The granularity in  $\phi$ - and pseudo-rapidity space is  $(0.03-0.06) \times (0.03-0.06)$ .

4) Muon chamber.

In the central region (pseudo-rapidity 0-1.5), there are two different proposals: one of them proposes to use the toroidal field of 2 Tesla, the other proposes to use only the return field. The required resolution is analysed in more detail later. The forward region is covered by the magnetized iron, as the solenoidal coil gives a poor bending power in this region.

Heavy quark reconstruction.

In Fig.2, a copy of transparency is shown. This is from the talk at KEK on Sept. 7 '88. In this analysis, heavy down-quark of mass  $500 \text{ GeV}/c^2$  is reconstructed through top and W reconstruction by reconstructing 8 jets, without lepton detection. The contents are hopefully self-explanatory. The cut conditions which were used in the reconstruction are:

\* $n_{\text{jet}}$  is greater than 6.

\* $P_{\text{jet}}$  is greater than  $30 \text{ GeV}/c$ .

\* $E_{T\text{jet}}$  is greater than  $20 \text{ GeV}$ .

\* $\cos(\text{jet opening angle from W})$  is greater than 0.7.

\*reconstructed  $m_W$  is between  $40$  and  $100 \text{ GeV}/c^2$ .

\*reconstructed  $P_W$  is greater than  $1000 \text{ GeV}/c$ .

\* $\cos(\text{jet opening angle from top})$  is greater than 0.8.

\*reconstructed top mass is between  $20$  and  $60 \text{ GeV}/c^2$ .

Note that the mass of top quark was assumed to be  $40 \text{ GeV}/c^2$  at that time. Later it was assumed to be  $200 \text{ GeV}/c^2$ .

In Fig.3, assumed detector resolution and conclusions are shown.

The  $t\bar{t}$ -bar background is further analysed in the talk at Sendai on Nov. 15 '88. In Fig.4, the number of  $t\bar{t}$ -bar events

which survived the cut conditions above at each cut point is shown as a function of  $P_T$  of the top quarks at generation time. The number drops abruptly for  $P_T$  below 300 GeV/c. A large contribution comes from the number of jets which survive the cut conditions. The requirement that the number of 'good jets' must be larger than 6 reduces the QCD background substantially, especially for lower  $P_T$  events, as Fig.5 shows. Fig.6 shows the  $t\bar{t}$  jet cross-section as a function of  $P_T$  of the top quark at generation time. Fig.7 is a copy of transparency showing a conclusion about ( $t\bar{t}$ ) jet background.

#### Higgs reconstruction through detection of 4 leptons.

Neutral Higgs particle of 500 GeV/c<sup>2</sup> decaying into two Z's is reconstructed by detecting 4 muons. This study is done in order to make clear the requirements (resolution and coverage) for the muon detector. The cut conditions on the muons are as follows:

- (1) Pseudo rapidity must be within  $\pm 3$ .
- (2)  $P_T$  must be larger than 20 GeV/c.
- (3)  $P$  must be larger than 30 GeV/c.
- (4) The number of muons which pass the conditions above must be greater than or equal to 4.
- (5) Cosine of the opening angle between a pair of muons with opposite charge must be greater than 0.
- (6) There must be at least two such pairs whose invariant mass lie within 88 and 100 GeV/c<sup>2</sup> (Z0 mass is assumed to be 94 GeV/c<sup>2</sup>).
- (7) The Higgs mass is reconstructed out of all combinations of the above pairs.

Fig.8 shows the reconstruction efficiency as a function of the type of resolution which is listed below. In the figure, the efficiency in the case of detecting two mu's and two jets is also shown for the resolution type 10, 11 and 'mixed' case. The 'mixed' case means the muon momentum is measured by the tracking chamber up to pseudo-rapidity  $y_p=1.6$ , and it is measured by the iron-toroid magnet for the pseudo-rapidity between 1.6 and 3. This clearly demonstrates that the tracking chamber is very important in the central region, and that the mu-mu+2 jet mode is more advantageous in reconstruction efficiency.

Fig.9 shows the reconstructed Higgs mass in the case of a perfect

detector resolution. Fig.10 shows the reconstructed Z0 mass and Higgs mass in the case of detector resolution type 6 (like the central tracker at 90 degree). Fig.11 shows the same things in the case of detector resolution type 12 (20%).

Fig.12 shows the reconstructed Higgs mass on the Z0-pair background for the integrated luminosity of  $10^{40}$ , in the cases of detector resolution type 1 (perfect resolution) and 4 ( $0.3 P_T$ ). Fig.13 shows the same things in the cases of detector resolution type 10 (iron toroid --- muon momentum vector is measured at the exit of iron filter) and 11 (iron toroid --- muon momentum vector is not measured at the exit of iron filter). It is clear that the resolution is very important, and that the luminosity is essential. Fig.14 shows the Higgs-reconstruction efficiency as a function of pseudo-rapidity coverage of the central tracker. This shows that the central tracker coverage up to  $y=1.5 - 2$  is very important.

