

T. Aaltonen,²⁴ J. Adelman,¹⁴ B. Álvarez González^{w,12} S. Amerio^{ee,44} D. Amidei,³⁵ A. Anastassov,³⁹ A. Annovi,²⁰ J. Antos,¹⁵ G. Apollinari,¹⁸ J. Appel,¹⁸ A. Apresyan,⁴⁹ T. Arisawa,⁵⁸ A. Artikov,¹⁶ J. Asaadi,⁵⁴ W. Ashmanskas,¹⁸ A. Attal,⁴ A. Aurisano,⁵⁴ F. Azfar,⁴³ W. Badgett,¹⁸ A. Barbaro-Galtieri,²⁹ V.E. Barnes,⁴⁹ B.A. Barnett,²⁶ P. Barria^{gg,47} P. Bartos,¹⁵ G. Bauer,³³ P.-H. Beauchemin,³⁴ F. Bedeschi,⁴⁷ D. Beecher,³¹ S. Behari,²⁶ G. Bellettini^{ff,47} J. Bellinger,⁶⁰ D. Benjamin,¹⁷ A. Beretvas,¹⁸ A. Bhatti,⁵¹ M. Binkley*,¹⁸ D. Bisello^{ee,44} I. Bizjak^{kk,31} R.E. Blair,² C. Blocker,⁷ B. Blumenfeld,²⁶ A. Bocci,¹⁷ A. Bodek,⁵⁰ V. Boisvert,⁵⁰ D. Bortoletto,⁴⁹ J. Boudreau,⁴⁸ A. Boveia,¹¹ B. Brau^{a,11} A. Bridgeman,²⁵ L. Brigliadori^{dd,6} C. Bromberg,³⁶ E. Brubaker,¹⁴ J. Budagov,¹⁶ H.S. Budd,⁵⁰ S. Budd,²⁵ K. Burkett,¹⁸ G. Busetto^{ee,44} P. Bussey,²² A. Buzatu,³⁴ K. L. Byrum,² S. Cabrera^{y,17} C. Calancha,³² S. Camarda,⁴ M. Campanelli,³¹ M. Campbell,³⁵ F. Canelli^{14,18} A. Canepa,⁴⁶ B. Carls,²⁵ D. Carlsmith,⁶⁰ R. Carosi,⁴⁷ S. Carrillo^{n,19} S. Carron,¹⁸ B. Casal,¹² M. Casarsa,¹⁸ A. Castro^{dd,6} P. Catastini^{gg,47} D. Cauz,⁵⁵ V. Cavaliere^{gg,47} M. Cavalli-Sforza,⁴ A. Cerri,²⁹ L. Cerrito^{q,31} S.H. Chang,²⁸ Y.C. Chen,¹ M. Chertok,⁸ G. Chiarelli,⁴⁷ G. Chlachidze,¹⁸ F. Chlebana,¹⁸ K. Cho,²⁸ D. Chokheli,¹⁶ J.P. Chou,²³ K. Chung^{o,18} W.H. Chung,⁶⁰ Y.S. Chung,⁵⁰ T. Chwalek,²⁷ C.I. Ciobanu,⁴⁵ M.A. Ciocci^{gg,47} A. Clark,²¹ D. Clark,⁷ G. Compostella,⁴⁴ M.E. Convery,¹⁸ J. Conway,⁸ M. Corbo,⁴⁵ M. Cordelli,²⁰ C.A. Cox,⁸ D.J. Cox,⁸ F. Crescioli^{ff,47} C. Cuenca Almenar,⁶¹ J. Cuevas^{w,12} R. Culbertson,¹⁸ J.C. Cully,³⁵ D. Dagenhart,¹⁸ N. d'Ascenzo^{v,45} M. Datta,¹⁸ T. Davies,²² P. de Barbaro,⁵⁰ S. De Cecco,⁵² A. Deisher,²⁹ G. De Lorenzo,⁴ M. Dell'Orso^{ff,47} C. Deluca,⁴ L. Demortier,⁵¹ J. Deng^{f,17} M. Deninno,⁶ M. d'Errico^{ee,44} A. Di Canto^{ff,47} B. Di Ruzza,⁴⁷ J.R. Dittmann,⁵ M. D'Onofrio,⁴ S. Donati^{ff,47} P. Dong,¹⁸ T. Dorigo,⁴⁴ S. Dube,⁵³ K. Ebina,⁵⁸ A. Elagin,⁵⁴ R. Erbacher,⁸ D. Errede,²⁵ S. Errede,²⁵ N. Ershaidat^{cc,45} R. Eusebi,⁵⁴ H.C. Fang,²⁹ S. Farrington,⁴³ W.T. Fedorko,¹⁴ R.G. Feild,⁶¹ M. Feindt,²⁷ J.P. Fernandez,³² C. Ferrazza^{hh,47} R. Field,¹⁹ G. Flanagan^{s,49} R. Forrest,⁸ M.J. Frank,⁵ M. Franklin,²³ J.C. Freeman,¹⁸ I. Furic,¹⁹ M. Gallinaro,⁵¹ J. Galyardt,¹³ F. Garbersson,¹¹ J.E. Garcia,²¹ A.F. Garfinkel,⁴⁹ P. Garosi^{gg,47} H. Gerberich,²⁵ D. Gerdes,³⁵ A. Gessler,²⁷ S. Giagu^{ii,52} V. Giakoumopoulou,³ P. Giannetti,⁴⁷ K. Gibson,⁴⁸ J.L. Gimmell,⁵⁰ C.M. Ginsburg,¹⁸ N. Giokaris,³ M. Giordani^{jj,55} P. Giromini,²⁰ M. Giunta,⁴⁷ G. Giurgiu,²⁶ V. Glagolev,¹⁶ D. Glenzinski,¹⁸ M. Gold,³⁸ N. Goldschmidt,¹⁹ A. Golossanov,¹⁸ G. Gomez,¹² G. Gomez-Ceballos,³³ M. Goncharov,³³ O. González,³² I. Gorelov,³⁸ A.T. Goshaw,¹⁷ K. Goulianos,⁵¹ A. Gresele^{ee,44} S. Grinstein,⁴ C. Grosso-Pilcher,¹⁴ R.C. Group,¹⁸ U. Grundler,²⁵ J. Guimaraes da Costa,²³ Z. Gunay-Unalan,³⁶ C. Haber,²⁹ S.R. Hahn,¹⁸ E. Halkiadakis,⁵³ B.