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Search for Anomalous $ZZ\gamma$ and $Z\gamma\gamma$ Couplings with $D\bar{O}$

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Abstract

A direct test of the Standard Model by searching for anomalous $ZZ\gamma$ and $Z\gamma\gamma$ couplings is presented. We analyze $p\bar{p} \rightarrow \ell\bar{\ell}\gamma + X$, ($\ell = e, \mu$) events at $\sqrt{s} = 1.8$ TeV with the $D\bar{O}$ detector at the Fermilab Tevatron Collider. A fit to the transverse energy spectrum of the photon in the signal events, based on the data set corresponding to an integrated luminosity of 13.9 pb^{-1} (13.3 pb^{-1}) for the electron (muon) channel, yields the following 95% CL limits on the anomalous CP -conserving $ZZ\gamma$ couplings: $-1.9 < h_{30}^Z < 1.8$ ($h_{40}^Z = 0$), and $-0.5 < h_{40}^Z < 0.5$ ($h_{30}^Z = 0$), for a form-factor scale $\Lambda = 500$ GeV. Limits on the $Z\gamma\gamma$ couplings and CP -violating couplings are also discussed.

Direct measurement of the $ZV\gamma$ ($V = Z, \gamma$) trilinear gauge boson couplings is possible by studying $Z\gamma$ production in $p\bar{p}$ collisions at the Tevatron ($\sqrt{s} = 1.8$ TeV). The most general Lorentz and gauge invariant $ZV\gamma$ vertex is described by four coupling parameters, h_i^V , ($i = 1\dots 4$) [1]. Combinations of the CP -conserving (CP -violating) parameters h_3^V and h_4^V (h_1^V and h_2^V) correspond to the electric (magnetic) dipole and magnetic (electric) quadrupole transition moments of the $ZV\gamma$ vertex. In the Standard Model (SM), all the $ZV\gamma$ couplings vanish at the tree level. Non-zero (i.e. *anomalous*) values of the h_i^V couplings result in an increase of the $Z\gamma$ production cross section and change the kinematic distribution of the final state particles [2]. Partial wave unitarity of the general $f\bar{f} \rightarrow Z\gamma$ process restricts the $ZV\gamma$ couplings uniquely to their vanishing SM values at asymptotically high energies [3]. Therefore, the coupling parameters have to be modified by form-factors $h_i^V = h_{i0}^V / (1 + \hat{s}/\Lambda^2)^n$, where \hat{s} is the square of the invariant mass of the $Z\gamma$ system, Λ is the form-factor scale, and h_{i0}^V are coupling values at the low energy limit ($\hat{s} \approx 0$) [2]. Following Ref. [2] we assume $n = 3$ for $h_{1,3}^V$ and $n = 4$ for $h_{2,4}^V$. Such a choice yields the same asymptotic energy behavior for all the couplings. Unlike $W\gamma$ production where the form-factor effects do not play a crucial role, the Λ -dependent effects cannot be ignored in $Z\gamma$ production at Tevatron energies. This is due to the higher power of \hat{s} in the vertex function, a direct consequence of the additional Bose-Einstein symmetry of the $ZV\gamma$ vertices [2].

This paper describes a measurement [4] of the $ZV\gamma$ couplings using $p\bar{p} \rightarrow \ell\bar{\ell}\gamma + X$ ($\ell = e, \mu$) events. Similar measurements were recently performed by CDF [5] and L3 [6]. The data used was observed with the $D\bar{O}$ detector during the 1992–1993 run, corresponding to an integrated luminosity of $13.9 \pm 1.7 \text{ pb}^{-1}$ ($13.3 \pm 1.6 \text{ pb}^{-1}$) for the electron (muon) data. Since the initial state radiation and radiative Z decays can also contribute to the $\ell\bar{\ell}\gamma$ final state, one would expect a non-zero signal even for the SM (i.e., vanishing) $ZV\gamma$ couplings.

The DØ detector [7] consists of three main systems: (i) uranium-liquid argon sampling calorimeter with coverage in pseudorapidity (η) for $|\eta| \leq 4.4$. and energy resolution of $15\%/\sqrt{E/\text{GeV}}$ for electrons and $50\%/\sqrt{E/\text{GeV}}$ for isolated pions [8]; (ii) central and forward drift chambers with a combined coverage of $|\eta| \leq 3.2$; and (iii) muon system providing coverage for $|\eta| \leq 3.3$ and momentum resolution for central ($|\eta| < 1.0$) muons of $\delta(1/p)/(1/p) = 0.18(p-2)/p \oplus 0.008p$ (p in GeV/ c).

