

Heavy Diboson Production at the Tevatron

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Tremendous progress has been made in recent years in the measurement of heavy boson pair-production at the Tevatron collider. I will briefly review the current status of heavy diboson production measurements at CDF and DØ with special emphasis on the recently released ZZ cross section measurements.

1 Introduction

Leading order Feynman diagrams for WW , WZ and ZZ production include t -channel quark exchange and s -channel vector boson exchange diagrams. Measurements of the cross-sections for these processes test QCD at NLO and, through the presence of triple-gauge coupling vertices in the s -channel diagrams, test fundamental predictions of the EWK Standard Model. Experimentally, it is important to verify the rate of production of vector boson pairs since they are significant and often irreducible backgrounds to searches for new physics, most notably in the Higgs sector. Technically, these measurements are also informing Higgs boson search strategies at CDF and DØ .

WW ([1, 2]) and WZ production ([3, 4]) are firmly established in fully leptonic decay modes and have been used to place limits on anomalous triple-gauge couplings that are complementary to those obtained at other colliders [5, 4, 6, 7]. Interestingly, a first tentative signal has been observed in the semi-leptonic mode $WW/WZ \rightarrow l\nu jj$ by the CDF experiment [8].

2 ZZ Cross Section Measurements

2.1 CDF

CDF have performed searches in the modes $ZZ \rightarrow llll$ and $ZZ \rightarrow ll\nu\nu$ in 1.9 fb^{-1} of data [9]. *A posteriori*, the 4-lepton channel has the greatest sensitivity. 2 events are observed in lepton channels with a Z + jets background expectation of only $0.014^{+0.010}_{-0.008}$ and a further single event in a channel with 6 times larger background. The complete breakdown is given in table 1 and the event kinematics are compared to signal and background expectations in figure 1.

The CDF search in the $ZZ \rightarrow ll\nu\nu$ channel uses a matrix-element based likelihood method to distinguish the ZZ signal from the largely irreducible WW background. The final result, combining both $llll$ and $ll\nu\nu$ channels, is a measured cross section of $\sigma(ZZ) = 1.4^{+0.7}_{-0.6}$ (stat.+syst.) pb, in good agreement with the prediction $\sigma(ZZ)_{\text{NLO}} = 1.4 \pm 0.1$ pb. The probability of the background alone fluctuating to give the observed events in both channels is 5.1×10^{-6} .

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| Category | Candidates without a trackless electron | Candidates with a trackless electron |
|-------------------|--|---|
| ZZ | $1.990 \pm 0.013 \pm 0.210$ | $0.278 \pm 0.005 \pm 0.029$ |
| $Z + \text{jets}$ | $0.014^{+0.010}_{-0.007} \pm 0.003$ | $0.082^{+0.089}_{-0.060} \pm 0.016$ |
| Total | $2.004^{+0.016}_{-0.015} \pm 0.210$ | $0.360^{+0.089}_{-0.060} \pm 0.033$ |
| Observed | 2 | 1 |

Table 1: The breakdown of expected and observed numbers of events in 2 separate event classes in the CDF search for $ZZ \rightarrow llll$. Note that events containing a forward electron with no matching track suffer a significantly higher background contamination. Overall sensitivity is optimised by keeping the two sub-sets separate throughout the analysis.

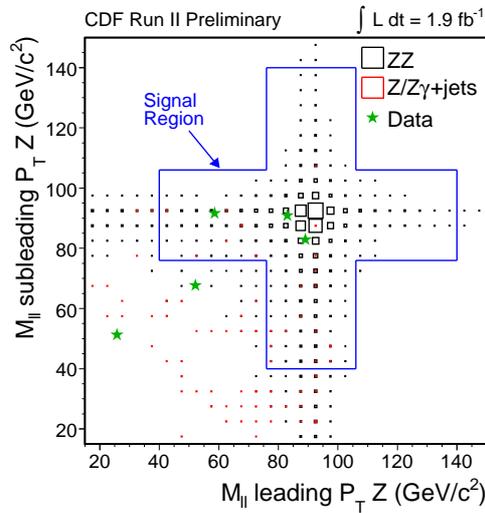


Figure 1: The distribution of event candidates in CDF's search for $ZZ \rightarrow llll$ events. The three stars inside the signal region are the final candidate events, compared to the expected distributions of both signal and background processes.

