



**Fermi National Accelerator Laboratory**

**FERMILAB-Conf-96/372-E**

**CDF**

## **Search for Squarks and Gluinos Using Dileptons**

**J. P. Done**  
For the CDF Collaboration  
*Department of Physics, Texas A&M University*  
*College Station, Texas 77843-4242*

*Fermi National Accelerator Laboratory*  
*P.O. Box 500, Batavia, Illinois 60510*

October 1996

Published Proceedings of the 1996 Annual Divisional Meeting (DPF 96) of the Division of Particles and Fields of the American Physical Society, Minneapolis, Minnesota, August 10-15, 1996.



## **Disclaimer**

*This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness or usefulness of any information, apparatus, product or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process or service by trade name, trademark, manufacturer or otherwise, does not necessarily constitute or imply its endorsement, recommendation or favoring by the United States Government or any agency thereof. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.*

## **Distribution**

*Approved for public release: further dissemination unlimited.*

**SEARCH FOR SQUARKS AND GLUINOS USING DILEPTONS**JAMES P. DONE<sup>a</sup>*Department of Physics, Texas A&M University  
College Station, TX 77843-4242, USA*

Gluinos and squarks can be produced in  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8$  TeV and decay via charginos and neutralinos to final states containing two (or more) leptons. There is no charge correlation between the leptons of charginos from gluino decays as gluinos are Majorana particles. The strategy of this analysis is to search for an excess of events containing two like-sign isolated leptons, missing energy, and jets. The analysis is based on  $81 \text{ pb}^{-1}$  of 1994-5 data recorded at CDF.

**1 Introduction**

Supersymmetry<sup>1</sup> is a promising theory that can allow for grand unification. The Minimal Supersymmetric Standard Model (MSSM)<sup>2</sup> is a supersymmetrized Standard Model (SM) with two Higgs doublets. Since the SUSY mass scale is expected to be of the order of the electroweak scale, data from the CDF<sup>3</sup> experiment can be used to search for SUSY.

**2 SUSY Signature and Data Analysis**

The assumption that R-parity (a multiplicative quantum number which differentiates particles from superparticles) is conserved implies: (1) SUSY objects are pair-produced and (2) a stable LSP exists. Cosmological constraints require that the LSP be the lightest neutralino ( $\tilde{\chi}_1^0$ )<sup>4</sup>. In our analysis, slepton and sneutrino masses are related to gaugino masses<sup>3</sup>. At hadron colliders, the dominant production mechanism for gluinos is  $gg \rightarrow \tilde{g}\tilde{g}$  via  $s$ -channel gluon and  $t$ - and  $u$ -channel gluino exchange. For heavy gluinos, cascade decays into charginos ( $\tilde{\chi}_1^\pm$ ) dominate (e.g.,  $\tilde{g} \rightarrow q\bar{q}'\tilde{\chi}_1^\pm$ )<sup>5</sup>. The charginos decay into leptons along with an LSP which carries away energy. Since the gluino is a Majorana fermion, it has the distinctive property of decaying with equal probability

---

<sup>a</sup>Representing the CDF Collaboration.

into fermions and antifermions. Since there are many SM processes which contribute opposite charge dileptons but very few that contribute those of the same sign, an excellent signature for pair production of gluinos results from events in which both gluinos decay to a chargino of the same sign, yielding like-sign dileptons in the final state<sup>6,7</sup>. We choose isolated dilepton ( $ee$ ,  $e\mu$ , and  $\mu\mu$ ) events from the 1994-5 Tevatron data recorded at CDF<sup>8</sup>. In order to choose leptons that passed our trigger with a high efficiency, we selected primary electrons (muons) with  $|\eta^e| < 1.1$  ( $|\eta^\mu| < 0.6$ ) and  $E_T^e > 11$  GeV ( $p_T^\mu > 11$  GeV/c). We require additional electrons (muons) to be within  $|\eta^e| < 2.4$  ( $|\eta^\mu| < 1.0$ ) and  $E_T^e > 5$  GeV ( $p_T^\mu > 5$  GeV/c) as indicated by SUSY MC studies. We require these events to have two or more jets with  $E_T > 15$  GeV and  $|\eta| < 2.4$  but at least one with  $|\eta| < 1.1$ . Jets and leptons are required to be well separated. A large missing energy cut ( $\cancel{E}_T > 25$  GeV) is made in order to identify LSPs in the gluino decay products. We then make our like-sign dilepton cut to reduce SM dilepton background. After all cuts, we observe two candidates.

