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ABSTRACT

We report preliminary results on the production of electron-positron pairs in the mass range 2.5 to 4 GeV in 400 GeV p-Be interactions. Production cross-sections of the  $\psi(3100)$  near  $x = 0$  as a function of  $p_t$ ,  $x$  and the decay angle are presented and implications of this new data for single direct leptons are discussed. A 90% confidence level upper limit on  $\sigma_B(\psi(3700))/\sigma_B(\psi(3100))$  of .02 is measured.

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We have reported the observation of high mass ( $\geq 4$  GeV)  $e^+e^-$  pairs produced in 400 GeV proton-beryllium collisions.<sup>1</sup> In the same experiment we have observed a strong  $J/\psi(3100)$  signal. In this letter we present the details of  $J/\psi$  production cross-sections as a function of  $p_T$ ,  $x$  and the decay angle.

The two-arm magnetic spectrometer used to make these measurements is described in detail in Ref. 1. The techniques described in Refs. 1 and 2 of comparing the energy deposited in a lead-glass array with the measured momentum and the use of the longitudinal shower development in the lead-glass array permit rejection against hadrons of  $4 \times 10^{-4}$  per arm while retaining an efficiency for electrons of  $.80 \pm .08$ .

The observed mass spectrum  $d\sigma/dm$  for electron-positron pairs is presented in Fig. 1. Outside the  $\psi$  peak the data is consistent with accidental background. To obtain cross-sections we assumed that the production at all masses is similar to that observed here for the  $\psi(3100)$  and specified below in Eq. 3. Cross-sections reported here have been corrected for trigger and reconstruction efficiencies which lead to a combined efficiency before electron cuts of  $.8 \pm .1$  for electron pair events. Corrections have also been made for electron cut inefficiency, target absorption of the proton beam, and for events lost due to bremsstrahlung of the electrons (estimated to be  $.075 \pm .025$  of the events at 3 GeV).

The  $\psi(3100)$  resonance is observed with a width of 40 MeV (rms) which is in good agreement with the calculated resolution of

the apparatus. We measure a mass value of  $3.096 \pm .030$  GeV. The cross sections for  $\psi$  production times its branching ratio to  $e^+e^-$  are shown in Figures 2a and b plotted versus  $x = p_{\parallel} \text{ cm}/p_{\text{max}}$  and  $p_t$ . No dip or significant variation in the invariant cross-section is observed in our narrow range of x-acceptance near  $x = 0$ . The average invariant cross-section on beryllium for  $x$  near 0 is found to be well described by a quadratic exponential in  $p_t$ :

$$E \frac{d^3\sigma}{dp^3} \cdot B = (1.7 \pm .4) \times 10^{-32} e^{-ap_t^2} \text{ cm}^2/\text{GeV}^2/\text{Be Nucleus} \quad (1)$$

with  $a = 1.1 \pm .35 \text{ GeV}^{-2}$ ,

where B is the branching ratio of  $\psi(3100)$  to  $e^+e^-$  and the  $p_t$  range measured is  $0 < p_t < 2.0$  GeV. An equally valid description of our data is given by:

$$E \frac{d^3\sigma}{dp^3} \cdot B = (2.5 \pm .6) \times 10^{-32} e^{-bp_t} \text{ cm}^2/\text{GeV}^2/\text{Be Nucleus} \quad (2)$$

with  $b = 1.6 \pm .35 \text{ GeV}^{-1}$ .

The errors on the cross-sections presented here are dominated by the systematic uncertainty in the absolute normalization. The decay angle distribution in the  $\psi$  rest system (helicity frame) is shown in Fig. 2c to favor  $1 + \cos^2\theta^*$  (confidence level .20).<sup>3</sup> The approximately 5% of background<sup>1</sup> under the  $\psi$  has not been subtracted in Fig. 2, but this does not significantly change any of these distributions.

The differential cross sections are measured to be:

$$\left(\frac{d\sigma}{dx}\right)_{x=0} \cdot B = (2.6 \pm .7) \times 10^{-31} \text{ cm}^2/\text{Be Nucleus} \quad -.06 \leq x \leq .08$$

$$\left(\frac{d\sigma}{dy}\right)_{y=0} \cdot B = (6.4 \pm 1.6) \times 10^{-32} \text{ cm}^2/\text{Be Nucleus} \quad -.28 \leq y \leq .32$$

where  $y$  is the center of mass rapidity. To obtain the total cross section, a distribution in  $x$  away from the  $x = 0$  region must be assumed. Using

$$E \frac{d^3\sigma}{dp^3} \propto (1 - |x|)^{4.3} \cdot e^{-1.6pt} \quad -1 \leq x \leq 1 \quad (3)$$

which characterizes data from this and other experiments at high  $x^4,^5$  we obtain a total cross section times branching ratio on beryllium of

$$\sigma \cdot B = (1.00 \pm .25) \times 10^{-31} \text{ cm}^2/\text{Be Nucleus}$$

