

# Electroweak Physics with CDF

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**Abstract.** The CDF experiment at the Tevatron has used  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV to perform electroweak physics measurements. A program of *precision* electroweak tests of SM started measuring  $W$  and  $Z$  bosons cross section using different leptonic final states, evaluating dielectron Forward-Backward Asymmetry  $A_{FB}$  and di-boson cross section production.

**PACS.** 13.38.Be Decays of  $W$  bosons – 13.38.Dg Decays of  $Z$  bosons

## 1 Introduction

The Collider Detector at Fermilab (CDF) is a general purpose detector located in one of the interaction regions at the Tevatron collider.  $p\bar{p}$  collisions at the Tevatron reach an energy in the center of mass of  $\sqrt{s} = 1.96$  TeV. We are reporting here the electroweak physics measurements performed using the first physics-quality data of RunII taken from March 2002 to January 2003.

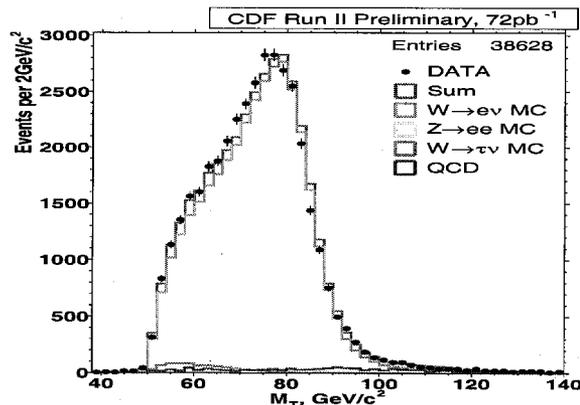
## 2 $W$ and $Z$ Cross Sections

$W$  and  $Z$  bosons are produced by  $q\bar{q}$  annihilation. Due to the large QCD background, decay channels of bosons involving quarks are difficult to identify; therefore  $W$  and  $Z$  bosons are identified through their leptonic decays.

### 2.1 Measuring $W$ Cross Section

The signature for a leptonic  $W$  boson decay is a high momentum isolated lepton with missing transverse energy accounting for the undetected neutrino.

$W \rightarrow e\nu$  candidates are collected with a trigger selecting high- $E_T$  central electron candidates; after requiring one tight electron with  $E_T > 25$  GeV matched to a track with  $P_T > 10$  GeV/c and missing transverse energy  $\cancel{E}_T > 25$  GeV, 38628 events are left in data. The transverse mass distribution of the candidate events is reported in Fig.1. Background from QCD dijets is estimated from data assuming that its distribution is flat with respect to the missing transverse energy. Background contamination from other electroweak processes like  $W \rightarrow \mu\nu$ ,  $Z \rightarrow ee$  and  $W \rightarrow \tau\nu$  is estimated from MC events after a detailed simulation through our detector.



**Fig. 1.** Transverse Mass distribution of  $W \rightarrow e\nu$  events collected by the CDF experiment

$W \rightarrow \mu\nu$  candidates are collected by a high- $P_T$  muon trigger. After requiring an isolated muon with  $P_T > 20$  GeV/c and  $\cancel{E}_T > 20$  GeV, 21599  $W$  candidates remain. The main background contamination comes from  $Z \rightarrow \mu\mu$ ,  $W \rightarrow \tau\nu$ , cosmic rays and dijet QCD events.

Specific triggers for selecting a sample enriched with  $\tau$  events decaying into hadrons have been designed for RunII [1]. Events with  $\cancel{E}_T > 25$  GeV and with the topology of the  $\tau$  decaying hadronically are selected at trigger level.  $W \rightarrow e\nu$  events are explicitly removed. 2345 events pass the selection with an estimated background mainly from QCD of  $612 \pm 61$  events. Using the removed  $W \rightarrow e\nu$  events collected in the same data sample, the ratio of the electroweak coupling constant has been measured:

