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NEUTRINO-ELECTRON SCATTERING AT NAL

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October 15, 1973

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Proposal

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Submitted by

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I. INTRODUCTION

The reactions of scattering neutrinos from electrons,

$$\nu_{\mu} + e^{-} \rightarrow \nu_{\mu} + e^{-},$$

$$\nu_{e} + e^{-} \rightarrow \nu_{e} + e^{-}$$

and their anti-neutrino counter parts, have not been observed experimentally. These reactions are predicted in the conventional V-A theory of weak interactions.¹ They also play an important role in some cosmological theories.

During the past few years, a great deal of theoretical effort has been concentrated on the development of a renormalizable gauge field theory which manifestly unifies the electromagnetic and the weak interactions.² In the context of this theory, there should exist a massive neutral boson, Z^0 , mediating the weak interactions in the same manner as the photons for the electromagnetic interactions. Both the V-A theory and the gauge theory give definite predictions on the cross sections of neutrino-electron scatterings. Specifically, the Z^0 -exchange makes the νe scatterings a first order process; while the W-boson exchange makes the reaction a second-order process.

The latest experimental development in high energy physics is the observation of events in neutrino-nucleus interactions in which no muons are produced.³ This important discovery indicates the existence of hadronic weak currents.

These important developments in physics give us a strong incentive to study neutrino-electron scattering at NAL where strong sources of high energy neutrinos are available. Such experiments will answer the following fundamental questions:

1. The existence of lepton-lepton current interactions.
2. The distinction between various weak interactions theories, such as;
 - a. The 4 fermion local V-A theory.

- b. Gauge theories.
 - c. The W-boson theory.
3. The existence of a lepton neutral current.

II. EXPERIMENTAL DESIGN

The basic experimental arrangement is shown in Figure 1. It is conceptually simple and novel. The main elements are specially designed multiple wire proportional chambers interlaced with steel radiators of ~ 1 mm in thickness. Neutrino-electron interaction can be uniquely identified by the observation of an electron shower with no hadrons present. We expect to measure the shower energy with about 15% accuracy and the shower mean direction to about 1 mr. The whole arrangement presents several hadron interaction lengths to the beam so that events with hadrons can be isolated. The accuracy of the shower direction measurement will enable us to sort out other types of backgrounds.

Two main types of background are:

1. π^0 production, and
2. $\nu n \rightarrow ep$.

The first one can be discriminated against by the presence of accompanying hadrons. The second process gives the same order of magnitude as the νe rate but such events can be recognized by the broader angular distribution.

A. The MWPC

The construction of the MWPC's is shown in Figure 2. In this construction, the anode wires are connected together to give information on the pulse-height induced by the electron shower. Signals at the anode wires will also provide the trigger. The two cathode planes are used to measure the position

of the electron shower. The "induced" signal on the cathode planes will be read out by a novel "delay-line" technique.⁴ The active area of the chamber is 1m x 1m. The mechanical construction is made easier by the use of aluminum plates to support the cathode-wires, which also serve as part of the radiators. The cathode-wires are glued on and insulated from the aluminum plates.

Our laboratory has already built 12 MWPC's 1m x 1m in size for the μ -p inelastic scattering experiment at NAL. The required electronics and mechanical techniques for building these chambers are now fully developed. We estimate the cost of the chambers for the proposed ν -e experiment at \$3000 each.

B. The Target

The steel plates between the chambers and the aluminum walls of the MWPC's will serve both as the target for the neutrino interactions and as the radiators for the recoil electrons. The total thickness will be 1 r.l. per cell.

Initially, we would like to build 50 cells, made of one MWPC and one steel plate each, to do the experiment. Later, this number can be increased to 100. If economy permits, we will replace a few steel plates with Pb-glasses. This should enable us to double check the electron energy measurement of the MWPC's.

C. The Trigger

Since the chambers are operating in the proportional mode we will use the anode signals to sample the electron shower energy. If the anode signals from several adjacent chambers show the characteristics of such an electron shower this signal will be used as a trigger to latch the information

in all the equipment preparatory to data readout.

Details of electronic logic, cosmic ray vetoes, etc., will not be discussed here.

D. Counting Rate

To estimate the event rate we use the neutrino spectrum shown in Figure 3 and a basic cross section

$$\sigma_0(E_\nu) = 10^{-41} E_\nu(\text{in GeV}) \text{ cm}^2.$$

For 50 cells 1m x 1m in size each with 1.38 cm of steel and 2 cm of aluminum

$$N_e = 6 \times 10^{23} \times \left(\frac{10.8 \times 26}{56} + \frac{5.4 \times 13}{27} \right) \times 50 = 3 \times 10^{26} \text{ electrons/cm}^2.$$

From Figure 3, the integral

$$\int_5^{120} N_\nu(E) E \, dE = 1.7 \times 10^{11} \text{ neutrinos/m}^2.$$

The counting rate is given by

$$N = N_e \int_5^{120} N_\nu(E) \sigma_0(E) \, dE = 0.3 \text{ events/hour,}$$

for 10^{13} protons of 350 GeV on target, 600 pulses per hour. At 1000 hours running, ~~this event~~ provide 137 events.

The expected cross-sections from different theories are shown in Figures 4 and 5. It is quite clear that 1000 hours running should provide a very significant amount of information in this totally unknown area of physics.

E. Beam

The neutrino beam at NAL can be reused many times. This proposed experiment can be set up behind Exp. 1A and run in parallel with other neutrino experiments under any beam conditions. ,

III. EXPERIMENTAL PLAN AND REQUEST TO NAL

Upon approval of this experimental proposal, we can immediately start the construction work of the MWPC's. The job can be completed within a year and the experiment can be started 3 months after.

The cost for the MWPC's and the associated electronics will be borne by the investigators from the Enrico Fermi Institute of the University of Chicago. We anticipate requesting some supplementary equipment funding from the NSF.

We request NAL to supply the following items:

- 1) 50 steel plates, 1m x 1m x 1.4 cm in size.
- 2) A suitable enclosure for the equipment.
- 3) Fast trigger electronics from PREP.
- 4) A PDP-11 computer for data acquisition.

Finally, we request 1,000 hours of running time, all parasitic.

REFERENCES

1. R. P. Feynman and M. Gell-Mann, Phys. Rev. 109, 193 (1958).
2. The theoretical review on this subject and original references can be found in the Rapporteur's talk given by B. W. Lee, and the experimental review by D. H. Perkins; Proceedings of the XVI International Conference on High Energy Physics at Chicago and Batavia, U. S. A., 1972.
3. The experimental results from CERN and NAL were recently announced at the International Symposium on Photons and Electrons, Bonn, Germany, 1973.
4. This technique has been developed by T. A. Nunamaker at the Enrico Fermi Institute.

FIGURE CAPTIONS

- Figure 1: Schematic arrangement of the experimental setup.
- Figure 2: Schematic diagram of the construction of the multiple-wire proportional chamber.
- Figure 3: Neutrino spectrum at NAL.
- Figure 4: Cross sections for ν_e scatterings given by different theories. The parameter e^2/g^2 is defined as $\sin^2\theta$ in Weinberg's theory.
- Figure 5: The energy dependence of ν_e scattering cross sections.
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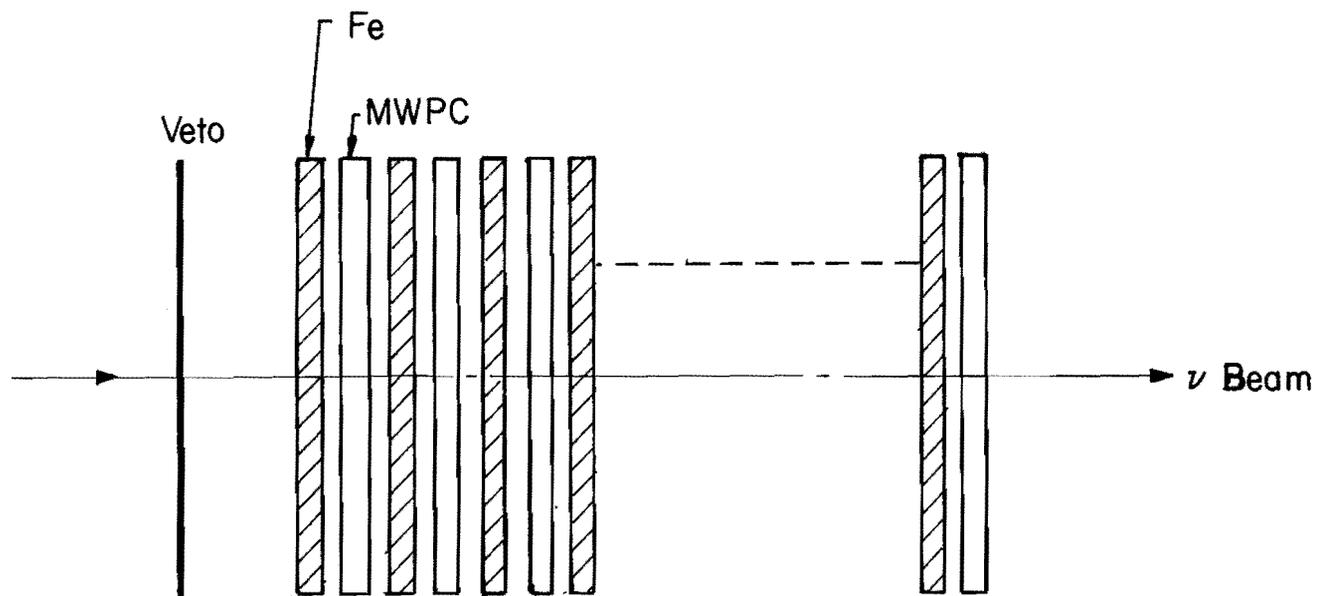