-Y. Han,⁵⁰ J.Y. Han,⁵⁰ F. Happacher,²⁰ K. Hara,⁵⁶ D. Hare,⁵³ M. Hare,⁵⁷ R.F. Harr,⁵⁹ M. Hartz,⁴⁸ K. Hatakeyama,⁵ C. Hays,⁴³ M. Heck,²⁷ J. Heinrich,⁴⁶ M. Herndon,⁶⁰ J. Heuser,²⁷ S. Hewamanage,⁵ D. Hidas,⁵³ C.S. Hill^{c,11} D. Hirschbuehl,²⁷ A. Hocker,¹⁸ S. Hou,¹ M. Houlden,³⁰ S.-C. Hsu,²⁹ R.E. Hughes,⁴⁰ M. Hurwitz,¹⁴ U. Husemann,⁶¹ M. Hussein,³⁶ J. Huston,³⁶ J. Incandela,¹¹ G. Introzzi,⁴⁷ M. Iori^{ii,52} A. Ivanov^{p,8} E. James,¹⁸ D. Jang,¹³ B. Jayatilaka,¹⁷ E.J. Jeon,²⁸ M.K. Jha,⁶ S. Jindariani,¹⁸ W. Johnson,⁸ M. Jones,⁴⁹ K.K. Joo,²⁸ S.Y. Jun,¹³ J.E. Jung,²⁸ T.R. Junk,¹⁸ T. Kamon,⁵⁴ D. Kar,¹⁹ P.E. Karchin,⁵⁹ Y. Kato^{m,42} R. Kephart,¹⁸ W. Ketchum,¹⁴ J. Keung,⁴⁶ V. Khotilovich,⁵⁴ B. Kilminster,¹⁸ D.H. Kim,²⁸ H.S. Kim,²⁸ H.W. Kim,²⁸ J.E. Kim,²⁸ M.J. Kim,²⁰ S.B. Kim,²⁸ S.H. Kim,⁵⁶ Y.K. Kim,¹⁴ N. Kimura,⁵⁸ L. Kirsch,⁷ S. Klimenko,¹⁹ B. R. Ko,¹⁷ K. Kondo,⁵⁸ D.J. Kong,²⁸ J. Konigsberg,¹⁹ A. Korytov,¹⁹ A.V. Kotwal,¹⁷ M. Kreps,²⁷ J. Kroll,⁴⁶ D. Krop,¹⁴ N. Krumnack,⁵ M. Kruse,¹⁷ V. Krutelyov,¹¹ T. Kuhr,²⁷ N.P. Kulkarni,⁵⁹ M. Kurata,⁵⁶ S. Kwang,¹⁴ A.T. Laasanen,⁴⁹ S. Lami,⁴⁷ S. Lammel,¹⁸ M. Lancaster,³¹ R.L. Lander,⁸ K. Lannon^{u,40} A. Lath,⁵³ G. Latino^{gg,47} I. Lazzizzera^{ee,44} T. LeCompte,² E. Lee,⁵⁴ H.S. Lee,¹⁴ J.S. Lee,²⁸ S.W. Lee^{x,54} S. Leone,⁴⁷ J.D. Lewis,¹⁸ C.-J. Lin,²⁹ J. Linacre,⁴³ M. Lindgren,¹⁸ E. Lipeles,⁴⁶ A. Lister,²¹ D.O. Litvintsev,¹⁸ C. Liu,⁴⁸ T. Liu,¹⁸ N.S. Lockyer,⁴⁶ A. Loginov,⁶¹ L. Lovas,¹⁵ D. Lucchesi^{ee,44} J. Lueck,²⁷ P. Lujan,²⁹ P. Lukens,¹⁸ G. Lungu,⁵¹ J. Lys,²⁹ R. Lysak,¹⁵ D. MacQueen,³⁴ R. Madrak,¹⁸ K. Maeshima,¹⁸ K. Makhoul,³³ P. Maksimovic,²⁶ S. Malde,⁴³ S. Malik,³¹ G. Manca^{e,30} A. Manousakis-Katsikakis,³ F. Margaroli,⁴⁹ C. Marino,²⁷ C.P. Marino,²⁵ A. Martin,⁶¹ V. Martin^{k,22} M. Martínez,⁴ R. Martínez-Ballarín,³² P. Mastrandrea,⁵² M. Mathis,²⁶ M.E. Mattson,⁵⁹ P. Mazzanti,⁶ K.S. McFarland,⁵⁰ P. McIntyre,⁵⁴ R. McNulty^{j,30} A. Mehta,³⁰ P. Mehtala,²⁴ A. Menzione,⁴⁷ C. Mesropian,⁵¹ T. Miao,¹⁸ D. Mietlicki,³⁵ N. Miladinovic,⁷ R. Miller,³⁶ C. Mills,²³ M. Milnik,²⁷ A. Mitra,¹ G. Mitselmakher,¹⁹ H. Miyake,⁵⁶ S. Moed,²³ N. Moggi,⁶ M.N. Mondragon^{n,18} C.S. Moon,²⁸ R. Moore,¹⁸ M.J. Morello,⁴⁷ J. Morlock,²⁷ P. Movilla Fernandez,¹⁸ J. Müllenstädt,²⁹ A. Mukherjee,¹⁸ Th. Muller,²⁷

