Electron neutrino and antineutrino appearance in the full MINOS data sample


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We report the results of a search for $\nu_e$ and $\bar{\nu}_e$ appearance in $\nu_\mu$ and $\bar{\nu}_\mu$ beams using the full MINOS data sample. This analysis uses an exposure of $10.6 \times 10^{20}$ $\text{cm}^2 \text{sr} \text{MeV}$ protons-on-target taken with a $\nu$ ($\bar{\nu}$) beam mode and is the first accelerator long-baseline search for $\nu_\mu \rightarrow \nu_e$. With the $\nu_e$ and $\bar{\nu}_e$ data, we probe $\theta_{13}$, $\delta$, and the mass hierarchy. Assuming a normal (inverted) mass hierarchy, $\delta = 0$, and $\theta_{23} < \frac{\pi}{4}$, we set a constraint of 0.01 (0.03) $< 2 \sin^2(2\theta_{13}) \sin^2(\theta_{23}) < 0.12$ (0.18) at 90% C.L. with the best-fit value of $2 \sin^2(2\theta_{13}) \sin^2(\theta_{23}) = 0.051_{-0.030}^{+0.038} (0.093_{-0.045})$. 

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The neutrino oscillation phenomenon is successfully modeled by a theory of massive neutrino eigenstates that are different from the neutrino flavor eigenstates. These sets of eigenstates are related by the PMNS matrix [1] which is commonly parameterized by three angles, $\theta_{ij}$, and a CP-violating phase, $\delta$.

The values of $\theta_{12}$ and $\theta_{23}$ have been measured [2-4] with indications that $\theta_{23}$ is not maximal [5-7]. The final angle, $\theta_{13}$, is now known to have a nonzero value from measurements by reactor experiments [8-10], the measurement by the T2K [11] accelerator experiment, and from earlier MINOS results [12, 13].

Despite these accomplishments, the value of $\delta$ is still unknown, as is the ordering of the neutrino masses, which is referred to as the neutrino mass hierarchy. Much of the attention in the neutrino community is now focused on resolving these unknowns. The mass hierarchy is not only a fundamental property of neutrinos but also has a direct impact on the ability of neutrinoless double beta decay searches to state definitively whether the neutrino is its own antiparticle [14]. Reactor experiments make a pure measurement of $\theta_{13}$, whereas the $\nu_e \rightarrow \nu_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ appearance probabilities measured by accelerator experiments such as MINOS depend on the value of $\delta$. In addition, the long-baseline of MINOS means that interactions between neutrinos and the matter of the Earth make the appearance probabilities dependent on the neutrino mass hierarchy [15, 16].

We report the result from the search for $\nu_e$ ($\bar{\nu}_e$) appearance in a $\nu_\mu$ ($\bar{\nu}_\mu$) beam using the full MINOS data sample. This result uses an exposure of $10.6 \times 10^{20}$ protons-on-target taken with a $\nu$ beam and an exposure of $3.3 \times 10^{20}$ protons-on-target taken with a $\bar{\nu}$ beam. The neutrino sample is $30\%$ larger than the sample used for the previous MINOS results on this topic [13]. This analysis represents the first long-baseline search for $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ appearance, places new constraints on $\delta$, and has a dependence on the value of $\delta$, $\theta_{23}$, and the neutrino mass hierarchy.

In the MINOS experiment [17], neutrino oscillation is studied with the NuMI beamline [18] by measuring neutrino interactions in two detectors. The Near Detector (ND), which has a fiducial mass of 29 tons, is at a distance of 1.04 km from the production target and is used to determine the composition of the beam before the neutrinos have oscillated. The Far Detector (FD), which has a fiducial mass of 3.8 kilotons, is at a distance of 735 km from the production target and is used to measure the change in the neutrino flavor composition of the beam. In both the $\nu$ and $\bar{\nu}$ beam modes, the NuMI beam has an energy spectrum that is peaked at 3 GeV. At the ND, the neutrino flavor composition of the neutrino interactions, as determined by a combination of simulation and measurement, is found to be $91.7\% \nu_\mu$, $7.0\% \bar{\nu}_\mu$, and $1.3\% \nu_\tau$ and $\bar{\nu}_\tau$ for the $\nu$ beam mode and $58.1\% \nu_\mu$, $39.9\% \bar{\nu}_\mu$, and $2.0\% \nu_\tau$ and $\bar{\nu}_\tau$ for the $\bar{\nu}$ beam mode.

Both detectors are magnetized tracking calorimeters consisting of alternating planes of 2.54 cm thick steel and 1 cm thick scintillating plastic [17]. The scintillator planes are segmented into 4.1 cm wide strips with wavelength-shifting fibers embedded in the strips to collect light for readout by multi-anode photomultiplier tubes.