Figure 1 Schematic Arrangement of the Experimental Setup

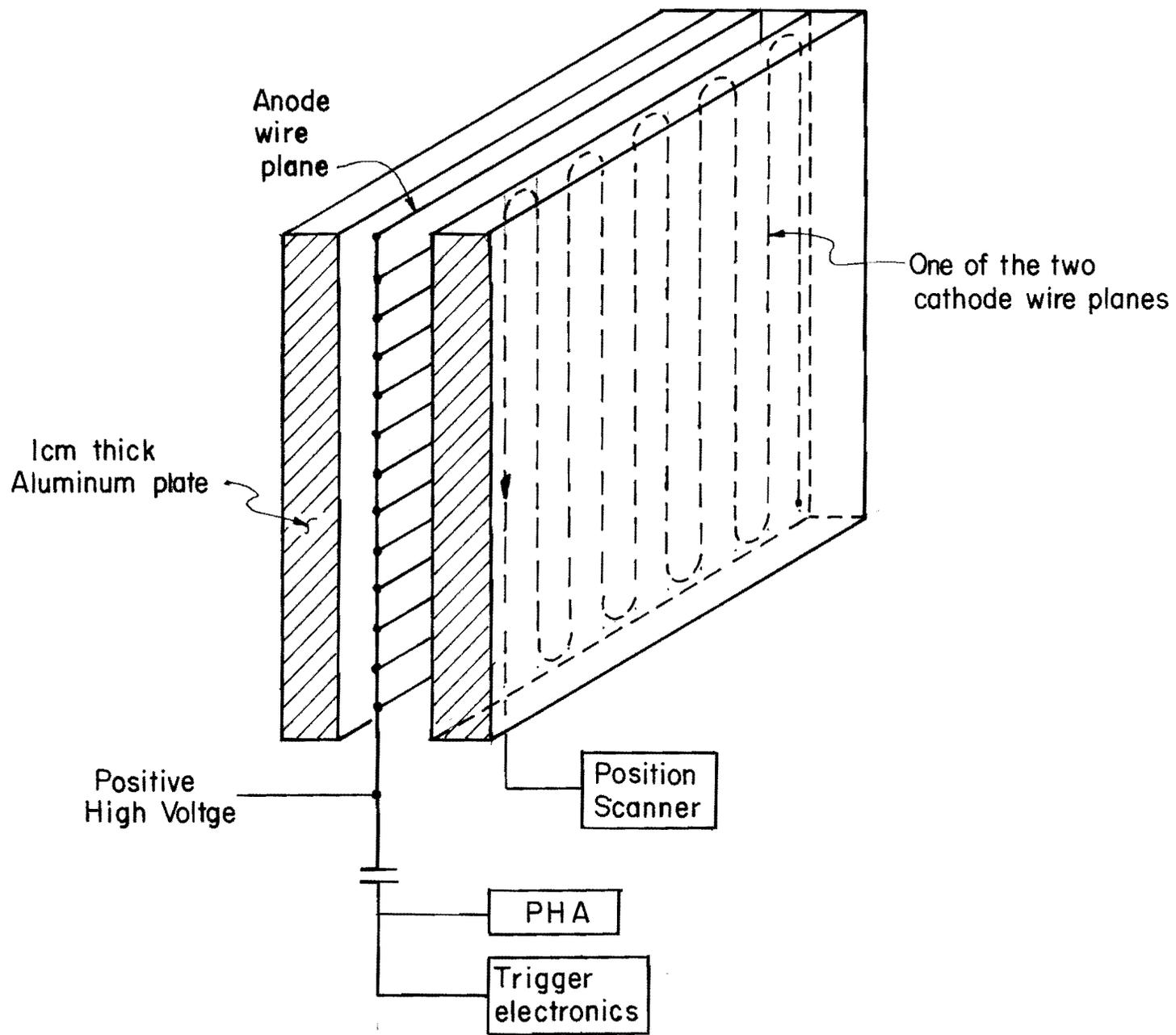


Figure 2 Schematic diagram of the MWPC construction

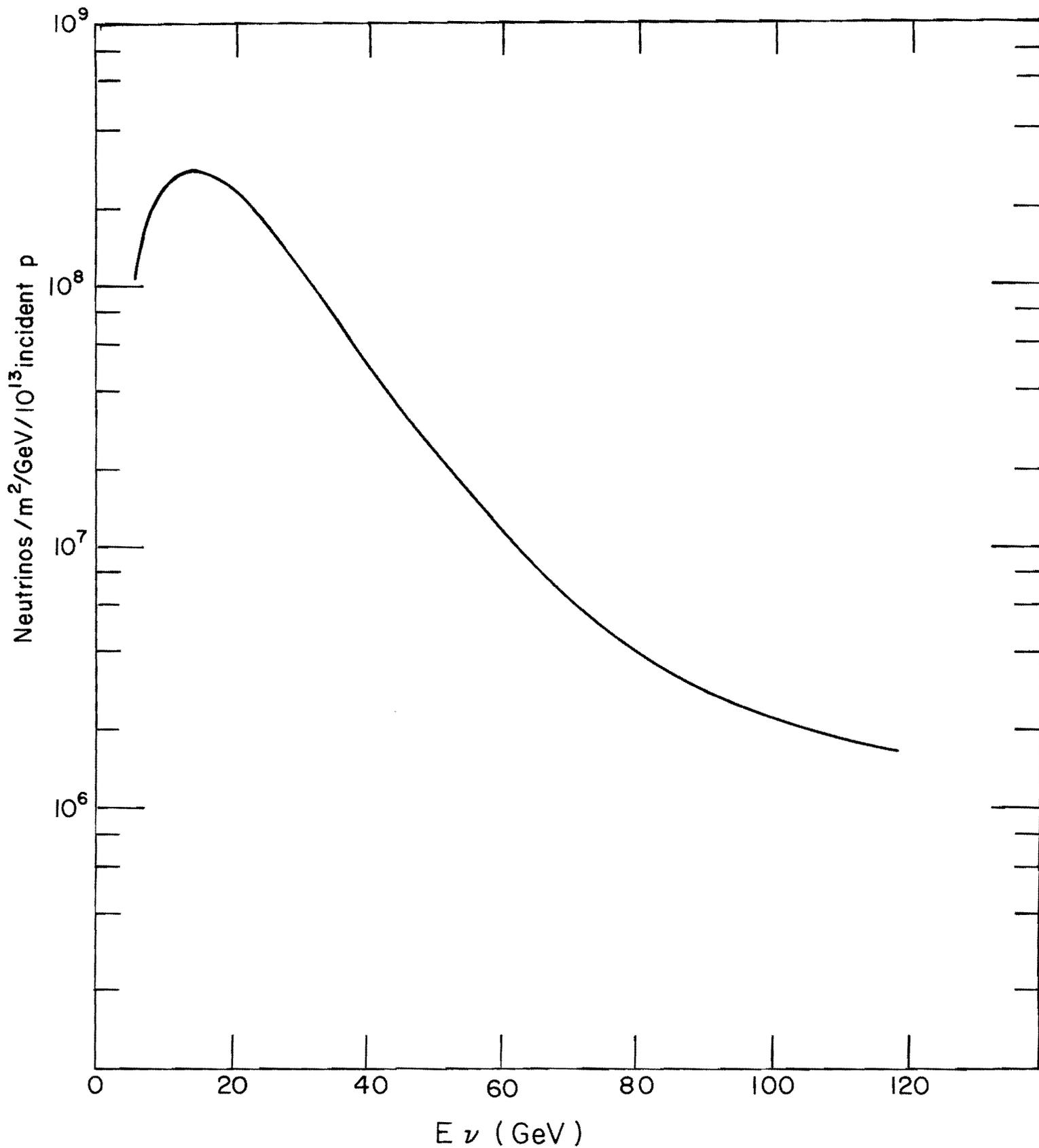
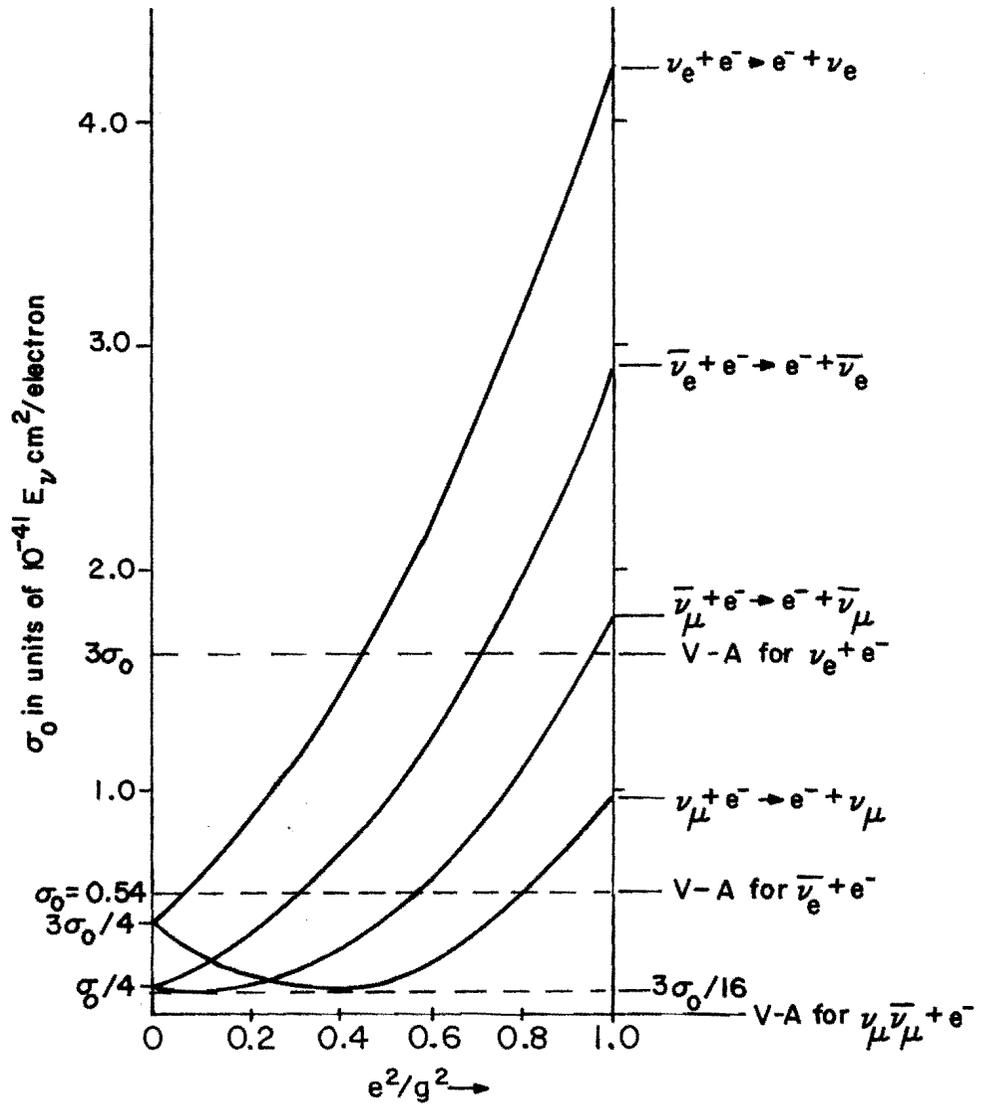


