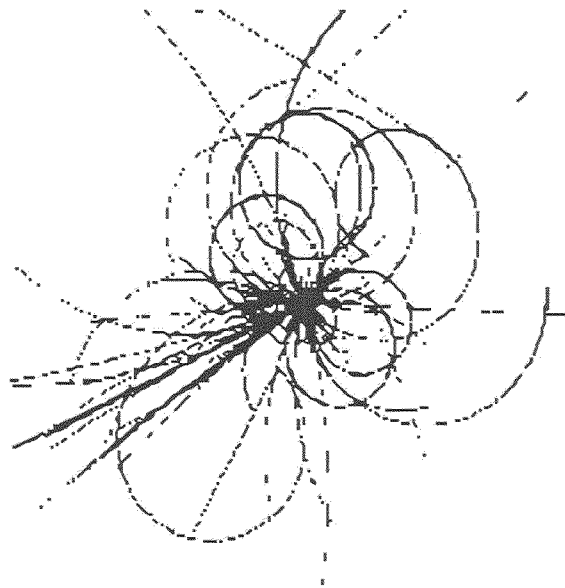


THE SUPERCONDUCTING SUPER COLLIDER LABORATORY



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NOTES

GENERALIZED FORMULA FOR POWER LOSS
OF A GAUSSIAN BUNCH IN RESONATORS

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ABSTRACT

Recently Furman derived a practical formula for computing the power lost by a gaussian bunch in periodic orbit traversing a resonator with quality factor $Q > 1/2$. We give a generalized expression valid for all Q values.

A practical formula to compute the power loss by a short gaussian bunch traversing a resonator was derived by Furman¹ recently. His formula is valid for resonator quality factor $Q > 1/2$. With a slightly different formulation, we obtain a general expression valid for all physical Q values. We shall take Eq. (10) in Furman's note as the starting point, which states that, for a gaussian bunch of total charge $N_B e$ and longitudinal width σ_z , moving in a periodic orbit of length $2\pi R$ with a frequency $f_0 = \omega_0/2\pi$, the power loss is given by

$$P = R_S (N_B e f_0)^2 (R/\sigma_z) \int_{-\infty}^{\infty} \frac{e^{-s^2} ds}{1 + Q^2 (s/\alpha - \alpha/s)^2} \quad (1)$$

where $\alpha = (\omega_R/\omega_0)(\sigma_z/R)$, and R_S , Q , and ω_R are the shunt impedance, quality factor, and resonance frequency of the resonator, respectively. The conditions which justify the integral representation Eq. (1) are $\omega_R/\omega_0 \gg Q$ and $(\sigma_z/R)^2 \gg 1$. The discussion of these conditions is given in Ref. 1.

Consider the identity

$$\begin{aligned} \frac{1}{1 + iQ(x-1/x)} &= -\frac{i x}{Q} \left[\frac{1}{(x^2-1)-ix/Q} \right] \\ &= -\left(\frac{i x}{Q}\right) \frac{1}{(x-x_1)(x-x_2)} \\ &= -\left(\frac{i}{Q}\right) \frac{1}{(x_1-x_2)} \left[\frac{x_1}{x-x_1} - \frac{x_2}{x-x_2} \right], \end{aligned} \quad (2)$$

where

$$x_{1,2} = \frac{1}{2Q} [i \pm \sqrt{4Q^2-1}] \quad (3)$$

Replacing Q by $-Q$ in (2), we get

$$\frac{1}{1-iQ(x-1/x)} = \left(\frac{i}{Q}\right) \frac{1}{(x_1-x_2)} \left[\frac{x_1}{x+x_1} - \frac{x_2}{x+x_2} \right] \quad (4)$$

It follows that

$$\begin{aligned} \frac{1}{1+Q^2(x-1/x)^2} &= \frac{1}{2} \left[\frac{1}{1+iQ(x-1/x)} + \frac{1}{1-iQ(x-1/x)} \right] \\ &= \frac{i}{Q} \frac{1}{(x_1-x_2)} \left[\frac{x_1^2}{(x_1^2-x^2)} - \frac{x_2^2}{(x_2^2-x^2)} \right] \end{aligned} \quad (5)$$

Let $x = \alpha x$, and $s_{1,2} = \alpha x_{1,2}$, the integral in Eq. (1) can be written as

$$\begin{aligned} I(Q) &\equiv \int_{-\infty}^{\infty} \frac{e^{-s^2} ds}{1+Q^2(x-1/x)^2} \\ &= \frac{i}{Q} \frac{1}{(x_1-x_2)} \int_{-\infty}^{\infty} ds e^{-s^2} \left[\frac{s_1^2}{s_1^2-s^2} - \frac{s_2^2}{s_2^2-s^2} \right] \\ &= \frac{\alpha\pi}{Q} [s_1 w(s_1) - s_2 w(s_2)] / (s_1 - s_2) \end{aligned} \quad (6)$$

where $w(z)$ is the complex error function² defined by

$$w(z) = \frac{iz}{\pi} \int_{-\infty}^{\infty} \frac{e^{-s^2}}{z^2 - s^2} ds \quad (\text{Im } z > 0)$$

From (3) $\text{Im } x_{1,2} > 0$ for $Q > 0$, clearly Eq. (6) is well-defined for all nonzero Q values. In particular, for $Q > 1/2$ we can write

$$I(Q) = \frac{2\pi}{\sqrt{4Q^2-1}} \text{Re}[s_1 w(s_1)] \quad (7)$$

by using the facts that $s_2 = -s_1^*$ and $w(-s_1^*) = [w(s_1)]^*$. Expression (7) is just Furman's result.

The function $I(Q)$ given by (6) is obviously continuous at $Q = 1/2$, we can easily evaluate it as

$$I\left(\frac{1}{2}\right) = 2\alpha\pi \left. \frac{d}{ds} [sw(s)] \right|_{s=i\alpha}$$

Using the identity²

$$\frac{dw(s)}{ds} = \frac{2i}{\sqrt{\pi}} - 2sw(s) ,$$

we get $I(\frac{1}{2}) = 2\pi\alpha[(1+2\alpha^2)w(i\alpha) - \frac{2\alpha}{\sqrt{\pi}}]$. The positivity of $I(1/2)$ is insured by the fact $w(i\alpha) \geq \frac{1}{\sqrt{\pi}}$.

REFERENCES

1. M. Furman, SSC-N-142, May 1986.
2. Handbook of Mathematical Functions, ed. by M. Abramowitz and I.A. Stegun, Dover (1965), p.297-298.

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