Measurement of the effective weak mixing angle in $p\bar{p} \to Z/\gamma^* \to e^+e^-$ events


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We present a measurement of the fundamental parameter of the standard model, the weak mixing angle, in $p\bar{p}\to Z/\gamma^*\to e^+e^-$ events at a center of mass energy of 1.96 TeV, using data corresponding to 9.7 fb$^{-1}$ of integrated luminosity collected by the D0 detector at the Fermilab Tevatron. The effective weak mixing angle is extracted from the forward-backward charge asymmetry as a function of the invariant mass around the $Z$ boson pole. The measured value of $\sin^2\theta_{\text{eff}} = 0.23146 \pm 0.00047$ is the most precise measurement from light quark interactions to date, with a precision close to the best LEP and SLD results.

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The weak mixing angle $\sin^2\theta_W$ is one of the fundamental parameters of the standard model (SM). It describes the relative strength of the axial-vector couplings $g_A^f$ to the vector couplings $g_f^f$ in neutral-current interactions of a $Z$ boson to fermions $f$ with Lagrangian

$$\mathcal{L} = -i\frac{g}{2\cos\theta_W} f\gamma^\mu \left(g_f^f - g_A^f \gamma_5\right) f Z^\mu, \quad (1)$$

with $g_f^f = I_f^f - 2Q_f \cdot \sin^2\theta_W$, $g_A^f = I_f^f$, where $I_f^f$ and $Q_f$ are the weak isospin component and the charge of the fermion. At tree level and in all orders of the on-shell renormalization scheme, the weak mixing angle can be written in terms of the $W$ and $Z$ boson masses as $\sin^2\theta_W = 1 - M_W^2/M_Z^2$. To include higher order electroweak radiative corrections, effective weak mixing angles are defined as

$$\sin^2\theta_{\text{eff}}^f = \frac{1}{4|Q_f|} \left( 1 - \frac{g_f^f}{g_A^f} \right), \quad (2)$$

for each fermion flavor.

It is customary to quote the charged lepton effective weak mixing angle $\sin^2\theta_{\text{eff}}^l$, determined by measurements of observables around the $Z$ boson pole. There is tension between the two most precise measurements of $\sin^2\theta_{\text{eff}}^l$, which are $0.23221 \pm 0.00029$ from the combined LEP measurement using the charge asymmetry $A_{FB}^b$ for $b$ quark production and $0.23098 \pm 0.00026$ from the SLD measurement of the $e^+e^-$ left-right polarization asymmetry $A_{l\nu}$ [1]. An independent determination of the effective weak mixing angle is therefore an important precision test of the SM electroweak breaking mechanism.

At the Tevatron, the mixing angle can be measured in the Drell-Yan process $p\bar{p}\to Z/\gamma^*\to \ell^+\ell^-$, through a forward-backward charge asymmetry in the distribution of the emission angle $\theta^*$ of the negatively charged lepton momentum relative to the incoming quark momentum, defined in the Collins-Soper frame [2]. Events with $\cos \theta^* > 0$ are classified as forward ($F$), and those with $\cos \theta^* < 0$ as backward ($B$). The forward-backward charge asymmetry, $A_{FB}$, is defined by

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}, \quad (3)$$

where $N_F$ and $N_B$ are the numbers of forward and backward events. The asymmetry arises from the interference between vector and axial vector coupling terms.
The asymmetry $A_{FB}$ can be measured as a function of the invariant mass of the dilepton pair ($M_{ee}$). The presence of both vector and axial-vector couplings of the $Z$ boson to fermions gives the most significant variation of $A_{FB}$ in vicinity of the $Z$ boson pole, which is sensitive to the effective weak mixing angle.

Measurements of $\sin^2 \theta_W$ have been reported previously by the CDF Collaboration in the $Z \rightarrow e^+e^-$ [3, 4] and $Z \rightarrow \mu^+\mu^-$ [5] channels, and the D0 Collaboration in the $Z \rightarrow e^+e^-$ channel [6, 7]. The angle $\sin^2 \theta_W$ has also been measured at the LHC in $pp$ collisions by the CMS Collaboration in the $Z \rightarrow \mu^+\mu^-$ channel at $\sqrt{s} = 7$ TeV [8].

