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Angular Distribution of the $J/\psi \pi^0$ events ON and OFF the 1P_1 Resonance

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Abstract

The angular distribution of the $J/\psi \pi^0$ signal in the 1P_1 energy region is compared to the distribution of the background events (off resonance and in the η_c' region)

Introduction

We are studying the reaction:

$$p\bar{p} \longrightarrow ^1P_1 \longrightarrow J/\psi + \pi^0 \quad (1)$$

to be compared with the background :

$$p\bar{p} \longrightarrow J/\psi + \pi^0 \quad (2)$$

(the J/ψ is detected through its e^+e^- decay).

The formation of the 1P_1 as intermediate state in (1), constrains the orbital angular momentum both in the initial and final state. Given the 1P_1 quantum numbers:

$$J = 1 \quad P = +1 \quad C = -1,$$

the $p\bar{p}$ initial state must have:

$$P(p\bar{p}) = (-1)^{L+1} = +1 \quad C(p\bar{p}) = (-1)^{L+S} = -1$$

This implies:

$$L = \text{odd} \quad \text{with} \quad S = 0;$$

in particular $L = 1$ to have $J = 1$. If the z-axis is taken along the incident \vec{p} direction, the projection of J along this axis should be:

$$J_z = M = L_z + S_z = 0$$

since $L_z = 0$. Then:

$$|initial\rangle = |JM\rangle = |10\rangle$$

To match the 1P_1 quantum numbers, the final state $J/\psi \pi^0$ has an orbital angular momentum $l = \text{even}$:

$$P = P(J/\psi) \times P(\pi^0) \times (-1)^l = (-1)(-1)(-1)^l = (-1)^l$$

$$P = (-1)^l = +1 \Rightarrow l = \text{even}$$

$$C = C(J/\psi) \times C(\pi^0) = (-1)(+1) = -1.$$

The angular momentum conservation in the decay:

$$^1P_1 \longrightarrow J/\psi + \pi^0$$

$J_{1P_1} = 1 = l + J_{J/\psi} = l + 1$ constrains the possible values of $l = 0, 2$.

The final state can then be a mixture of a S and D wave:

$$|final\rangle = A |Swave\rangle + B |Dwave\rangle = |JM\rangle$$

1.

pure S wave

$$|JM\rangle = |10\rangle = A Y_0^0 |10\rangle_{J/\psi}$$

The decay amplitude is then $T = A Y_0^0$, yielding a flat angular distribution. Using the helicity formalism language, this corresponds to the case where the 3 helicity states for the J/ψ are equally fed by the 1P_1 decay: in this case the two independent helicity amplitudes $f_{\pm 1,0}^1$ and $f_{0,0}^1$ are equal.

2.

pure D wave

$$|JM\rangle = |10\rangle = B (\sqrt{(3/10)} Y_2^1 |1-1\rangle_{J/\psi} - \sqrt{(2/5)} Y_2^0 |10\rangle_{J/\psi} + \sqrt{(3/10)} Y_2^{-1} |11\rangle_{J/\psi})$$

Assuming $A=0$ the angular distribution is:

$$W(\theta) \propto \frac{2}{5} |Y_2^0|^2 + \frac{3}{10} (|Y_2^1|^2 + |Y_2^{-1}|^2) = \frac{1}{8\pi} (1 + 3 \cos^2 \theta)$$

This corresponds to have $f_{0,0}^1 = -2 f_{\pm 1,0}^1$

3. In the general case, both waves can contribute to the angular distribution.

In the case of reaction (2), where no intermediate state constrains the initial or final angular momentum, we can have:

$$L = \text{odd} \quad S = 0 \quad \Rightarrow \quad l = \text{even} \quad (a)$$

$$L = \text{even} \quad S = 1 \quad \Rightarrow \quad l = \text{odd} \quad (b)$$

This is easily obtained by conserving P and C quantum numbers:

$$P(p\bar{p}) = (-1)^{L+1} \quad P(J/\psi\pi^0) = (-1)^l$$

$$C(p\bar{p}) = (-1)^{L+S} \quad C(J/\psi\pi^0) = -1.$$

$$J = L + S = l + J_{J/\psi} = l + 1$$

$$M = L_z + S_z = 0 \quad L = \text{odd}, \quad S = 0$$

$$M = L_z + S_z = 0, \pm 1 \quad L = \text{even}, \quad S = 1.$$

To study the background we should consider that in the final state, besides the S and D wave (already discussed for the 1P_1 formation, with $L = 1$ and $S = 0$), an extra P wave should be taken into account.