P. Murat,¹⁸ M. Mussini^{dd,6} J. Nachtman^{o,18} Y. Nagai,⁵⁶ J. Naganoma,⁵⁶ K. Nakamura,⁵⁶ I. Nakano,⁴¹ A. Napier,⁵⁷ J. Nett,⁶⁰ C. Neu^{aa,46} M.S. Neubauer,²⁵ S. Neubauer,²⁷ J. Nielsen^{g,29} L. Nodulman,² M. Norman,¹⁰ O. Norriella,²⁵ E. Nurse,³¹ L. Oakes,⁴³ S.H. Oh,¹⁷ Y.D. Oh,²⁸ I. Oksuzian,¹⁹ T. Okusawa,⁴² R. Orava,²⁴ K. Osterberg,²⁴ S. Pagan Griso^{ee,44} C. Pagliarone,⁵⁵ E. Palencia,¹⁸ V. Papadimitriou,¹⁸ A. Papaikonomou,²⁷ A.A. Paramanov,² B. Parks,⁴⁰ S. Pashapour,³⁴ J. Patrick,¹⁸ G. Pauletta^{jj,55} M. Paulini,¹³ C. Paus,³³ T. Peiffer,²⁷ D.E. Pellett,⁸ A. Penzo,⁵⁵ T.J. Phillips,¹⁷ G. Piacentino,⁴⁷ E. Pianori,⁴⁶ L. Pinera,¹⁹ K. Pitts,²⁵ C. Plager,⁹ L. Pondrom,⁶⁰ K. Potamianos,⁴⁹ O. Poukhov^{*,16} F. Prokoshin^{z,16} A. Pronko,¹⁸ F. Ptohos^{t,18} E. Pueschel,¹³ G. Punzi^{ff,47} J. Pursley,⁶⁰ J. Rademacker^{c,43} A. Rahaman,⁴⁸ V. Ramakrishnan,⁶⁰ N. Ranjan,⁴⁹ I. Redondo,³² P. Renton,⁴³ M. Renz,²⁷ M. Rescigno,⁵² S. Richter,²⁷ F. Rimondi^{dd,6} L. Ristori,⁴⁷ A. Robson,²² T. Rodrigo,¹² T. Rodriguez,⁴⁶ E. Rogers,²⁵ S. Rolli,⁵⁷ R. Roser,¹⁸ M. Rossi,⁵⁵ R. Rossin,¹¹ P. Roy,³⁴ A. Ruiz,¹² J. Russ,¹³ V. Rusu,¹⁸ B. Rutherford,¹⁸ H. Saarikko,²⁴ A. Safonov,⁵⁴ W.K. Sakumoto,⁵⁰ L. Santi^{jj,55} L. Sartori,⁴⁷ K. Sato,⁵⁶ V. Saveliev^{v,45} A. Savoy-Navarro,⁴⁵ P. Schlabach,¹⁸ A. Schmidt,²⁷ E.E. Schmidt,¹⁸ M.A. Schmidt,¹⁴ M.P. Schmidt^{*,61} M. Schmitt,³⁹ T. Schwarz,⁸ L. Scodellaro,¹² A. Scribano^{gg,47} F. Scuri,⁴⁷ A. Sedov,⁴⁹ S. Seidel,³⁸ Y. Seiya,⁴² A. Semenov,¹⁶ L. Sexton-Kennedy,¹⁸ F. Sforza^{ff,47} A. Sfyrla,²⁵ S.Z. Shalhout,⁵⁹ T. Shears,³⁰ P.F. Shepard,⁴⁸ M. Shimojima^{t,56} S. Shiraishi,¹⁴ M. Shochet,¹⁴ Y. Shon,⁶⁰ I. Shreyber,³⁷ A. Simonenko,¹⁶ P. Sinervo,³⁴ A. Sisakyan,¹⁶ A.J. Slaughter,¹⁸ J. Slaunwhite,⁴⁰ K. Sliwa,⁵⁷ J.R. Smith,⁸ F.D. Snider,¹⁸ R. Snihur,³⁴ A. Soha,¹⁸ S. Somalwar,⁵³ V. Sorin,⁴ P. Squillacioti^{gg,47} M. Stanitzki,⁶¹ R. St. Denis,²² B. Stelzer,³⁴ O. Stelzer-Chilton,³⁴ D. Stentz,³⁹ J. Strologas,³⁸ G.L. Strycker,³⁵ J.S. Suh,²⁸ A. Sukhanov,¹⁹ I. Suslov,¹⁶ A. Taffard^{f,25} R. Takashima,⁴¹ Y. Takeuchi,⁵⁶ R. Tanaka,⁴¹ J. Tang,¹⁴ M. Tecchio,³⁵ P.K. Teng,¹ J. Thom^{h,18} J. Thome,¹³ G.A. Thompson,²⁵ E. Thomson,⁴⁶ P. Tipton,⁶¹ P. Ttito-Guzmán,³² S. Tkaczyk,¹⁸ D. Toback,⁵⁴ S. Tokar,¹⁵ K. Tollefson,³⁶ T. Tomura,⁵⁶ D. Tonelli,¹⁸ S. Torre,²⁰ D. Torretta,¹⁸ P. Totaro^{jj,55} M. Trovato^{hh,47} S.-Y. Tsai,¹ Y. Tu,⁴⁶ N. Turini^{gg,47} F. Ukegawa,⁵⁶ S. Uozumi,²⁸ N. van Remortel^{b,24} A. Varganov,³⁵ E. Vataga^{hh,47} F. Vázquez^{n,19} G. Velev,¹⁸ C. Vellidis,³ M. Vidal,³² I. Vila,¹² R. Vilar,¹² M. Vogel,³⁸ I. Volobouev^{x,29} G. Volpi^{ff,47} P. Wagner,⁴⁶ R.G. Wagner,² R.L. Wagner,¹⁸ W. Wagner^{bb,27} J. Wagner-Kuhr,²⁷ T. Wakisaka,⁴² R. Wallny,⁹ S.M. Wang,¹ A. Warburton,³⁴ D. Waters,³¹ M. Weinberger,⁵⁴ J. Weinelt,²⁷ W.C. Wester III,¹⁸ B. Whitehouse,⁵⁷ D. Whiteson^{f,46} A.B. Wicklund,² E. Wicklund,¹⁸ S. Wilbur,¹⁴ G. Williams,³⁴ H.H. Williams,⁴⁶ P. Wilson,¹⁸ B.L. Winer,⁴⁰ P. Wittich^{h,18} S. Wolbers,¹⁸ C. Wolfe,¹⁴ H. Wolfe,⁴⁰ T. Wright,³⁵ X. Wu,²¹ F. Würthwein,¹⁰ A. Yagil,¹⁰ K. Yamamoto,⁴² J. Yamaoka,¹⁷ U.K. Yang^{r,14} Y.C. Yang,⁶² W.M. Yao,²⁹ G.P. Yeh,¹⁸ K. Yi^{o,18} J. Yoh,¹⁸ K. Yorita,⁵⁸ T. Yoshida^{l,42} G.B. Yu,¹⁷ I. Yu,²⁸ S.S. Yu,¹⁸ J.C. Yun,¹⁸ A. Zanetti,⁵⁵ Y. Zeng,¹⁷ X. Zhang,²⁵ Y. Zheng^{d,9} and S. Zucchelli^{dd6}