This letter reports a measurement of the effective weak mixing angle from the $A_{FB}$ distribution using 9.7 fb$^{-1}$ of integrated luminosity collected with the D0 detector at the Fermilab Tevatron collider. The precision of the previous D0 measurement using 5 fb$^{-1}$ of data [7], $\sin^2 \theta_W = 0.2309 \pm 0.0008$ (stat.) $\pm 0.0006$ (syst.), was dominated by available statistics and the uncertainty on the electron energy scale. The analysis of the full 9.7 fb$^{-1}$ data set presented here features an extended acceptance and a new electron energy calibration method providing substantially improved accuracy.

The D0 detector comprises a central tracking system, a calorimeter and a muon system [9–11]. The central tracking system consists of a silicon microstrip tracker (SMT) and a scintillating central fiber tracker (CFT), both located within a 2 T superconducting solenoidal magnet. The tracking system consists of a silicon microstrip tracker (SMT) and a scintillating central fiber tracker (CFT), both located within a 2 T superconducting solenoidal magnet [9–11]. The central tracking system comprises a silicon microstrip tracker (SMT) and a scintillating central fiber tracker (CFT), both located within a 2 T superconducting solenoidal magnet.

The MC Drell-Yan $Z/\gamma^* \rightarrow e^+e^-$ sample is generated using the D0 simulation software, based on the leading-order PYTHIA generator [14] with the NNPDF2.3 [15] parton distribution functions (PDFs), followed by a GEANT-based simulation [16] of the D0 detector. The PYTHIA MC samples, with data events from random beam crossings overlaid, are mainly used to understand the geometric acceptance, and the energy scale and resolution of electrons in the calorimeter.

A new method of electron energy calibration is developed and applied to both the data and MC, which significantly reduces the systematic uncertainty due to the electron energy measurement. The weak mixing angle, which is extracted from $A_{FB}$ as a function of $M_{ee}$, depends strongly on the position of the peak value of $M_{ee}$. Therefore, it is critical to have a precise electron energy measurement and a consistent measured peak value of $M_{ee}$ from different regions of the detector across various Tevatron running conditions. In Ref. [7], an overall scale factor is applied to simulations to model the detector response for electron energy depositions, where the scale factor is determined by comparing the $M_{ee}$ spectrum in data and MC, yielding a large uncertainty due to background estimation and detector resolution. In this analysis, a new energy calibration method is applied to the data and the MC separately. For CC electrons, an instantaneous luminosity-dependent scale factor ($\alpha_{CC}^{\text{inst}}$) and an $\eta_{det}$-dependent scale factor ($\alpha_{\eta}^{CC}$) are applied to the electron energy. For EC electrons in addition to the scale factors requiring at least two electromagnetic (EM) clusters reconstructed near the boundaries of CC calorimeter modules [9] ($\phi$-mod boundary) are included. The geometric acceptance is further extended compared with previous D0 results [7] from $|\eta_{det}| < 1.0$ to $|\eta_{det}| < 1.1$ for the CC, and from $1.5 < |\eta_{det}| < 2.5$ to $1.5 < |\eta_{det}| < 3.2$ for the EC. In addition, EC-EC events, which were excluded due to their poorer track reconstruction and calorimeter energy resolution, are now included. These extensions in $\eta_{det}$ and $\phi$-mod acceptance give a 70% increase in the number of events over what would be expected from the increase in integrated luminosity. An additional 15% increase is gained from the improvements in the track reconstruction algorithm. The number of $Z \rightarrow e^+e^-$ candidate events in the data sample is 560,267 which includes 248,380 CC-CC events, 240,593 CC-EC events and 71,294 EC-EC events.
factors $\alpha_L^{EC}$ and $\alpha_T^{EC}$, an $\eta_{det}$-dependent offset $\beta_\eta^{EC}$ is introduced to model the $\eta_{det}$ dependence of the electron energy. All correction factors are determined by scaling the peak of the $M_{ee}$ distribution as a function of instantaneous luminosity and $\eta_{det}$ to the Z boson mass measured by LEP ($M_Z = 91.1875$ GeV) [1]. The CC correction factors are tuned with the CC-CC events. To remove one degree of freedom, $\beta_\eta^{EC}$ is expressed as a function of $\alpha_\eta^{EC}$, and the relationship is measured with the CC-EC events. The values of $\alpha_\eta^{EC}$ and $\beta_\eta^{EC}(\alpha_\eta^{EC})$ are fitted with the EC-EC events. After the calibration, the standard deviation $\delta M$ of the $M_{ee}$ peak value from the LEP $M_Z$ value is $\approx 20$ MeV. Various closure tests are performed to check the validity of the calibration procedure. For example, an extra $\eta_{det}$-dependent offset is applied to the corrected energy and fixed by performing the calibration again. The extra offset is found to be consistent with $\delta M$. The ratio of $\delta M$ to $M_Z$ is propagated into the uncertainty of the $\sin^2\theta_W$ measurement to estimate the systematic uncertainty arising from the energy calibration.

