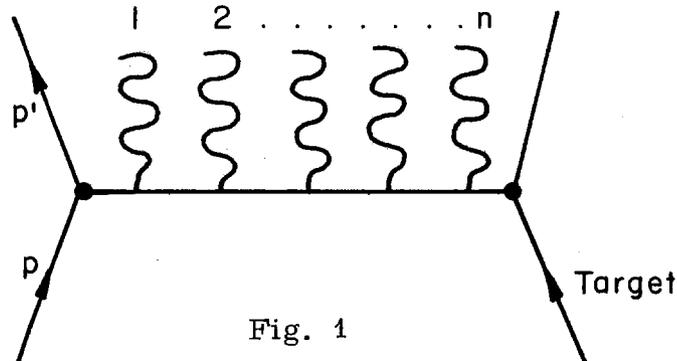


A SIMPLE FORMULA FOR THE DISTRIBUTION OF ENERGETIC
SECONDARY BARYONS FROM PROTON-INITIATED COLLISIONS

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A slight generalization of the multiperipheral model developed by authors in UCRL-18275 leads to a simple formula for the energy-angular distribution of the most energetic secondary baryons from proton-initiated collisions. The multiperipheral diagram is that of Fig. 1, corresponding to the production of n pions together with



the energetic non-strange baryon p , which may be either a nucleon or a low-mass resonance.

In the model of UCRL-18275 let us integrate over all variables except the two referring to the left-most link in Fig. 1. These two remaining variables we here call x and t , rather than x_1 and t_1 . The differential cross section is then $(e^{X_0} = E_{\text{lab}}/M_p)$

$$\frac{\partial^2 \sigma_n^{p \rightarrow p'}}{\partial x \partial t} \approx \sigma_{\text{tot}}^{p \rightarrow p'} \left(g^2 \right)^n \frac{\left(X_0 - x \right)^{n-1}}{(n-1)!} \bar{a} e^{at} e^{-g^2 X_0}, \quad (1)$$

where we have dropped the M subscript of UCRL-18275, restricting ourselves to the simple model of Sec. E of that paper, where all trajectories are of the "meson" class, none of the Pomeranchuk. We have here inserted a simple exponential t-dependence with a width determined by the parameter a, which may be taken as a constant--or in a better approximation, equal to $a_0 + 2\alpha'x$ where $\alpha' \approx 1 \text{ GeV}^{-2}$ is the slope of the average trajectory. Note that g^2 determines both the energy and multiplicity dependence of the cross section. In particular, $\bar{n} = g^2 X_0$. Note also that if we integrate Formula (1) over dx (from 0 to X_0) and over dt (from 0 to $-\infty$), we find

$$\sigma_n \approx \sigma_{\text{tot}}^{p \rightarrow p'} \frac{\left(g^2 X_0 \right)^n}{n!} e^{-g^2 X_0}. \quad (2)$$

Finally, if we sum Formula (2) over all n, we get $\sigma_{\text{tot}}^{p \rightarrow p'}$ as the energy-independent total cross section for the proton to produce the baryon p' . Alternately we may first sum Formula (1) over n to obtain

$$\frac{\partial^2 \sigma^{p \rightarrow p'}}{\partial x \partial t} \approx \sigma_{\text{tot}}^{p \rightarrow p'} g^2 e^{-g^2 x} \bar{a} e^{at}. \quad (3)$$

If the x-dependence of a is neglected, we see that the dependence on t and the dependence on x are independent of each other.

The variable t is approximately related to the transverse momentum of the final baryon p' :

$$\left. \begin{aligned} p'_{\perp}{}^2 &\approx -(t - t_{\min}) \\ -t_{\min} &\approx \left(m_{p'}^2 - m_p^2 \right) \frac{s_r}{s} \end{aligned} \right\} \quad (4)$$

where

if s_r is the square of the invariant mass of all outgoing particles except p' . The longitudinal momentum of the final baryon, either in the lab or c.m. systems, is given by

$$p'_{\parallel} \approx p_a \left(1 - \frac{s_r}{s} \right).$$

What we need, finally, is the relation between s_r and x . This is given by

$$\frac{s_r}{s} \approx k e^{-x}, \quad (5)$$

where k is a constant near unity that depends on dynamical details of the multiperipheral chain. The model fails for x very close to zero, so we cannot determine k by normalization at $x = 0$. In terms of p'_{\perp} and p'_{\parallel} , the distributions are

$$\frac{\partial \sigma}{\partial (p'_{\perp})^2} \sim e^{-a (p'_{\perp})^2},$$

$$\frac{\partial \sigma}{\partial p'_{\parallel}} \sim (p_a - p'_{\parallel}) g^2 - 1, \quad \frac{p'_{\parallel}}{p_a} > 1 - k.$$

The model thus contains four parameters, $\sigma_{\text{tot}}^{p \rightarrow p'}$, g^2 , a_0 , and k , but g^2 is a universal constant which has been determined to be approximately 1.5 by fitting overall energy multiplicity data. We may estimate a_0 by realizing that the average value of x is $X_0/n + 1$, while the average value of $(t - t_{\text{min}})$ is a^{-1} . For $n = 3$ at 28.5 GeV, and thus $\langle x \rangle \approx 1$, \bar{p}_{\perp} for outgoing protons is ≈ 0.34 GeV/c. Thus

$$\bar{a} = a_0 + 2a' \approx 9 \text{ GeV}^{-2}, \text{ or } a_0 \approx 7 \text{ GeV}^{-2}. \quad (6)$$

The value of k can be obtained from the experimental mean inelasticity factor

$$\left\langle \frac{p'_{\parallel}}{p_a} \right\rangle \approx 0.6, \quad (7)$$

which from (5) and (6) tells us that

$$k e^{-\langle x \rangle} \approx 0.4. \quad (8)$$

Now the average n is $g^2 X_o$, so the overall average value of x is roughly independent of X_o :

$$\langle x \rangle = \frac{X_o}{\langle n \rangle + 1} = \frac{X_o}{g^2 X_o + 1} \approx \frac{1}{g^2} \approx 0.67 . \quad (9)$$

Thus

$$k \approx 0.4 e^{-0.67} \approx 0.8 . \quad (10)$$

The coefficient $\sigma_{\text{tot}}^{p \rightarrow p'}$ can be determined by measuring at any given energy the total cross section for producing energetic baryons of the type p' .

The spectrum of secondary pions is also well defined in the multi-peripheral model but cannot be so easily expressed as in the simple Formula (1). It could easily be calculated on a numerical basis.