Assuming a linear  $A$ -dependence, the total cross section per nucleon is

$$\sigma \cdot B = (1.1 \pm .3) \times 10^{-32} \text{ cm}^2/\text{nucleon.}$$

No evidence is seen for production of the  $\psi(3700)$ . Taking 5 events as the 90% confidence level leads to the upper limit

$$\frac{\sigma \cdot B(\psi(3700))}{\sigma \cdot B(\psi(3100))} < .02.$$

Table I shows a comparison of the results of this experiment with other measurements of  $\psi(3100)$  production in hadron collisions.<sup>5</sup> Although some additional assumptions, indicated in the table, are necessary to compare these experiments, the agreement seems to be quite good with experiments with  $\sqrt{s} > 20$  GeV. Below that value, a substantial s-dependence is apparent: there is a factor of 10 between Fermilab and Serpukhov energies and a factor of 100 between Fermilab and BNL energies.

Finally, we calculate the yield of single direct leptons resulting from the  $\psi(3100)$  production cross-section measured here. Figure 3 shows a comparison of this yield, using Eq. (1), with the dashed curve representing  $\frac{\pi^+ + \pi^-}{2} \times 10^{-4}$  which is consistent with the direct lepton yield measured previously by this and other groups.<sup>2,7</sup> We conclude that for single lepton  $p_t$  between 1.6 GeV and 2.3 GeV, a substantial fraction, if not all, of the single lepton cross-section comes from the  $\psi(3100)$ . However, the drop in  $e/\pi$  below 1.5 GeV which results from this source is not observed.<sup>8</sup> Above 2.3 GeV, the contribution depends on an extrapolation of the measured  $\psi$  yield. The quadratic form (Eq. 1) suggests that even at higher  $p_t$ , the  $\psi(3100)$  makes a major contribution, and the linear form (Eq. 2) gives more direct leptons at high  $p_t$  than are observed.<sup>9</sup>

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- <sup>2</sup>J. A. Appel, et al., Phys. Rev. Lett. 33, 722 (1974).
- <sup>3</sup>An isotropic distribution with confidence level .05 cannot be ruled out and would lead to an increase of 20% for all cross sections reported here.
- <sup>4</sup>An x-distribution given by  $E \frac{d^3\sigma}{dp^3} \propto e^{-6|x|} e^{-1.6p_t}$  for  $|x| > .3$  and flat in x for  $|x| < .3$  also seems to describe the data in Ref. 5. Assuming such a distribution would decrease the total cross sections quoted here by 20%. The differential cross sections, however, are not significantly affected.
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- <sup>8</sup>F. W. Büsser, Direct Leptons at the I.S.R. II, Proc. of the International Conf. on Lepton and Photon Interactions, Stanford (1975).

<sup>9</sup>For the  $\psi$  to be the dominant source of direct leptons at all  $s$ , the  $s$ -dependence of  $\psi$  production must be similar to that of direct leptons. We note that direct leptons behave (see Ref. 10) with  $s$  much as do pions, i.e.,  $\sim e^{-25 p_t/\sqrt{s}}$ . Our comparison with  $\psi$  yield is made near  $\sqrt{s} = 27$  GeV. The  $s$ -dependence of  $\psi$ 's is as yet poorly known, but is not inconsistent with the above form.

<sup>10</sup>F. W. Büsser, Direct Leptons at the I.S.R. I, Loc. Cit.

