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Search for Scalar Top

Christopher M. Holck
For the CDF Collaboration
University of Pennsylvania

*Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510*

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Christopher M. Holck
University of Pennsylvania

for the

CDF Collaboration

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Abstract

We report results of a search for direct production of scalar top quark decaying to charm + LSP. The observed number of events is consistent with expectations from Standard Model processes. Assuming $BR=100\%$ we set a 95% CL in the $m_{\tilde{t}}-m_{\tilde{\chi}_1^0}$ parameter space.

Supersymmetry (SUSY) assigns to every fermionic SM particle a bosonic superpartner and vice versa. Therefore, the SM quark helicity states q_L and q_R acquire scalar SUSY partners \tilde{q}_L and \tilde{q}_R which are also the mass eigenstates for the first two generations to a good approximation. However, a large mixing can occur in the third generation leading to a large splitting between the mass eigenstates [1]. This can lead to a stop quark which is not only the lightest squark but also lighter than the top quark.

At the Tevatron, stop is produced in pairs via gg and $q\bar{q}$ fusion diagrams. The cross section depends only on stop quark mass at leading order [2]. The dominant next-to-leading order SUSY corrections depend on the other squark masses and are small ($\sim 1\%$). For $m_{\tilde{t}}=115 \text{ GeV}/c^2$, $\sigma(\tilde{t}\bar{\tilde{t}}) = 6.8 \text{ pb}$ in the next-to-leading order.

Whenever kinematically allowed, $\tilde{t} \rightarrow b\tilde{\chi}_1^+$. If this channel is closed but sneutrino ($\tilde{\nu}$) is light, then $\tilde{t} \rightarrow b\tilde{\nu}$ dominates. If neither of these channels is allowed, then stop decays via a one-loop diagram to charm+neutralino: $\tilde{t} \rightarrow c\tilde{\chi}_1^0$, which is the mode considered for this search. We assume that $\tilde{\chi}_1^0$ is the lightest SUSY particle and that R-parity is conserved. Hence the $\tilde{\chi}_1^0$ is undetected and causes an imbalance of energy. In this scenario, the signature of $\tilde{t}\bar{\tilde{t}}$ is 2 acolinear charm jets, significant missing transverse energy (\cancel{E}_T), and no lepton(s) in the final state.

We have searched data corresponding to a total integrated luminosity of $88 \pm 6 \text{ pb}^{-1}$ collected using the CDF detector during the 1993-95 Tevatron run.

Sample	N_{exp}
$W^\pm(\rightarrow \mu^\pm \nu_\mu) + \geq 2$ jets	1.5 ± 0.8
$W^\pm(\rightarrow \tau^\pm \nu_\tau) + \geq 1$ jet	6.9 ± 2.6
$Z^0(\rightarrow \nu\bar{\nu}) + \geq 2$ jets	1.3 ± 0.6
$t\bar{t}$	0.7 ± 0.4
total W/Z/Top bkg	10.5 ± 3.6
QCD	3 ± 2.5
total background	13.4 ± 4.3
Data	11

Table 1: The number of data and expected background events after final cuts.

CDF is a general purpose detector consisting of tracking, vertexing, calorimeter components and a muon system. Events from this analysis were collected using a trigger which required $\cancel{E}_T > 35$ GeV.

We select events with 2 or 3 jets which have $E_T \geq 15$ GeV and $|\eta| \leq 2$. The \cancel{E}_T cut is increased beyond trigger threshold to 40 GeV and we require that \cancel{E}_T is neither parallel nor antiparallel to any of the jets in the event to reduce the contribution from the processes where missing energy comes from jet energy mismeasurement: $\min \Delta\phi(\cancel{E}_T, j) > 45^\circ$, $\Delta\phi(\cancel{E}_T, j_1) < 165^\circ$, and $45^\circ < \Delta\phi(j_1, j_2) < 165^\circ$. The jet indices are ordered by decreasing E_T . We reject events with an identified electron or muon.