$Z\gamma$ candidates are selected by searching for events containing two isolated electrons (muons) with high transverse energy E_T (transverse momentum p_T), and an isolated photon. The $ee\gamma$ sample is selected from a trigger requiring two isolated EM clusters, each with $E_T \geq 20$ GeV. An electron cluster is required to be within the fiducial region of the calorimeter ($|\eta| \leq 1.1$ in the central calorimeter (CC), or $1.5 \leq |\eta| \leq 2.5$ in the end calorimeters (EC)). Offline electron identification requirements are: (i) the ratio of the EM energy to the total shower energy must be > 0.9 ; (ii) the lateral and longitudinal shower shape must be consistent with an electron shower [9]; (iii) the isolation variable of the cluster (I) must be < 0.1 , where I is defined as $I = [E_{\text{tot}}(0.4) - E_{\text{EM}}(0.2)]/E_{\text{EM}}(0.2)$, $E_{\text{tot}}(0.4)$ is the total shower energy inside a cone defined by $\mathcal{R} = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$, and $E_{\text{EM}}(0.2)$ is the EM energy inside a cone of $\mathcal{R} = 0.2$; (iv) at least one of the two electron clusters must have a matching track in the drift chambers; and (v) $E_T > 25$ GeV for both electrons.

The $\mu\mu\gamma$ sample is selected from a trigger requiring an EM cluster with $E_T > 7$ GeV and a muon track with $p_T > 5$ GeV/ c . A muon track is required to have $|\eta| \leq 1.0$ and must have: (i) hits in the inner drift-tube layer; (ii) a good overall track fit; (iii) bend view impact parameter < 22 cm; (iv) a matching track in the central drift chambers; and (v) minimum energy deposition of 1 GeV in the calorimeter along the muon path. The muon must be isolated from a nearby jet ($\mathcal{R}_{\mu\text{-jet}} > 0.5$). At least one of the muon tracks is required to traverse a minimum length of magnetized iron ($\int Bdl > 1.9$ Tm); it is also required that $p_T^{\mu_1} > 15$ GeV/ c and $p_T^{\mu_2} > 8$ GeV/ c .

The requirements for photon identification are common to both electron and muon samples. We require a photon transverse energy $E_T^\gamma > 10$ GeV and the same quality cuts as those on the electron, except that there must be no track pointing toward the calorimeter cluster. Additionally, we require that the separation between a photon and both leptons be $\Delta\mathcal{R}_{\ell\gamma} > 0.7$. This cut suppresses the contribution of the radiative $Z \rightarrow \ell\ell\gamma$ decays [2]. The above selection criteria yield four $ee\gamma$ and two $\mu\mu\gamma$ candidates. Figure 1 shows the E_T^γ distribution for these events. Three $ee\gamma$ and both $\mu\mu\gamma$ candidates have a three body invariant mass close to that of the Z and low separation between the photon and one of the leptons, consistent with the interpretation of these events as radiative $Z \rightarrow \ell\ell \rightarrow \ell\ell\gamma$ decays. The remaining candidate in electron channel has a dielectron mass compatible with that of the Z and a photon well separated from the leptons, an event topology typical for direct $Z\gamma$ production in which a photon is radiated from one of the interacting partons [2].

The estimated background includes contributions from (i) $Z + \text{jet}(s)$ production where one of the jets fakes a photon or an electron (the latter case corresponds to the $ee\gamma$ signature if additionally one of the electrons from the $Z \rightarrow ee$ decay is not detected in a tracking chamber); (ii) QCD multijet production with jets being misidentified as electrons or photons; (iii) $\tau\tau\gamma$ production followed by decay of each τ to $\ell\bar{\nu}_\ell\nu_\tau$.

