

Latest Results on Bottom Spectroscopy and Production with CDF

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Abstract

Using data collected with the CDF Run II detector, new measurements on bottom production cross-sections are presented. The latest achievements in bottom hadron spectroscopy are discussed. The results are based on a large sample of semileptonic and hadronic decays of bottom states made available by triggers based on the precise CDF tracking system.

1 First Observation of the Baryons Σ_b and Σ_b^* in CDF

The bottom $\Sigma_b^{(*)}$ states decay strongly into Λ_b^0 by emitting soft pion as shown in Figure 1. Our results are based on data collected with the CDF II detector [2] and corresponding to an integrated luminosity of $\sim 1.1 \text{ fb}^{-1}$. The trigger used in this study is based on displaced tracks. It reconstructs with the central tracker a pair of $p_T \gtrsim 2.0 \text{ GeV}/c$ tracks at Level 1 and enables secondary vertex selection at Level 2 requiring each of these tracks to have impact parameter measured by the CDF silicon detector SVX II larger than $120 \mu\text{m}$. The signals of $\Sigma_b^{(*)\pm}$ states were sought in the decay chain $\Sigma_b^{(*)\pm} \rightarrow \Lambda_b^0 \pi_{soft}^\pm$, $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$, $\Lambda_c^+ \rightarrow p K^- \pi^+$ ^a. To remove the contribution due to a mass resolution

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^aUnless otherwise stated all references to the specific charge combination imply the charge conjugate combination as well.

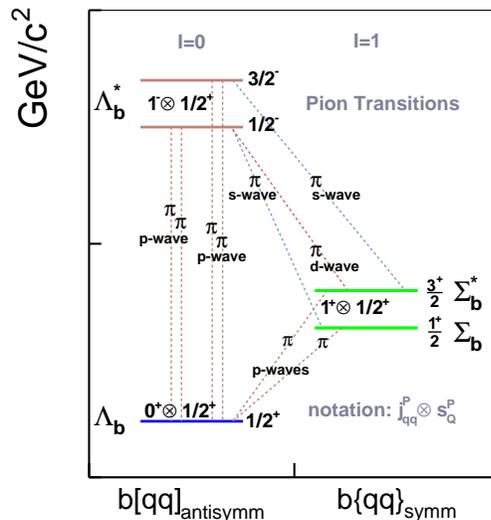


Figure 1: The low lying Σ^- - and Λ^- -like b -baryons and their strong decays with pion emissions.

of each Λ_b^0 candidate and to avoid absolute mass scale systematic uncertainties, the $\Sigma_b^{(*)\pm}$ candidates were reconstructed in the mass difference Q -value spectra defined as $Q = M(\Lambda_b^0 \pi_{soft}^\pm) - M(\Lambda_b^0) - M_{\text{PDG}}(\pi^\pm)$ for every charge state of $\Sigma_b^{(*)\pm}$ candidates. Here we assume also that the width of the weakly decaying Λ_b^0 candidate is determined by the corresponding detector mass resolution. The fitted experimental spectra are shown at Figure 2, and fit results are summarized in Tables 1 and 2 [3].

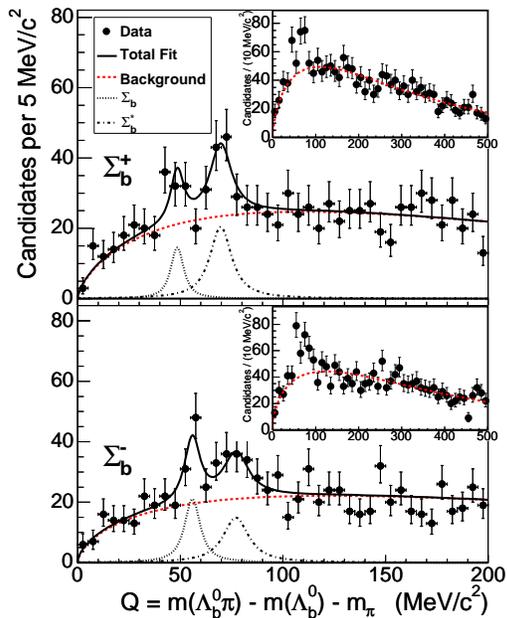


Figure 2: The experimental mass difference spectra [3] for the candidates of both charged partners, $\Sigma_b^{(*)\pm}$. Double peak signatures are observed in every case.