(CDF Collaboration[†])

¹*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

²*Argonne National Laboratory, Argonne, Illinois 60439, USA*

³*University of Athens, 157 71 Athens, Greece*

⁴*Institut de Fisica d'Altes Energies, Universitat Autònoma de Barcelona, E-08193, Bellaterra (Barcelona), Spain*

⁵*Baylor University, Waco, Texas 76798, USA*

⁶*Istituto Nazionale di Fisica Nucleare Bologna, ^{dd}University of Bologna, I-40127 Bologna, Italy*

⁷*Brandeis University, Waltham, Massachusetts 02254, USA*

⁸*University of California, Davis, Davis, California 95616, USA*

⁹*University of California, Los Angeles, Los Angeles, California 90024, USA*

¹⁰*University of California, San Diego, La Jolla, California 92093, USA*

¹¹*University of California, Santa Barbara, Santa Barbara, California 93106, USA*

¹²*Instituto de Fisica de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain*

¹³*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

¹⁴*Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA*

¹⁵*Comenius University, 842 48 Bratislava, Slovakia; Institute of Experimental Physics, 040 01 Kosice, Slovakia*

¹⁶*Joint Institute for Nuclear Research, RU-141980 Dubna, Russia*

¹⁷*Duke University, Durham, North Carolina 27708, USA*

¹⁸*Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA*

¹⁹*University of Florida, Gainesville, Florida 32611, USA*

²⁰*Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy*

²¹*University of Geneva, CH-1211 Geneva 4, Switzerland*

²²*Glasgow University, Glasgow G12 8QQ, United Kingdom*

²³*Harvard University, Cambridge, Massachusetts 02138, USA*

²⁴*Division of High Energy Physics, Department of Physics, University of Helsinki and Helsinki Institute of Physics, FIN-00014, Helsinki, Finland*

- ²⁵University of Illinois, Urbana, Illinois 61801, USA
- ²⁶The Johns Hopkins University, Baltimore, Maryland 21218, USA
- ²⁷Institut für Experimentelle Kernphysik, Karlsruhe Institute of Technology, D-76131 Karlsruhe, Germany
- ²⁸Center for High Energy Physics: Kyungpook National University, Daegu 702-701, Korea; Seoul National University, Seoul 151-742, Korea; Sungkyunkwan University, Suwon 440-746, Korea; Korea Institute of Science and Technology Information, Daejeon 305-806, Korea; Chonnam National University, Gwangju 500-757, Korea; Chonbuk National University, Jeonju 561-756, Korea
- ²⁹Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA
- ³⁰University of Liverpool, Liverpool L69 7ZE, United Kingdom
- ³¹University College London, London WC1E 6BT, United Kingdom
- ³²Centro de Investigaciones Energeticas Medioambientales y Tecnologicas, E-28040 Madrid, Spain
- ³³Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA
- ³⁴Institute of Particle Physics: McGill University, Montréal, Québec, Canada H3A 2T8; Simon Fraser University, Burnaby, British Columbia, Canada V5A 1S6; University of Toronto, Toronto, Ontario, Canada M5S 1A7; and TRIUMF, Vancouver, British Columbia, Canada V6T 2A3
- ³⁵University of Michigan, Ann Arbor, Michigan 48109, USA
- ³⁶Michigan State University, East Lansing, Michigan 48824, USA
- ³⁷Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia
- ³⁸University of New Mexico, Albuquerque, New Mexico 87131, USA
- ³⁹Northwestern University, Evanston, Illinois 60208, USA
- ⁴⁰The Ohio State University, Columbus, Ohio 43210, USA
- ⁴¹Okayama University, Okayama 700-8530, Japan
- ⁴²Osaka City University, Osaka 588, Japan
- ⁴³University of Oxford, Oxford OX1 3RH, United Kingdom
- ⁴⁴Istituto Nazionale di Fisica Nucleare, Sezione di Padova-Trento, ^{ee}University of Padova, I-35131 Padova, Italy
- ⁴⁵LPNHE, Universite Pierre et Marie Curie/IN2P3-CNRS, UMR7585, Paris, F-75252 France
- ⁴⁶University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA
- ⁴⁷Istituto Nazionale di Fisica Nucleare Pisa, ^{ff}University of Pisa, ^{gg}University of Siena and ^{hh}Scuola Normale Superiore, I-56127 Pisa, Italy
- ⁴⁸University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA
- ⁴⁹Purdue University, West Lafayette, Indiana 47907, USA
- ⁵⁰University of Rochester, Rochester, New York 14627, USA
- ⁵¹The Rockefeller University, New York, New York 10021, USA
- ⁵²Istituto Nazionale di Fisica Nucleare, Sezione di Roma 1, ⁱⁱSapienza Università di Roma, I-00185 Roma, Italy
- ⁵³Rutgers University, Piscataway, New Jersey 08855, USA
- ⁵⁴Texas A&M University, College Station, Texas 77843, USA
- ⁵⁵Istituto Nazionale di Fisica Nucleare Trieste/Udine, I-34100 Trieste, ^{jj}University of Trieste/Udine, I-33100 Udine, Italy
- ⁵⁶University of Tsukuba, Tsukuba, Ibaraki 305, Japan
- ⁵⁷Tufts University, Medford, Massachusetts 02155, USA
- ⁵⁸Waseda University, Tokyo 169, Japan
- ⁵⁹Wayne State University, Detroit, Michigan 48201, USA
- ⁶⁰University of Wisconsin, Madison, Wisconsin 53706, USA
- ⁶¹Yale University, New Haven, Connecticut 06520, USA
- ⁶²Center for High Energy Physics: Kyungpook National University, Daegu 702-701, Korea; Seoul National University, Seoul 151-742, Korea; Sungkyunkwan University, Suwon 440-746, Korea; Korea Institute of Science and Technology Information, College London Daejeon 305-806, Korea; Chonnam National University, Gwangju 500-757, Korea; Chonbuk National University, Jeonju 561-756, Korea