After the electron energy calibration, an additional electron energy resolution smearing is derived and applied to the MC to achieve agreement of the width of the $M_{ee}$ distribution in data. For the CC $\phi$-mod boundary electrons, the resolution smearing is modeled with a Crystal Ball function [17]. For other electrons, the smearing is modeled with a Gaussian function.

Additional corrections and reweightings are applied to the MC simulation to improve the agreement with data. The scale factors of the electron identification efficiency between the MC and the data are measured using the tag-probe method [18] and applied to the MC distribution as functions of $p_T$ and $\eta_{det}$. The simulation is further corrected for higher-order effects not included in PYTHIA by reweighting the MC events at the generator level in two dimensions ($p_T$ and rapidity $y$ of the Z boson) to match RESBOS [19] predictions. In addition, next-to-next-leading order QCD corrections are applied as a function of $M_Z$ [19, 20]. The distribution of the instantaneous luminosity and the $z$ coordinate of the $p\bar{p}$ collision vertices are also weighted to match those in the data. Since $A_{FB}$ is defined as a ratio of numbers of events, many small uncertainties cancel out. Only the electron selection efficiency scale factor in these additional corrections contributes significantly to the final uncertainty.

The charge of the particle track matched to the EM cluster is used to determine if the EM cluster is associated to an electron or positron and to classify the event as forward or backward. The charge misidentification probability $f_q$ is given by

$$f_q = \frac{1}{2} \frac{N_{SS}}{(N_{SS} + N_{OS})}, \quad (4)$$

where $N_{SS}$ ($N_{OS}$) is the total number of $p\bar{p} \rightarrow Z/\gamma^* \rightarrow e^+e^-$ events reconstructed with same-sign (opposite-sign) electrons. The probabilities are measured in data and MC using tag-and-probe method. The charge of electrons and positrons reconstructed in the MC is randomly changed to match the misidentification probability in the data averaged over $p_T$ spectrum of electrons. In the CC region the average charge misidentification rate in data is about 0.3%, whereas in the EC region it varies from 1% at $|\eta_{det}| = 1.5$ to 10% at $|\eta_{det}| = 3.0$. The statistical uncertainty of the measured charge misidentification rate is included as a systematic uncertainty.

The background is suppressed by the strict requirements on the quality of the matched track. The main contribution is from multijet events, in which jets are misreconstructed as electrons, and is estimated from data. Multijet events are selected by reversing part of the electron selections to study the differential distributions of the multijet background, which are different from the real multijet background that passes all the electron selections. Therefore, a correction factor is applied as a function of electron $p_T$, given by the ratios of the efficiencies for EM-like jets (which are selected in a multijet-enriched data sample and pass all the electron selections) and “reverse-selected” jets. The normalization of the multijet background is determined by fitting the sum of the $M_{ee}$ distributions of multijet events and the signal MC events to the distribution from the selected data events. The $W^\pm$+jets, $Z/\gamma^* \rightarrow \tau\tau$, di-boson (WW and $WZ$) and $t\bar{t}$ backgrounds are estimated using the PYTHIA MC simulation and found to be negligible around the $Z$ boson pole. The total number of background events is found to be less than 1% in CC-CC events and 4% in CC-EC and EC-EC events. The $M_{ee}$ and $\cos \theta^* \gamma^*$ distributions of data and of the sum of signal MC and background expectations are in good agreement with the SM predictions.