Figure 3 Neutrino Flux Produced
by 350 GeV Protons at NAL



**Figure 4 Neutrino-Electron Scattering
Cross-Sections in Weinberg Model**

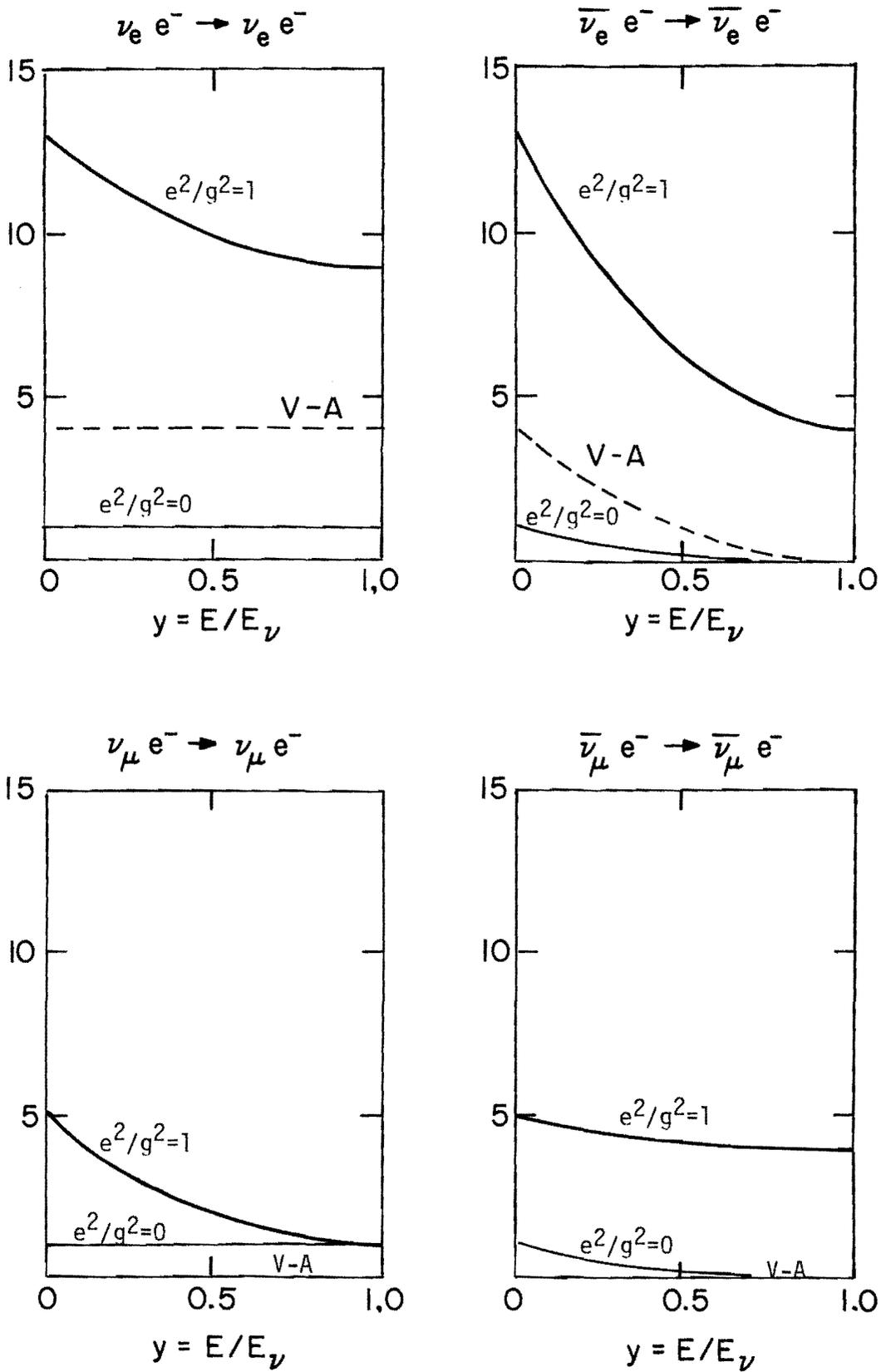


Figure 5 $d\sigma/dE$ for ν - e scattering
(units are $0.4 \times 10^{-41} \text{ cm}^2/\text{GeV}$)

Addendum to Proposal 253

Detection Technique

In this addendum, we describe the technique to be used in detecting the recoil electrons from high energy ν_e scattering. Experience at SLAC is reported, as well as the test we recently made at the Wilson Synchrotron Laboratory of Cornell University.

A. Principle of Detection

To identify ν_e scatters, events which occur at the 10^{-3} level of the dominant νN interactions, a detector which unambiguously singles out electromagnetic interactions is clearly necessary. Moreover, since hadronic reactions produce electromagnetic showers via neutral pions, the extreme forward going character of the electrons from ν_e scatterings must be recognized by this detector.

It is well established now that the energies of high energy electrons or photons can be measured very well by an electromagnetic shower detector of approximately 20 radiation lengths in thickness. The energy of the incident particle is fully contained in the detector. However, such a detector is not thick enough to contain the high energy hadronic cascade because 20 r.l. is equivalent to approximately only 2 strong interaction collision lengths. In the initial 3 r.l. of the detector, the

energy loss of the electromagnetic particles is several times that of the minimum ionizing particles. Therefore, an EM shower detector in conjunction with a $(\frac{dE}{dX})$ measurement at its front end can determine the energy of an electron or photon and at the same time reject hadrons.

In Proposal #253, we try to apply this simple principle to detect the recoil electrons from ν_e scattering. The detector we plan to use is sensitive only to the electromagnetic showers induced by high energy electrons or photons. It should be "blind" to the nuclear scattering events because it capitalizes on the radical differences between electromagnetic and hadronic shower development. Of course, as in all neutrino experiments, the detector is part of the target. After a description of a SLAC detector in (B), our solution is described in (C).

B. Experience at SLAC

At SLAC, two large size shower detectors were built for the 8- and 20-GeV/c spectrometers. Each one of these detectors was 20 r.l. thick. It consisted of 20 slabs of Pb-plate and lucite sandwich, 1 r.l. thick each. The first 3 slabs were used also for $(\frac{dE}{dX})$ measurement. Signals from the whole assembly were electronically added to provide the total energy determination. The arrangement is shown schematically in Figure 1.

For the SLAC deep inelastic electron scattering experiment, the π/e ratio accepted by the heavily shielded spectrometer was at times as bad as 400/1. Therefore, pion rejection was very important to the success of the experiment. For a typical run, the raw pulse height spectrum is shown in Figure 2. The minimum ionization peak was greatly suppressed by strong attenuation, but $\sim 30\%$ still remained. The pulse height spectrum of the front (dE/dX) detector is shown in Figure 3. If one demands a (dE/dX) cut, as shown in the Figure, then all the pion contamination vanishes in the total energy spectrum. The resultant total energy spectrum of the electrons is shown in Figure 4. We found that pions could be rejected to the level of $\sim 10^{-3}$, with an overall electron detection efficiency ranging from 67% at 2 GeV to 89% at energies above 4 GeV.

It should be remarked that the SLAC shower detectors were built in 1965 and have successfully operated since 1966. Although no details were ever published, the technique has been widely adopted in high energy physics. For example, a similar system was used at SPEAR for the Ψ discovery.

C. ν_e Detection at Fermilab

In order to do the ν_e scattering experiment at the Fermilab successfully, one has to measure (1) the recoil electron energy,

(2) the recoil electron direction. Also one has to be able to reject the nuclear scattering events which will be $\sim 10^3$ times more abundant as mentioned before. The Pb-lucite sandwich described above is a device much less expensive than Pb-glass or NaI(Tl) crystals. However, the cost for using it to do a ν_e scattering experiment is still formidable. In addition, it does not have good spatial resolution. The economic consideration led us to the development of a shower detector as described in Proposal #253. It basically works the same way as a Pb-lucite sandwich does. The only difference is that instead of the lucite-photomultiplier tube assembly, less expensive proportional chambers are used to allow accurate shower direction measurements. In each module, the signal from the anode will be used for pulse height analysis. The delay-line readout from the cathode planes will determine the centroid position of the electron shower.

The signature of a possible ν_e scattering event is also provided by the anode signals of the proportional chambers. The trigger scheme is illustrated in Figure 5. The analog signals from every ten or so successive anode planes are added together. The triggering threshold is adjusted by the calibrated attenuator in front of the discriminator which will accept electrons above a certain energy and reject most of the signals caused by nuclear scatterings. We plan to build the electronic logic required for this purpose at the Enrico Fermi Institute.

D. The Test at Cornell University

During November, 1974, five members of the collaboration (R. Fine, R. Heisterberg, L. Mo, T. Nunamaker, and S. C. Wright) did a test run with 5 MWPC's at the Wilson Synchrotron Laboratory, Cornell University.

The test set-up at Cornell University is shown in Figure 6. The bremsstrahlung from the 12 GeV electron synchrotron was incident upon a thin converter in front of a dipole magnet. The asymmetrically produced electron pair was energy selected by a magnet and a collimation channel to an energy resolution of $\sim 10\%$. This electron beam was used to test the shower detector we intend to use for the ν_e scattering experiment at Fermilab.

We installed five 1m X 1m MWPC's exactly the way we intent to do the ν_e scattering experiment as shown in Figure 7. Since there were only five chambers, the assembly was not thick enough to contain all the electron shower energy. Hence, the energy resolution was consequently not ideal. We recorded the pulse height spectrum of each individual chamber. The total energy spectrum was obtained by software addition. Figure 8 shows the pulse height spectrum at 8 GeV from the first chamber. Figure 9 shows the total energy spectrum at 8 GeV from the sum of five chambers. Figure 10 gives the energy calibration from 2.5 through 9.5 GeV.

As shown in Figure 11, on each cathode plane there are five tap points to measure the particle trajectory position. Because of time and financial limitations, we did not build the electronic centroid finder. Instead we use standard commercial start-stop type, time-to-digital converters (TDC) to measure the transit time. By doing it this way, the positions measured were actually the edges, instead of the center, of the shower core. A typical picture of a shower induced by 8 GeV electrons is shown in Figure 12. Figure 13 shows the spatial resolution which was obtained by taking the average values as the centroid positions and fit with a straight line.

E. Conclusion

The test run at Cornell University indicated that the proposed detector is a sound one. One can measure the electron energy and determine simultaneously the shower direction to better than 8 mr. If the centroid finders are used instead of the TDC's, we anticipate a factor of ~ 4 improvement in direction determination.