We present a search for τ sneutrino production using the Tevatron $p\bar{p}$ collision data collected with the CDF II detector and corresponding to an integrated luminosity of 1 fb^{-1} . We focus on the scenarios predicted by the R-parity violating (RPV) supersymmetric models in which τ sneutrinos decay to two charged leptons of different flavor. With the data consistent with the standard model expectations, we set the upper limits on $\sigma(p\bar{p} \rightarrow \tilde{\nu}_\tau) \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu, \mu\tau, e\tau)$ and use these results to constrain the RPV couplings as a function of τ sneutrino mass.

Supersymmetric (SUSY) extensions of the standard model (SM) are among the leading candidates for a theory which can solve the hierarchy problem and provide a framework for unifying particle interactions [1]. Gauge-invariant and renormalizable interactions introduced in the SUSY models can violate the conservation of baryon (B) and lepton (L) number and lead to a proton lifetime shorter than the current experimental limits [2]. This problem is usually solved by postulating conservation of an additional quantum number, R-parity $R_p = (-1)^{3(B-L)+2s}$, where s is the particle spin [3]. However, models with R-parity-violating (RPV) interactions conserving spin and either B or L can also avoid direct contradiction with the proton lifetime upper limits [4]. Such models have the advantage that they naturally introduce lepton flavor violation and can generate non-zero neutrino masses and angles [5] consistent with neutrino-oscillation data [6]. They can also explain the recently reported anomalous phase of the b to s transition [7]. From an experimental standpoint, RPV interactions allow for single production of supersymmetric particles (sparticles) in high-energy particle collisions and for sparticles to decay directly into SM particles only; this makes the lightest sparticle unstable and critically affects the experimental strategy of the SUSY searches. Due to their clean final-state signatures, processes of single slepton production followed by decay to a pair of SM charged leptons become promising search channels for R-parity violating SUSY particles [8].

In this Letter we report a search for a heavy τ sneu-

trino $\tilde{\nu}_\tau$ produced in quark-antiquark $d\bar{d}$ annihilation and decaying via lepton-flavor-violating interactions into $e\mu/\mu\tau/e\tau$ final states. The search is performed using data corresponding to an integrated luminosity of 1 fb^{-1} collected in $p\bar{p}$ collisions at $\sqrt{s} = 1.96 \text{ TeV}$ by the CDF II detector at the Tevatron. The results are analysed in the framework of the minimal supersymmetric extensions of the SM [1], where the RPV part of the superpotential relevant to the τ sneutrino production and decay can be written as

$$W_{RPV} = \lambda'_{ijk} \mathbf{L}_i \mathbf{Q}_j \bar{\mathbf{d}}_k + \frac{1}{2} \lambda_{ijk} \mathbf{L}_i \mathbf{L}_j \bar{\mathbf{e}}_k \quad (1)$$

\mathbf{L} and \mathbf{Q} in Eq.(1) are the $SU(2)$ doublet superfields of leptons and quarks; $\bar{\mathbf{e}}$, $\bar{\mathbf{u}}$, and $\bar{\mathbf{d}}$ are the $SU(2)$ singlet superfields of leptons, u-type and d-type quarks; λ' and λ are the Yukawa couplings at the production and decay vertex respectively; the indices i, j , and k denote the fermion generations. We assume single-coupling dominance and the third super-generation to be the lightest. The couplings $\lambda'_{311} = 0.10$ and $\lambda_{i3k} = 0.05$, compatible with the current indirect limits [9], are chosen as a benchmark point. Heavy sneutrinos have been extensively searched for at LEP [9]. Recently searches for heavy sneutrinos decaying into the $e\mu$ final state have been performed by the CDF [10] and D0 collaborations [11]. The results in this paper supersede [10]. This analysis also represents the first search for lepton-flavor-violating decays of heavy sneutrinos into final states involving a third generation lepton, the τ , at the Tevatron.