2 Observation and Mass Measurement of the Baryon Ξ_b

The bottom cascade baryons Ξ_b consist of a single bottom quark, one strange quark and one light quark. Theoretical predictions for these heavy baryons are outlined in Table 3 [4]. We consider the lowest lying Ξ_b states that decay weakly and the Ξ_b' states that decay radiatively or strongly via pion emission. The Ξ_b candidates are reconstructed in the decay chain $\Xi_b^- \rightarrow J/\psi \Xi^-$ with secondary states $J/\psi \rightarrow \mu^+ \mu^-$ and $\Xi^- \rightarrow \Lambda^0 \pi^-$, $\Lambda^0 \rightarrow p \pi^-$ (see Figure 3). Since experiments with bubble chambers the strange cascade, given its long decay path of $c \cdot \tau = 4.91$ cm [5], is identified as a charged track with a 1-track decay vertex at the end formed by a kinked soft pion track as shown at Figure 3. The subse-

State	Q or $\Delta_{\Sigma_b^*}$ (MeV/c ²)	Mass (MeV/c ²)
Σ_b^+	$Q_{\Sigma_b^+} = 48.5^{+2.0+0.2}_{-2.2-0.3}$	$5807.8^{+2.0}_{-2.2} \pm 1.7$
Σ_b^-	$Q_{\Sigma_b^-} = 55.9 \pm 1.0 \pm 0.2$	$5815.2 \pm 1.0 \pm 1.7$
Σ_b^{*+}	$\Delta_{\Sigma_b^*} = 21.2^{+2.0+0.4}_{-1.9-0.3}$	$5829.0^{+1.6+1.7}_{-1.8-1.8}$
Σ_b^{*-}		$5836.4 \pm 2.0^{+1.8}_{-1.7}$

Table 1: The masses resulting from the simultaneous fit of both spectra [3].

Yields of the signals			
Σ_b^+	Σ_b^-	Σ_b^{*+}	Σ_b^{*-}
32^{+13+5}_{-12-3}	59^{+15+9}_{-14-4}	77^{+17+10}_{-16-6}	69^{+18+16}_{-17-5}

Table 2: The fitted yields [3] of the identified $\Sigma_b^{(*)\pm}$ states. The combined significance of all four peaks relative to the null hypothesis well exceeds 5 Gaussian standard deviations.

quent V^0 decay vertex of the Λ^0 is associated with the 1-track vertex and included in a two-vertex kinematic fit. The key technique in this analysis is the tracking algorithm developed to reconstruct Ξ^- tracks leaving hits in the CDF silicon vertex tracker SVX II. A finest tracking resolution [2] coupled with the custom software provide a clean signal for Ξ^- , see Figure 4. The analysis [6] uses a data sample of integrated $\mathcal{L} = 1.9 \text{ fb}^{-1}$ collected by the CDF dimuon trigger [2] which saves events with two oppositely charged tracks reconstructed in the CDF cen-

State	b sq	J^P	I_3	j_{sq}	$M, \text{GeV}/c^2$
Ξ_b^0	$b\{su\}$	$1/2^+$	$1/2$	0	5.80
Ξ_b^-	$b\{sd\}$	$1/2^+$	$-1/2$	0	5.80
$\Xi_b^{0'}$	$b\{su\}$	$1/2^+$	$1/2$	1	5.94
$\Xi_b^{-'}$	$b\{sd\}$	$1/2^+$	$-1/2$	1	5.94

Table 3: Theoretical expectations for properties of bottom cascade baryons containing a single b -quark [4]. The lowest lying states have a light quark pair with momentum $j_{sq} = 0$ while the next ones have light quarks aligned with $j_{sq} = 1$.