In the MINOS data sample, the flavor of a neutrino is determined only for charged-current (CC) interactions. $\nu_\tau$-CC and $\bar{\nu}_\tau$-CC interactions are identified by the presence of a long muon track that extends beyond a cluster of energy depositions that are consistent with hadronic activity at the interaction vertex. Neutral-current (NC) interactions are identified by the energy depositions associated with hadronic activity. $\nu_e$-CC and $\bar{\nu}_e$-CC interactions produce an electromagnetic shower that typically leaves a compact cluster of energy depositions within 6 to 12 planes. This analysis does not distinguish between $\nu_e$-CC and $\bar{\nu}_e$-CC interactions.

The sample of events classified as $\nu_e$-CC and $\bar{\nu}_e$-CC interactions contains a background of interactions that produce a compact cluster of energy depositions in the scintillator. NC interactions with a significant electromagnetic component and $\nu_\mu$-CC or $\bar{\nu}_\mu$-CC interactions in which the muon track is not easily identified make up the majority of the background. Smaller contributions to the background arise from $\nu_\tau$-CC and $\bar{\nu}_\tau$-CC interactions. In addition to backgrounds that mimic $\nu_\tau$-CC and $\bar{\nu}_\tau$-CC event topologies, intrinsic $\nu_e$ and $\bar{\nu}_e$ components of the NuMI beam must be taken into account.

Candidate $\nu_e$-CC and $\bar{\nu}_e$-CC events are required to fall within a fiducial volume and to be coincident in time and direction with the NuMI beam. We require the events to have shower-like topologies by rejecting events with tracks that are longer than 25 planes or extend more than 15 planes from a shower edge. In addition, reconstructed events must have at least five consecutive planes with deposited energy above a threshold; this threshold is defined as half of the energy deposited by a minimum ionizing particle. We require the events to have a reconstructed energy between 1 and 8 GeV where most of the $\nu_e$ and $\bar{\nu}_e$ appearance is expected.

We further classify the events in this pre-selected sample of shower-like events by using a library-event-matching (LEM) algorithm [19, 20]. Within the LEM algorithm, the topology of energy depositions that characterize the event is compared to a library of simulated signal and background events. Separate libraries are used for the events in the $\nu$ beam mode and $\bar{\nu}$ beam mode. The 50 best-matching events in the library are collected and
used to produce three variables. These variables are the fraction of best-matching library events that are $\nu_e$-CC or $\bar{\nu}_e$-CC, the average inelasticity of the best-matching $\nu_e$-CC or $\bar{\nu}_e$-CC library events, and the average fraction of the energy depositions that overlap between the test event and the best-matching $\nu_e$-CC or $\bar{\nu}_e$-CC library events. These three variables and the reconstructed neutrino energy of the test event are then used as an input into an artificial neural network. The output value from the neural network is used to discriminate between signal and background events. This discriminant variable is referred to as $\alpha_{LEM}$ and is shown in Fig. 1. Signal events have a value near one, while background events cluster near zero. The maximum sensitivity to $\nu_e$ and $\bar{\nu}_e$ appearance is obtained by analyzing events with $\alpha_{LEM} > 0.6$.

Following the selection of $\nu_e$-CC and $\bar{\nu}_e$-CC candidate events, the ND data is used to study the rate of background from NC, $\nu_\mu$-CC and $\bar{\nu}_\mu$-CC and intrinsic beam $\nu_e$-CC and $\bar{\nu}_e$-CC interactions. The NuMI beam can be tuned to produce different energy spectra. Among these different beam configurations, the relative contributions of the various backgrounds change in a well-understood way. By measuring the total of the three backgrounds in three different beam configurations, the relative amounts of the individual backgrounds can be deduced [21].

We use the measurement of the ND backgrounds to derive the FD background predictions for the data samples in the $\nu$ beam mode and in the $\bar{\nu}$ beam mode. For each sample, we divide simulated FD events into bins of energy and $\alpha_{LEM}$ and correct each FD background component, bin-by-bin, by multiplying it by the measured ND ratio of data to simulated events for that background. Since the ND data sample does not contain $\nu_e$-CC or $\bar{\nu}_e$-CC events from oscillation, we estimate the FD contribution from this small background component through simulation and a correction based on the observed ND $\nu_e$-CC and $\bar{\nu}_e$-CC spectra.