We estimate the QCD background from data using the probability, $P(\text{jet} \rightarrow e/\gamma)$, for a jet

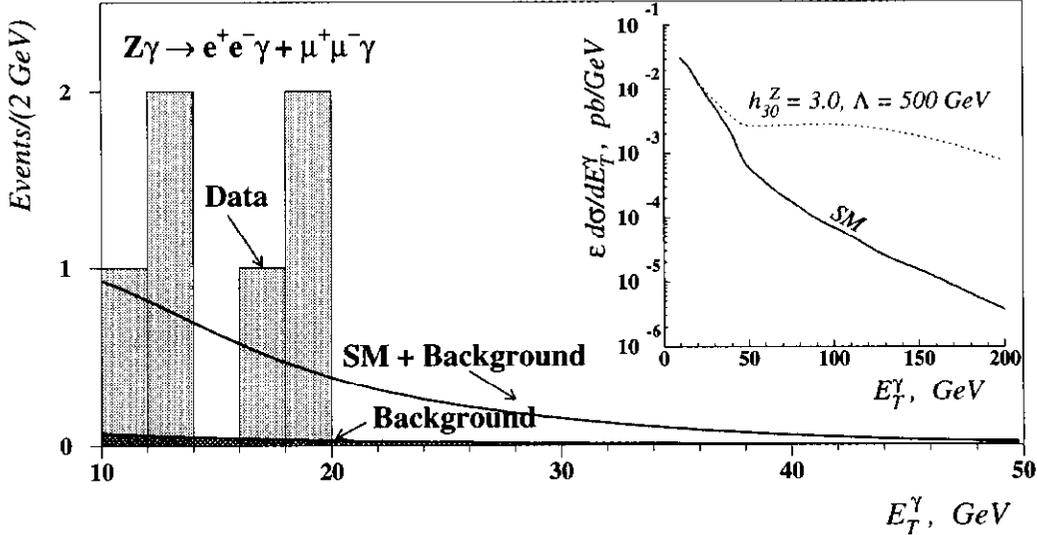


FIG. 1. Photon transverse energy spectrum for $ee\gamma$ and $\mu\mu\gamma$ events. The shadowed bars correspond to data, the hatched curve represents the total background, and the solid line shows the sum of the SM predictions and the background. The insert shows $d\sigma/dE_T^\gamma$ folded with the efficiencies for SM and anomalous ($h_{30}^Z = 3.0$) couplings.

to be misidentified as an electron/photon. This probability is determined by measuring the fraction of non-leading jets in samples of QCD multijet events that pass our electron/photon identification cuts, and takes into account a 0.25 ± 0.25 fraction of direct photon events in the multijet sample [10]. We find the probabilities $P(\text{jet} \rightarrow e/\gamma)$ to be $\sim 10^{-3}$ in the typical E_T ranges for the electrons and photons of between 10 and 50 GeV. We find the background from $Z+\text{jet}(s)$ and QCD multijet events in the electron channel by applying misidentification probabilities to the jet E_T spectrum of the inclusive $ee + \text{jet}(s)$ and $e\gamma + \text{jet}(s)$ data. The background is 0.43 ± 0.06 events. For the muon channel the QCD background is estimated by applying the misidentification probability to the inclusive $\mu\mu + \text{jet}(s)$ spectrum. The estimation of the QCD background from data in the muon case also accounts for cosmic ray background. The combined background from QCD multijet and cosmic ray events in the muon channel is found to be 0.02 ± 0.01 events.

The $\tau\tau\gamma$ background is estimated using the ISAJET Monte Carlo event generator [11] followed by a full simulation of the DØ detector, resulting in 0.004 ± 0.002 events for $ee\gamma$ and 0.03 ± 0.01 events for $\mu\mu\gamma$ channels.

Subtracting the estimated backgrounds from the observed number of events, the signal is $3.57_{-1.91}^{+3.15} \pm 0.06$ for the $ee\gamma$ channel and $1.95_{-1.29}^{+2.62} \pm 0.01$ for the $\mu\mu\gamma$ channel, where the first and dominant uncertainty is due to Poisson statistics, and the second is due to the systematic error of the background estimate.