CDF II is a general-purpose particle detector, described in detail elsewhere [12]. This measurement uses information from the central tracker [13], calorimeters [14, 15], and muon detectors [16] for charged lepton reconstruction and identification. Reconstruction of photons and π^0 mesons makes extensive use of the CES, the central shower maximum detector which is embedded at a depth of six radiation lengths within the electromagnetic calorimeter [14]. The luminosity is measured by a hodoscopic system of Cherenkov counters [17]. The event geometry and kinematics are described using the azimuthal angle ϕ around the beamline and the pseudorapidity $\eta = -\ln \tan \frac{\theta}{2}$, where θ is the polar angle with respect to the beamline. The transverse energy and momentum of the reconstructed particles are defined as: $E_T = E \sin \theta$, $p_T = p \sin \theta$, where E is the energy and p is the momentum.

The data used in this measurement are collected using inclusive high- p_T electron and muon triggers which select high- p_T electron and muon candidates with $|\eta| \lesssim 1.0$. After event reconstruction, electron and muon candidates with $p_T \geq 20 \text{ GeV}/c$ are identified using the procedures described in [18]. In addition we use independent measurements of the electron energy in the CES to improve the overall electron selection efficiency and identification of electron candidates radiating sig-

*Deceased

†With visitors from ^aUniversity of Massachusetts Amherst, Amherst, Massachusetts 01003, ^bUniversiteit Antwerpen, B-2610 Antwerp, Belgium, ^cUniversity of Bristol, Bristol BS8 1TL, United Kingdom, ^dChinese Academy of Sciences, Beijing 100864, China, ^eIstituto Nazionale di Fisica Nucleare, Sezione di Cagliari, 09042 Monserrato (Cagliari), Italy, ^fUniversity of California Irvine, Irvine, CA 92697, ^gUniversity of California Santa Cruz, Santa Cruz, CA 95064, ^hCornell University, Ithaca, NY 14853, ⁱUniversity of Cyprus, Nicosia CY-1678, Cyprus, ^jUniversity College Dublin, Dublin 4, Ireland, ^kUniversity of Edinburgh, Edinburgh EH9 3JZ, United Kingdom, ^lUniversity of Fukui, Fukui City, Fukui Prefecture, Japan 910-0017, ^mKinki University, Higashi-Osaka City, Japan 577-8502, ⁿUniversidad Iberoamericana, Mexico D.F., Mexico, ^oUniversity of Iowa, Iowa City, IA 52242, ^pKansas State University, Manhattan, KS 66506, ^qQueen Mary, University of London, London, E1 4NS, England, ^rUniversity of Manchester, Manchester M13 9PL, England, ^sMuons, Inc., Batavia, IL 60510, ^tNagasaki Institute of Applied Science, Nagasaki, Japan, ^uUniversity of Notre Dame, Notre Dame, IN 46556, ^vObninsk State University, Obninsk, Russia, ^wUniversity de Oviedo, E-33007 Oviedo, Spain, ^xTexas Tech University, Lubbock, TX 79609, ^yIFIC(CSIC-Universitat de Valencia), 56071 Valencia, Spain, ^zUniversidad Tecnica Federico Santa Maria, 110v Valparaiso, Chile, ^{aa}University of Virginia, Charlottesville, VA 22906, ^{bb}Bergische Universität Wuppertal, 42097 Wuppertal, Germany, ^{cc}Yarmouk University, Irbid 211-63, Jordan, ^{kk}On leave from J. Stefan Institute, Ljubljana, Slovenia,

nificant energy due to the bremsstrahlung. The τ leptons are identified via their hadronic decays as narrow calorimeter clusters associated with one or three charged tracks [19]. As the neutrino from the τ decay escapes detection, the “visible” four-momentum of a τ candidate, p_τ^{vis} , is reconstructed summing the four-momenta of charged particle tracks and neutral particle CES showers with a pion mass hypotheses. The resolution in p_τ^{vis} is further improved by combining measurements of the track momenta and energies of the CES showers with the energy measurements in the calorimeter. A reconstructed τ candidate is required to have the visible transverse energy, E_T^{vis} , greater than 25 GeV and its most energetic track must have $p_T > 10$ GeV/ c . The invariant mass of its decay products, $M_\tau^{\text{vis}} = \sqrt{p_\tau^{\text{vis}2}}$, is required to be consistent with the τ lepton decay: $M_\tau^{\text{vis}} < (1.8 + 0.0455 \times (E_T^{\text{vis}}/\text{GeV} - 20)) \text{ GeV}/c^2$, where the second term in the formula accounts for a degradation of the resolution in M_τ^{vis} at high energy. Reconstructed τ candidates within $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2} < 0.4$ of an identified electron or of a loosely-identified muon with large electromagnetic energy deposit in the calorimeter are excluded from the analysis.

Events selected for the analysis are required to have two identified central ($|\eta| < 1$) lepton candidates of different flavor and opposite electric charge. The leptons have to be isolated: the extra energy measured within a cone of radius $\Delta R \leq 0.4$ surrounding the leptons must be less than 10% of the lepton energy. Events with leptons consistent with a photon conversion or a cosmic ray hypothesis are removed from the analysis sample [18].

Signal and background studies are performed using Monte Carlo (MC) samples generated by PYTHIA [20] with the Tune A of CTEQ5L parton distribution functions [21]. The detector response is simulated with a GEANT3-based package [22]. The trigger, reconstruction and identification efficiencies are measured using Z events as calibration samples [18].

The predicted yield of signal events is calculated using the next-to-leading order (NLO) $p\bar{p} \rightarrow \tilde{\nu}_\tau$ production cross section [23]. The total $\tilde{\nu}_\tau$ width is defined by the $d\bar{d}$ and $l_i l_k$ decay modes as $\Gamma_{\tilde{\nu}_\tau} = (3\lambda_{311}^2 + 2\lambda_{i3k}^2)M_{\tilde{\nu}_\tau}/16\pi$, where $M_{\tilde{\nu}_\tau}$ is $\tilde{\nu}_\tau$ mass.