TABLE I

$\psi(3100)$ : COMPARISON WITH OTHER EXPERIMENTS

EXPERIMENT	$\sqrt{s}$	CROSS-SECTION	ADDITIONAL ASSUMPTIONS NECESSARY
This Expt. Busser et al.	27.4 mostly 52	$\frac{d\sigma}{dy} \Big _{y=0} \cdot B$ 7.1 $\pm$ 1.8 x 10 <sup>-33</sup> cm <sup>2</sup> /Nucleon 7.5 $\pm$ 2.5 x 10 <sup>-33</sup> cm <sup>2</sup> /Proton	A <sup>1</sup> -----
This Expt. Knapp et al.	27.4 21.7	$\sigma \cdot B$ 1.1 $\pm$ .3 x 10 <sup>-32</sup> cm <sup>2</sup> /Nucleon ~1.0 x 10 <sup>-32</sup> cm <sup>2</sup> /Nucleon	A <sup>1</sup> , (1 -  x ) <sup>4.3</sup>
This Expt. Anderson et al.	27.4 16.8	$\frac{d\sigma}{dy} \Big _{y=0} \cdot B$ 6.4 $\pm$ 1.6 x 10 <sup>-32</sup> cm <sup>2</sup> /Be Nucleus 5.9 x 10 <sup>-32</sup> cm <sup>2</sup> /Be Nucleus	----- corrected for 1 + COS <sup>2</sup> $\theta^*$ decay distribution
This Expt. Antipov et al.	27.4 12.0	$\sigma \cdot B$ 1.0 $\pm$ .25 x 10 <sup>-31</sup> cm <sup>2</sup> /Be Nucleus 9.5 $\pm$ 2.5 x 10 <sup>-33</sup> cm <sup>2</sup> /Be Nucleus	(1 -  x ) <sup>4.3</sup> -----
This Expt. Aubert et al.	27.4 5.1	$\sigma \cdot B$ 1.1 x 10 <sup>-32</sup> cm <sup>2</sup> /Nucleon ~10 <sup>-34</sup> cm <sup>2</sup> /Nucleon	A <sup>1</sup> , (1 -  x ) <sup>4.3</sup> e <sup>-6pt</sup> , P <sub>  </sub> independent

FIGURE CAPTIONS

- Fig. 1: Electron-Positron Mass Spectrum:  $d\sigma/dm$  per nucleon versus the effective mass. A linear A-dependence is assumed.
- Fig. 2:  $\psi(3100)$  Cross-Sections: a) Invariant cross-section versus  $x \equiv P_{\parallel}^{\text{cm}}/P_{\text{max}}^{\text{cm}}$  (integrated over  $p_t^2$ ).  
b) Invariant cross-section versus  $p_t$  near  $x = 0$ .  
c) Decay angle distribution in the helicity frame of the  $\psi$ . (Plotted errors are statistical only).
- Fig. 3: Single direct leptons coming from the  $\psi(3100)$ . The dotted and dot-dashed curves are two possible extrapolations beyond  $p_t$  values for which this cross-section is determined in this experiment.

