Prospects of Physics at CDF with the SVX

The CDF Collaboration

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Abstract
During next physics run CDF will strongly enhance its heavy flavor tagging capabilities with the installation of a silicon vertex detector (SVX), that will allow precise measurements of secondary decay vertices in the plane transverse to the beam (impact parameter resolution \( \approx 12 \, \mu m \)). We expect this detector to have a significant impact on b-physics (\( c\tau_B \approx 350 \, \mu m \)) and top search. In the following we will discuss CDF prospects for top search and for CP violation asymmetry measurements in the B-sector.

1 Physics Motivations for Vertex Detection

CDF is in a privileged position to perform top searches [1] and study b-physics [2]. In fact, the TEVATRON provides pp collisions at the highest \( \sqrt{s} \) currently available and this results into high cross sections both for top and beauty production. Several CDF upgrades are being commissioned to increase the efficiency for tagging heavy flavors; primarily the extension of the muon detector coverage and the installation of a silicon vertex detector.

With an integrated luminosity of 25 \( pb^{-1} \) (the goal for the 1992 run), approximately \( 5 \times 10^8 \) b's will be produced. At CDF the total inelastic cross section, \( \sigma_{inel} \), is of the order of 50 mb and the b-production cross section is \( \sigma_b \approx 20 \, \mu b \) for \( \left| y_b \right| \leq 1.0 \) [3], which gives a ratio \( \sigma_b/\sigma_{inel} \approx 4 \times 10^{-4} \). The secondary vertex detection provided by the SVX ([4,5]), in association with the good CDF lepton identification, will be essential for the required physical background rejection. We estimate that this can be done by preserving a good signal detection efficiency for B meson tagging. The SVX will allow detecting B-decay modes (leptonic and non-leptonic) with an impact parameter (D) resolution \( \sigma_D \approx 12 \, \mu m \), that is fully adequate for B-tagging, given typical D values of B sons (\( c\tau_B \approx 350 \, \mu m \)). This will be extremely useful in particular in the search for top, since top quark events contain always two b's (section 3), and for CP violation studies, where a neutral B is selected to decay in an exclusive mode (section 4).

2 SVX Performance

The SVX (fig.1,2) is built to provide high resolution r-\( \phi \) tracking in a pseudorapidity interval \( \left| \eta \right| < 1.9 \).

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Figure 1: Isometric view of one of the two SVX barrels
Tagging the b's with the SVX will be essential, especially for the lepton + jets decays, because of the low rate of dileptons. CDF already tried to tag b's in the lepton + jets sample, by searching for low $P_t$ muons from sequential decays. In spite of its low efficiency, the low-$P_t$ $\mu$ search, combined with results from the other modes, provided an increase of a couple of GeV in the lower limit on $M_{top}$ [1]. By using the SVX for tagging b's we expect to get a larger improvement.

### 3.1 B-tag Efficiency in Top Events

Preliminary studies on B meson tagging efficiency in top decays have been performed using the ISAJET Monte-carlo program as event generator and a simplified SVX followed by detector simulation [10,11]. We assume the $\sigma_D$ curve of fig. 4 and we say that a B has been tagged when there are at least 3 charged prongs within SVX+CTC acceptance and with $|D| > 3 \sigma_D$. Table 1 shows the efficiencies for tagging at least 1 B (B1) and 2 B's (B2) in the event for three values of the top mass. The event $z$-vertex smearing gives $\sim 60\%$ contribution to B1 and B2.

<table>
<thead>
<tr>
<th>$M_{top}$ (GeV/$c^2$)</th>
<th>B1 (%)</th>
<th>B2 (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>100</td>
<td>19</td>
<td>2</td>
</tr>
<tr>
<td>150</td>
<td>28</td>
<td>4</td>
</tr>
<tr>
<td>200</td>
<td>30</td>
<td>5</td>
</tr>
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</table>

More detailed studies are in progress in which we select a sample of $t\bar{t}$ events passing the cuts of the lepton + jets analysis described in [1,8,9] and exploit correlations between impact parameter, $D$, and azimuthal production angle, $\phi$, of tracks coming from a secondary vertex (for a description of $D - \phi$ correlations see [12]). Preliminary results indicate that efficiencies increase with this method.

### 3.2 S/N Improvement with the B-tag

The physical background for the lepton + jets mode is given by direct $b\bar{b}$ production, gluon splitting ($g \rightarrow b\bar{b}$), $\gamma$ conversions and production of W bosons associated to QCD jets. To get rid of the background we require a high $P_t$ isolated lepton, large missing transverse energy ($E_t$) and we cut on the associated jet $E_t$'s. After this selection W+jets is the only non-negligible background left. This is no surprise, since the top decay is

$$t\bar{t} \rightarrow W + bW - b,$$

followed by

$$W \rightarrow l + E_t, \quad W \rightarrow jet jet,$$

which is to be compared with

$$W + jets \rightarrow l + E_t + jets,$$

which basically leads to the same topology. However, a clear difference is given by the presence of the two b's in $t\bar{t}$ events. On the basis of present top analysis results, we think to be able to achieve a S/N slightly better than 1 for $M_{top} = 150$ GeV/$c^2$, without SVX information.