There are several sources of background events that pass our analysis selections. We classify these contributions based on whether the lepton candidates reconstructed in these events originated from a “real” lepton (produced from a W or Z decay) or were a result of a hadron being misidentified as a lepton, lepton flavor mis-assignment or a secondary lepton inside a jet. We collectively refer to the lepton candidates of the second category as “fakes” and classify each contributing background process into Type I, II and III according to the typical number of real and fake leptons reconstructed. Type I contains events with two real leptons and includes

$Z/\gamma^* \rightarrow \tau\tau$, diboson (WW , WZ , ZZ) and $t\bar{t}$ events. Type I is therefore called physics background. Type II includes events with one reconstructed fake lepton. They come from either (i) the $W/Z/\gamma^* + \text{jet(s)}$ events where one of the reconstructed leptons is in fact a jet misidentified as a lepton or (ii) $Z/\gamma^* \rightarrow ee/\mu\mu$ events with one of the leptons misidentified as a lepton of a different flavor. The backgrounds in Type I and II are estimated using MC, and their expected event yields are normalized to the NLO cross sections [24–26]. Events with two fake leptons (Type III) are dominated by multi-jet events with two jets misidentified as leptons and $\gamma + \text{jets}$ events; in the latter case, a converted photon is not identified as such and gets reconstructed as an electron and a jet is misidentified as a μ or a τ . The contribution of the processes in Type III is estimated using a data sample with two leptons of the same-charge and assuming no charge correlation between the two misidentified leptons. Both Type II and III are called fake background.

The systematic uncertainties in this search arise from a number of sources. The uncertainty on the luminosity measurement is 6% [27]. Uncertainties on lepton identification efficiency are 3% for τ 's, 1% for electrons, and 1% for muons [18]. The jet-to- τ misidentification probability is measured with an accuracy of 15%. Uncertainties in the parton distribution functions (PDF) result in the systematic error on the predicted signal cross section, which varies from 4% to 20% and increases with the $\tilde{\nu}_\tau$ mass. Variations of the signal acceptance due to PDF uncertainties are less than 1%.

We search for a signal from $\tilde{\nu}_\tau$ decays into lepton pairs in the distributions for dilepton invariant mass, M_{ll} (note that in the case of a hadronically decaying τ , it is the τ visible energy that is used to calculate the mass). The low mass region, $50 \text{ GeV}/c^2 < M_{ll} < 110 \text{ GeV}/c^2$, is used to validate the event selection and the background normalization. The observed and expected event yields in this region are in good agreement, as summarized in Table I. With the normalization fixed, the backgrounds are extrapolated into the region $M_{ll} > 100 \text{ GeV}/c^2$, where the search is performed as a “blind” counting experiment. Figure 1 compares data distributions in M_{ll} to the SM expectations for each of the three channels. With no statistically significant excesses observed, we use the data to set upper limits on $\sigma(p\bar{p} \rightarrow \tilde{\nu}_\tau) \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$.

TABLE I: The observed and predicted event yields in the control region. Uncertainties on the predicted yields include both statistical and systematic contributions.

| Control Region | $50 \text{ GeV}/c^2 < M_{ll} < 110 \text{ GeV}/c^2$ | | |
|---------------------|---|------------------|------------------|
| Channel | $e\mu$ | $\mu\tau$ | $e\tau$ |
| Physics Backgrounds | 100.2 ± 6.5 | 262.4 ± 21.0 | 309.6 ± 24.7 |
| Fake Backgrounds | 9.4 ± 3.3 | 222.6 ± 31.7 | 577.8 ± 37.0 |
| Total SM | 109.6 ± 7.7 | 485.0 ± 40.9 | 887.4 ± 51.0 |
| Observed | 105 | 477 | 894 |

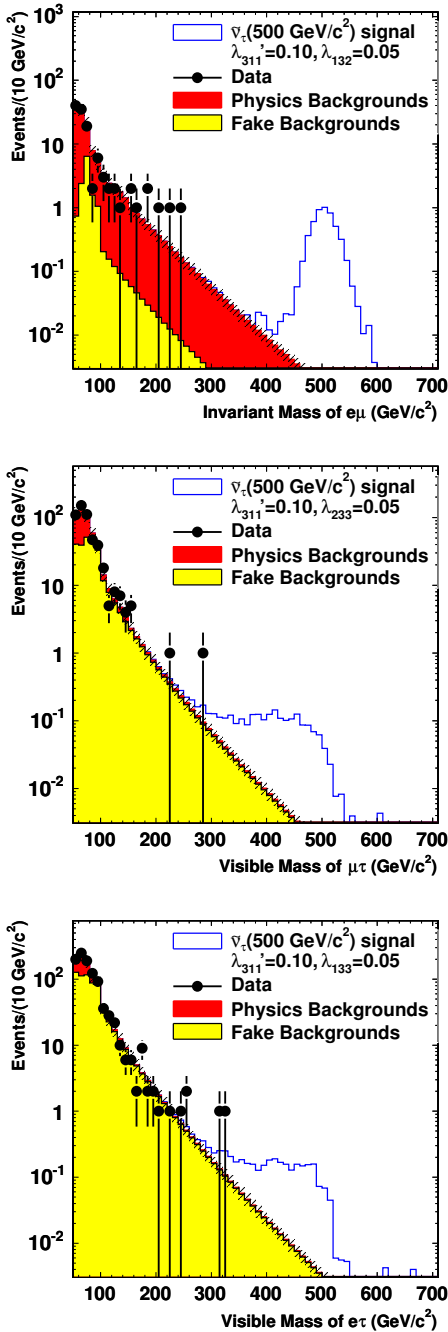


FIG. 1: Expected SM and observed distributions in M_{II} for $e\mu$, $\mu\tau$ and $e\tau$ channels. Also shown is an expected $\tilde{\nu}_\tau$ signal for $M_{\tilde{\nu}_\tau} = 500 \text{ GeV}/c^2$ and RPV couplings $\lambda'_{311} = 0.10$ and $\lambda_{i3k} = 0.05$.