Although Z0-pair background is seen to be not serious as in the figures 12 and 13 in the case of detecting 4 leptons, the QCD-background may be serious. Some of the cuts which suppress this background is investigated. Fig.15 shows the cumulative percentage of transverse energy deposited within the cone of  $\Delta R=0.6$  from a muon track (lepton isolation). It is seen that Higgs and Z0-pair events are very isolated, but the  $t\bar{t}$  events (here top mass of  $200 \text{ GeV}/c^2$  is assumed) are not, especially when the generated top quark has a large transverse momentum (i.e., when the collision is hard). The rejection factor of 300 is expected. Fig.16 shows the distribution of the number of muons passing cuts in an event. The decay in flight of charged kaons and pions are included. Higgs events have 60 % efficiency if the cut is made at 4 (pseudo-rapidity coverage up to 3). This efficiency drops to 1/3 if the coverage is limited to  $y=1.6$ . This demonstrates the necessity of forward muon detector. On the other hand, top-pair events have only 0.6% of efficiency; thus the rejection of 100 is expected.

#### Higgs reconstruction through 2 leptons and 2 jets.

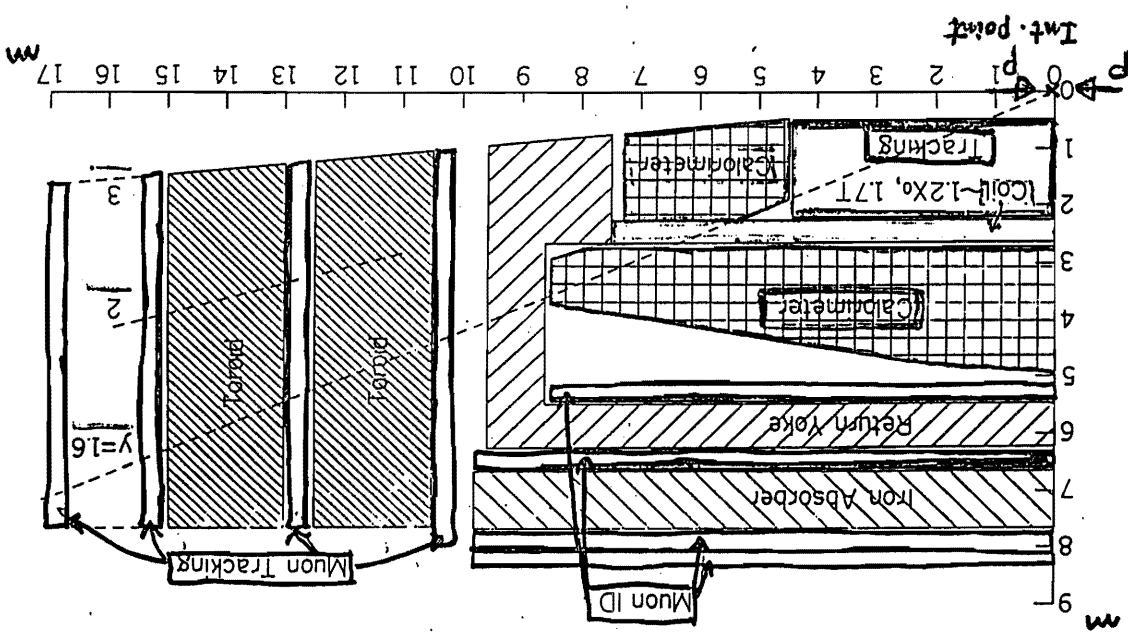
Neutral Higgs decaying into 2 Z0's is reconstructed by detecting 2 muons and 2 jets. Fig.17 shows Z0 reconstruction by 2 jets. Fig. 18 shows the reconstructed Higgs for the integrated luminosity of  $5 \times 10^{39}$ . A comparison with Fig.12 shows

the advantage in efficiency. Fig.19 shows the W-pair background for the luminosity of  $10^{38}$ . If more cuts are imposed, this background is not serious. Other QCD background needs more study.

Conclusion.

- 1) Lepton id is very important.  
Good resolution up to a few Tev and at low momentua.  
Good coverage up to pseudo-rapidity 3.
- 2) Lepton isolation, number of jets, particle multiplicity etc, are very powerful cuts.
- 3) To study QCD background and jets, a reliable and fast event generator is essentially needed.
- 4) For neutral Higgs reconstruction, a good detection efficiency and high luminosity are needed in 4-lepton mode. In 2-lepton + 2 jet mode, QCD background needs more study, but very promising efficiencywise.