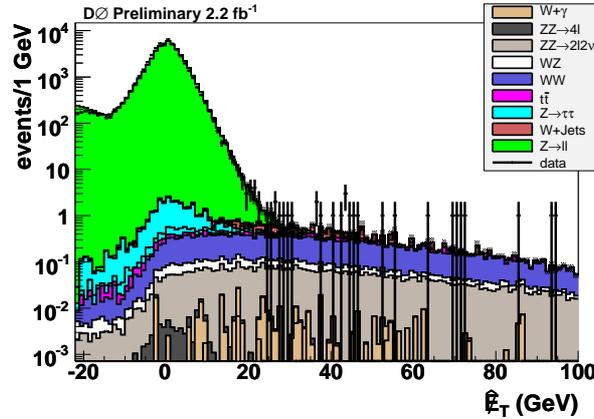


Figure 2: The distribution of effective missing- E_T for $ZZ \rightarrow eev\nu$ signal and backgrounds. A cut is made requiring this quantity to be larger than 30 GeV, effectively selecting a sample of WW and ZZ events that are then separated statistically.

2.2 $D\bar{O}$

$D\bar{O}$ have searched for ZZ production in the 4-lepton mode in an earlier data set and used the resulting upper limit to place constraints on non-Standard Model $ZZ(Z/\gamma^*)$ couplings [7].

Recently, $D\bar{O}$ have searched for ZZ production in the mode $ZZ \rightarrow \ell\nu\nu$ in 2.2 fb^{-1} of data [10]. They construct an effective missing- E_T object, which gives extra weight to the most reliable components of the missing- E_T measurement. The distribution of this quantity in the di-electron channel is shown in figure 2. As can be seen, a cut can be made which selects a relatively pure sample of WW and ZZ events. The events are separated statistically by forming a likelihood discriminant based on the (di-)lepton kinematics. The probability for both ee and $\mu\mu$ samples to be comprised entirely of non- ZZ events is 8.2×10^{-3} . Equivalently, a tentative cross section can be quoted: $\sigma(ZZ) = 2.1 \pm 1.2$ (stat.+syst.) pb. $D\bar{O}$ quote a theoretical cross section for comparison of $\sigma(ZZ)_{\text{NLO}} = 1.6 \pm 0.1$ pb. Note that the CDF quoted cross section has been computed in the zero-width approximation while the $D\bar{O}$ number includes lower mass γ^* contributions; however in both cases the measured cross sections have been corrected appropriately to provide a true comparison with theory.

3 Conclusions and Outlook

Figure 3 summarises recent cross section measurements in the electroweak sector at the Tevatron. Precision measurements are now being made of many diboson final states, and the last 2 years have seen both WZ and ZZ signals firmly established for the first time. So far, measured cross sections are in line with Standard Model expectations, confirming important predictions and placing further constraints on non-Standard Model couplings.

Excitingly, increased luminosity together with technical advances made partly through the measurement of heavy diboson production will soon yield experimental sensitivity to Standard Model Higgs production in certain mass ranges, as indicated in figure 3.

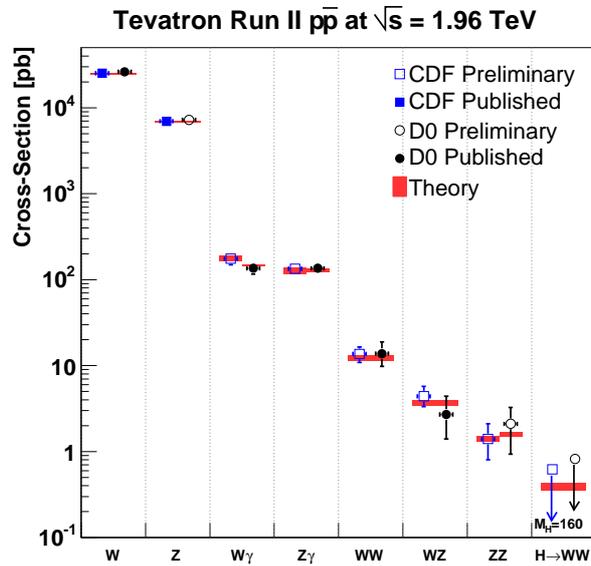


Figure 3: A summary of electroweak cross section measurements at Fermilab's Tevatron collider. The new ZZ cross section measurements described here are indicated second from right.

References

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