

## 2.5 JETS AND MASSIVE DILEPTONS AT LARGE $p_T$

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We have investigated the possibility of studying massive dileptons recoiling against large  $p_T$  hadron jets. While experiments of this type are of general theoretical interest, the experimental problems are formidable. The physics involved in the production of dileptons at large  $p_T$  is very similar to that discussed above for direct photons at large  $p_T$  and has been discussed by several authors.<sup>1-4</sup> Additional interest arises from the possibility of observing large  $p_T$   $J/\psi$ 's recoiling against jets.<sup>5</sup> Common to these authors' discussions is the desirability of studying the correlation of the jets and dimuons.

Experiments have been performed at Fermilab, ISR, and SPS which show a rather large  $\langle p_T \rangle$  for massive dimuons.<sup>5-9</sup> Studying both the dilepton and the jet are possible at the Tevatron because of the increased rates at some moderate mass and  $p_T$ , say  $M = 4$  and  $p_T = 4$ . The distinct advantage of the Tevatron is the variation of incident beams which can probe different production mechanisms and hadron structure functions.

We have considered experiments with either dimuons or dielectrons in the final state opposite high  $p_T$  jets.

### Dimuons

We considered an open geometry with a large solid-angle calorimeter (e.g., E-609 or E-557) augmented by a downstream muon spectrometer as sketched in Fig. 1. The background from pion decay is enormous. Our crude estimate indicates that there will be approximately  $5 \times 10^5$  dijet events (of which some  $\lesssim 1\%$  will contain dimuons from  $\pi$  decay) for every dimuon event at a given  $p_T$ ; however, it is possible that with a sophisticated trigger processor which required massive dimuons at large  $p_T$  and a veto on the same side jet that a signal might be observable.

We also considered a half-closed geometry as shown in Fig. 2. This geometry has the advantage of reducing the pion decay background. This geometry, however, has a greatly reduced acceptance for the dimuons, and because half of the detector is open for jet observation, it is not possible to increase the flux by large factors to overcome the acceptance loss. Another drawback occurs due to multiple scattering. This may not be a serious effect depending on the absorber and particular spectrometer arrangement.

There is an inherent incompatibility in the relatively large magnetic field required for good muon mass resolution and the relative small magnetic field required to reduce smearing of the jet  $p_T$ ; perhaps a compromise can be achieved.

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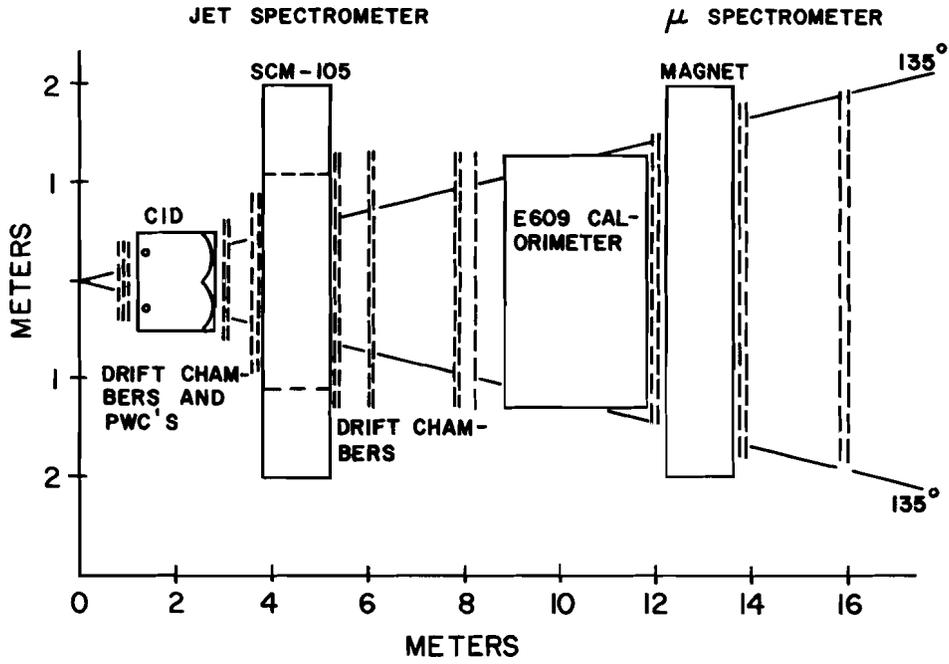


Fig. 1. Open geometry for detection of dimuons and jets.