The acceptance of the DØ detector for the $ee\gamma$ and $\mu\mu\gamma$ final states was studied using the leading order event generator of Baur and Berger [2] followed by a fast detector simulation program which takes into account effects of the electromagnetic and missing transverse energy resolutions, muon momentum resolution, variations in position of the vertex along

the beam-axis, and trigger and offline efficiencies. These efficiencies are estimated using $Z \rightarrow ee$ data for the electron channel. The muon trigger efficiency is estimated from the $e\mu$ data selected using non-muon triggers. The offline efficiency for the muon channel is calculated based on $e\mu$ and $Z \rightarrow \mu\mu$ samples. The trigger efficiency for $ee\gamma$ is 0.98 ± 0.01 while the efficiency of offline dielectron identification is 0.64 ± 0.02 in the CC and 0.56 ± 0.03 in the EC. For the muon channel the trigger efficiency is $0.94_{-0.09}^{+0.06}$, and the offline dimuon identification efficiency is 0.54 ± 0.04 . The photon efficiency depends slightly on E_T^γ due to the calorimeter cluster shape algorithm and the isolation cut, and accounts for loss of the photon due to a random track overlap (which results in misidentification of the photon as an electron) and the photon conversion into an e^+e^- pair before the outermost tracking chamber. The average photon efficiency is 0.53 ± 0.05 . The geometrical acceptance for the electron (muon) channel is 53% (20%) for the SM case and increases slightly for non-zero anomalous couplings. The overall efficiency for the electron (muon) channel for SM couplings is 0.17 ± 0.02 (0.06 ± 0.01). The MRSD--' [12] set of structure functions is used in the calculations. The uncertainties due to the choice of structure function (6%, as determined by variation of the results for different sets) are included in the systematic error of the Monte Carlo calculation. The effect of higher order QCD corrections are accounted for by multiplying the rates by a constant factor $k = 1.34$ [2].

The observed number of events is compared with the SM expectation using the estimated efficiency and acceptance. We expect the signal in the e and μ channels for SM couplings to be: $S_{ee\gamma}^{SM} = 2.7 \pm 0.3$ (sys) ± 0.3 (lum) and $S_{\mu\mu\gamma}^{SM} = 2.2 \pm 0.4$ (sys) ± 0.3 (lum) events, where the first error is due to the uncertainty in the Monte Carlo modelling, and the second reflects the uncertainty in the integrated luminosity calculation. Our observed signal agrees within errors with the SM prediction for both channels.

To set limits on the anomalous coupling parameters, we fit the observed E_T spectrum of the photon (E_T^γ) with the Monte Carlo predictions plus the estimated background, combining the information in the spectrum shape and the event rate. The fit is performed for the $ee\gamma$ and $\mu\mu\gamma$ samples, using a binned likelihood method [13], including constraints to account for our understanding of luminosity and efficiency uncertainties. Because the contribution of the anomalous couplings is concentrated in the high E_T^γ region, the differential distribution $d\sigma/dE_T^\gamma$ is more sensitive to the anomalous couplings than a total cross section (see insert in Fig. 1 and Ref. [2]). To optimize the sensitivity of the experiment for the low statistics, we assume Poisson statistics for each E_T^γ bin and use the maximum likelihood method to fit the experimental data. To exploit the fact that anomalous coupling contributions lead to an excess of events at high transverse energy of the photon, a high- E_T^γ bin, in which we observe no events is explicitly used in the histogram [13]. The results were cross-checked using an unbinned likelihood fit which yields similar results.

Figure 1 shows the observed E_T^γ spectrum with the SM prediction plus the estimated background for the $e + \mu$ combined sample. The 95% confidence level (CL) limit contour for anomalous coupling parameters h_{30}^Z and h_{40}^Z (h_{10}^Z and h_{20}^Z) is shown in Fig. 2. Also shown in this figure are the recent CDF [5] and L3 [6] limits. A form-factor scale $\Lambda = 500$ GeV is used for the calculations of the limits and partial wave unitarity constraints. We obtain the following 95% CL limits for the CP -conserving $ZZ\gamma$ and $Z\gamma\gamma$ couplings (in the assumption that all couplings except one are at the SM values, i.e. zeros):