The weak mixing angle is extracted from the background-subtracted raw $A_{FB}$ spectrum in the regions $75 < M_{ee} < 115$ GeV for CC-CC and CC-EC events, and $81 < M_{ee} < 97$ GeV for EC-EC events by comparing the data to simulated $A_{FB}$ templates corresponding to different input values of $\sin^2\theta_W$. The mass window for EC-EC events is narrower to take into account the different track reconstruction and energy measurement. The templates are obtained by reweighting $M_Z$ and $\cos \theta^* \gamma^*$ distributions at the generator level to different Born-level $\sin^2\theta_W$ predictions. The $A_{FB}$ distribution is negligibly sensitive to the effect of QED final state radiation because most of these radiated photons are emitted co-linearly with the electron and are reconstructed as one single EM object by the detector. The background-subtracted raw $A_{FB}$ distribution and the PYTHIA prediction with the fitted $\sin^2\theta_W$ are shown in Fig. 1.

The results of the fits for different event categories, with statistical and systematic uncertainties, are listed in Table 1.
The PDF uncertainty is estimated by reweighting the PDF set in the MC simulations to different sets of the NNPDF2.3 parameterization [15].

To have a consistent SM definition and make our result comparable with previous measurements, the PYTHIA interpretation of the weak mixing angle is compared to the predictions from modified RESBOS with CTEQ6.6 PDF set [21], which uses different values of effective weak mixing angle for leptons and up or down quarks [22]. A constant 0.00008 positive shift in $\sin^2 \theta_W$ in the prediction from RESBOS relative to the LO prediction from PYTHIA changes the measured leptonic effective weak mixing angle to $\sin^2 \theta_{\ell e} = 0.23146 \pm 0.00047$, with the same breakdown of uncertainties as above. The comparison between our measurement and other experimental results is shown in Fig. 2. Our measurement is consistent with the current world average.

In conclusion, we have measured the effective weak mixing angle $\sin^2 \theta_{\ell e}$ from the distribution of the forward-backward charge asymmetry $A_{FB}$ in the process $p\bar{p} \rightarrow Z/\gamma^* \rightarrow e^+e^-$ at the Tevatron. This measurement, which supersedes that reported in [7], uses nearly twice the integrated luminosity and significantly extends the electron acceptance. The primary systematic uncertainty is reduced by introducing a new electron energy calibration method. The final result from 9.7 fb$^{-1}$ of integrated luminosity is $\sin^2 \theta_{\ell e} = 0.23146 \pm 0.00047$, and is the most precise measurement from light quark interactions, and is close to the precision of the world’s best measurements performed by the LEP and SLD Collaborations.
\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure2.png}
\caption{(color online). Comparison of measured $\sin^2 \theta_{\text{eff}}$ with results from other experiments. The average is a combination of $A_{\phi,1}^b$, $A_{\phi}^b$, $A_{\phi}^c$, $Q^\text{had}$, $A_{\eta}^\text{SLD}$ (CDF), and $A_{\eta}^\text{SLD}$ (D0) measurements from the LEP and SLD Collaborations [1].}
\end{figure}

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[12] D0 uses a cylindrical coordinate system with the z axis along the beam axis in the proton direction. Angles $\theta$ and $\phi$ are the polar and azimuthal angles, respectively. Pseudorapidity is defined as $\eta = -\ln(\tan(\theta/2))$ where $\theta$ is measured with respect to the interaction vertex. In the massless limit, $\eta$ is equivalent to the rapidity $y = (1/2) \ln((E+p_z)/(E-p_z))$, and $\eta_{\text{det}}$ is the pseudorapidity measured with respect to the center of the detector.