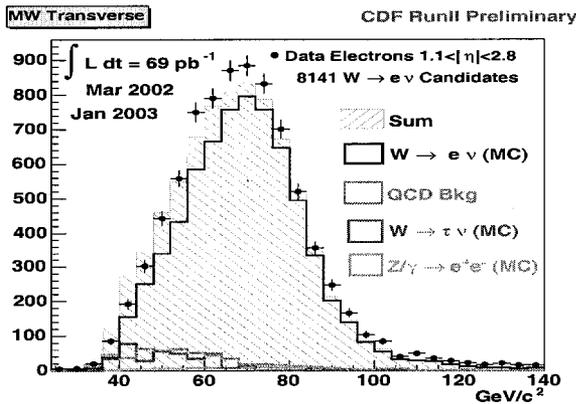
$$g_\tau/g_e = 0.99 \pm 0.04(\text{stat.}) \pm 0.07(\text{sys.}), \quad (1)$$

consistent with SM expectations.

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**Table 1.** Yields of  $W$  boson events in the different leptonic decay channels with the measured  $\sigma \times BR(W \rightarrow \ell\nu)$  ( $\mathcal{L} = 72\text{pb}^{-1}$ ). Quoted uncertainties are respectively for statistics, systematics and luminosity.

Channel	Events	Bkg. (%)	$\sigma \times BR(W \rightarrow \ell\nu)$ (nb)
$e\nu_e$	38625	6	$2.64 \pm 0.02 \pm 0.09 \pm 0.16$
$\mu\nu_\mu$	21599	11	$2.64 \pm 0.02 \pm 0.12 \pm 0.16$
$\tau\nu_\tau$	2345	26	$2.62 \pm 0.07 \pm 0.21 \pm 0.16$
$e\nu_e, \mu\nu_\mu$	combined		$2.640 \pm 0.012 \pm 0.093 \pm 0.158$



**Fig. 2.** Transverse Mass distribution of  $W \rightarrow e\nu$  with electrons detected by forward calorimeters ( $1.1 \leq |\eta| \leq 2.8$ )

Estimating the acceptance of the selection in the different decay channels, the production cross section times branching ratio is measured and reported in Tab.1.

The measured values are in agreement with the predicted theoretical values[2] (NNLO) of  $2.731 \pm 0.002$  nb.

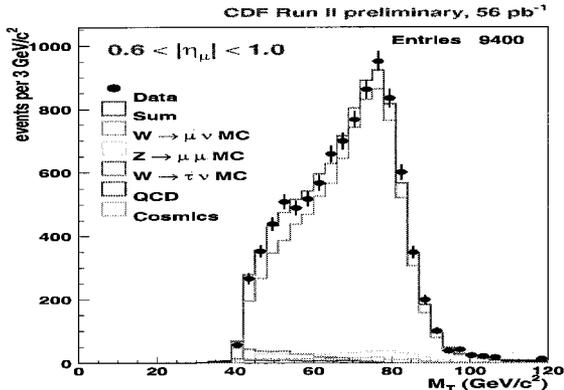
### 2.1.1 Extending Detector Acceptance

A substantial part of the RunII upgrade was devoted to extend the pseudorapidity coverage of lepton identification[3].

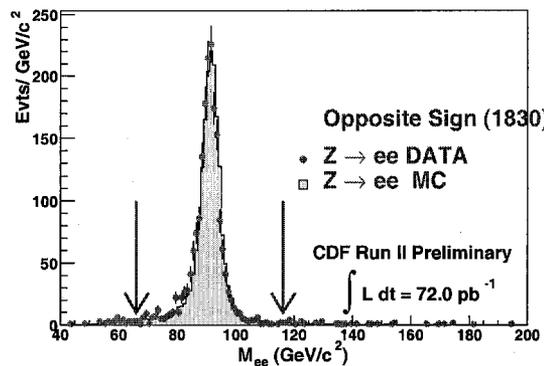
$W \rightarrow e\nu$  candidates are selected using a trigger looking for an electromagnetic cluster in the plug calorimeter that covers the pseudorapidity region  $1 < |\eta| < 3.6$ . Events are selected with  $\cancel{E}_T > 20$  GeV,  $E_T > 20$  GeV and matching the cluster with a track. As pseudorapidity increases, the acceptance of the Central Outer Tracker [6] gets smaller; it is therefore necessary to perform tracking with the silicon detectors[7,8]. Transverse mass of  $W \rightarrow e\nu$  candidates is shown in Fig.2. RunII muon detection system covers the region  $|\eta| \leq 1.5$ [4]. Transverse mass of  $W \rightarrow \mu\nu$  candidates for  $0.6 \leq |\eta| \leq 1.0$  is shown in Fig.3.