We wish to thank Dr. R. Humphrey, Dr. M. Tigner, Professor B. McDaniel, and the staff of the Wilson Synchrotron Laboratory for their help and hospitality. It has always been a pleasure to visit Cornell University.

Figure Captions

- Figure 1 Electron shower detector used in the 20-GeV/c spectrometer at SLAC. The detector is a Pb-lucite sandwich.
- Figure 2 Pulse height spectrum of the SLAC shower detector with a mixture of pions and scattered electrons of 6.4 GeV.
- Figure 3 Pulse height spectrum of the third slab of the SLAC shower detector at 6.4 GeV. The cut used for electron-pion separation is also shown.
- Figure 4 Pulse height spectrum of the SLAC shower detector at 6.4 GeV. Pion contaminations has been removed by imposing a (dE/dX) cut as shown in Figure 3.
- Figure 5 Trigger scheme for the proposed ν_e scattering experiment at Fermilab.
- Figure 6 Arrangement of the test setup at the Wilson Synchrotron Laboratory of Cornell University.
- Figure 7 Details of the test setup. The steel radiator is $3/8$ in thick each. The distance between chambers is 13". Each chamber consists of two aluminum plates, $1/4$ in thick each.
- Figure 8 Pulse height spectrum of the first proportional chamber with 8 GeV electrons.
- Figure 9 Total energy spectrum given by 5 MWPC's and 8 GeV electrons. This spectrum was obtained by software addition of spectra from five chambers.
- Figure 10 Energy calibration curve of the 5-MWPC assembly. The channel numbers are those of the centroids of the total energy spectra as shown in Figure 9.
- Figure 11 Schematic of the delay-line readout on the cathode plane. The positions seen by the TDC's are the edges of the shower.

Figure 12

A typical picture from the cathode plane delay-line readout. In this illustration, the triangles indicate the outer edges of the shower measured by the TDC's, and the crosses their average values which were taken as the centroids.

Figure 13

Typical spatial resolution given by the electron shower.

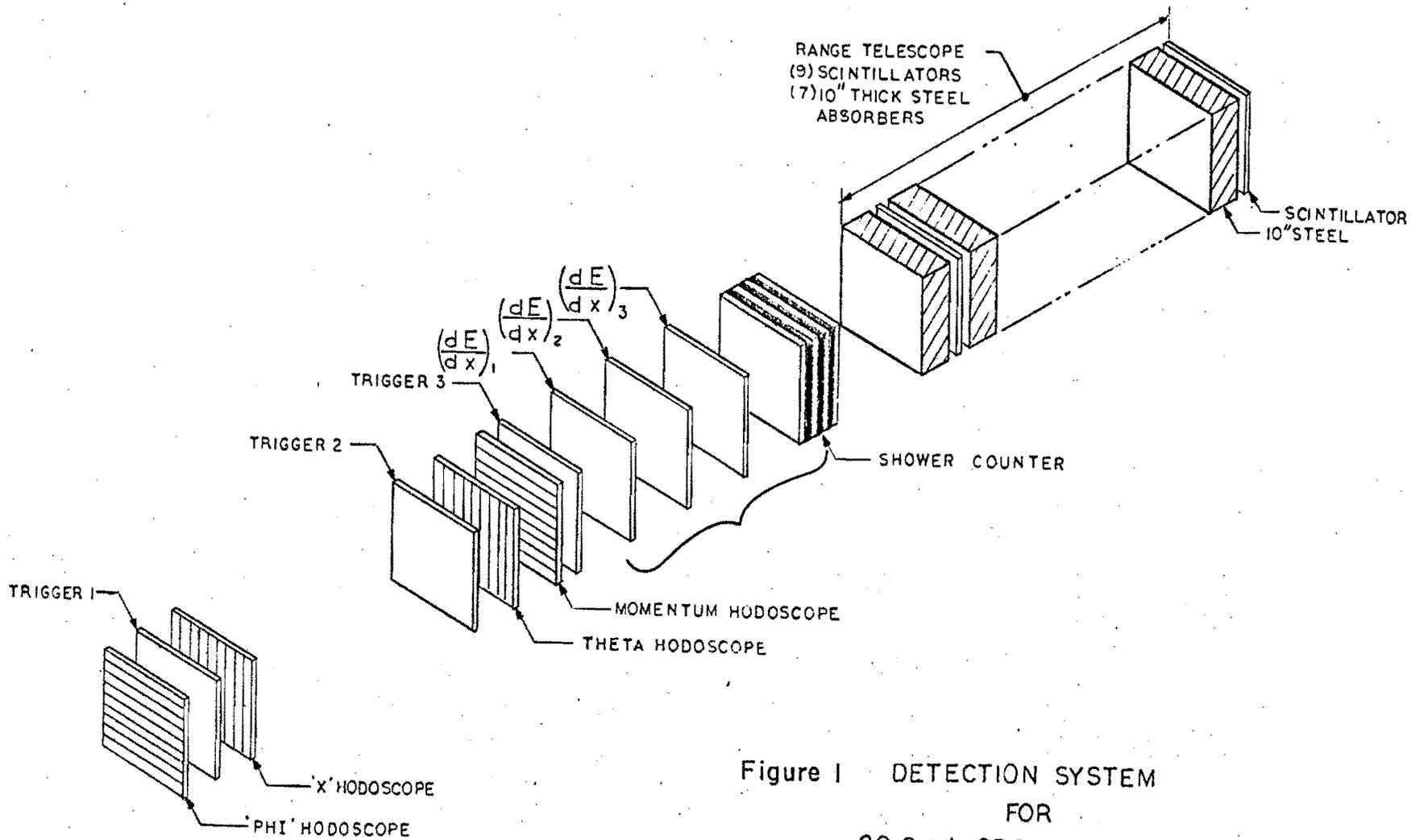
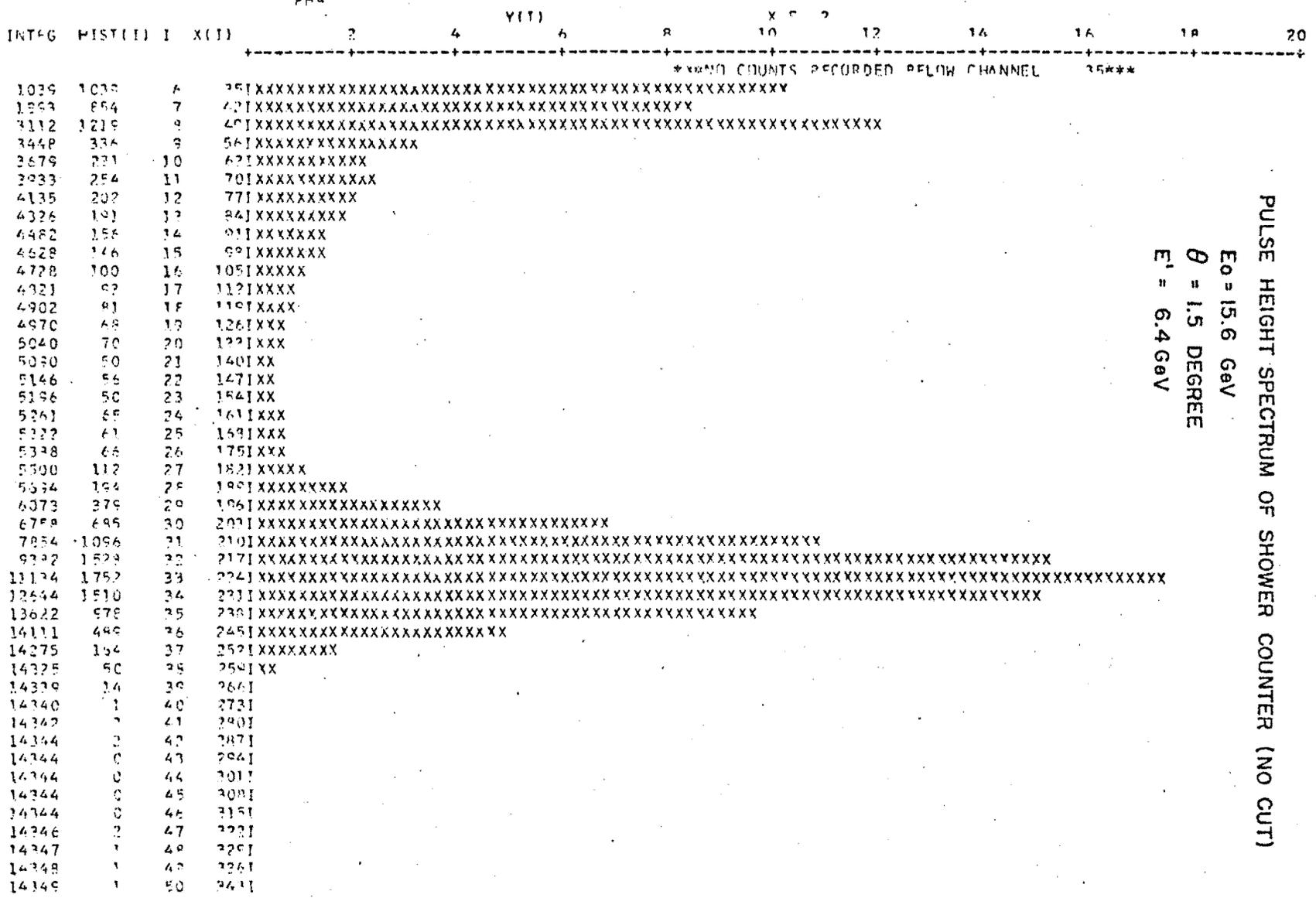


Figure 1 DETECTION SYSTEM
 FOR
 20 GeV/c SPECTROMETER



PULSE HEIGHT SPECTRUM OF SHOWER COUNTER (NO CUT)
 $E_0 = 15.6 \text{ GeV}$
 $\theta = 1.5 \text{ DEGREE}$
 $E_1 = 6.4 \text{ GeV}$

Figure 2

14374 14349 0 0 0
 OBSERVED PEDESTAL LIPS AT CHANNEL NO. = 35
 TA-CUT= 176.6449 ,ICUT= 7 ,BY MOMENTS (ON CHANNELS 176 TO 301); PEAK CENTRE= 220, SIGMA= 15.94 PEAK TOTAL (OBS)= 9083
 SHOWER COUNT# FOR CONDITION 9 = ALL EVENTS INCLUDING PIONS