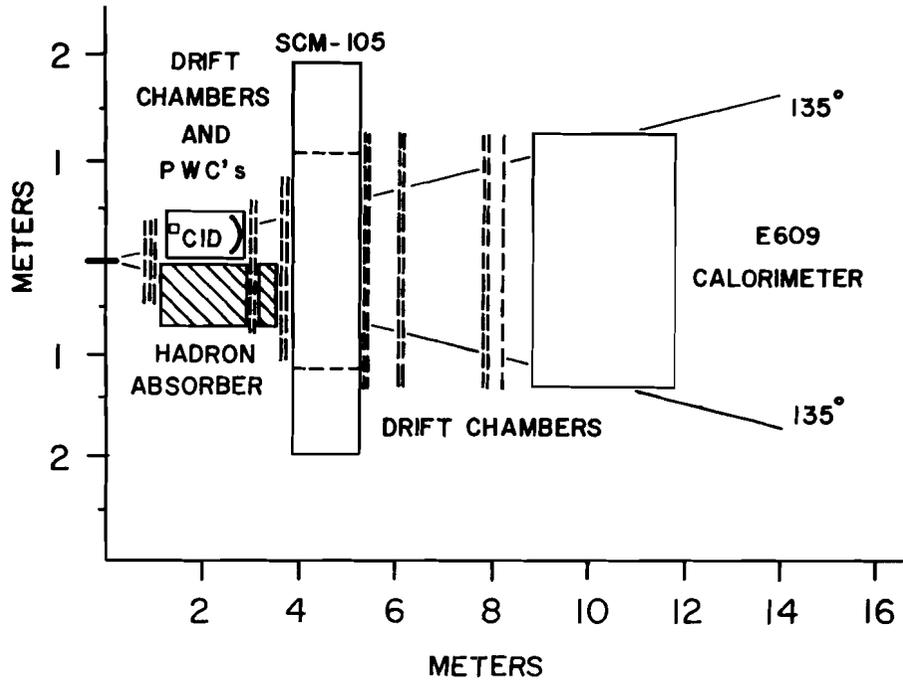


Fig. 2. "Half closed geometry" for detection of dimuons and jets.

The trigger for this experiment will still require a large  $p_T$  jet on one side along with mass requirement on the dimuons.

### Dielectrons

For studying massive dielectrons at high  $p_T$  we have considered a geometry of the type shown in Fig. 3. Backgrounds of a different variety are inherent in electron-triggered experiments. The primary backgrounds are from hadrons simulating electrons in the electromagnetic detector and from the Dalitz decays and externally converted electromagnetic decays of  $\pi^0$ ,  $\eta$ ,  $\eta'$ , ...

The signal at a given  $p_T$  is again approximately one part in  $5 \times 10^5$  of the dijet cross section. The probability for  $\pi^0$  conversion or for hadron simulation of electrons is about 1% each, and the probability for two such background electron "signals" is of the order of  $10^{-3}$ . Consequently, a further rejection factor on the order of  $10^3$  is required in the trigger by requiring massive dielectrons at large  $p_T$  and the absence of a second high  $p_T$  jet. We remark, with guarded optimism, that such a trigger may be feasible but a detailed study would be required.

Finally, we show in Fig. 4 predicted rates in a high  $p_T$  dilepton trigger compared to the expected rates for a direct photon experiment. The rate calculations are based on a typical run of 100 hrs. with a flux of  $10^7/s$  incident on 0.45 cm LH2 target.

We observe (as does McLeod in Section 2.5) that an apparatus of the type shown in Fig. 3 could simultaneously measure high  $p_T$  jets, direct photons, and jets, and perhaps dielectrons and jets.

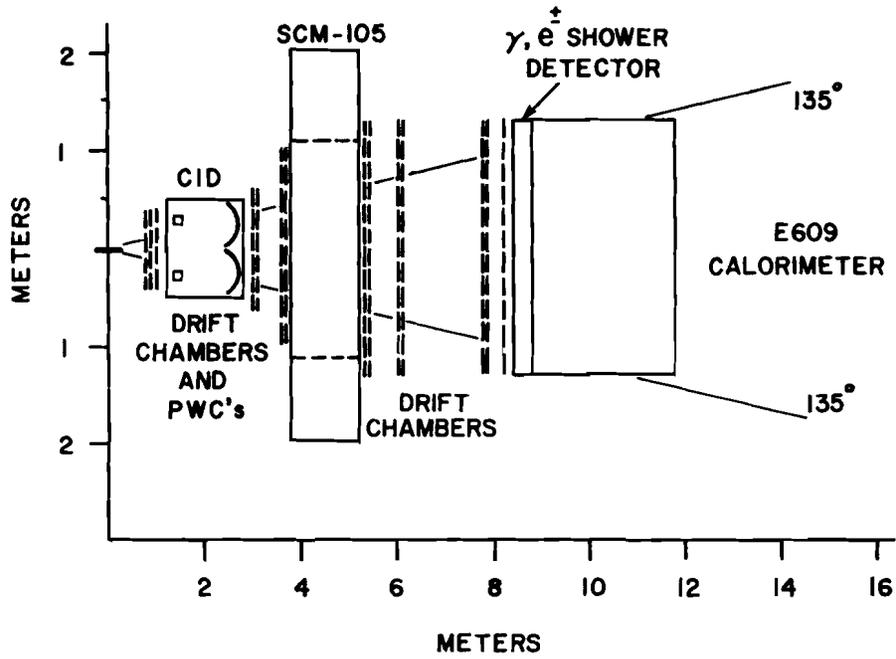


Fig. 3. Detector for study of dielectrons and high  $p_T$  jets.