$$l = 1 \quad P \text{ wave}$$

$$J = l + J_{J/\psi} = 1 + 1 = 0, 1, 2.$$

Assuming $L = 0$, for the $p\bar{p}$ state:

$$J = L + S = 1 \quad M = 0, \pm 1$$

The possible initial state $|JM\rangle$ are then: $|10\rangle$, $|11\rangle$ and $|1-1\rangle$. The corresponding decay amplitudes are:

$$T(|10\rangle) = \sqrt{\frac{1}{2}}Y_1^{-1}|11\rangle - \sqrt{\frac{1}{2}}Y_1^1|1-1\rangle$$

$$T(|11\rangle) = \sqrt{\frac{1}{2}}Y_1^1|10\rangle - \sqrt{\frac{1}{2}}Y_1^0|11\rangle$$

$$T(|1-1\rangle) = -\sqrt{\frac{1}{2}}Y_1^{-1}|10\rangle + \sqrt{\frac{1}{2}}Y_1^0|1-1\rangle$$

Assuming to sum over the three polarization of the initial state with the J/ψ helicity states equally populated we expect a flat angular distribution, being:

$$T^2(|1\pm 1\rangle) = \frac{3}{16\pi}(1 + \cos^2\theta)$$

$$T^2(|10\rangle) = \frac{3}{8\pi}\sin^2\theta$$

Assuming $L = 2$ for the $p\bar{p}$ state:

$$J = L + S = 2 + 1 = 1, 2, 3 \quad M = 0, \pm 1$$

we can have $J = 1, 2$ with $M = 0, \pm 1$, compatible with the P wave ($l = 1$) in the final state.

If $J = 1$ the result is a flat angular distribution (as discussed above).

If $J = 2$, summing over the three polarization of the initial state, we obtain an angular distribution:

$$W(\theta) \propto T^2(|20\rangle) + T^2(|21\rangle) + T^2(|2-1\rangle)$$

$$W(\theta) \propto \frac{1}{4\pi}(2 + 3\cos^2\theta)$$

initial>			final>	$W(\theta)$
L	S	J	l	
1	0	1	0	$1/4\pi$ (A)
1	0	1	2	$\frac{1}{8\pi}(1 + 3\cos^2\theta)$ (B)
0	1	1	1	$1/4\pi$
2	1	1	1	$1/4\pi$
2	1	2	1	$\frac{1}{4\pi}(2 + 3\cos^2\theta)$ (C)

Table 1: Expected angular distribution as a function of the initial and final quantum numbers

Data analysis

Figure 1 shows the $\cos\theta^*$ of the π^0 (c.m. frame) as calculated by the fit for data selected by the analysis as $J/\psi \pi^0$ candidates (see 1P_1 paper). In particular fig. 1a and 1d are respectively the distribution of the events at the 1P_1 resonance and off (1P_1 and η'_c energies).

To study if there is any evidence of a preferred distribution of the signal events respect to the background, we used a Monte Carlo that generates the events with the standard CERN package GENBOD.

As discussed before we made the assumption that the helicity states of the J/ψ are equally populated. This results in the further assumption that the final decay $J/\psi \rightarrow e^+e^-$ has a flat distribution. Choosing for this decay a $(1 + \cos^2\theta^*)$ distribution in the other extreme hypothesis doesn't imply any relevant difference to the final conclusions.

The three distributions A, B and C are used to generate $J/\psi \pi^0$ events at the c.m. energy of the 1P_1 (the distributions doesn't change drastically fixing the c.m. energy at the η'_c mass - 3610 MeV -).

The geometrical acceptance of our detector was imposed using the following cuts:

$$15^\circ < \theta_{e_{1,2}} < 60^\circ$$

$$2^\circ < \theta_{\gamma_{1,2}} < 68^\circ$$

In figure 2 the continuous and dashed line are respectively the data and the MC (normalized to data) distribution for each of the three hypothesis A , B and C .

The following table summarizes the results of the χ^2 calculation, where:

$$\chi^2 = \sum_{i=1}^{10} \frac{(n_i^{exp} - n_i^{th})^2}{\sigma_i^2}$$

with:

n_i^{exp} number of events in the i -th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i -th bin

σ_i poissonian errors of n_i^{exp}

	A	B	C
ON	1.19	0.51	-
OFF	0.42	0.73	0.46

Table 2: $\chi^2/\# d.o.f.$; $\# d.o.f. = N_{bin} - 1 = 10 - 1$

Hp.: The J/ψ decay has a flat distribution.