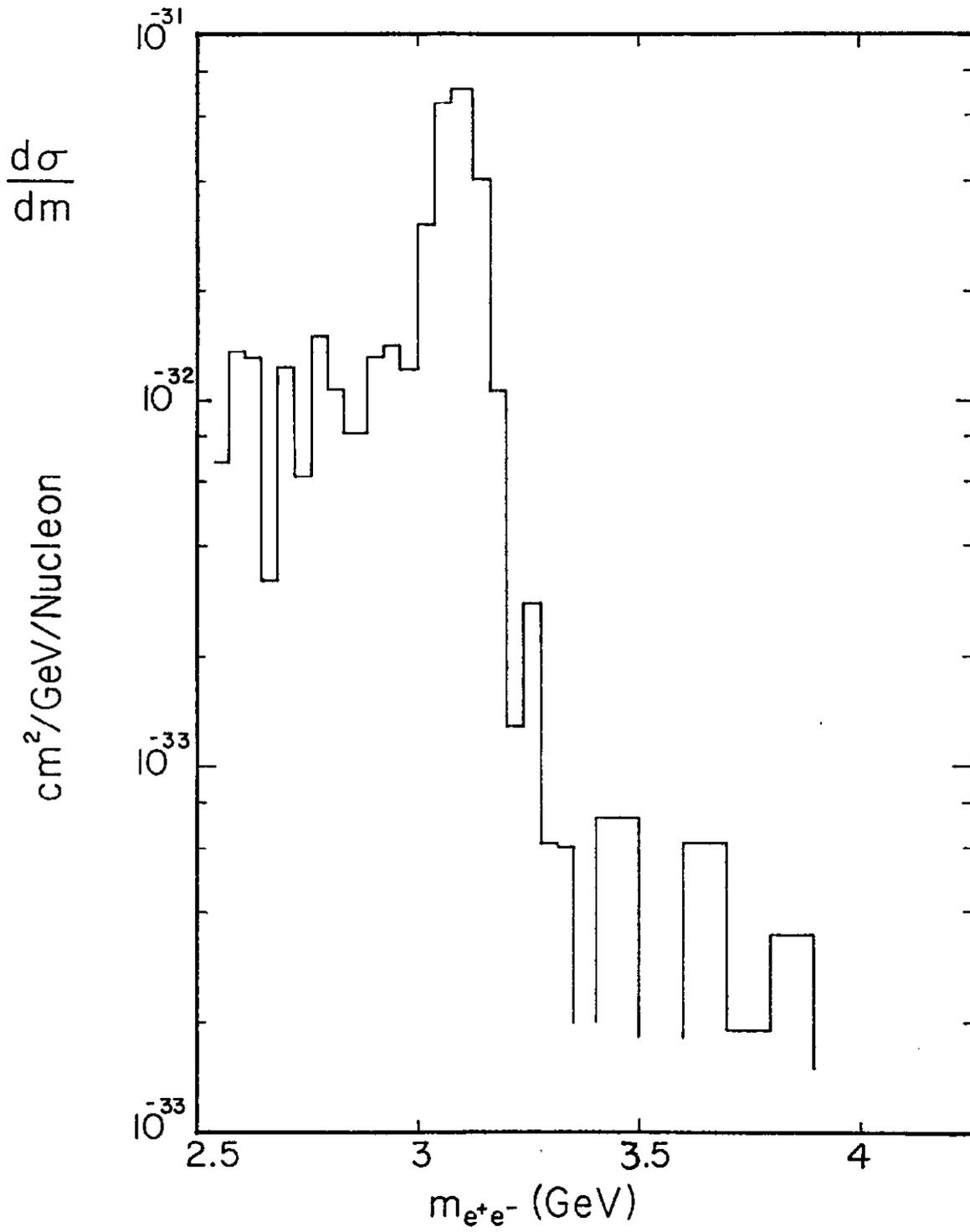


Fig. 1

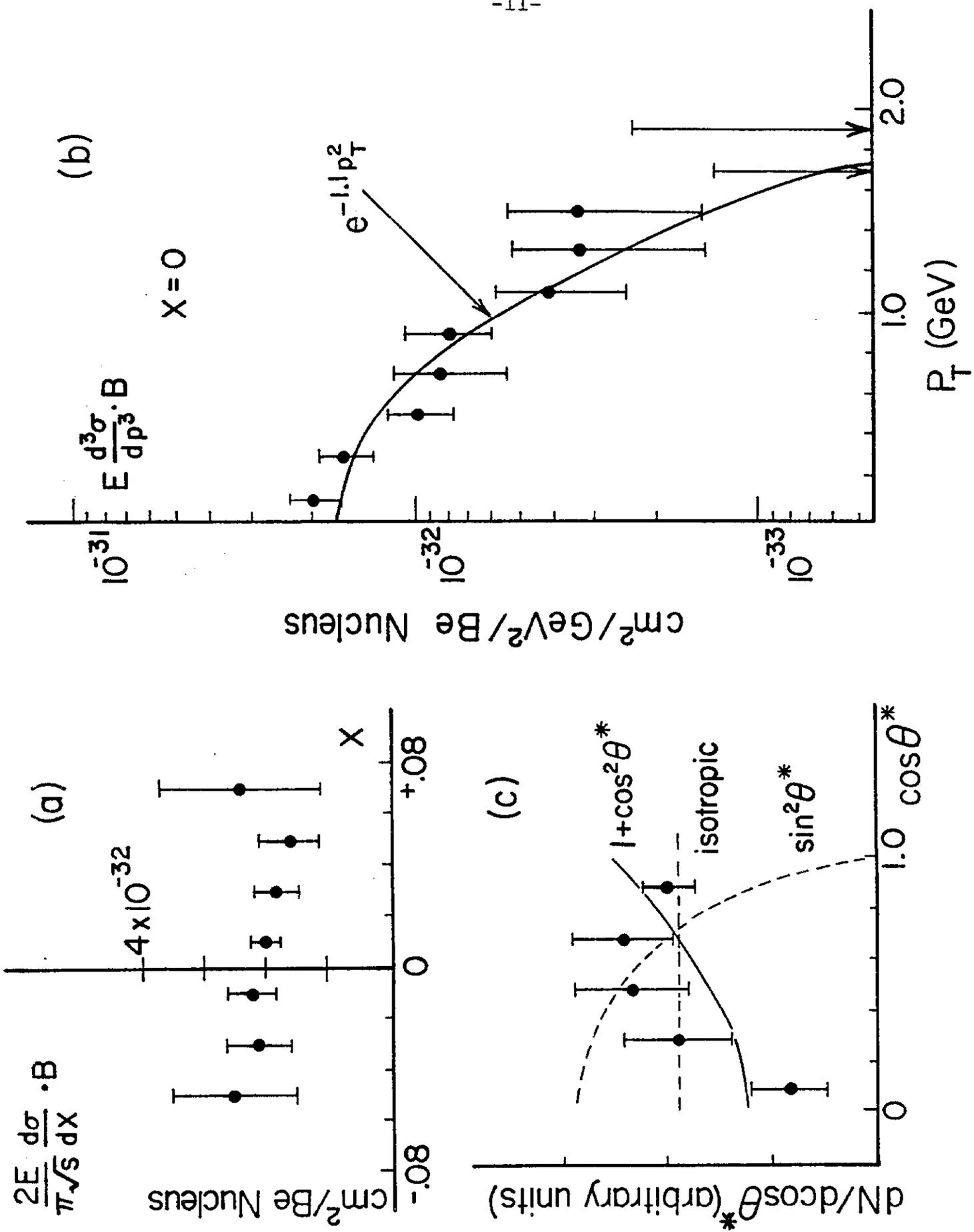