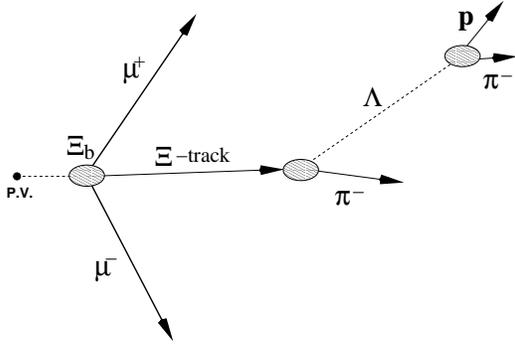


Figure 3: Topology of the $\Xi_b^- \rightarrow J/\psi \Xi^-$ decay.

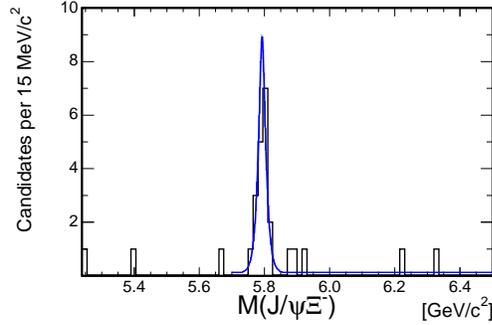


Figure 5: The invariant mass distribution of $J/\psi \Xi^-$ candidates after optimized selection criteria have been applied. The profile of the unbinned fit is superimposed. A clear signal is observed [6].

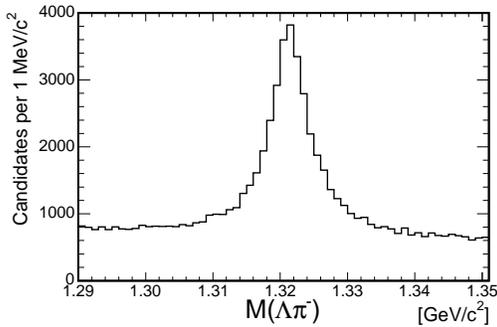


Figure 4: The Ξ^- signal [6] when the cascade track has at least 2 hits in the CDF SVX II tracker.

tral tracker, matched to hits in the CDF muon chambers and selected in the mass window $M(\mu^+\mu^-) \in [2.7, 4.0] \text{ GeV}/c^2$ around the mass of the J/ψ [5]. The sample yields $\sim 15 \times 10^6$ J/ψ and ~ 23500 Ξ^- candidates. The final selection criteria for Ξ_b^- candidates have been studied using ~ 31000 B -mesons in the mode $B^+ \rightarrow J/\psi K^+$ as a control sample assuming very similar decay kinematics. The invariant mass of selected $J/\psi \Xi^-$ candidates is shown in Figure 5. An unbinned likelihood fit finds [6] $17.5 \pm 4.3(\text{stat})$ Ξ_b^- candidates at a mass of $5792.9 \pm 2.5(\text{stat}) \pm 1.7(\text{syst}) \text{ MeV}/c^2$ and with a significance of 7.7 of Gaussian standard de-

viations. The results [6] are in good agreement with theoretical predictions and with the observation made by the DØ Collaboration [7].