The sources of systematic uncertainty that affect the background prediction are given in Table I. The effect of each source of uncertainty is evaluated by producing simulated ND and FD event samples that are modified according to the estimated size of each systematic effect. These modified samples are used to produce an altered FD background prediction for the systematic effect in question. We take the resulting difference between the nominal and modified prediction as the systematic uncertainty on the background prediction. The systematic effect that results in the largest reduction in sensitivity is a 2.0% uncertainty on the relative energy scale between the ND and FD.

With the absence of a $\nu_e$-CC and $\bar{\nu}_e$-CC signal in the ND, the signal selection efficiency cannot be extrapolated from the ND events in the same way as the background estimate. Therefore, to evaluate the signal efficiency, we select a sample of well identified $\nu_\mu$-CC events [22, 23], remove the energy depositions that are associated with the muon track [24], and insert the simulated energy depositions of an electron with an identical three-momentum [25]. This method effectively turns a well identified sample of $\nu_\mu$-CC and $\bar{\nu}_\mu$-CC data events into a sample of $\nu_e$-CC and $\bar{\nu}_e$-CC data events. For the $\nu$ beam mode ($\bar{\nu}$ beam mode) data sample, we find the expected number of FD signal events with $\alpha_{LEM} > 0.6$ and the associated systematic uncertainty to be $33.7 \pm 1.9$ (3.9 $\pm$ 0.2) assuming $\sin^2(2\theta_{13}) = 0.1$, $\delta = 0$, $\theta_{23} = \pi/4$, and a normal mass hierarchy. This corresponds to an identification efficiency of $(57.4 \pm 2.8)\%$ for the $\nu$ beam mode and of $(63.3 \pm 3.1)\%$ for the $\bar{\nu}$ beam mode. The systematic uncertainties are evaluated in a way that is similar to the evaluation of the background systematics.
Events with $\alpha_{\text{LEM}} < 0.5$ are insensitive to $\nu_e$ and $\bar{\nu}_e$ appearance. These events are therefore used in a separate study to validate the analysis procedure. ND events with $\alpha_{\text{LEM}} < 0.5$ are used to predict FD event yields, which are found to agree with the FD data to within 0.3 (0.6) standard deviations of the statistical uncertainty for the data sample in the $\nu$ ($\bar{\nu}$) beam mode.

Events with $\alpha_{\text{LEM}} > 0.6$ are selected for further analysis in the $\nu$ beam mode and in the $\bar{\nu}$ beam mode. The expected and observed event counts in these samples are shown in Table I. The observed FD reconstructed energy spectra, in bins of $\alpha_{\text{LEM}}$, are shown for the candidate events in Fig. 2. Assuming a three-flavor neutrino oscillation probability that includes matter effects [16], we simultaneously fit the data from the $\nu$ beam mode and $\bar{\nu}$ beam mode samples for the value of $2\sin^2(2\theta_{23})\sin^2(\delta_{23})$ while the value of the mass hierarchy and $\delta$ are held fixed. The fit is performed using the 15 bins formed by three bins of $\alpha_{\text{LEM}}$ and five bins of energy. This procedure is performed for all values of delta and both mass hierarchies, and the resulting confidence intervals, calculated using the Feldman-Cousins technique [26], are shown in Fig. 3. The values of the oscillation parameters used in the fit are taken from previous measurements [2, 4] and are set to $\sin^2(\theta_{13}) = 0.051\pm0.035$, $|\Delta m^2_{32}| = (2.39^{+0.98}_{-1.01}) \times 10^{-3}$ eV$^2$, $\theta_{12} = 60.0 \pm 0.2\degree$, and $\Delta m^2_{21} = (7.59^{+0.19}_{-0.21}) \times 10^{-5}$ eV$^2$. The full set of statistical and systematic uncertainties on the prediction are taken into account when calculating the likelihood, as are the uncertainties on the oscillation parameters. The resulting curve of twice the negative log likelihood difference ($-2\Delta \ln L$) for the possible values of $\delta$ is shown in Fig. 4. The difference is taken with respect to the best-fitting solution. Assuming $\theta_{23} > \pi/4$ ($\theta_{23} < \pi/4$), the inverted hierarchy is preferred at 0.631 (0.041) units of $-2\Delta \ln L$.

In conclusion, we have presented a search for $\nu_\mu$ and $\bar{\nu}_e$ appearance in $\nu_\mu$ and $\bar{\nu}_\mu$ beams from the full MINOS data sample. We have used these data to place new constraints on the mixing angle $\theta_{13}$. This Letter also represents the first long-baseline search for $\bar{\nu}_\mu \rightarrow \nu_e$ appearance, as part of an analysis which is dependent on the CP-violating phase $\delta$ and the neutrino mass hierarchy.
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[23] P. Adamson et al. (MINOS), Phys. Rev. Lett. 106,