## 2.2 Measuring Z Cross Section

$Z/\gamma^* \rightarrow \ell\ell$  events are selected requesting two central high- $P_T$  isolated leptons with opposite charge.  $Z/\gamma^* \rightarrow ee$  events are required to have two central electrons with opposite charge and with  $E_T > 25$  GeV and  $P_T > 10$  GeV/c.



**Fig. 3.** Transverse Mass distribution of  $W \rightarrow \mu\nu$  with muons identified for  $0.6 < |\eta| < 1.0$ .



**Fig. 4.** Dielectron invariant mass distribution of  $Z \rightarrow e^+e^-$  candidates.

**Table 2.** Yields of events in the different leptonic decay channels of  $Z$  boson with the measured  $\sigma \times BR(Z \rightarrow \ell^+\ell^-)$  ( $\mathcal{L} = 72\text{pb}^{-1}$ ). Quoted uncertainties are respectively for statistics, systematics and luminosity.

Channel	Events	Bkg. (%)	$\sigma \times BR(Z \rightarrow \ell^+\ell^-)$ (nb).
$e^+e^-$	1830	0.6	$267 \pm 6 \pm 15 \pm 16$
$\mu^+\mu^-$	1631	0.9	$246 \pm 6 \pm 12 \pm 15$
$e^+e^-, \mu^+\mu^-$	combined		$251.5 \pm 4.3 \pm 10.6 \pm 15.1$

The invariant mass is required to be between 66 and 116  $\text{GeV}/c^2$ . The total yield of candidates is 1830 with an estimated background of  $10 \pm 5$  events. The invariant mass of the dielectron pair is reported in Fig.4.

The  $Z/\gamma^* \rightarrow \mu\mu$  candidates are selected requiring one isolated central muon with  $P_T > 20$   $\text{GeV}/c$  and a second isolated high- $P_T$  track passing minimum ionizing energy requirements. The production cross section times branching fraction has been measured; the values are reported in Tab.2 and are in agreement with the predicted  $Z$  boson production (NNLO)  $250.5 \pm 3.8$   $\text{pb}$  [5].

$Z \rightarrow \tau^+\tau^-$  candidates are selected requiring one  $\tau$  identified from its electronic decay and the other one from its hadronic decay. To increase the purity we require

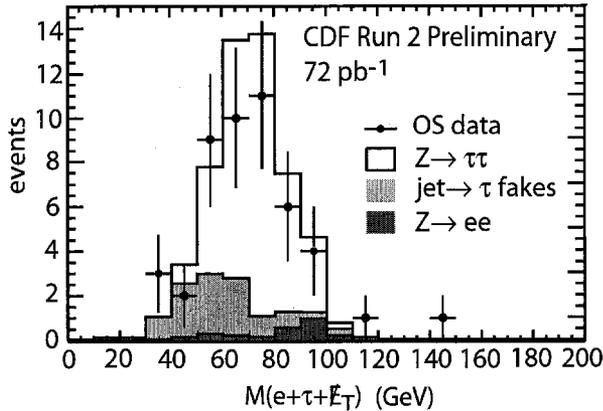


Fig. 5. Invariant Mass distribution of  $Z \rightarrow \tau\tau$  candidates.

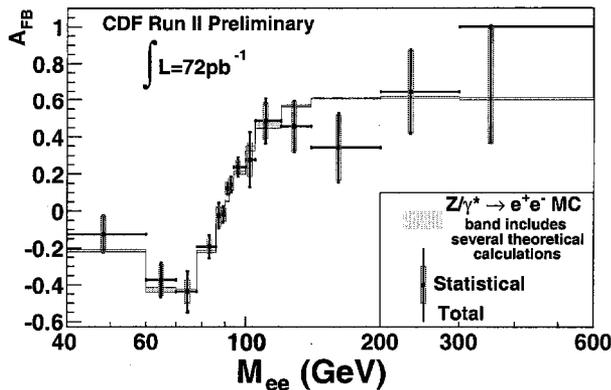


Fig. 6. Forward Backward asymmetry for dielectron pairs compared with theoretical predictions.