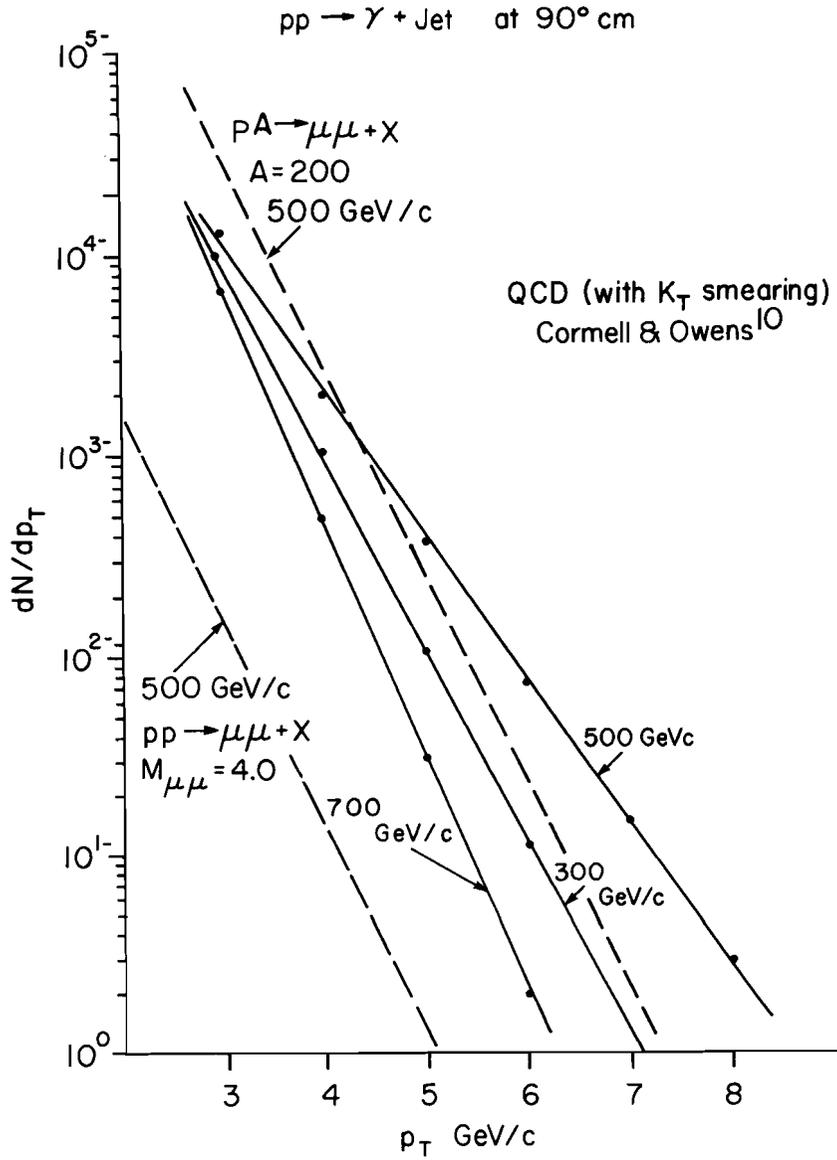


Fig. 4. Predicted rates for massive dimuons at large  $p_T$  compared with the predicted rates of direct photons and jets at  $p_{beam} = 500 \text{ GeV}/c^{10}$ . The rates are based on a run of 100 hrs at  $10^7$  proton/sec incident on a 10% target.

### References

1. H. Fritzsch and P. Minkowski, Phys. Lett. **73B**, 80 (1978).
  2. G. Altarelli, G. Parisi, and R. Petronzio, Phys. Lett. **76B**, 351 (1978).
  3. K. Kajantie, J. Lindfors, and R. Raitio, Nucl. Phys. **B144**, 422 (1978).
  4. Z. Kunszt, E. Pietarinen, and E. Reya, Phys. Rev. **D3**, 733 (1980).
  5. J. Collins, proceedings of this workshop.
  6. L. Lederman, XIX International Conf. on High Energy Physics, Tokyo, 1978.
  7. R. Barate et al., Phys. Rev. Lett. **43**, 1541 (1979).
  8. K. J. Anderson et al., Phys. Rev. Lett. **42**, 944 (1979),  
G. E. Hogan et al., Phys. Rev. Lett. **42**, 948 (1979).
  9. W. Kienzle et al., in Proceedings of the International Symposium on Lepton and Photon Interactions at High Energies, Fermi National Accelerator Laboratory, August 1979.
  10. L. Cormell and J. F. Owens, Phys. Rev. **D22**, 1609 (1980).
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