Fig. 1



August 1989 SM

An example of SSC detector (quadrant)

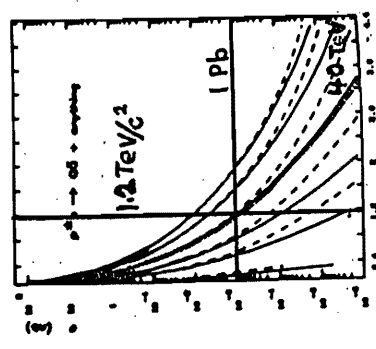
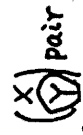
~ 30,000 ton

Fig. 2

Heavy Quark Detection at SSC

Sept. 7 '88  
Yuzo Asano  
Univ. of Tsukuba

From unitarity,  $M_F \lesssim 550 \text{ GeV}/c^2$   
for  $M_X \approx M_Y = M_F \cdot \frac{1}{2}$   
 $\beta = 1.02 \pm 0.02 \rightarrow |M_X^2 - M_Y^2| \lesssim 350 \text{ GeV}/c^2$



CROSS SECTION →

$\approx 1.2 \text{ TeV}/c^2$  for 1 Pb

If  $\int L = 10^{40} \text{ cm}^{-2}$

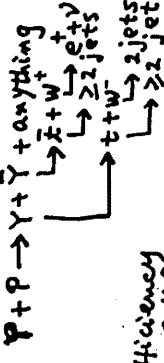
number of quarks  $\sim 10^4$

FIG. 143 Integrated cross sections for pair production of heavy quarks in proton-proton fixed and colliding beam experiments. Theoretical curves are shown for the case of independent standard model couplings, according to the predictions of the 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100.

EHLQ (1984)

Assume,  $m_Y = 500 \text{ GeV}/c^2$ ,  $m_X = 40 \text{ GeV}/c^2$

• Kim and Takikawa



efficiency  $\sim 0.4\%$

• Similar channel with lepton detection  $\sim 0.2\%$

TRIED,

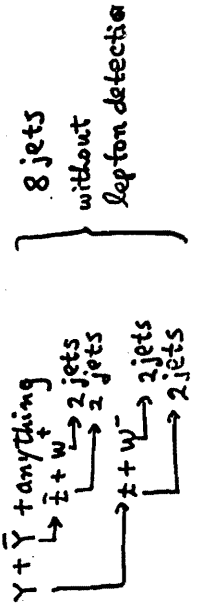


Fig. 3

Assumed Detector:  $\frac{\Delta E}{E} = 20\%/E$  for e and T  $\Delta\eta = 0.06$   $|\eta| \leq 3$   
 $\frac{\Delta E}{E} = 50\%/E$  for hadrons  $\Delta\phi = 0.06$   
 jet finding radius  $R_{jet} = 0.6$

Reconstruct W  $\rightarrow$  Reconstruct  $\bar{t}$   $\rightarrow$  Reconstruct Y  
 from (hadron) calorimeter only.  
 cuts, as follows...

Conclusion

Signal seen : efficiency  $\sim 2.5\%$   
 (accurate mass determination is difficult)

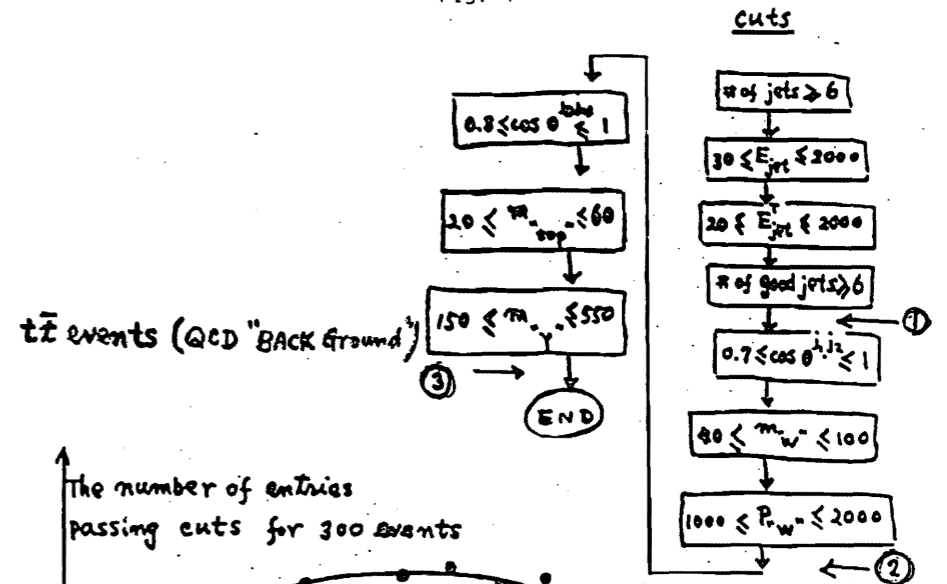
Back ground :

To be done  $\left\{ \begin{array}{l} W^+W^- \text{ OK.} \\ t\bar{t} \text{ probably OK} \\ \text{min. bias} \text{ but, can't tell definitely.} \\ \text{others To be done} \end{array} \right.$