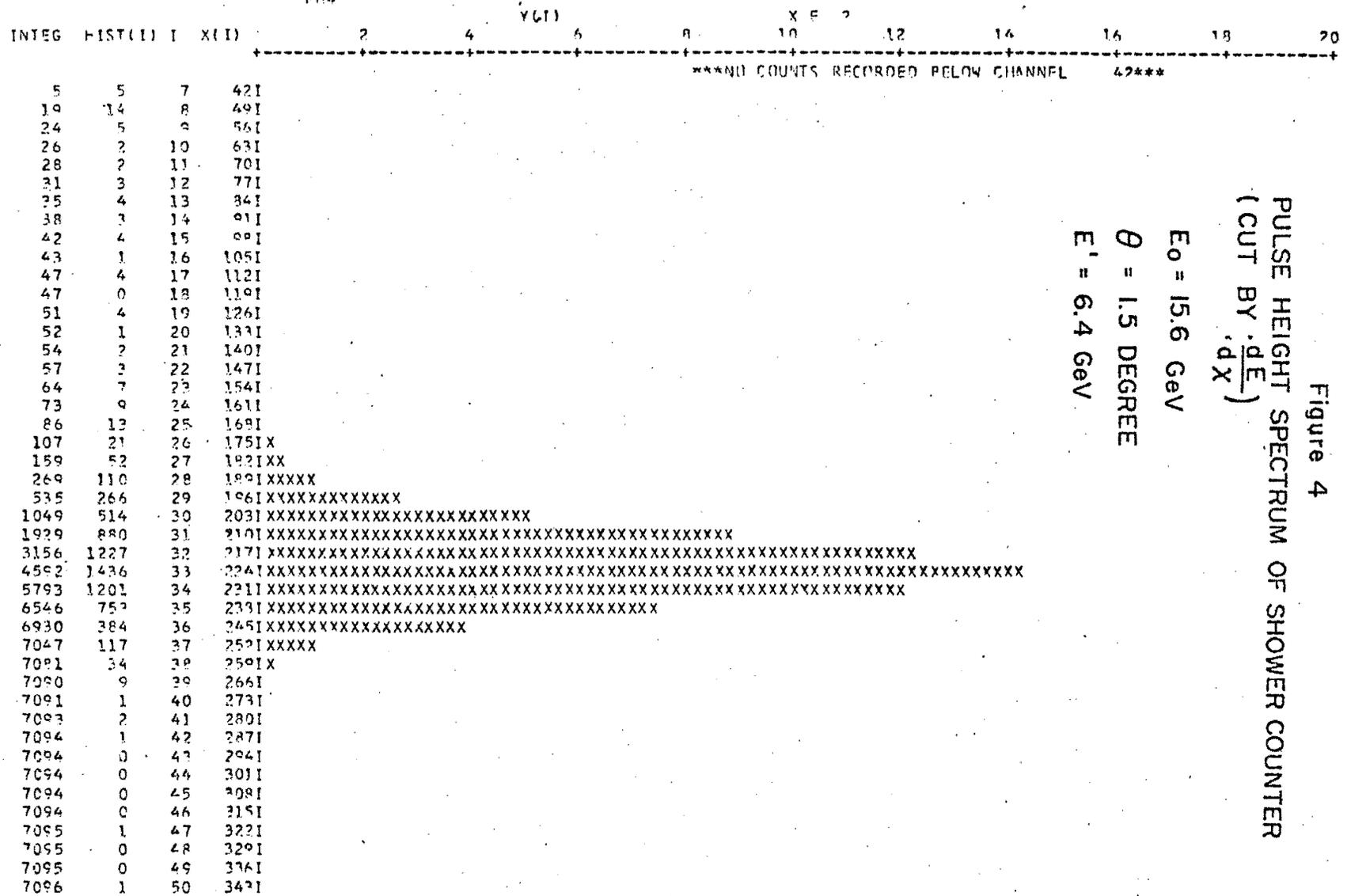
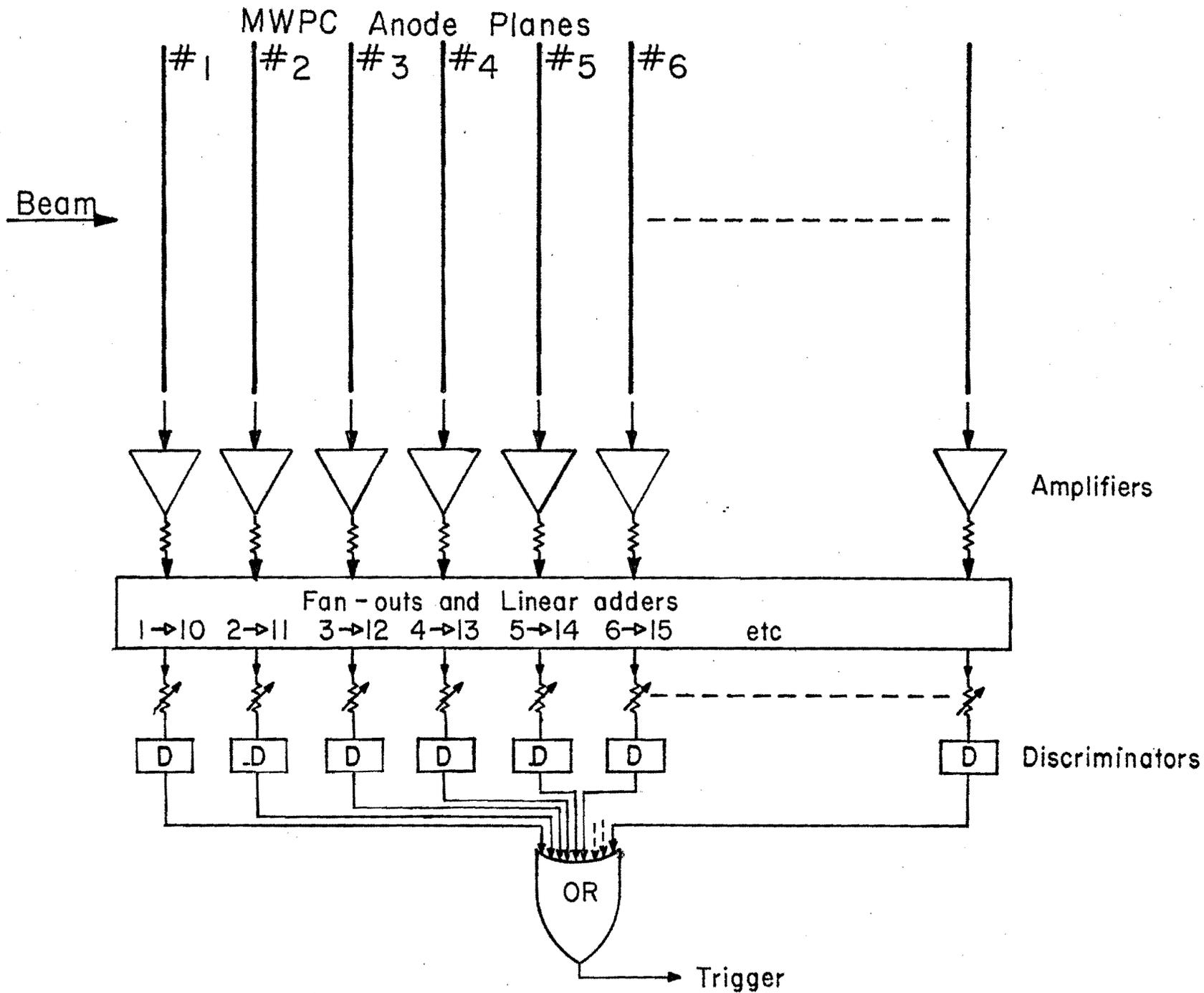


Figure 4
 PULSE HEIGHT SPECTRUM OF SHOWER COUNTER
 (CUT BY $\frac{dE}{dX}$)
 $E_0 = 15.6$ GeV
 $\theta = 1.5$ DEGREE
 $E' = 6.4$ GeV

7112 7096 0 0 0
 TA-CUT= 176.6449 ,ICUT= 7 .BY MOMENTS(ON CHANNELS 176 TO 201);PEAK CENTRE= 221.SIGMA= 14.59 PEAK TOTAL(CRS)= 7021
 SHOWER COUNTER FOR CONDITION 3 = YES CCL YES DE/DX AND GOOD P-T CODES