To select events with a charm jet, we determine the probability that the ensemble of tracks within a jet is consistent with coming from the primary vertex. We require that at least one jet has a probability of less than 5%.

The dominant source of background for this analysis is W/Z+jets production where the vector boson decays to a lepton (e/ μ) that is not identified or to a τ lepton which decays hadronically. There is also a small contribution from QCD multijet production.

We observe 11 events which is consistent with 13.4 ± 4.3 events from Standard Model processes (see Table 1). We interpret this as an excluded region in the $m_{\tilde{t}} - m_{\tilde{\chi}_1^0}$ parameter space as shown in figure 1. The maximum $m_{\tilde{t}}$ excluded is $122 \text{ GeV}/c^2$ for $m_{\tilde{\chi}_1^0} = 37 \text{ GeV}/c^2$.

References

- [1] see for example, H. Baer *et al*, *Phys. Rev. D* **44**, 725 (1991); H. Baer *et al*, *Phys. Rev. D* **50**, 4517 (1994).
- [2] W. Beenakker *et al.*, *Nucl. Phys.* **B515**, 3 (1998)
- [3] S. Abachi *et al.*, *Phys. Rev. Lett.* **76**, 2222 (1996)
- [4] ALEPH collaboration, CERN EP/98-76 submitted to *Phys. Lett. B*.

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$$\text{BR}(\tilde{t}_1 \rightarrow c + \tilde{\chi}_1^0) = 100\%$$

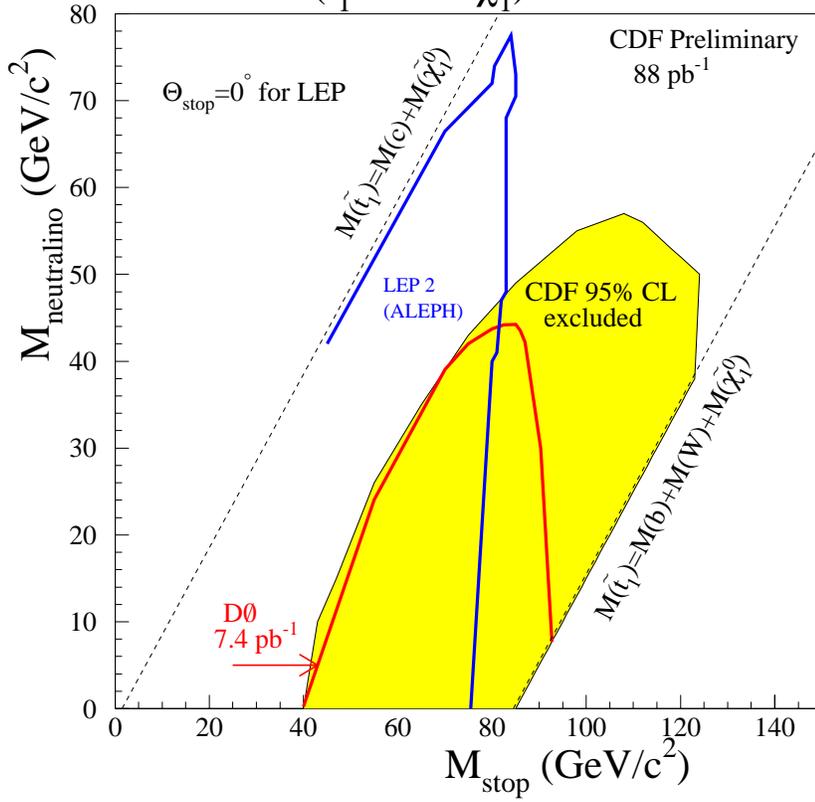


Figure 1: 95% CL limit for $\tilde{t} \rightarrow c\tilde{\chi}_1^0$. Also shown are results from D0 [3] and from ALEPH [4]