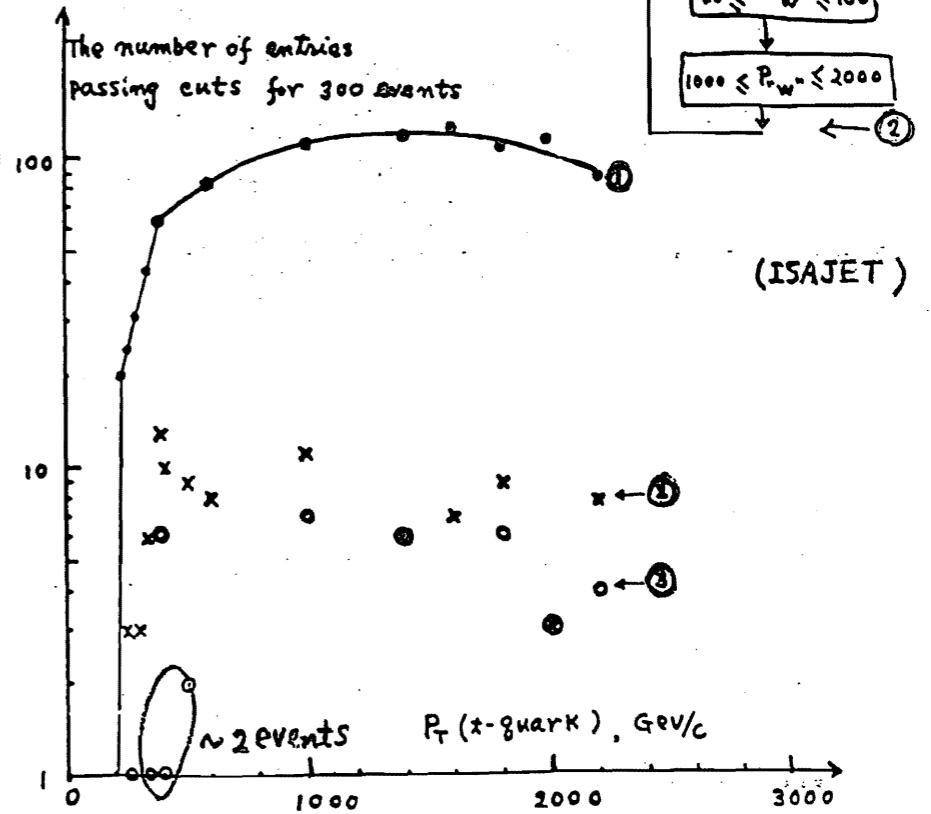
how to cope with difference in  $\int L$  as much as  
 $\times 3000$   $t\bar{t}$   
 $\times 10^9$  min. bias ??

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Fig. 4



$t\bar{t}$  events (QCD "BACK ground")



(ISAJET)