In each of the channels ($e\mu$, $\mu\tau$, $e\tau$), the expected and observed upper limits are calculated using a Bayesian technique [28] at 95% credibility level (C.L.) as a function of $M_{\tilde{\nu}_\tau}$. For a given $M_{\tilde{\nu}_\tau}$, the limits are calculated by integrating the differential cross section $d\sigma/dM_{II}$ over the region $M_{II} > M_{II}^{\min}$, where M_{II}^{\min} optimizes the search

sensitivity for a selected $M_{\tilde{\nu}_\tau}$. The search results for $M_{\tilde{\nu}_\tau} = 500 \text{ GeV}/c^2$ are summarized in Table II. Figure 2 shows the expected and observed 95% C.L. upper limits on $\sigma(p\bar{p} \rightarrow \tilde{\nu}_\tau) \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$ as a function of $M_{\tilde{\nu}_\tau}$. We also set 95% C.L. upper limits on $\lambda'_{311}{}^2 \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$ as shown in Table III.

TABLE II: Expected and observed number of events in $e\mu$, $\mu\tau$, and $e\tau$ channels. The expected yields of $\tilde{\nu}_\tau$ events are calculated for $M_{\tilde{\nu}_\tau} = 500 \text{ GeV}/c^2$ and RPV couplings $\lambda'_{311} = 0.10$ and $\lambda_{i3k} = 0.05$.

| Channel | $e\mu$ | $\mu\tau$ | $e\tau$ |
|----------------------------------|-----------------|----------------|----------------|
| $M_{II}^{\min} (\text{GeV}/c^2)$ | 440 | 300 | 310 |
| Physics Backgrounds | 0.03 ± 0.01 | 0.1 ± 0.02 | 0.2 ± 0.03 |
| Fake Backgrounds | 0.01 ± 0.01 | 0.3 ± 0.1 | 0.6 ± 0.1 |
| Total SM background | 0.04 ± 0.01 | 0.4 ± 0.1 | 0.9 ± 0.1 |
| Expected signal | 5.9 ± 0.1 | 2.0 ± 0.1 | 2.7 ± 0.1 |
| Observed | 0 | 0 | 2 |

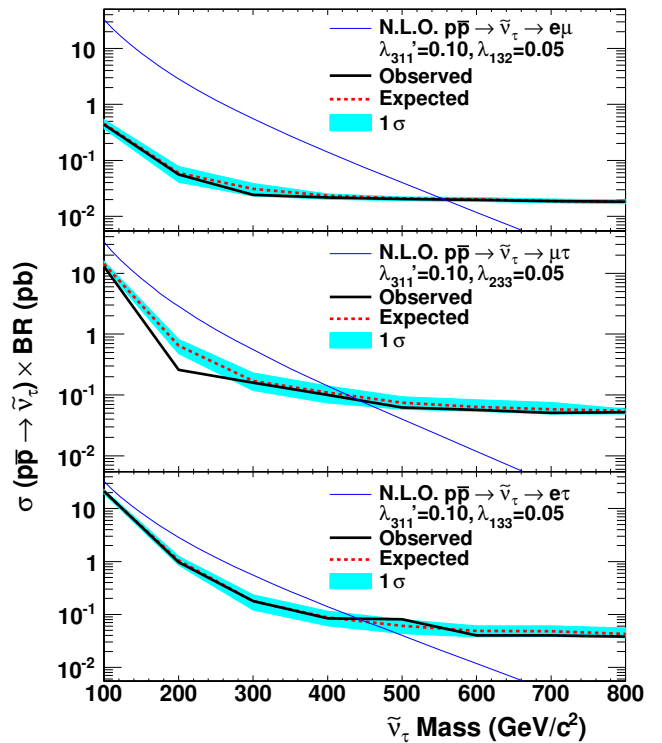


FIG. 2: The expected and observed 95% C.L. upper limits on $\sigma(p\bar{p} \rightarrow \tilde{\nu}_\tau) \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$ as a function of $M_{\tilde{\nu}_\tau}$.

In conclusion, we have searched for production of a massive sneutrino decaying to $e\mu$, $\mu\tau$, or $e\tau$ final states via R-parity violating interactions. We find the data consistent with the SM predictions and calculate the 95% C.L. upper limits on the $\sigma(p\bar{p} \rightarrow \tilde{\nu}_\tau) \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$ in the mass range up to $800 \text{ GeV}/c^2$. Using

TABLE III: 95% C.L. upper limits on $\lambda'_{311}{}^2 \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$.

| $M_{\tilde{\nu}_\tau}$ (GeV/ c^2) | $e\mu$ | $\mu\tau$ | $e\tau$ |
|---------------------------------------|--------------------|--------------------|--------------------|
| 300 | 6×10^{-5} | 4×10^{-4} | 5×10^{-4} |
| 400 | 2×10^{-4} | 1×10^{-3} | 9×10^{-4} |
| 500 | 7×10^{-4} | 2×10^{-3} | 3×10^{-3} |
| 600 | 2×10^{-3} | 7×10^{-3} | 5×10^{-3} |
| 700 | 8×10^{-3} | 2×10^{-2} | 2×10^{-2} |

these cross section limits, we constrain $\lambda'_{311}{}^2 \times \text{BR}(\tilde{\nu}_\tau \rightarrow e\mu/\mu\tau/e\tau)$ as a function of $M_{\tilde{\nu}_\tau}$. This analysis sets to the first Tevatron limits for lepton-flavor violating decays of heavy sneutrinos into final states involving a third generation lepton. For the RPV couplings $\lambda'_{311} = 0.10$ and $\lambda_{i3k} = 0.05$ the observed 95% C.L. lower limits on $\tilde{\nu}_\tau$ mass are 558 GeV/ c^2 in the $e\mu$ channel, 441 GeV/ c^2 in the $\mu\tau$ channel, and 442 GeV/ c^2 in the $e\tau$ channel.

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