1

The following table summarizes the results of the χ^2 calculation, where:

$$\chi^2 = \sum_{i=1}^{10} \frac{(n_i^{exp} - n_i^{th})^2}{\sigma_i^2}$$

with:

n_i^{exp} number of events in the i-th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i-th bin

σ_i poissonian errors of n_i^{exp}

	A	B	C
ON	1.19	0.51	0.68
$-\ln(lik)$	19.85	18.68	18.52
OFF	0.42	0.73	0.46
$-\ln(lik)$	15.97	17.41	16.56

Table 1: $\chi^2/\# d.o.f.$; $\# d.o.f. = N_{bin} - 1 = 10 - 1$

N.B. $-\ln(lik)$ is the same for all the χ^2 calculations.

2

The following table summarizes the results of the χ^2 calculation, where:

$$\chi^2 = \sum_{i=1}^{10} \frac{(n_i^{exp} - n_i^{th})^2}{\sigma_i^2}$$

with:

n_i^{exp} number of events in the i-th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i-th bin

σ_i errors as $\sqrt{n_i^{th}}$

3

The following table summarizes the results of the χ^2 calculation, taking the formula of the

DPG for bins with few events, where:

$$\chi^2 = \sum_{i=1}^{10} 2(n_i^{th} - n_i^{exp}) + 2n_i^{exp} \ln(n_i^{exp}/n_i^{th})$$

	A	B	C
ON	0.85	0.69	0.59
OFF	0.42	0.75	0.57

Table 2: $\chi^2/\# \text{ d.o.f.}$; $\# \text{ d.o.f.} = N_{bin} - 1 = 10 - 1$

with:

n_i^{exp} number of events in the i-th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i-th bin

	A	B	C
ON	0.90	0.64	0.61
OFF	0.43	0.75	0.56

Table 3: $\chi^2/\# \text{ d.o.f.}$; $\# \text{ d.o.f.} = N_{bin} - 1 = 10 - 1$

Hp.: The J/ψ decay has a distribution $(1 + \cos^2 \theta^*)$.

1

The following table summarizes the results of the χ^2 calculation, where:

$$\chi^2 = \sum_{i=1}^{10} \frac{(n_i^{exp} - n_i^{th})^2}{\sigma_i^2}$$

with:

n_i^{exp} number of events in the i-th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i-th bin

σ_i poissonian errors of n_i^{exp}

	A	B	C
ON	1.46	0.57	0.67
$-\ln(lik)$	20.34	18.69	18.46
OFF	0.57	0.65	0.48
$-\ln(lik)$	16.63	17.11	16.48

Table 1: $\chi^2/\# d.o.f.$; $\# d.o.f. = N_{bin} - 1 = 10 - 1$

N.B. $-\ln(lik)$ is the same for all the χ^2 calculations.

2

The following table summarizes the results of the χ^2 calculation, where:

$$\chi^2 = \sum_{i=1}^{10} \frac{(n_i^{exp} - n_i^{th})^2}{\sigma_i^2}$$

with:

n_i^{exp} number of events in the i-th bin

n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the i-th bin

σ_i errors as $\sqrt{n_i^{th}}$

3

	<i>A</i>	<i>B</i>	<i>C</i>
ON	0.94	0.66	0.59
OFF	0.57	0.69	0.55

Table 2: $\chi^2/\# \text{ d.o.f.}$; $\# \text{ d.o.f.} = N_{bin} - 1 = 10 - 1$

The following table summarizes the results of the χ^2 calculation, taking the formula of the DPG for bins with few events, where:

$$\chi^2 = \sum_{i=1}^{10} 2(n_i^{th} - n_i^{exp}) + 2n_i^{exp} \ln(n_i^{exp}/n_i^{th})$$