Fig. 5

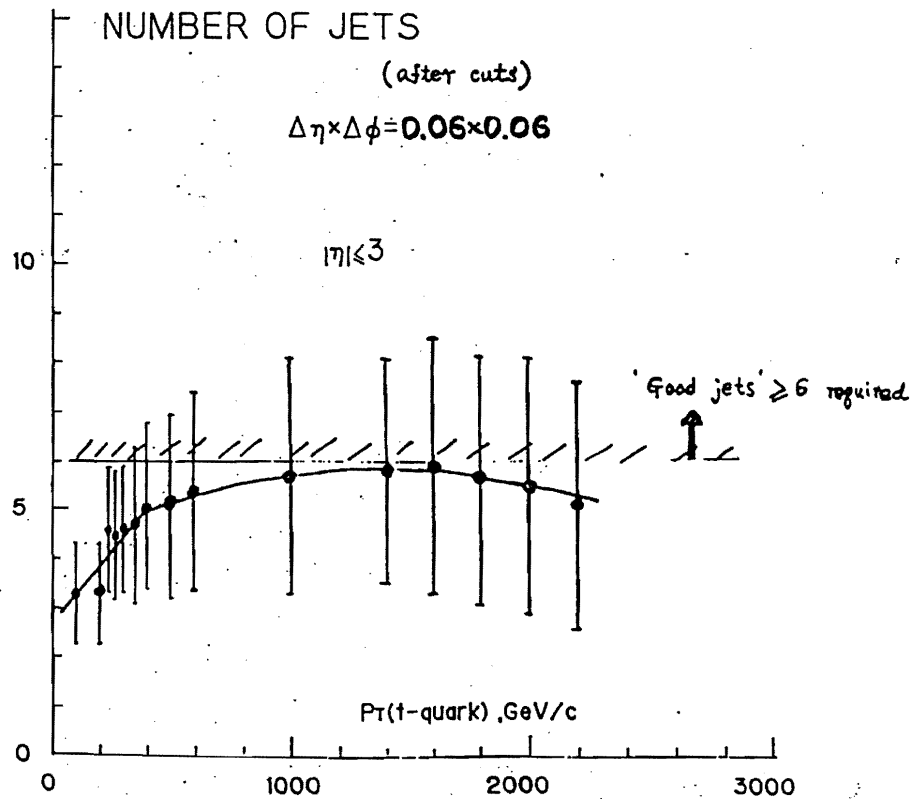


Fig. 6

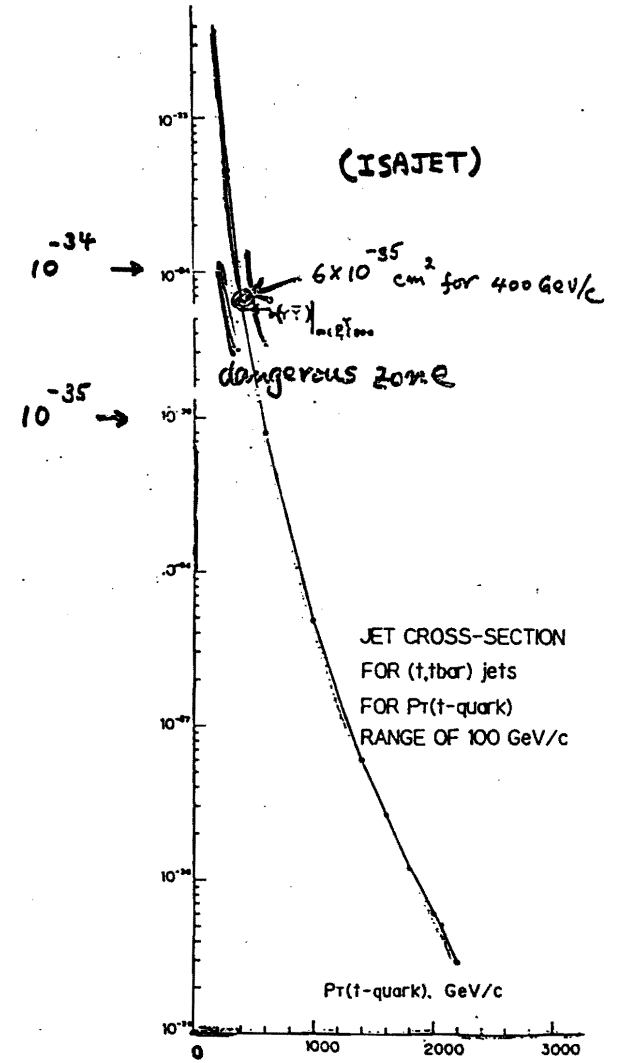




Fig. 7

( $t\bar{t}$ ) jet background

- $P_T \gg 400 \text{ GeV}/c$  : cross section too small
- $P_T \sim 400 \text{ GeV}/c$  :

$\int L = 10^{40}$  : 対する Background events

$$\sim \frac{2 \text{ events}}{300 \text{ trials}} \times (6 \times 10^{-35}) \times 10^{40} \times 2$$

$\uparrow$   $\sigma$  for  $\Delta P_T = 100$        $\uparrow$   $\int L$        $\uparrow$   $\Delta P_T = 200$

$\sim \underline{\underline{8 \times 10^3}}$

- $P_T \leq 200 \text{ GeV}/c$  :

{ # of 'good jets'  $> 6$  cut  $\rightarrow$  rejection  $\times 100$   
 cross section  $\times 100$

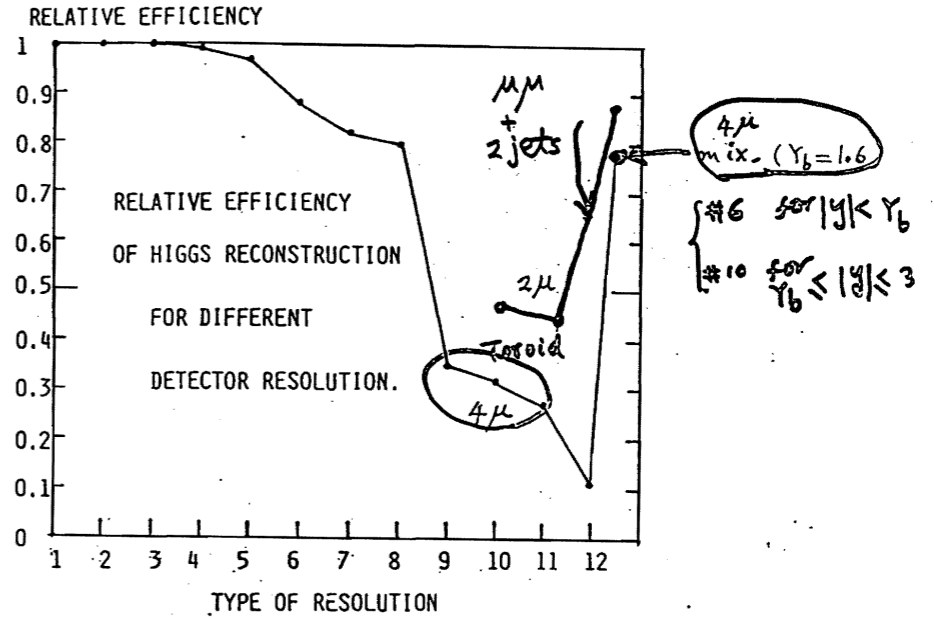
Signal  $\frac{15 \text{ events}}{5.6 \times 10^{36} \text{ cm}^{-2}}$  とほぼ cancel

$\rightarrow \frac{15 \times 10^{40}}{5.6 \times 10^{36}} = \underline{\underline{3 \times 10^4}}$