Fig. 2

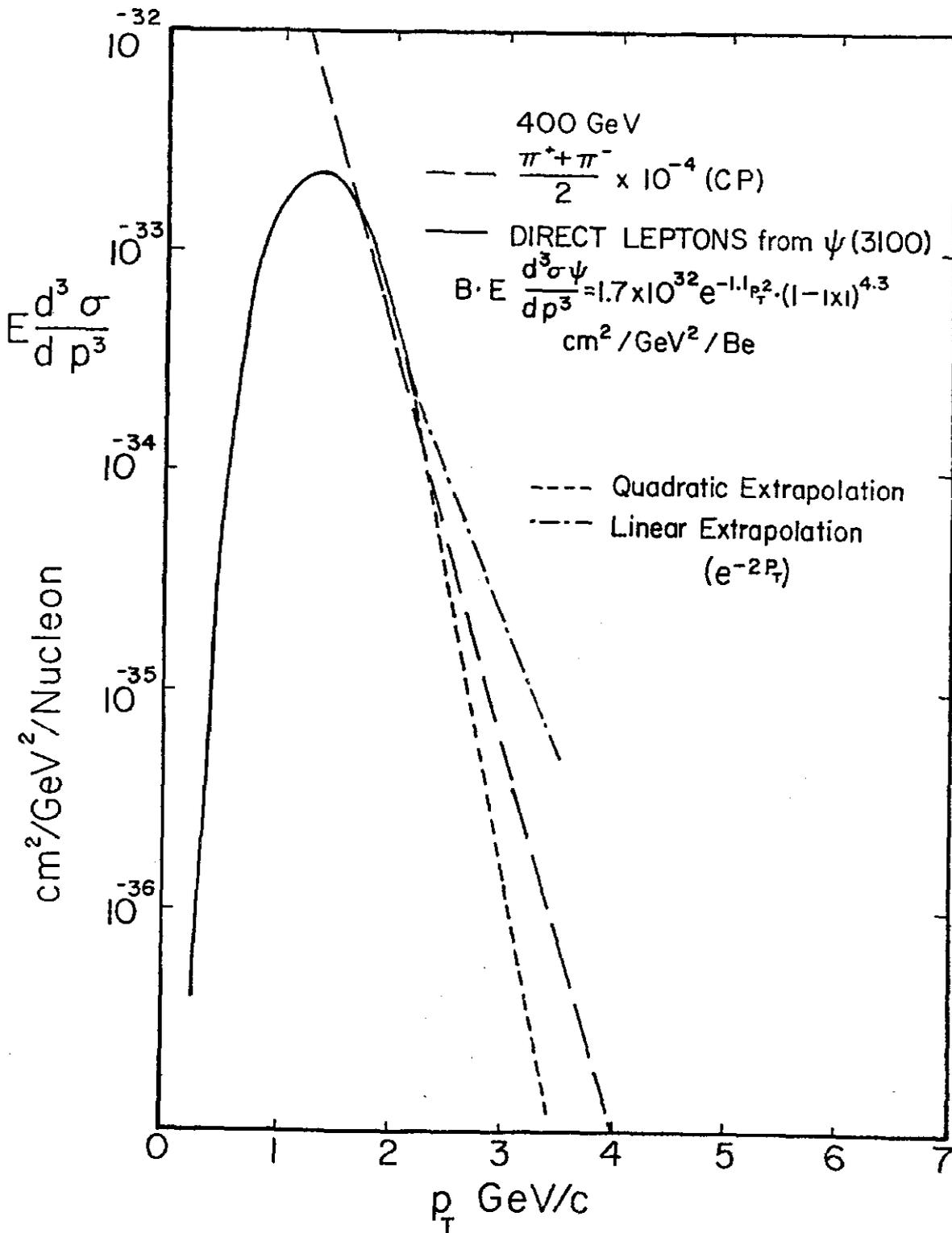


Fig. 3