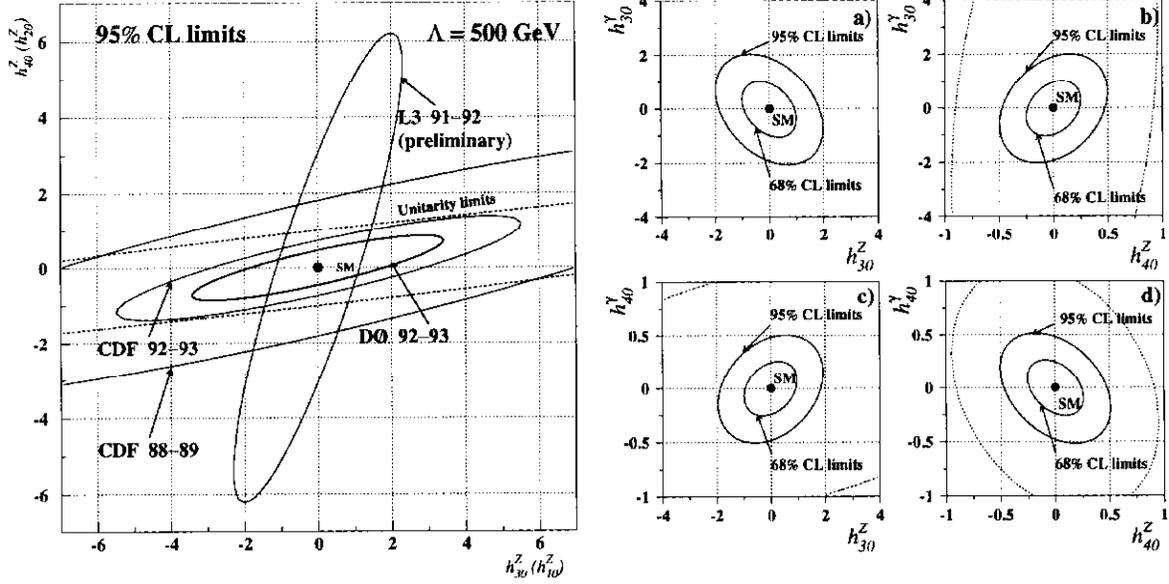


FIG. 2. Comparison of the limits on the correlated CP -conserving (CP -violating) anomalous $ZZ\gamma$ coupling parameters h_{30}^Z and h_{40}^Z (h_{10}^Z and h_{20}^Z). The solid ellipses represent 95% CL exclusion contours for DØ, CDF [5], and L3 [6] experiments. The dashed curve shows limits from partial wave unitarity for $\Lambda = 500$ GeV.

FIG. 3. Limits on the weakly correlated CP -conserving pairs of anomalous $ZV\gamma$ couplings: a) $(h_{30}^Z, h_{30}^\gamma)$, b) $(h_{40}^Z, h_{30}^\gamma)$, c) $(h_{30}^Z, h_{40}^\gamma)$, and d) $(h_{40}^Z, h_{40}^\gamma)$. The solid ellipses represent 68% and 95% CL exclusion contours. Dashed curves show limits from partial wave unitarity for $\Lambda = 500$ GeV.

$$\begin{aligned} -1.9 < h_{30}^Z < 1.8; & -0.5 < h_{40}^Z < 0.5 \\ -1.9 < h_{30}^\gamma < 1.9; & -0.5 < h_{40}^\gamma < 0.5 \end{aligned}$$

The correlated limits for pairs of couplings (h_{30}^V, h_{40}^V) are less stringent due to the strong interference between these couplings:

$$\begin{aligned} -3.3 < h_{30}^Z < 3.3; & -0.9 < h_{40}^Z < 0.9 \\ -3.5 < h_{30}^\gamma < 3.5; & -0.9 < h_{40}^\gamma < 0.9 \end{aligned}$$

Limits on the CP -violating $ZV\gamma$ couplings are numerically the same as those for the CP -conserving couplings. The limits on the h_{20}^Z , h_{40}^Z , and h_{10}^γ couplings are currently the most stringent.

Global limits on the anomalous couplings (i.e., limits independent of the values of other couplings) are close to the correlated limits for (h_{30}^V, h_{40}^V) and (h_{10}^V, h_{20}^V) pairs, since other possible combinations of couplings interfere with each other only at the level of 10%. This is illustrated in Fig. 3, which shows the limits for pairs of couplings of the same CP -parity (couplings with different CP -parity do not interfere with each other).

We also study the form-factor scale dependence of the results. The chosen value of the scale $\Lambda = 500$ GeV is close to the sensitivity limit of this experiment for the $h_{20,40}^V$ couplings: for larger values of the scale partial wave unitarity is violated for certain values of anomalous couplings allowed at 95% CL by this measurement.

The luminosity expected in the current Tevatron run may enable us to improve the limits

on the anomalous couplings by a factor of ~ 2 . Another way to improve the sensitivity toward the couplings is to include the $\nu\bar{\nu}$ decay channel of the Z in the above analysis. This study is in progress.

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[†] See, e.g. [7] for the list of the DØ Collaboration.

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