### 3 Signal Simulation and Background Studies

We simulate the signal via ISAJET<sup>11</sup> using MSSM parameters:  $\tan\beta = 4$ ,  $\mu = -400$  GeV/ $c^2$ ,  $A_t = -100$  GeV/ $c^2$ , and  $m_A = 500$  GeV/ $c^2$ . We take into account SUSY production via NLO cross-sections<sup>9,10</sup>. The trajectories and decays of these particles are evolved through our detector simulation. The main SM contributions to our analysis are due to  $t\bar{t}$  and Drell-Yan. Cross-sections for these processes were taken from CDF measured values<sup>12,13</sup>. The total number of expected background is  $1.3 \pm 0.6$  events.

### 4 Conclusion

After all cuts, we find two events consistent with expected background; hence there is no excess signal above background. We set limits at the 95 % C. L. on the mass of the gluino folding in a total systematic error of 35%. We set a preliminary limit for  $m_{\tilde{g}} > 230$  GeV/ $c^2$  where  $m_{\tilde{g}} \simeq m_{\tilde{q}}$  and a limit for  $m_{\tilde{g}} > 180$  GeV/ $c^2$  independent of  $m_{\tilde{q}}$  (see Figure 1).

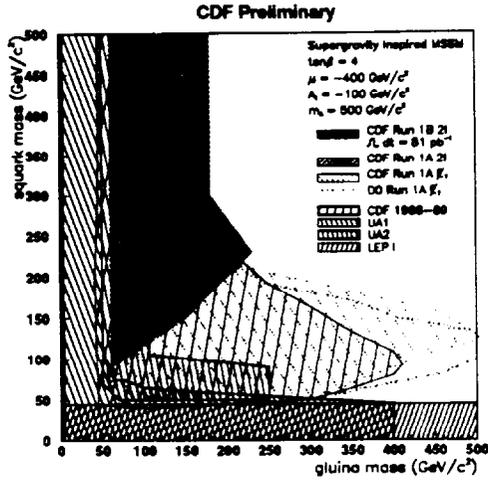


Figure 1: SUSY Mass Limits.

### Acknowledgments

This work is supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Science and Culture of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the A. P. Sloan Foundation; and the Alexander von Humboldt-Stiftung.

### References

1. D.V. Volkov and V.P. Akulov, *Phys. Lett.* B**46**, 109 (1973).
2. H.E. Haber and G.L. Kane, *Phys. Rep.* **117**, 75 (1985).
3. A. H. Chamseddine, R. Arnowitt, and P. Nath, *Phys. Rev. Lett.* **49**, 970 (1982).
4. J. Ellis, *et al.*, *Nucl. Phys. B* **238**, 453 (1984).
5. H. Baer, *et al.*, *Phys. Rev. D* **48**, 2978 (1993).
6. R. Barnett *et al.*, *Phys. Lett.* 315 (1993) 349-54.
7. H. Haber, International Workshop on Supersymmetry and Unification of Fundamental Interactions, SUSY 93, (World Scientific, 1993), pp. 373-

90.

8. F. Abe *et al.*, *Nucl. Instrum. Methods* **A271** (1988) 387.
9. W. Beenakker, *et al.*, *Z. Phys. C* **69**, 163 (1995)
10. W. Beenakker, *et al.*, *Phys. Rev. Lett.* **74** (1995) 2905-8.
11. H. Baer, *et al.*, FSU-HEP-930329 and UH-511-764-93.
12. F. Abe, *et al.*, *Phys. Rev. Lett.* **74**, 2626 (1995).
13. F. Abe, *et al.*, *Phys. Rev. D* **49**, 1 (1994).