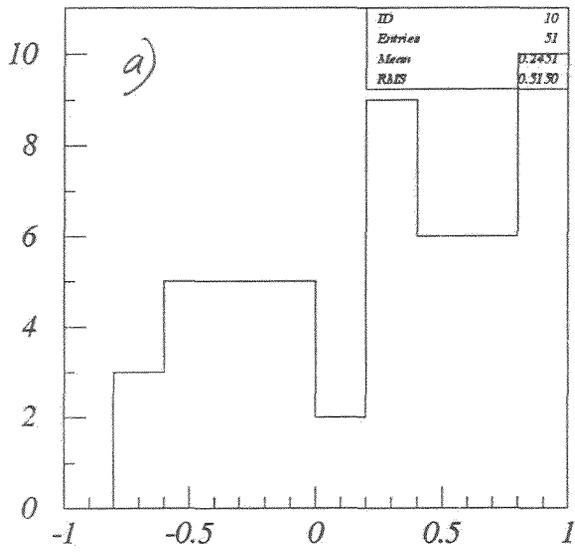
with:

n_i^{exp} number of events in the *i*-th bin

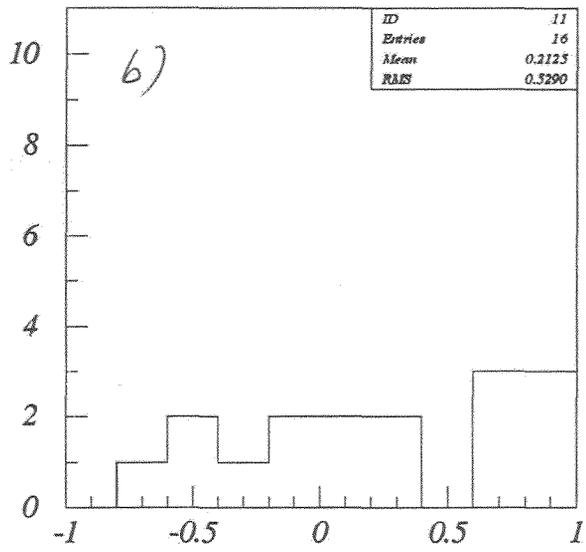
n_i^{th} number of MC events (normalized to $\sum_i n_i = n$) in the *i*-th bin

	<i>A</i>	<i>B</i>	<i>C</i>
ON	1.01	0.65	0.59
OFF	0.58	0.68	0.54

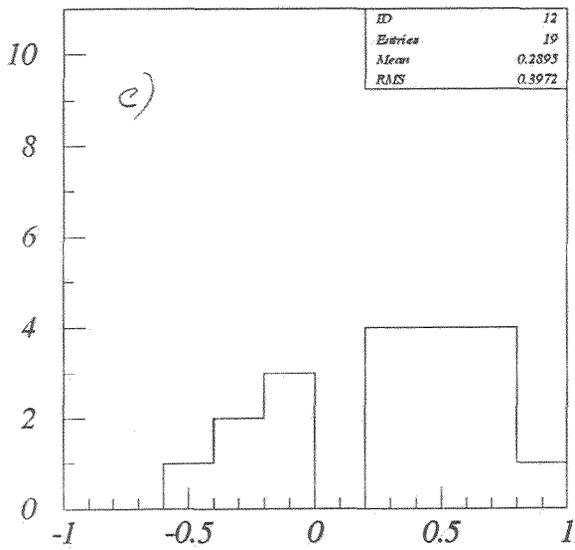
Table 3: $\chi^2/\# \text{ d.o.f.}$; $\# \text{ d.o.f.} = N_{bin} - 1 = 10 - 1$