**結論** QCD background は, signal とほぼ同数  
 従って, jet analysis [multiplicity, energy flow, etc.] により  
 更に数倍の rejection があれば,  
 十分 business になりうる。

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Fig. 8

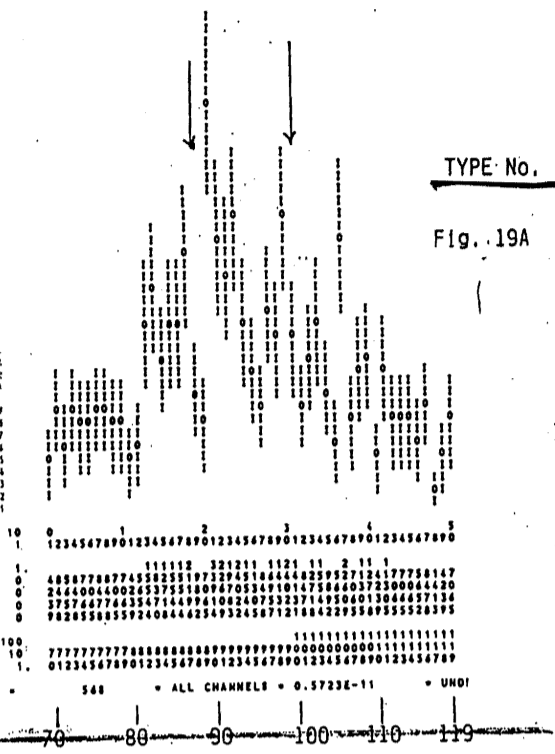


TYPE NO.	RESOLUTION $\Delta P/P$
1	Perfect
2	2%
3	$0.15/\sqrt{E}$
4	$0.3P_T$
5	4%
6	$0.14P + 0.01$
7	$0.54P_T$
8	$0.02$ for $ABS(y) < 1$ , $0.02 + (ABS(y) - 1) \cdot 0.06$ for $ y  > 1$ and $< 3$
9	$0.1 + 0.1P_T$ (TeV/c)
10	Toroid case 3 (Amako program #3...muon vector is measured at the exit of the filter.)
11	Toroid case 2 (Amako program #2...only position of the mu is measured at the exit of the filter.)
12	20%