The trigger schematic for νe scatterings

FIGURE 5

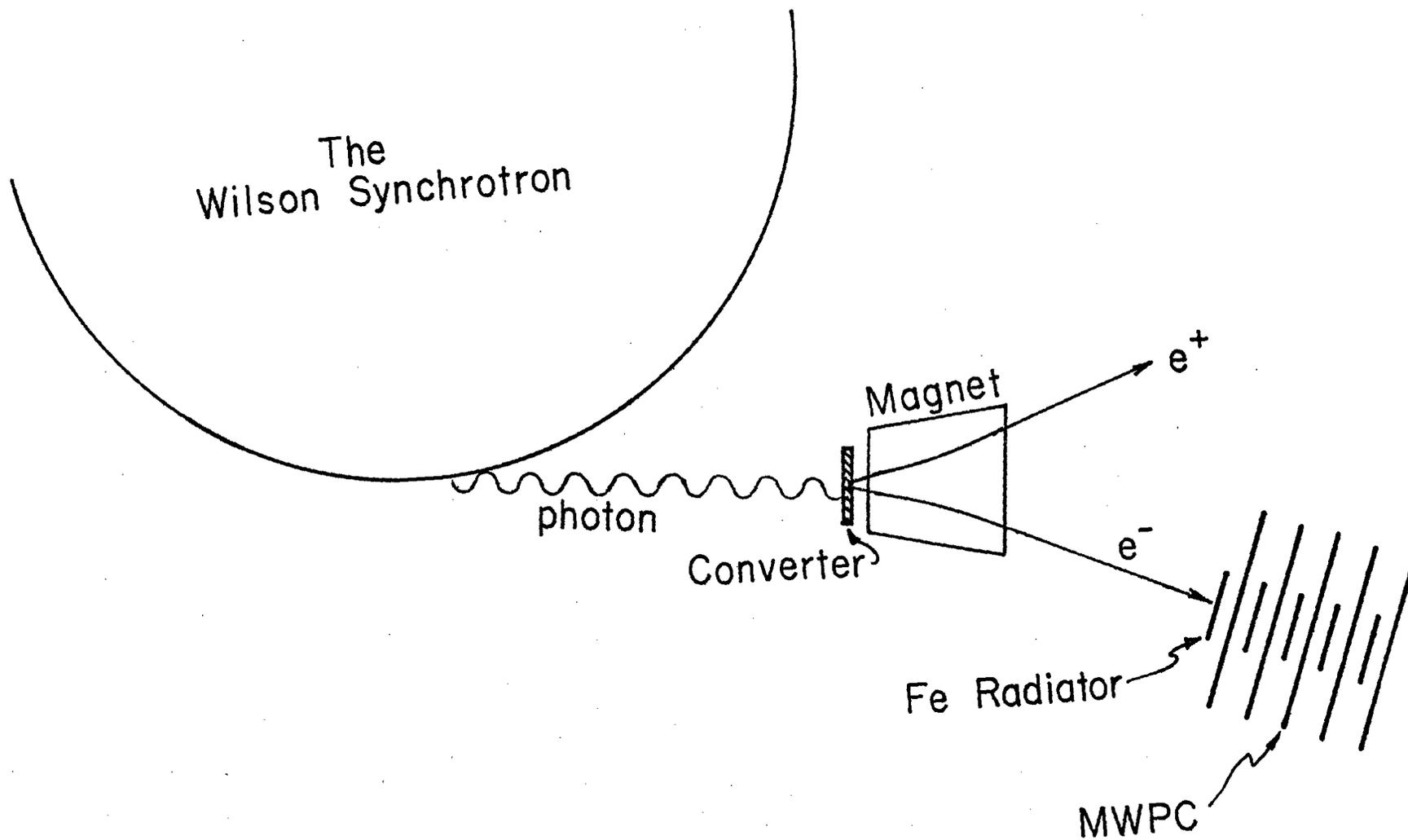


FIGURE 6

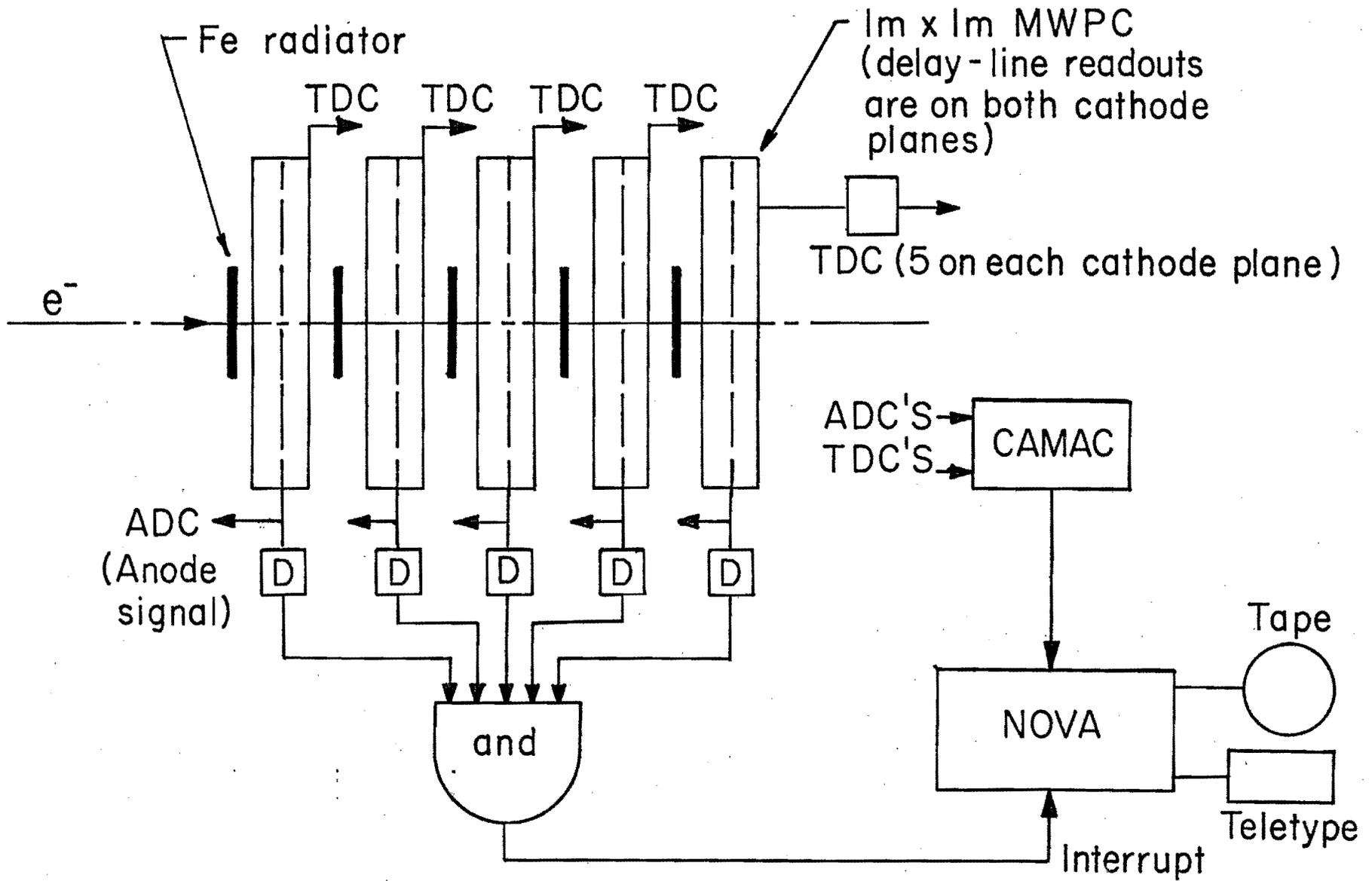


Figure 7

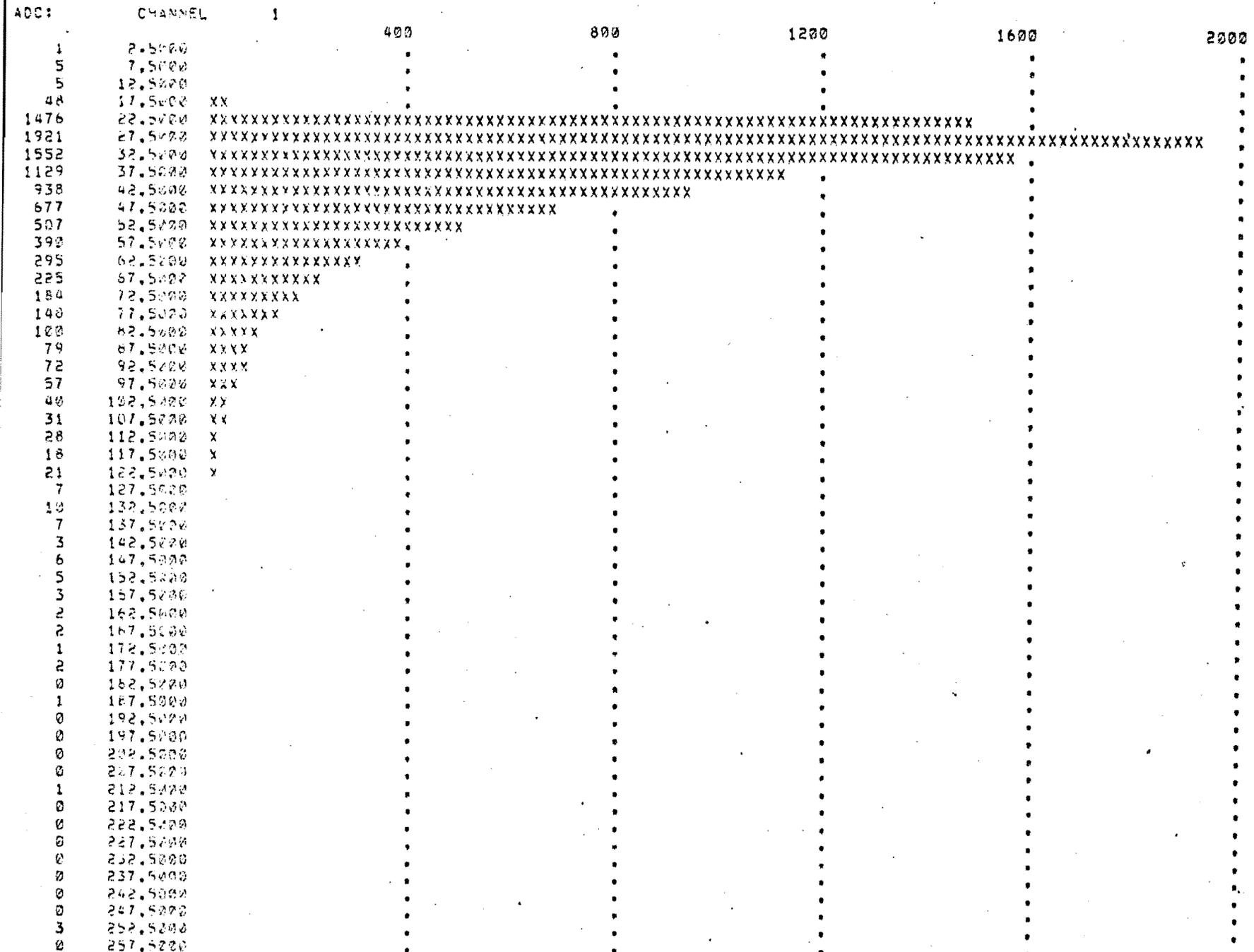


Figure 8
Pulse Height spectrum of the first
chamber with 8 GeV electrons

10000	EVENTS IN PLOT.	MEAN=	40.481248	AND SIGMA=	19.928125				
TOTAL	9660.	MEAN	38.679752 +-	0.152804	SIGMA	15.033975 +-	0.108065	CHIDF&NDF	1681.4280 15
TOTAL	9661.	MEAN	38.278044 +-	0.147226	SIGMA	14.425925 +-	0.104120	CHIDF&NDF	1634.0157 14
TOTAL	9661.	MEAN	38.278044 +-	0.147226	SIGMA	14.425925 +-	0.104120	CHIDF&NDF	1634.0157 14