$M_T(e, \cancel{E}_T) \leq 25 \text{ GeV}/c^2$  and  $P_T(e, \cancel{E}_T) \geq 25 \text{ GeV}/c$ . The invariant mass is reported in Fig.5.

### 3 Precision Electroweak Measurements

#### 3.1 Dielectron Forward Backward Asymmetry

With dielectron pairs created by Drell-Yan process it is possible to measure the Forward Backward asymmetry ( $A_{FB}$ ).  $A_{FB}$  is a probe of the strength of the vector and axial-vector couplings. What is unique at the Tevatron experiments is the possibility to probe  $A_{FB}$  not only at the  $Z$ -pole but up to a range of  $M_{ee}$  of  $600 \text{ GeV}/c^2$  as shown in Fig.6.

#### 3.2 Diboson production

Electroweak interaction allow bosons to self-interact. Therefore direct production of  $W\gamma$  and  $Z\gamma$  is searched. These processes probe anomalous couplings increasing therefore the sensitivity to physics beyond the Standard Model.  $W\gamma$  and  $Z\gamma$  events are searched starting from the samples described in previous section. Additionally a high energy

**Table 3.** Yield of events in the  $e$  and  $\mu$  decay channel of direct  $W\gamma$  production ( $\mathcal{L} = 72\text{pb}^{-1}$ ). Quoted uncertainties are respectively for statistics, systematics and luminosity.

Channel	Events	Bkg. (%)	$\sigma \times BR(W\gamma \rightarrow \ell\nu\gamma)$ (pb)
$e$	43	33	$17.2 \pm 3.8 \pm 2.8 \pm 1.0$
$\mu$	38	29	$19.8 \pm 4.5 \pm 2.4 \pm 1.2$

**Table 4.** Yields of events in the  $e$  and  $\mu$  decay channel of direct  $Z\gamma$  production ( $\mathcal{L} = 72\text{pb}^{-1}$ ). Quoted uncertainties are respectively for statistics, systematics and luminosity.

Channel	Events	Bkg. (%)	$\sigma \times BR(Z\gamma \rightarrow \ell^- \ell^+ \gamma)$ (pb)
$e$	11	4.6	$5.5 \pm 1.7 \pm 0.6 \pm 0.3$
$\mu$	14	4.0	$6.0 \pm 1.6 \pm 0.7 \pm 0.4$

photon with  $E_T > 7 \text{ GeV}$  is required with  $\Delta R_{(\gamma,\ell)} \geq 0.7$ . The yield of  $W\gamma$  and  $Z\gamma$  events are reported respectively in Tab.3 and Tab.4. The predicted production cross sections are respectively  $18.7 \pm 1.3 \text{ pb}$  and  $5.4 \pm 0.4 \text{ pb}$ .

$WW$  candidates are searched in the dilepton decay channel  $WW \rightarrow \ell\ell\nu\nu'$  ( $\ell = e, \mu$ ). Candidates events are selected requiring two high- $P_T$  isolated leptons with opposite charge and  $\cancel{E}_T > 25 \text{ GeV}$ . Additional requirements reject  $t\bar{t}$  dilepton and  $Z$  bosons candidates. Combining the  $ee$ ,  $\mu\mu$  and  $e\mu$  channels 2 events are observed with an estimated background of  $1.53 \pm 0.64$  and  $2.79 \pm 0.62$  signal events expected.

## 4 Conclusions

Electroweak measurements have been performed using the first *physics quality* data collected by the CDF experiment.  $W$  and  $Z$  production cross section values, dielectron forward backward asymmetry and diboson production have been established showing a substantial agreement with SM predictions. These latter measurements are statistically limited; larger data samples have been collected and are ready to be analyzed allowing a larger precision in electroweak measurements.

## References

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