12 FOR SHEAR NO. 1  
 D = 116  
 DATE 07/06/89

20 MASS



20-PAIR MASS FOR SHEAR NO. 1  
 HBOOK ID = 128

H MASS

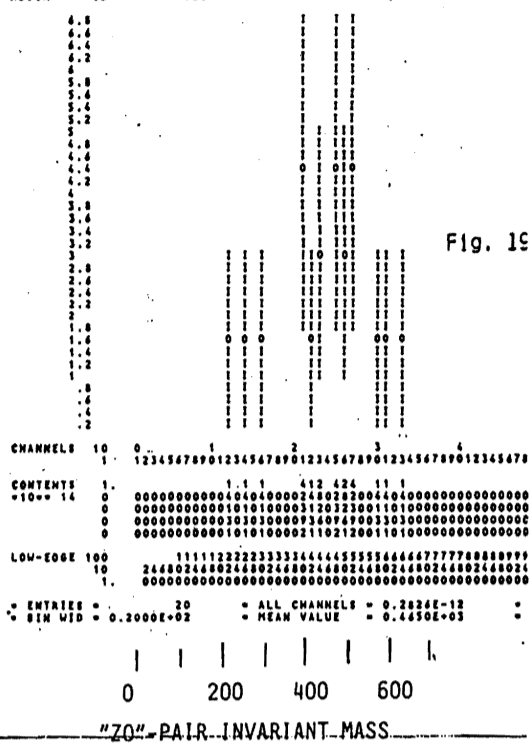


FIG. 11

MU-PAIR INVARIANT MASS

"Z0"-PAIR INVARIANT MASS

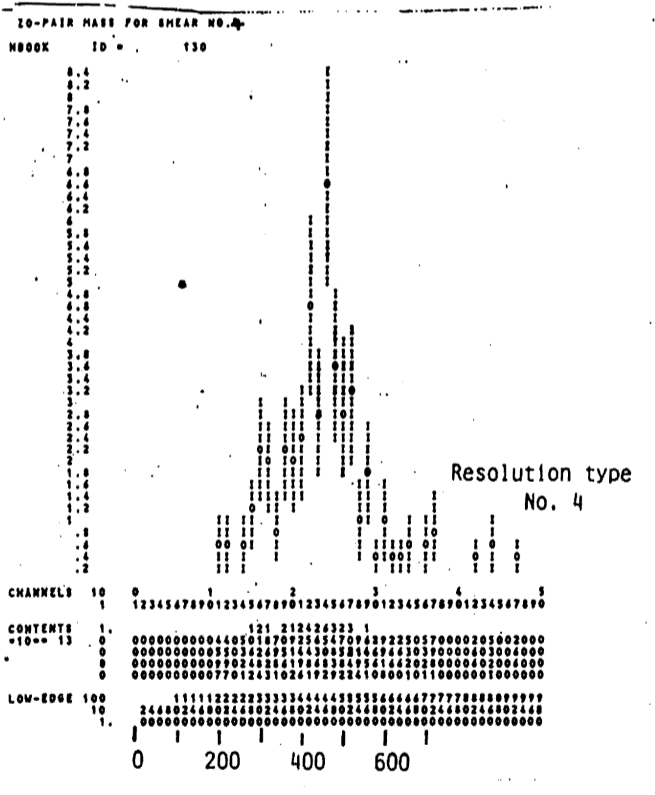
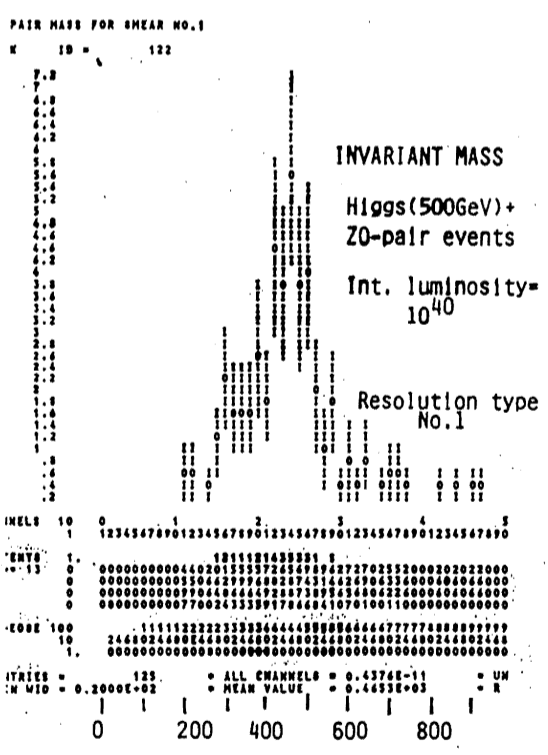


FIG. 12

J-PAIR MASS FOR SHEAR NO.10

30K ID = 133

Z0-PAIR MASS FOR SHEAR NO.11

HBOOK ID = 132

INVARIANT MASS  
Higgs+Z0-pair

Res. type 10

Res. type 11

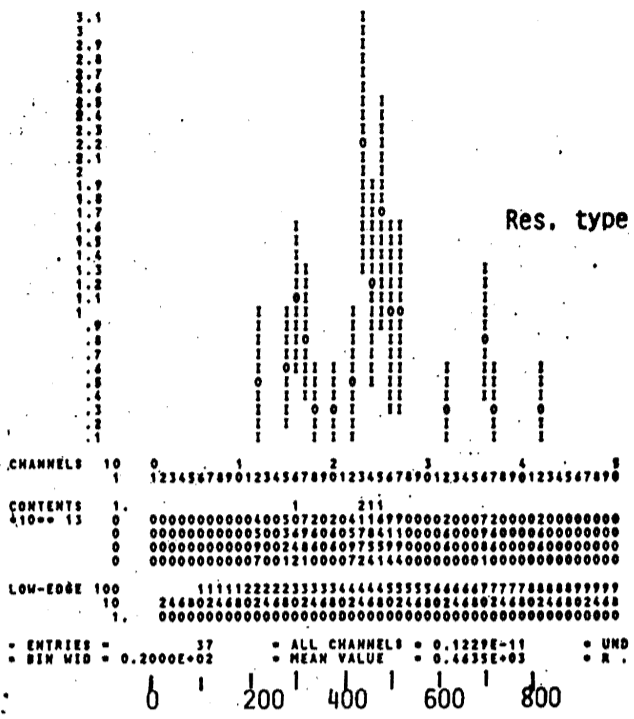
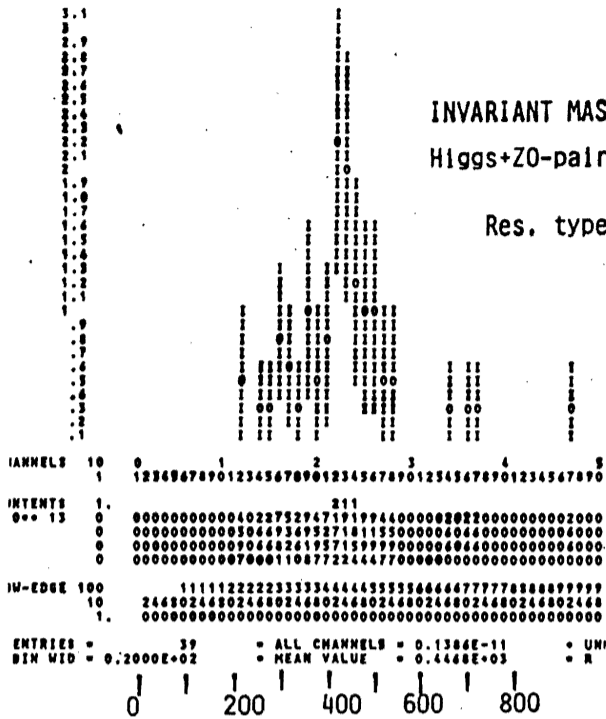