3 Correlated $b\bar{b}$ Production in CDF II Detector

In this chapter we cover briefly a unique analysis on a paired $b\bar{b}$ production measurement. As leading order (LO) processes dominate $b\bar{b}$ production, $\sigma_{b\bar{b}}$, while next-to-leading (NLO) processes are essential for inclusive σ_b studies, the measurement of $\sigma_{b\bar{b}}$ will help to disentangle LO and NLO contributions and to resolve the controversy between the Run I DØ and CDF measurements [8]. We select dimuon events with invariant masses $5 < M(\mu_1\mu_2) < 80 \text{ GeV}/c^2$, outside of the domain populated by sequential decays of single b -quarks and Z^0 modes, and extract $\sigma(b \rightarrow \mu^- + X, \bar{b} \rightarrow \mu^+ + X)$, subtracting contributions from $c\bar{c}$, prompt Drell-Yan pairs, c - and b -onium prompt decays, π^- , K -decays, and misidentified dimuon candidates. The signal and background contributions are determined by fitting the experimental 2-dimensional impact parameter $d_0(\mu_1), d_0(\mu_2)$

distribution to corresponding templates expected for various dimuon sources. The method exploits the fact that the shape of the $d_0(\mu)$ distribution is largely determined by the lifetime of its parent heavy hadron. The analysis is based on a data sample of total luminosity $\mathcal{L} = 740 \text{ pb}^{-1}$ collected with the CDF dimuon trigger [2] having no biases with respect to $d_0(\mu)$ distribution. The projection of the 2-dimensional fit onto $d_0(\mu)$ comprising various background contributions is shown in Figure 6. The extracted exper-

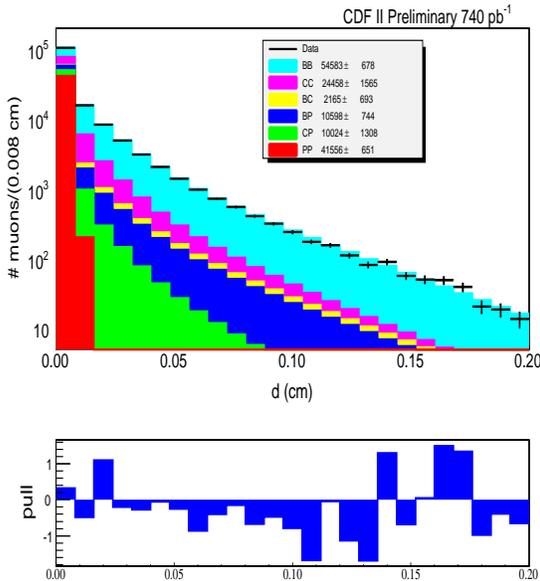


Figure 6: The projection of the 2-dimensional fit of $d_0(\mu_1), d_0(\mu_2)$ with background templates summed up and data superimposed. The notations used are “B” as b -source, “C” as c -source and “P” as the source of prompt muons.

imental cross-section is found to be $\sigma(b \rightarrow \mu^-, \bar{b} \rightarrow \mu^+) = 1549 \pm 133 \text{ pb}$. The exact NLO predictions are made using Herwig Monte-Carlo program [9], MNR code [10] running with EVTGEN generator [11], parton structure functions from MRST [12] fits and Peterson fragmentation function [13].

The ratio of data to NLO theoretical Monte-Carlo calculation is found to be $R2(b \rightarrow \mu^-, \bar{b} \rightarrow \mu^+) = 1.20 \pm 0.21$. The errors include statistical and systematic uncertainties added in quadratures. From this measurement we derive $\sigma(b\bar{b}, p_T \geq 6 \text{ GeV}/c, |y| \leq 1) = 1618 \pm 148 \text{ nb}$. The systematic uncertainty due to choice of the fragmentation model is $\sim 25\%$.

4 Summary

CDF announces the first observation of four bottom baryon $\Sigma_b^{(*)\pm}$ resonance states. CDF has also observed the strange bottom cascade baryon Ξ_b^- , and our measurements are in agreement with the DØ observation and with theoretical predictions. CDF II detector has measured the correlated production cross-section of $b\bar{b}$ pairs with b -quarks identified in their muonic semileptonic modes. The measurement is consistent with theoretical expectations. Using NLO Monte-Carlo cross-section calculations, the full $b\bar{b}$ production cross-section in the kinematic domain ($p_T \geq 6 \text{ GeV}/c, |y| \leq 1$) has been derived.

5 Acknowledgments

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