The upper limit to the S/N improvement achievable with the B-tag is set by the b-content in W+jets events. A conservative estimate of the upper limit can be done assuming that: a) $\sim 3\%$ of gluon jets ($P_t > 10$ GeV) produce b pairs [13], b) there are in average 4 gluons in W+jets events, that is $\sim 12\%$ of W+jets events have b pairs; c) the percentage of tagging errors (fake secondary vertices) is negligible compared to the fraction of real b's in W+jets events (confirmed by preliminary studies); d) the tagging efficiency for b's from gluons is the same as for b's from top decays. With these assumptions the upper limit to the S/N improvement is $100/12 \approx 8$. In
practice, we expect the efficiency for tagging b's from gluons to be lower because they have a lower $P_t$ than in top decays and therefore the improvement to be bigger. In this estimate we neglected other sources of physical background (like $c\bar{c}$), because we think that they are not large compared to the b-content in W+jets events; with next run's data we will have a good sample to measure them quantitatively. For example, the efficiency for tagging D mesons should be significantly lower than for B mesons, because of the lower D lifetime and because of the lower $P_t$ and multiplicity of D decays.

4 CP Violation in the B-sector

CDF has an extensive b-physics program, which may ultimately allow attacking the most ambitious goal, i.e. observing CP violation in the B-sector. We are presently evaluating our chances, by studying the asymmetry in the decay rates of $B^0$ and $B^0$ to the CP eigenstate $J/\Psi K_s$:

$$A = \frac{\Gamma(B^0 \rightarrow J/\Psi K_s) - \Gamma(B^0 \rightarrow J/\Psi K_s)}{\text{sum}}.$$  

The final state can be easily tagged by reconstructing $J/\Psi \rightarrow \mu^+\mu^-$ and $K_s \rightarrow \pi^+\pi^-$. The flavor of the neutral B meson decaying to the CP eigenstate can obtained by looking at the charge of the lepton coming from the semileptonic decay of the other b-hadron.

An important parameter in this study is the integrated luminosity, $L$, that is needed in order to see CP violation with the CDF detector. The measured CP violating asymmetry, $A$, is proportional to $\sin^2\beta$, the raw CP violating asymmetry from the CKM matrix [14,15]. The error on $\sin^2\beta$ can be expressed as

$$\delta(\sin^2\beta) = \left( L d_{\text{phys}} d_{\text{rec}} d_{\text{tag}}\right)^{-\frac{1}{2}},$$  

where $d_{\text{phys}}$ depends only on the physics process, $d_{\text{rec}}$ is a measure of how well $J/\Psi K_s$ events are reconstructed and $d_{\text{tag}}$ is a measure of how well the flavor of the other b-hadron is tagged. Theoretical predictions [16] give $\sin^2\beta = 0.34$ with an uncertainty $0.1 < \sin^2\beta < 1.0$.

Extrapolating from present data and assuming significant detector upgrades for future runs, we estimate [15] that CDF could reach $d_{\text{phys}} \times d_{\text{rec}} \approx 1.7$ pb and $d_{\text{tag}} \approx 0.01$ (or $0.005 < d_{\text{tag}} < 0.05$). Fig. 5 shows that in order to measure the raw CP violation with an error $\delta(\sin^2\beta) \leq \sin^2\beta$ we need $L > 500 \text{ pb}^{-1}$.

This integrated luminosity is not out of CDF's reach, if Fermilab accelerator will be upgraded with the construction of the new Main Injector (1995 or after). As far as the detector is concerned, in order to measure CP violation in the $J/\Psi K_s$ channel it is essential to have high reconstruction/tagging efficiency and good background rejection. The figures used in the previous estimates should be granted by the secondary vertex detection that will be provided by the SVX.

Figure 5: $L$ required to observe CP asymmetries

5 Conclusions

CDF expects a large physics outcome from 1992 run. In particular the enhanced tracking resolution provided by
the SVX will increase our chances to discover the top quark. A measurement of CP violation in the B-sector may not be out of reach in the long term future, but requires an integrated luminosity in excess of 500 pb⁻¹.

References

[1] 'Heavy Flavor Physics from Colliders (UA1, UA2, CDF)', L. Galtieri for the CDF Collaboration, these proc.


[10] 'Silicon Vertex Simulation, Tracking, Physics Estimates', F. Bedeschi et al, CDF internal note # 601

[11] FNAL proposal P-775: to upgrade CDF with a silicon vertex detector in order to tag long-lived heavy flavors.


[15] 'b-Physics at CDF', edited by P. Tipton, CDF internal note # 1482