Fig. 13

Fig. 14

HIGGS RECONSTRUCTION EFFICIENCY  
IN CENTRAL TRACKER + TOROID

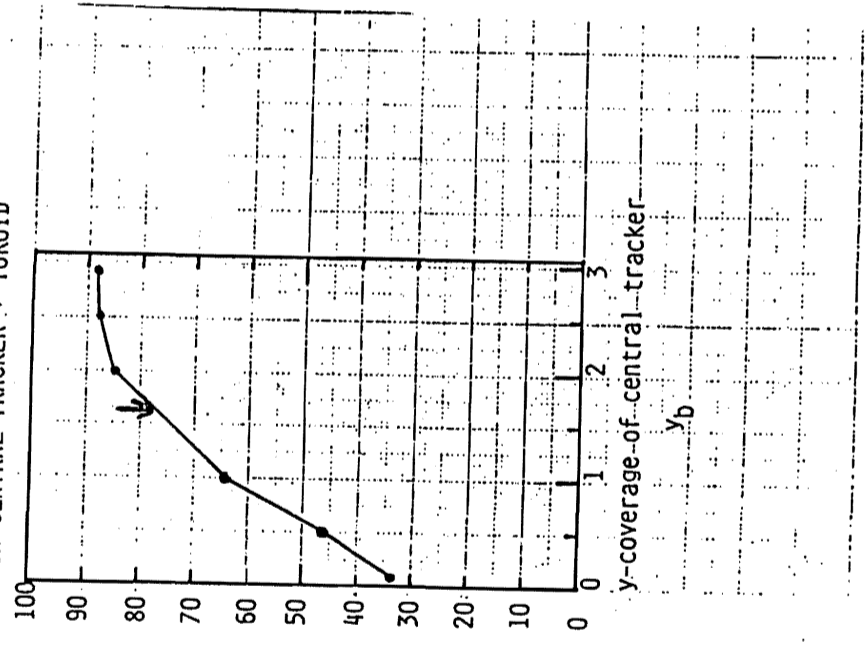
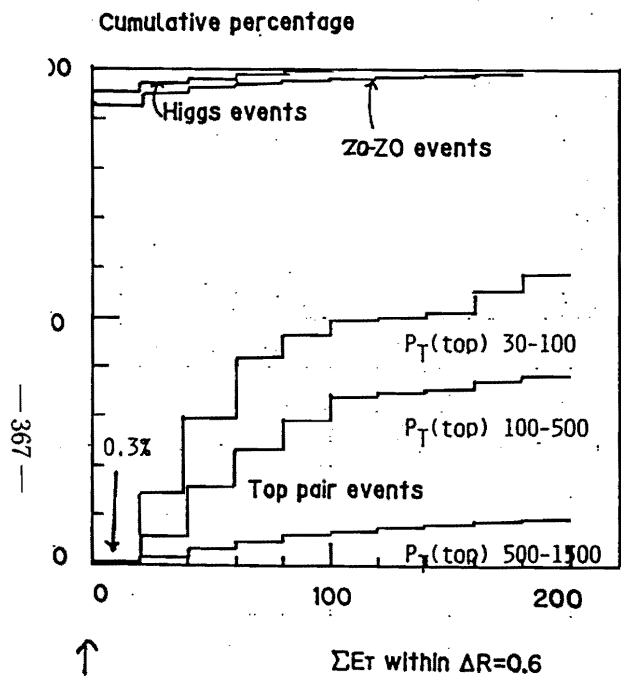


Fig. 15

MUON ISOLATION



$m_t = 200$

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↑  $\Sigma E_t$  within  $\Delta R=0.6$

$E_{T \text{ sum}}$  が  $\Delta R$  値より  
小さい確率

例えば、20 GeV cut  
をすれば、Higgsは95%  
とわかるが、 $Z^0$  eventsは、  
0.3%しか混入しない。

$$\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2} = 0.6$$

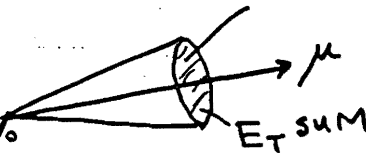
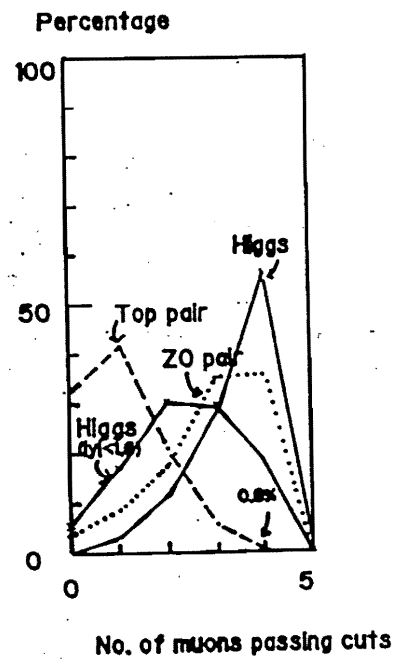


Fig. 16



$m_t = 200$   
Decay included.