ADC: TOTALS

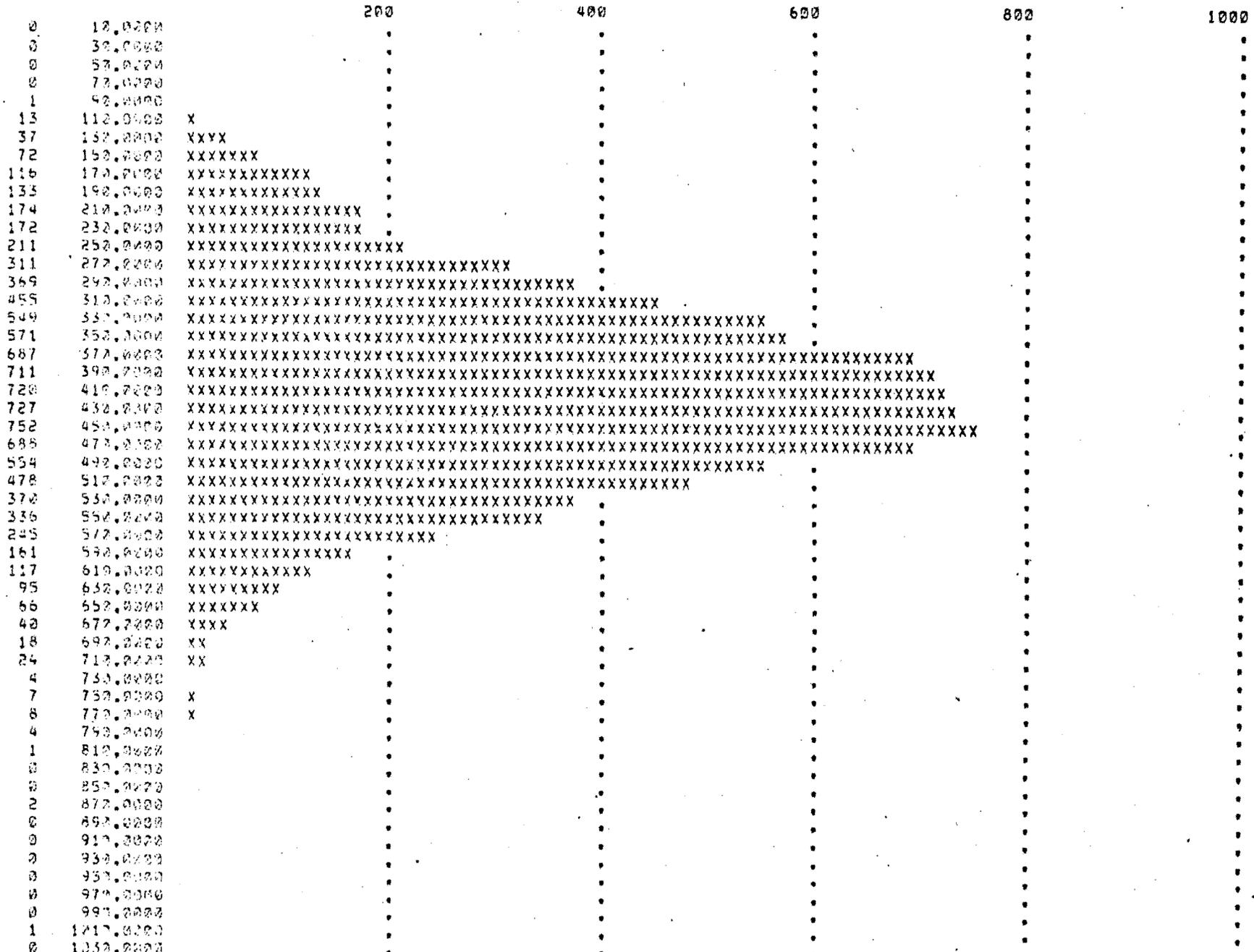


Figure 9
Total energy spectrum with 5 MWPC's
and 8 GeV electrons

10000 EVENTS IN PLOT.	MEAN=	408.231998	AND SIGMA=	110.247476					
TOTAL	9977.	MEAN	407.853062 +- 1.090729	SIGMA	108.947441 +- 0.771379	CHIDF&NDF	3.7045	30	
TOTAL	9977.	MEAN	407.853062 +- 1.090729	SIGMA	108.947441 +- 0.771379	CHIDF&NDF	3.7045	30	
TOTAL	9977.	MEAN	407.853062 +- 1.090729	SIGMA	108.947441 +- 0.771379	CHIDF&NDF	3.7045	30	

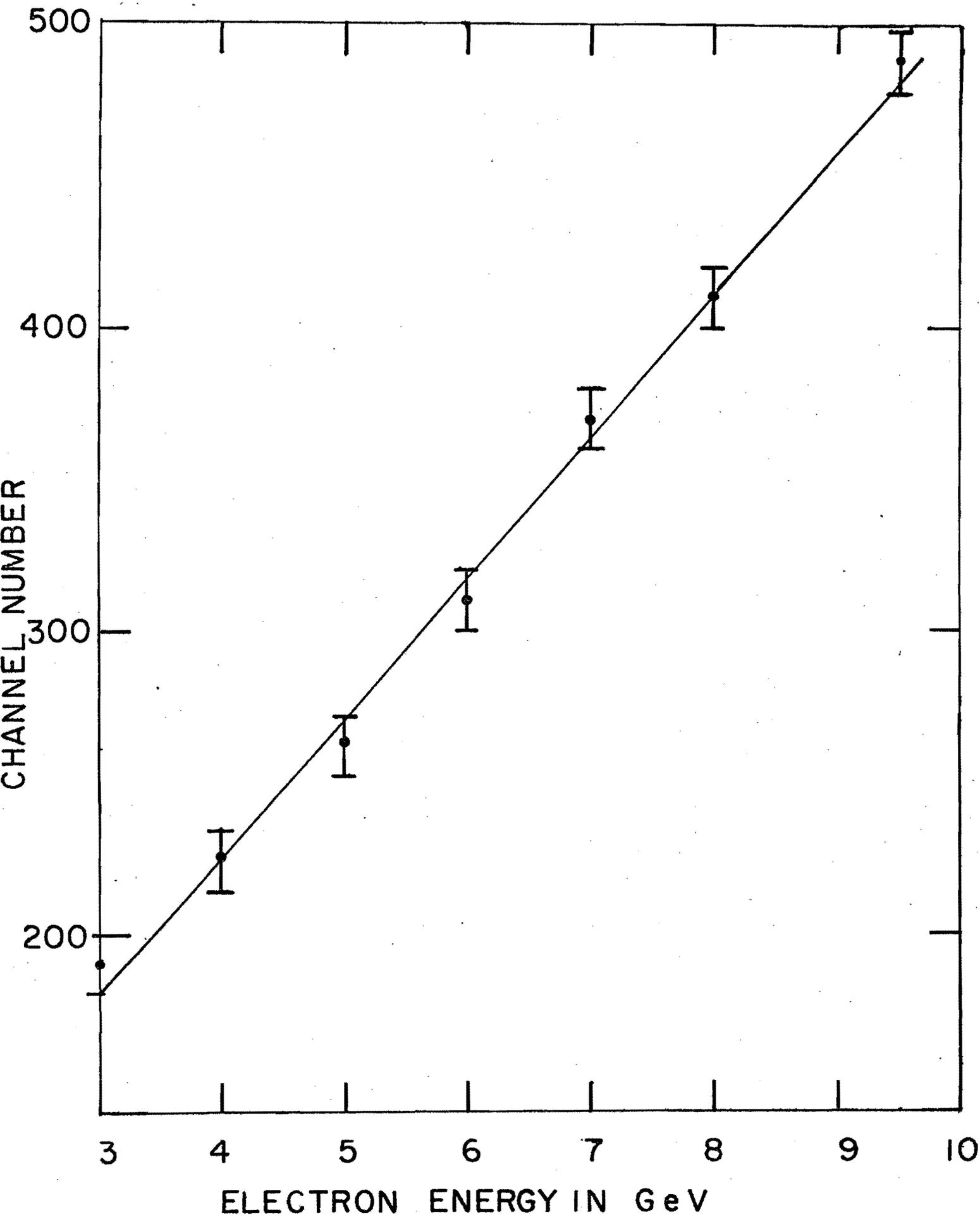


FIGURE 10

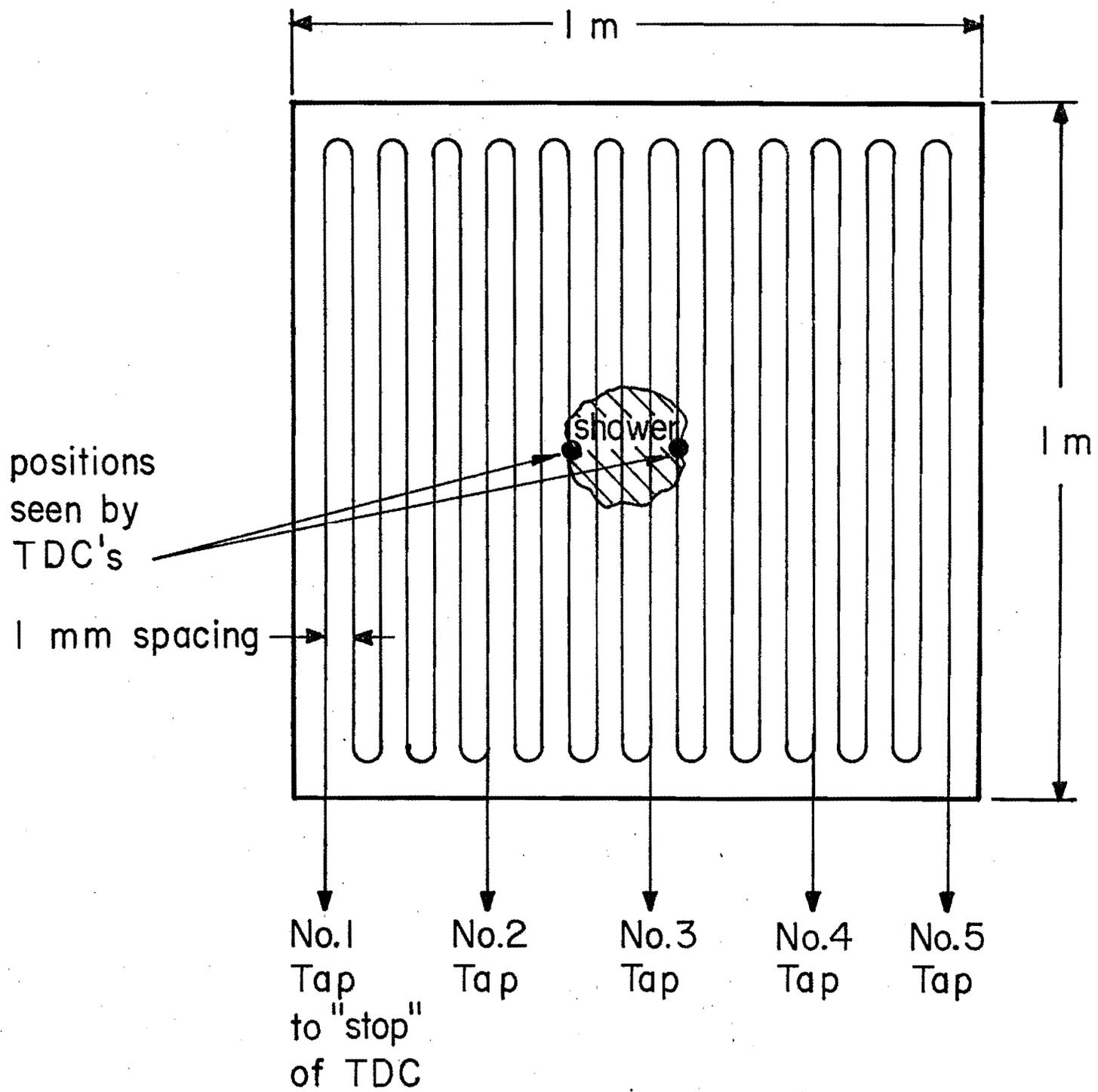


FIGURE II.