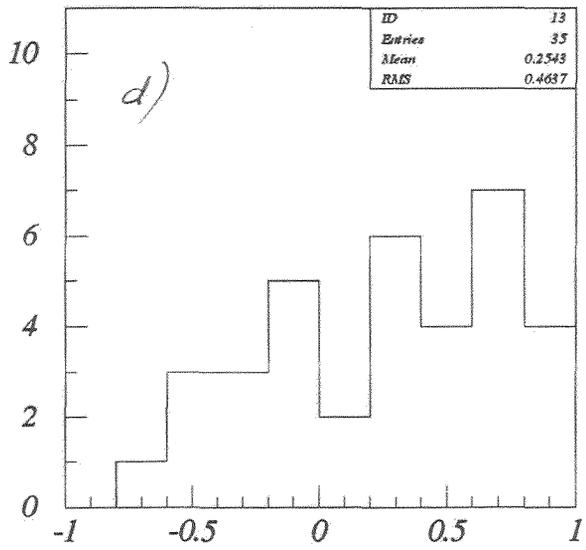
cosh pi0 in c.m. system ON IPI



cosh pi0 in c.m. system OFF IPI



cosh pi0 in c.m. system ETACP

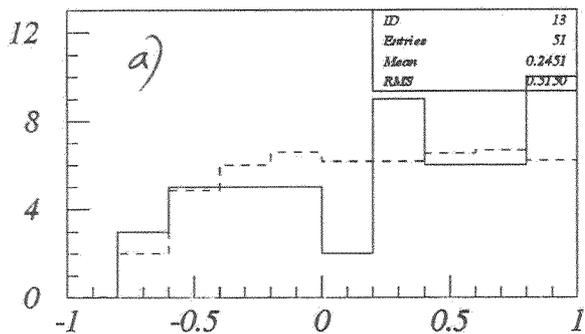


cosh pi0 in c.m. system OFF

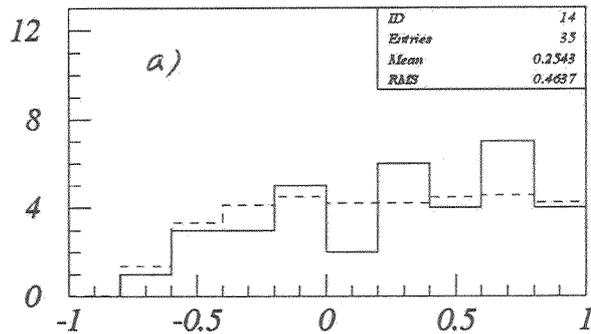
FIGURE 1

ON - ENTRIES = 51

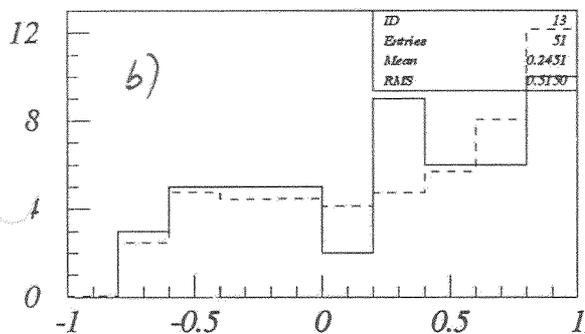
OFF - ENTRIES = 35



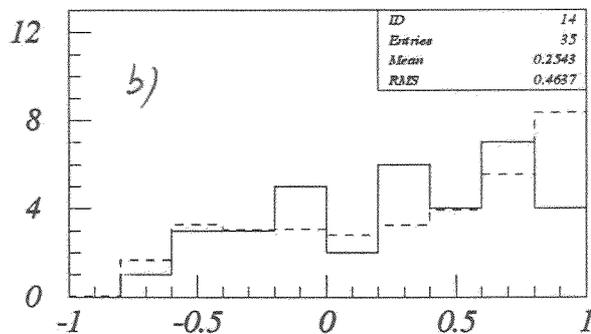
cosh_fit pi0 in c.m. system



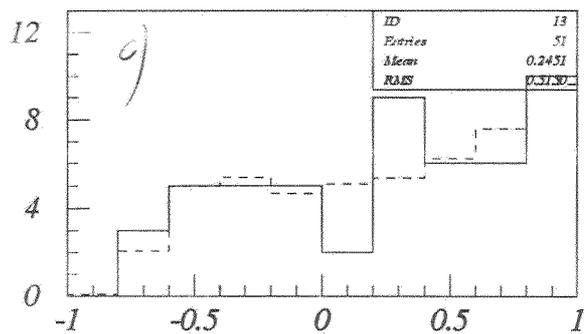
cosh_fit pi0 in c.m. system



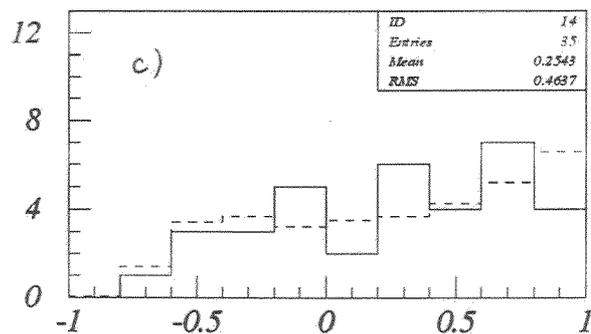
cosh_fit pi0 in c.m. system



cosh_fit pi0 in c.m. system



cosh_fit pi0 in c.m. system



cosh_fit pi0 in c.m. system

FIGURE 2