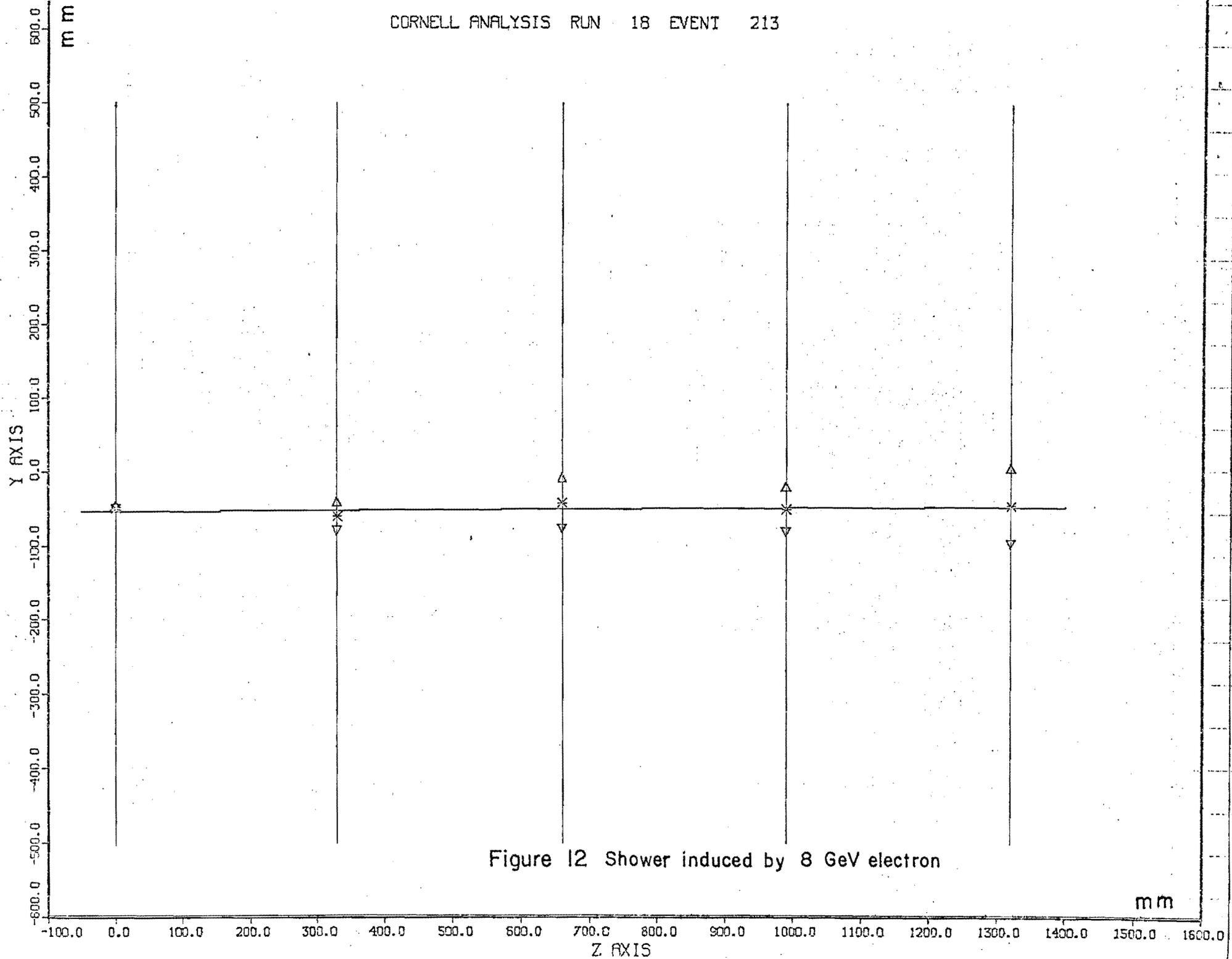


Figure 12 Shower induced by 8 GeV electron

HISTOGRAM # 1

DEVIATION HIST: GRAMS

SUM X10	Y X10	I	X	2	4	6	8	10	12	14	16	18	20
0.0	0.0	1	-50.00	:									
0.0	0.0	2	-48.00	:									
0.0	0.0	3	-46.00	:									
0.0	0.0	4	-44.00	:									
0.0	0.0	5	-42.00	:									
0.0	0.0	6	-40.00	:									
0.0	0.0	7	-38.00	:									
0.0	0.0	8	-36.00	:									
0.0	0.0	9	-34.00	:									
0.0	0.0	10	-32.00	:									
0.0	0.0	11	-30.00	:									
0.0	0.0	12	-28.00	:									
0.0	0.0	13	-26.00	:									
0.0	0.0	14	-24.00	:									
0.0	0.0	15	-22.00	:									
0.100	0.100	16	-20.00	:X									
0.200	0.100	17	-18.00	:X									
0.500	0.300	18	-16.00	:XX									
0.800	0.300	19	-14.00	:XX									
1.600	0.800	20	-12.00	:XXXX									
2.500	0.900	21	-10.00	:XXXXX									
5.000	2.500	22	-8.00	:XXXXXXXXXXXX									
10.400	5.400	23	-6.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
16.200	5.800	24	-4.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
23.300	7.100	25	-2.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
31.800	8.500	26	0.0	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
41.400	9.600	27	2.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
52.500	11.100	28	4.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
59.700	7.200	29	6.00	:XXXXXXXXXXXXXXXXXXXXXXXXXXXX									
62.600	2.900	30	8.00	:XXXXXXXXXXXXXXXXXXXX									
64.400	1.800	31	10.00	:XXXXXXXXXX									
65.300	0.900	32	12.00	:XXXXXX									
65.700	0.400	33	14.00	:XX									
66.100	0.400	34	16.00	:XX									
66.100	0.0	35	18.00	:									
66.100	0.0	36	20.00	:									
66.200	0.100	37	22.00	:X									
66.200	0.0	38	24.00	:									
66.200	0.0	39	26.00	:									
66.200	0.0	40	28.00	:									
66.200	0.0	41	30.00	:									
66.200	0.0	42	32.00	:									
66.200	0.0	43	34.00	:									
66.200	0.0	44	36.00	:									
66.200	0.0	45	38.00	:									
66.200	0.0	46	40.00	:mm									
66.200	0.0	47	42.00	:mm									
66.200	0.0	48	44.00	:									
66.200	0.0	49	46.00	:									
66.200	0.0	50	48.00	:									

Figure 13
Spatial resolution with
4 GeV electrons

UNDER: 0.0 < AVG > = 1.7556
OVER: 0.0 < SIG > = 5.5134

HISTOGRAM SHOWN FOR PLANE # 1

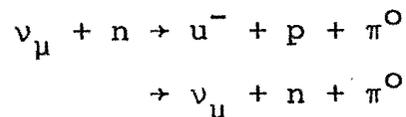
ADDENDUM III TO E-253

N. Booth, R. Heisterberg, L. Mo, A. Skuja

The setting up of E-253 is now almost complete in the Wonder Building and parasitic tuning is underway. The apparatus consists of 41 multi-wire proportional chambers, 1 m² in active area, with 1 r.l. of Al between successive chambers. In addition we have a scintillation hodoscope at the upstream end to veto/tag incoming muons and two hodoscopes at the downstream end separated by 4 ft of steel to veto/tag muons produced in the apparatus. At present we have no capability to measure the momenta of these muons. Now that we are ready to start our $\nu_{\mu}e$ scattering experiment we feel that it is time to reconsider the supplemental request which we made in Addendum II to our proposal.

We would like to supplement this setup with an iron magnet and drift chambers to measure the momenta of forward-going muons. Not only will this help us to better understand the background to our primary reaction of interest, $\nu_{\mu}e$ scattering, but will enable us to study as byproducts several interesting processes in a unique way. In particular, our apparatus is well suited to measure coincidences between electromagnetic showers and muons. With a small amount of additional equipment, we will be able to measure the following two reactions:

- 1) Single π^0 production by charged and neutral currents,



This provides information on the isospin structure of the weak current.

2) Inverse muon decay,

$$\nu_{\mu} + e^{-} \rightarrow \nu_e + \mu^{-}$$

$$\bar{\nu}_{\mu} + e^{-} \rightarrow \bar{\nu}_e + \mu$$

These reactions have not yet been observed and have rates comparable to $\nu_{\mu}e$ elastic scattering. They can only be done at Fermilab because the threshold energy is 10 GeV. Since the muons go very forward our apparatus has very good acceptance. The study of these reactions answers the question of the additive or multiplicative law of lepton conservation.

We wish to stress that these measurements should not cost us additional effort.

We should also mention that once we can measure muon momenta we can also examine events of the following types in our data:

a) Muon-electron pair production,

$$\nu_{\mu} + N \rightarrow \mu^{-} + e + X$$

b) Multi-lepton production,

$$\begin{aligned} \nu_{\mu} + N &\rightarrow \mu^{-} \mu^{+} \mu^{-} + X \\ &\rightarrow \mu^{-} e^{+} e^{-} + X \\ &\rightarrow e^{-} \mu^{+} \mu^{-} + X \\ &\rightarrow e^{-} e^{+} e^{-} + X \end{aligned}$$

We have begun to construct 10 drift chambers. We ask Fermilab to supply two iron dipole magnets each 1.5 m long and 8 ft X 8 ft in active cross-section. This requires an additional 15' of space in the Wonder Building. We do not envisage that this program will interfere with our approved $\nu_{\mu}e$ scattering experiment. Besides the additional physics which we can gain we consider it to be a necessary improvement to enable us to understand the many background processes which will trigger our apparatus.